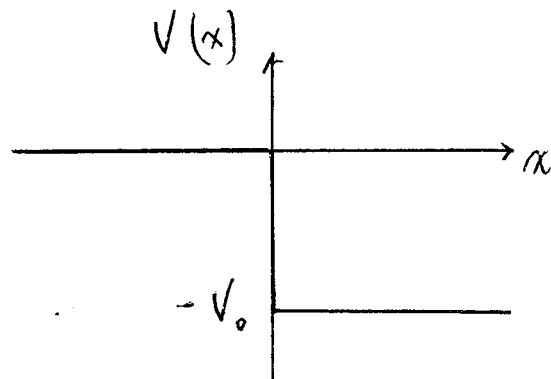


Physics 471: Homework 13 Solutions

1.

$$V(x) = \begin{cases} 0 & \text{for } x < 0 \\ -V_0 & \text{for } x > 0 \end{cases}$$



a) Classical physics:

Energy is conserved: $\frac{1}{2} m v_i^2 = \frac{1}{2} m v_f^2 - V_0$

$$v_f^2 = v_i^2 + \frac{2V_0}{m}$$

$$\underline{v_f = \sqrt{v_i^2 + \frac{2V_0}{m}}}$$

b) Energy eigenvalue eqn is $\hat{H} |E\rangle = E |E\rangle$

In position representation: $\left[-\frac{\hbar^2}{2m} \frac{d^2}{dx^2} + V(x) \right] \psi_E(x) = E \psi_E(x)$

$x < 0$: $\frac{d^2 \psi_E}{dx^2} = -\frac{2mE}{\hbar^2} \psi_E \equiv -k_1^2 \psi_E \quad k_1^2 = \frac{2mE}{\hbar^2}$

$$\psi_E^I(x) = A e^{ik_1 x} + B e^{-ik_1 x}$$

$x > 0$: $\frac{d^2 \psi_E}{dx^2} = -\frac{2m(E+V_0)}{\hbar^2} \psi_E \equiv -k_2^2 \psi_E \quad k_2^2 = \frac{2m(E+V_0)}{\hbar^2}$

$$\psi_E^II(x) = C e^{ik_2 x} \quad (D=0)$$

Boundary conditions at $x=0$:

ψ_E is continuous: $A + B = C$

$\frac{d\psi_E}{dx}$ is continuous: $ik_1(A - B) = ik_2 C$

c) Substitute C from (1) into (2): $ik_1(A - B) = ik_2(A + B)$

$$A(k_1 - k_2) = B(k_1 + k_2)$$

$$\frac{B}{A} = \frac{k_1 - k_2}{k_1 + k_2} = \frac{\sqrt{E} - \sqrt{E + V_0}}{\sqrt{E} + \sqrt{E + V_0}}$$

$$R = \left| \frac{B}{A} \right|^2 = \left(\frac{k_1 - k_2}{k_1 + k_2} \right)^2 = \left(\frac{\sqrt{E} - \sqrt{E + V_0}}{\sqrt{E} + \sqrt{E + V_0}} \right)^2$$

Classically, we expect $R=0$, so QM is different.

If E is small, then k_1 is small but k_2 is not small,

so $k_1 \ll k_2$ and $R \rightarrow 1$ total reflection!

If V_0 is small, then $k_2 \approx k_1$ and $R \rightarrow 0$.

If $E \gg V_0$, then $k_2 \approx k_1$ and $R \rightarrow 0$.

So the peculiarity of QM only appears for small E .

d) For plotting R vs. E/V_0 , re-write R :

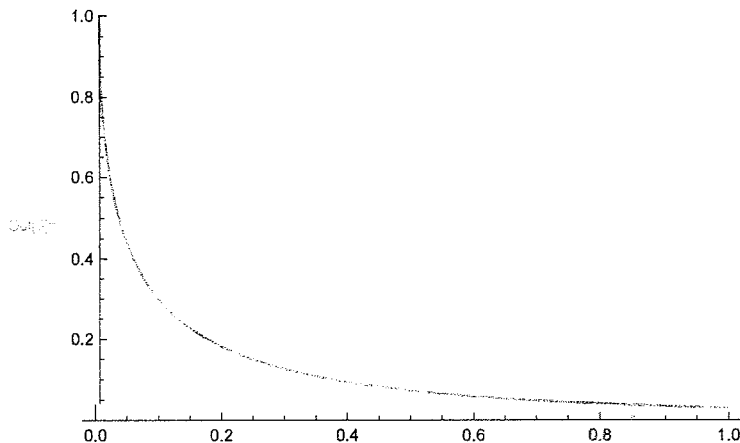
$$R = \left(\frac{\sqrt{E} - \sqrt{E+V_0}}{\sqrt{E} + \sqrt{E+V_0}} \right)^2 = \left(\frac{\sqrt{\frac{E}{V_0}} - \sqrt{\frac{E}{V_0} + 1}}{\sqrt{\frac{E}{V_0}} + \sqrt{\frac{E}{V_0} + 1}} \right)^2 = \left(\frac{\sqrt{x} - \sqrt{x+1}}{\sqrt{x} + \sqrt{x+1}} \right)^2$$

where $x = \frac{E}{V_0}$

In Mathematica I wrote R as

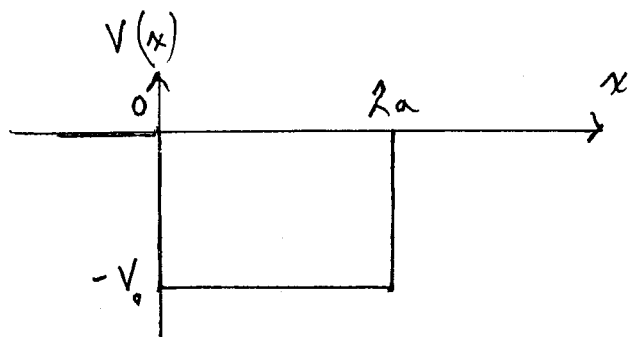
$$R = \left(\frac{1 - \sqrt{1 + \frac{1}{x}}}{1 + \sqrt{1 + \frac{1}{x}}} \right)^2 \quad \text{which is equivalent.}$$

`Plot[((1 - Sqrt[1 + 1/x]) / (1 + Sqrt[1 + 1/x]))^2, {x, 0, 1}, PlotRange -> {0, 1}]`



All my comments are in part (c).

$$2. \quad V(x) = \begin{cases} 0 & \text{for } x < 0 \\ -V_0 & \text{for } 0 \leq x \leq 2a \\ 0 & \text{for } x > 2a \end{cases}$$



a)

This is physically equivalent to the well potential discussed in Mc Intyre section 6.4. So we can use all the formulas from there.

I will re-write equation (6.94) because I don't like the \sin^2 term in the denominator of the denominator!

$$R = \frac{\frac{(k_2^2 - k_1^2)^2}{4k_1^2 k_2^2} \sin^2(2k_2 a)}{1 + \frac{(k_2^2 - k_1^2)^2}{4k_1^2 k_2^2} \sin^2(2k_2 a)}$$

$$k_1^2 = \frac{2mE}{\hbar^2}$$

$$k_2^2 = \frac{2m(E + V_0)}{\hbar^2}$$

$$\frac{k_2^2 - k_1^2}{4k_1^2 k_2^2} = \frac{V_0^2}{4E(E + V_0)}$$

$$R = \frac{\frac{V_0^2}{4E(E + V_0)} \sin^2\left(\frac{2a}{\hbar} \sqrt{2m(E + V_0)}\right)}{1 + \frac{V_0^2}{4E(E + V_0)} \sin^2\left(\frac{2a}{\hbar} \sqrt{2m(E + V_0)}\right)}$$

$$\text{With } E = \frac{V_0}{3}, \quad \frac{V_0^2}{4E(E + V_0)} = \frac{V_0^2}{\frac{4}{3}V_0 \cdot \frac{4}{3}V_0} = \frac{9}{16}$$

$$\frac{2a}{\hbar} \sqrt{2m(E + V_0)} = \frac{2a}{\hbar} \sqrt{2m \cdot \frac{4}{3}V_0} = \frac{2a}{\hbar} \sqrt{\frac{8}{3}mV_0}$$

2. a) continued

Case i) $\frac{2mV_0}{\hbar^2} a^2 \ll 1$ sin term is very small $\Rightarrow R \rightarrow 0$

Case ii) $\frac{2mV_0}{\hbar^2} a^2 = \frac{\pi^2}{48}$

Argument of sin is

$$\frac{2a}{\hbar^2} \sqrt{\frac{8}{3} m V_0} = \left(\frac{32}{3} \frac{m V_0 a^2}{\hbar^2} \right)^{\frac{1}{2}}$$

$$\sin\left(\frac{\pi}{3}\right) = \frac{\sqrt{3}}{2}$$

$$\sin^2\left(\frac{\pi}{3}\right) = \frac{3}{4}$$

$$= \left(\frac{16}{3} \cdot \frac{\pi^2}{48} \right)^{\frac{1}{2}} = \frac{\pi}{3}$$

$$R = \frac{\frac{9}{16} \cdot \frac{3}{4}}{1 + \frac{9}{16} \cdot \frac{3}{4}} = \frac{\frac{27}{64}}{1 + \frac{27}{64}} = \frac{\frac{27}{64}}{\frac{91}{64}} = \frac{27}{91} = 0.297 \approx \underline{\underline{0.30}}$$

Case iii) $\frac{2mV_0}{\hbar^2} a^2 = \frac{3\pi^2}{16}$

Argument of sin (from above) is

$$\left(\frac{16}{3} \cdot \frac{3\pi^2}{16} \right)^{\frac{1}{2}} = \pi$$

$$\sin(\pi) = 0$$

$$\Rightarrow R = 0, \text{ so } T = 1$$

This corresponds to a "transmission resonance" discussed in the book and on problem 4 of this homework.

b) V_0 fixed, $E \rightarrow 0$; $R \rightarrow 1$ similar to single step in problem 1

$E \gg V_0$: $R \rightarrow 0$ from numerator

at very high energies, the small ^{step} is hardly noticeable.

3. Tunneling

a) Start with McIntyre (6.93) for a potential well:

$$T = \frac{1}{1 + \frac{V_0^2}{4E(E+V_0)} \sin^2\left(\frac{2a}{\hbar} \sqrt{2m(E+V_0)}\right)}$$

Change well to bump: $V_0 \rightarrow -V_0$

$$T = \frac{1}{1 + \frac{V_0^2}{4E(E-V_0)} \sin^2\left(\frac{2a}{\hbar} \sqrt{2m(E-V_0)}\right)}$$

Now let $E < V_0$, so $E - V_0 < 0$

$$\sqrt{2m(E-V_0)} = i \sqrt{2m(V_0-E)}$$

We need

$$\begin{aligned} \sin(ix) &= \frac{e^{i(ix)} - e^{-i(ix)}}{2i} = \frac{e^{-x} - e^x}{2i} = i \cdot \frac{e^x - e^{-x}}{2} \\ &= i \sinh x \end{aligned}$$

$$T = \frac{1}{1 + \frac{V_0^2}{-4E(V_0-E)} \cdot i^2 \sinh^2\left(\frac{2a}{\hbar} \sqrt{2m(V_0-E)}\right)}$$

$$T = \frac{1}{1 + \frac{V_0^2}{4E(V_0-E)} \sinh^2\left(\frac{2a}{\hbar} \sqrt{2m(V_0-E)}\right)}$$

matches
(6.105)

$$R = 1 - T = \frac{\frac{V_0^2}{4E(V_0 - E)} \sinh^2\left(\frac{2a}{\hbar} \sqrt{2m(V_0 - E)}\right)}{1 + \frac{V_0^2}{4E(V_0 - E)} \sinh^2\left(\frac{2a}{\hbar} \sqrt{2m(V_0 - E)}\right)}$$

matches my version of (6.106)

b) Electron $m = 9.1 \cdot 10^{-31} \text{ kg}$

$$E = 0.5 \text{ eV}, \quad V_0 = 1.0 \text{ eV}, \quad V_0 - E = 0.5 \text{ eV}$$

$$\frac{V_0^2}{4E(V_0 - E)} = \frac{(1.0 \text{ eV})^2}{4 \cdot 0.5 \text{ eV} \cdot 0.5 \text{ eV}} = 1$$

Argument of \sinh : Use SI units $2a = 5 \text{ \AA} = 5 \cdot 10^{-10} \text{ m}$

$$V_0 - E = 0.5 \text{ eV} \cdot 1.6 \cdot 10^{-19} \frac{\text{J}}{\text{eV}} = 0.8 \cdot 10^{-19} \text{ J}$$

$$\sqrt{2m(V_0 - E)} = \sqrt{2 \cdot 9.1 \cdot 10^{-31} \text{ kg} \cdot 0.8 \cdot 10^{-19} \text{ J}} = 3.8 \cdot 10^{-25} \text{ kg} \cdot \text{m/s}$$

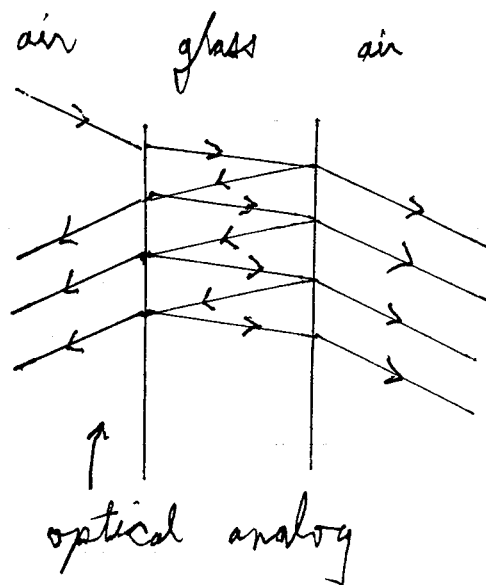
$$\frac{2a}{\hbar} \sqrt{2m(V_0 - E)} = \frac{5 \cdot 10^{-10} \text{ m}}{1.05 \cdot 10^{-34} \text{ J}\cdot\text{s}} \cdot 3.8 \cdot 10^{-25} \text{ kg} \cdot \text{m/s} = 1.82$$

$$\sinh^2(1.82) = 9.0$$

$$T = \frac{1}{1 + \sinh^2(1.82)} = \frac{1}{1 + 9.0} = \underline{\underline{0.10}} = 10 \%$$

(There may be some round-off error in my answer.)

4. Interference approach to understanding transmission resonances.



- a) In problem 16), I derived

$$r_1 = \frac{B}{A} = \frac{k_1 - k_2}{k_1 + k_2} \quad \text{for a single interface}$$

↑
optical analog

where $k_1 = \sqrt{\frac{2mE}{\hbar^2}}$ is the wave vector in free space

and $k_2 = \sqrt{\frac{2m(E+V_0)}{\hbar^2}}$ is the wave vector inside the "medium", i.e. the potential well between $x=0$ and $x=2a$.

$$C = A + B = A + A \frac{k_1 - k_2}{k_1 + k_2} = A \left[1 + \frac{k_1 - k_2}{k_1 + k_2} \right] = A \frac{2k_1}{k_1 + k_2}$$

$$t_1 = \frac{C}{A} = \frac{2k_1}{k_1 + k_2}$$

When we approach the same interface from the "medium" side, reverse the roles of k_1 and k_2 :

$$r_2 = \frac{k_2 - k_1}{k_2 + k_1} \quad t_2 = \frac{2k_2}{k_2 + k_1}$$

4. b) From the picture, sum over trajectories:

$$t = t_1 t_2 + t_1 r_2^2 t_2 + t_1 r_2^4 t_2 + \dots$$

$$= t_1 t_2 (1 + r_2^2 + r_2^4 + \dots)$$

$$c) \quad t = t_1 t_2 \cdot \frac{1}{1 - r_2^2} = \left(\frac{2k_1}{k_1 + k_2} \right) \left(\frac{2k_2}{k_2 + k_1} \right) \cdot \frac{1}{1 - \left(\frac{k_2 - k_1}{k_2 + k_1} \right)^2}$$

$$t = \frac{4k_1 k_2}{(k_1 + k_2)^2} \cdot \frac{(k_2 + k_1)^2}{[(k_2 + k_1)^2 - (k_2 - k_1)^2]} = \frac{4k_1 k_2}{4k_1 k_2} = 1$$

$$d) \quad r = r_1 + t_1 r_2 t_2 + t_1 r_2^3 t_2 + \dots$$

$$= r_1 + t_1 r_2 t_2 (1 + r_2^2 + r_2^4 + \dots)$$

$$= r_1 + t_1 r_2 t_2 \frac{1}{1 - r_2^2}$$

$$= \frac{k_1 - k_2}{k_1 + k_2} + \frac{2k_1 (k_2 - k_1) \cdot 2k_2}{(k_1 + k_2)^3} \cdot \frac{(k_1 + k_2)^2}{4k_1 k_2}$$

$$= \frac{k_1 - k_2}{k_1 + k_2} + \frac{k_2 - k_1}{k_1 + k_2} = 0$$

from part (c)

Extra credit:

$$\begin{aligned}
 t &= t_1 t_2 \left(1 + r_2^2 e^{4ik_2 a} + r_2^4 e^{8ik_2 a} + \dots \right) \\
 &= \frac{t_1 t_2}{1 - r_2^2 e^{4ik_2 a}} = \frac{4k_1 k_2}{(k_1 + k_2)^2} \cdot \frac{1}{1 - \left(\frac{k_2 - k_1}{k_2 + k_1} \right)^2 e^{4ik_2 a}} \\
 &= \frac{4k_1 k_2}{(k_1 + k_2)^2 - (k_2 - k_1)^2 e^{4ik_2 a}}
 \end{aligned}$$

Multiply numerator & denominator by $e^{-2ik_2 a}$:

$$\begin{aligned}
 t &= \frac{4k_1 k_2 e^{-2ik_2 a}}{(k_1^2 + 2k_1 k_2 + k_2^2) e^{-2ik_2 a} - (k_1^2 - 2k_1 k_2 + k_2^2) e^{2ik_2 a}} \\
 &= \frac{4k_1 k_2 e^{-2ik_2 a}}{(k_1^2 + k_2^2) \left(e^{-2ik_2 a} - e^{2ik_2 a} \right) + 2k_1 k_2 \left(e^{-2ik_2 a} + e^{2ik_2 a} \right)} \\
 &= \frac{4k_1 k_2 e^{-2ik_2 a}}{(k_1^2 + k_2^2) \cdot (-2i \sin(2k_2 a)) + 2k_1 k_2 \cdot 2 \cos(2k_2 a)} \\
 t &= \frac{e^{-2ik_2 a}}{\cos(2k_2 a) - i \frac{k_1^2 + k_2^2}{2k_1 k_2} \sin(2k_2 a)}
 \end{aligned}$$

Agrees w/ (6.91)
up to an
overall phase
factor!