

PHY 471: Homework 9 Solutions

1. $|\psi(t=0)\rangle = |E_2\rangle$

$$E_n = \frac{n^2 \pi^2 \hbar^2}{2mL^2}$$

a) $\langle \hat{H} \rangle = E_2 = \frac{2\pi^2 \hbar^2}{mL^2}$ at all times.

$$|\psi(t)\rangle = e^{-iE_2 t/\hbar} |E_2\rangle$$

An energy measurement will always yield E_2 .

b) $\psi(x,t) = e^{-iE_2 t/\hbar} \psi_{E_2}(x) = e^{-iE_2 t/\hbar} \sqrt{\frac{2}{L}} \sin\left(\frac{2\pi x}{L}\right)$
 McIntyre (5.66)

$$|\psi(x,t)|^2 = \frac{2}{L} \sin^2\left(\frac{2\pi x}{L}\right)$$

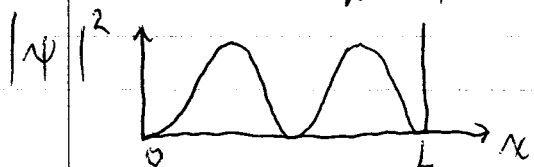
$$\langle x(t) \rangle = \int_0^L x |\psi(x,t)|^2 dx = \frac{2}{L} \int_0^L x \sin^2\left(\frac{2\pi x}{L}\right) dx$$

Method 1: Change to dimensionless variables $u = \frac{x}{L}$, $x = Lu$, $dx = L du$

$$\langle x(t) \rangle = \frac{2}{L} \int_0^1 (Lu) \sin^2(2\pi u) (L du) = 2L \int_0^1 u \sin^2(2\pi u) du$$

Mathematica evaluates dimensionless integral = $\frac{1}{4}$

$$\langle x(t) \rangle = \frac{2}{L} \cdot \frac{1}{4} = \underline{\underline{\frac{L}{2}}}$$



This makes sense because $|\psi|^2$ is symmetric about the middle of the well.

1.6) Method 2: Evaluate the integral yourself:

$$\begin{aligned} \int_0^1 u \sin^2(2\pi u) du & \quad \sin^2 x = \frac{1}{2}(1 - \cos 2x) \\ &= \frac{1}{2} \int_0^1 u (1 - \cos 4\pi u) du = \frac{1}{2} \left[\frac{1}{2} u^2 \Big|_0^1 - \int_0^1 u \cos(4\pi u) du \right] \\ \text{Integrate by parts: } dv &= \cos(4\pi u) du \quad v = \frac{1}{4\pi} \sin(4\pi u) \\ &= \frac{1}{4} - \frac{1}{2} \left[\underbrace{\frac{u}{4\pi} \sin(4\pi u)}_0 \Big|_0^1 - \frac{1}{4\pi} \underbrace{\int_0^1 \sin(4\pi u) du}_0 \right] \\ &= \frac{1}{4} \end{aligned}$$

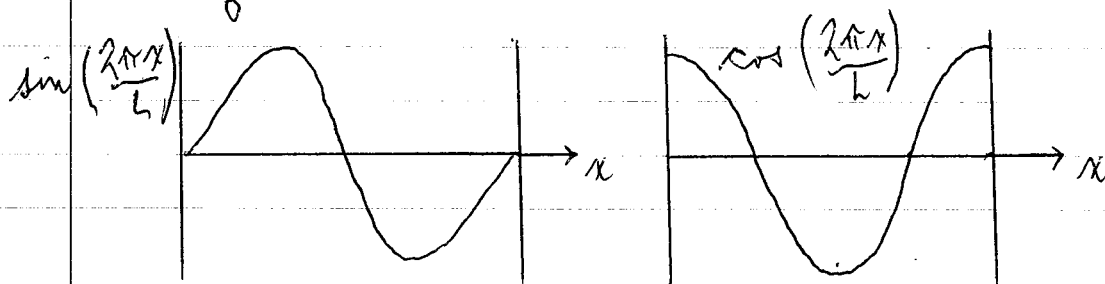
Method 3: Use Symbolab or Mathematica to give you an algebraic expression that contains L .

Then you don't have to change to dimensionless variables in the first place. (But it's a good trick to know!)

c) Calculate $\langle \hat{p}(t) \rangle$

Method 1: Brute force

$$\begin{aligned} \langle \hat{p}(t) \rangle &= \langle \psi(t) | \hat{p} | \psi(t) \rangle = \int_{-\infty}^{\infty} \psi^*(x, t) \left(-i\hbar \frac{d}{dx} \right) \psi(x, t) dx \\ &= \frac{2}{L} \int_0^L e^{+iE_2 t/\hbar} \sin\left(\frac{2\pi x}{L}\right) \cdot (-i\hbar) \cdot \frac{2\pi}{L} \cos\left(\frac{2\pi x}{L}\right) e^{-iE_2 t/\hbar} dx \end{aligned}$$



1.c) $\sin \frac{2\pi x}{L}$ is odd around $x = \frac{L}{2}$
 $\cos \frac{2\pi x}{L}$ is even around $x = \frac{L}{2}$ } Their product is odd,
 so the integral is 0.

If you aren't confident about that result, then continue

$$\begin{aligned}\langle \hat{p}(t) \rangle &= -i\hbar \frac{4\pi}{L^2} \int_0^L \sin\left(\frac{2\pi x}{L}\right) \cos\left(\frac{2\pi x}{L}\right) dx \\ &= -i\hbar \frac{4\pi}{L^2} \cdot \frac{L}{2\pi} \sin^2\left(\frac{2\pi x}{L}\right) \Big|_0^L = \underline{\underline{0}}\end{aligned}$$

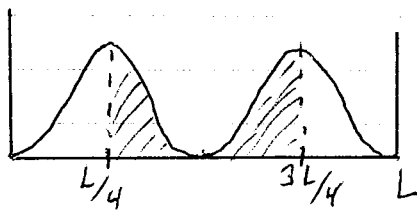
Method 2: Ehrenfest's Theorem: $\langle \hat{p}(t) \rangle = m \frac{d}{dt} \langle \hat{x}(t) \rangle$

But $\langle \hat{x}(t) \rangle = \frac{L}{2}$ is independent of time, so $\langle \hat{p} \rangle = 0$.

This makes sense because the system is in a stationary state!

$$\begin{aligned}d) P\left[\frac{L}{4} < x < \frac{3L}{4}\right] &= \int_{L/4}^{3L/4} |\psi(x, t)|^2 dx \\ &= \frac{2}{L} \int_{L/4}^{3L/4} \sin^2\left(\frac{2\pi x}{L}\right) dx = \frac{1}{L} \int_{L/4}^{3L/4} \left(1 - \cos\left(\frac{4\pi x}{L}\right)\right) dx \\ &= \frac{1}{L} \cdot \left[x - \frac{L}{4\pi} \sin\left(\frac{4\pi x}{L}\right) \right]_{L/4}^{3L/4} \\ &= \frac{1}{L} \left[\left(\frac{3L}{4} - \frac{L}{4}\right) - \frac{L}{4\pi} (\sin 3\pi - \sin \pi) \right] = \frac{1}{L} \cdot \frac{L}{2} = \underline{\underline{\frac{1}{2}}}\end{aligned}$$

We could have guessed this from $|\psi|^2$:



1. e) We already have $\langle \hat{x} \rangle$ and $\langle \hat{p} \rangle$, so we need $\langle \hat{x}^2 \rangle$ and $\langle \hat{p}^2 \rangle$

$$\langle x^2 \rangle = \int_0^L x^2 |\psi(x,t)|^2 dx = \frac{2}{L} \int_0^L x^2 \sin^2\left(\frac{2\pi x}{L}\right) dx$$

Again I'll use $u = \frac{x}{L}$:

$$\langle x^2 \rangle = 2L^2 \int_0^1 u^2 \sin^2(2\pi u) du = L^2 \left(\frac{1}{3} - \frac{1}{8\pi^2} \right)$$

For the integral, Symbolab gives $\frac{32\pi^3 - 12\pi}{192\pi^3} = \frac{1}{6} - \frac{1}{16\pi^2}$

$$\Delta x = \sqrt{\langle x^2 \rangle - \langle x \rangle^2} = \sqrt{L^2 \left(\frac{1}{3} - \frac{1}{8\pi^2} \right) - \frac{L^2}{4}} = \sqrt{L^2 \left(\frac{1}{12} - \frac{1}{8\pi^2} \right)}$$

$$= L \sqrt{0.0707...} \approx \underline{\underline{0.266 L}}$$

$$\langle p^2 \rangle = \int_{-\infty}^{\infty} \psi^* \left(-\hbar^2 \frac{d^2}{dx^2} \psi \right) dx = -\hbar^2 \frac{2}{L} \int_0^L \sin\left(\frac{2\pi x}{L}\right) \left(\frac{2\pi}{L}\right)^2 \left(-\sin\frac{2\pi x}{L}\right) dx$$

$$= \frac{8\pi^2}{L^3} \hbar^2 \underbrace{\int_0^L \sin^2\left(\frac{2\pi x}{L}\right) dx}_{= L/2} = 4\pi^2 \frac{\hbar^2}{L^2}$$

$$\Delta p = \sqrt{\langle p^2 \rangle - \langle p \rangle^2} = \sqrt{4\pi^2 \frac{\hbar^2}{L^2}} = \underline{\underline{\frac{2\pi\hbar}{L}}}$$

$$\Delta x \Delta p = 0.266 L \cdot \frac{2\pi\hbar}{L} = \underline{\underline{1.67 \hbar}}$$

Uncertainty Principle says $\Delta x \Delta p \geq \frac{1}{2} \hbar$, so this is consistent.

$$2. \quad |\psi(t=0)\rangle = A \left(|E_1\rangle + 3i|E_2\rangle \right)$$

$$a) \quad 1 = \langle \psi | \psi \rangle = |A|^2 (1+9) = 10 |A|^2 \rightarrow A = \frac{1}{\sqrt{10}}$$

Measure energy: results are E_1 or E_2

$$P(E_1) = \frac{1}{10} \quad P(E_2) = \frac{9}{10}$$

$$\begin{aligned} \langle \hat{H} \rangle &= \frac{1}{10} E_1 + \frac{9}{10} E_2 = \frac{1}{10} \cdot \frac{\pi^2 \hbar^2}{2mL^2} + \frac{9}{10} \cdot \frac{4\pi^2 \hbar^2}{2mL^2} \\ &= (0.1 + 3.6) \frac{\pi^2 \hbar^2}{2mL^2} = \underline{\underline{3.7 \cdot \frac{\pi^2 \hbar^2}{2mL^2}}} = \underline{\underline{3.7 E_1}} \end{aligned}$$

$$b) \quad |\psi(t)\rangle = \frac{1}{\sqrt{10}} e^{-iE_1 t/\hbar} |E_1\rangle + \frac{3i}{\sqrt{10}} e^{-iE_2 t/\hbar} |E_2\rangle$$

Energy measurements are the same as in part (a)!

$$c) \quad \psi(x,t) = \frac{1}{\sqrt{10}} e^{-iE_1 t/\hbar} \cdot \sqrt{\frac{2}{L}} \sin \frac{\pi x}{L} + \frac{3i}{\sqrt{10}} e^{-iE_2 t/\hbar} \sqrt{\frac{2}{L}} \sin \frac{2\pi x}{L}$$

$$\begin{aligned} |\psi(x,t)|^2 &= \frac{1}{10} \cdot \frac{2}{L} \left| e^{-iE_1 t/\hbar} \sin \frac{\pi x}{L} + 3i e^{-iE_2 t/\hbar} \sin \frac{2\pi x}{L} \right|^2 \\ &= \frac{1}{5L} \left(e^{iE_1 t/\hbar} \sin \frac{\pi x}{L} - 3i e^{iE_2 t/\hbar} \sin \frac{2\pi x}{L} \right) \left(e^{-iE_1 t/\hbar} \sin \frac{\pi x}{L} + 3i e^{-iE_2 t/\hbar} \sin \frac{2\pi x}{L} \right) \\ &= \frac{1}{5L} \left[\sin^2 \frac{\pi x}{L} + 9 \sin^2 \frac{2\pi x}{L} + 3i \sin \frac{\pi x}{L} \sin \frac{2\pi x}{L} \left(e^{i(E_1-E_2)t/\hbar} - e^{-i(E_2-E_1)t/\hbar} \right) \right] \\ E_2 - E_1 &= \frac{4\pi^2 \hbar^2}{2mL^2} - \frac{\pi^2 \hbar^2}{2mL^2} = \frac{3\pi^2 \hbar^2}{2mL^2} = \hbar \omega \end{aligned}$$

$$2i \sin \frac{(E_1 - E_2)t}{\hbar}$$

$$2. c) |\psi(x,t)|^2 = \frac{1}{5L} \left[\sin^2 \frac{\pi x}{L} + 9 \sin^2 \frac{2\pi x}{L} + 6 \sin \frac{\pi x}{L} \sin \frac{2\pi x}{L} \sin \omega t \right]$$

$$t_0 = 0 \rightarrow \sin \omega t = 0, \quad t_1 = \frac{\pi}{2\omega} \rightarrow \sin \omega t = \sin \frac{\pi}{2} = 1$$

$$t_2 = \frac{\pi}{\omega} \rightarrow \sin \omega t = \sin \pi = 0, \quad t_3 = \frac{3\pi}{2\omega} \rightarrow \sin \omega t = \sin \frac{3\pi}{2} = -1$$

$$t_4 = \frac{2\pi}{\omega} \rightarrow \sin \omega t = 0$$

\Rightarrow The plots for $t = t_0, t_2$, and t_4 will be identical.

See next page for plots made in Mathematica. The probability density sloshes back and forth in the well.

$$d) \langle x(t) \rangle = \int_0^L x |\psi(x,t)|^2 dx$$

Let's separate the integral into a time-independent part (first two terms) plus a time-dependent part (last term).

$$\frac{1}{5L} \int_0^L dx \left(x \sin^2 \frac{\pi x}{L} + 9 x \sin^2 \frac{2\pi x}{L} \right) = \frac{1}{5L} \left(\frac{L^2}{4} + 9 \frac{L^2}{4} \right) = \frac{L}{2}$$

$$\frac{1}{5L} \cdot \int_0^L 6 x \sin \frac{\pi x}{L} \sin \frac{2\pi x}{L} dx = \frac{6}{5L} \left(\frac{-8L^2}{9\pi^2} \right) = -\frac{16}{15\pi^2} L$$

$$\langle x(t) \rangle = \frac{L}{2} - \frac{16}{15\pi^2} L \sin(\omega t)$$

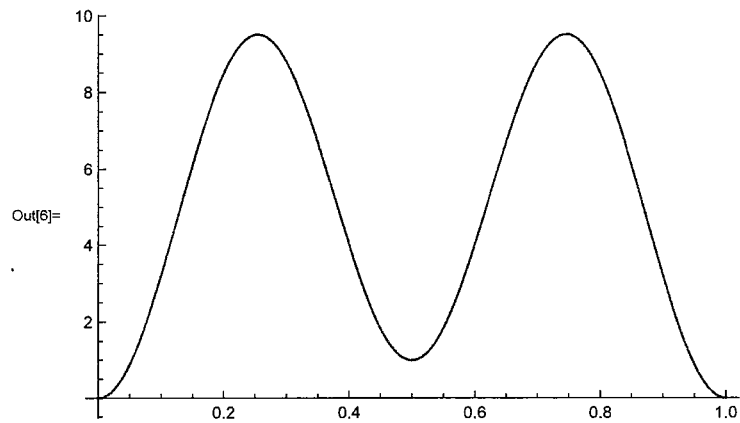
\nwarrow from Mathematica

$$x_{\max} = L \left(\frac{1}{2} + \frac{16}{15\pi^2} \right) = 0.608 L$$

$$x_{\min} = L \left(\frac{1}{2} - \frac{16}{15\pi^2} \right) = 0.392 L$$

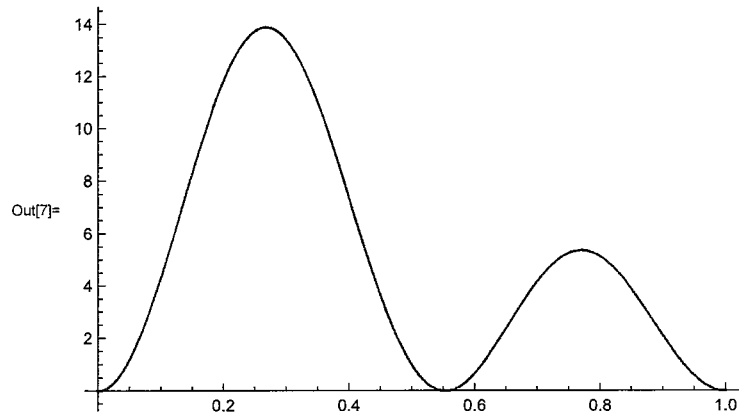
In problem 1,
 $\langle x(t) \rangle = \frac{L}{2}$
 always.

In[6]:= `Plot[Sin[π *x]^2 + 9*Sin[2* π *x]^2, {x, 0, 1}]`



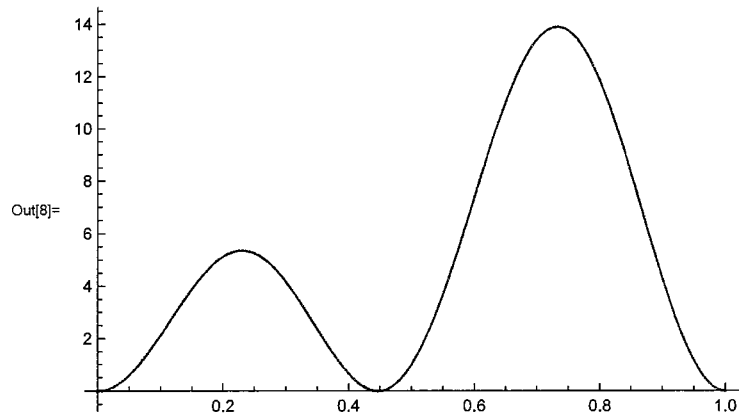
$$t = t_0, t_2, \text{ or } t_4$$

In[7]:= `Plot[Sin[π *x]^2 + 9*Sin[2* π *x]^2 + 6*Sin[π *x]*Sin[2* π *x], {x, 0, 1}]`



$$t = t_1$$

In[8]:= `Plot[Sin[π *x]^2 + 9*Sin[2* π *x]^2 - 6*Sin[π *x]*Sin[2* π *x], {x, 0, 1}]`



$$t = t_3$$

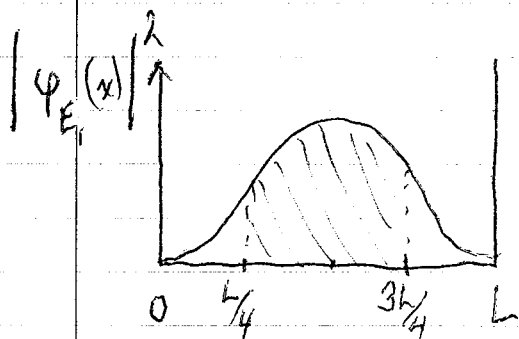
2. a) Measure energy, obtain result E_1 . The system is now in the state

$$|E_1\rangle \rightarrow \varphi_{E_1}(x) = \sqrt{\frac{2}{L}} \sin \frac{\pi x}{L}$$

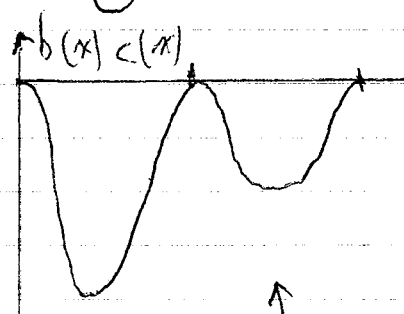
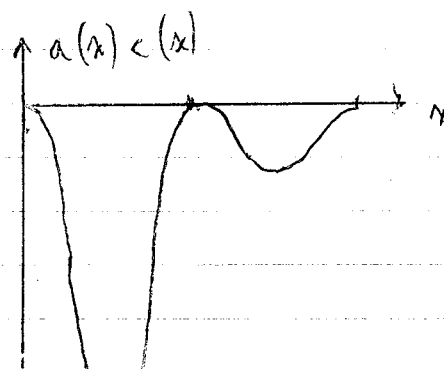
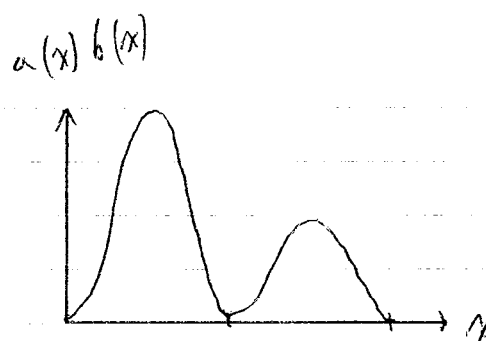
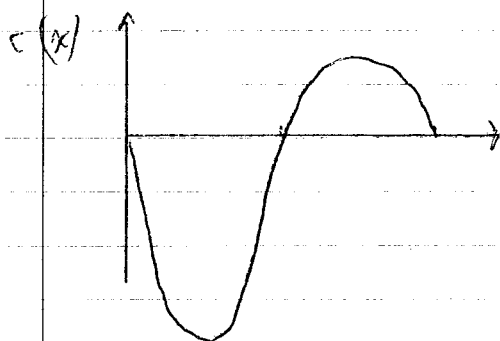
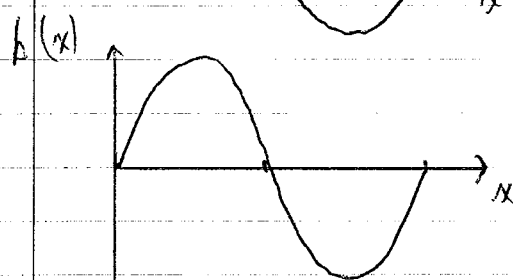
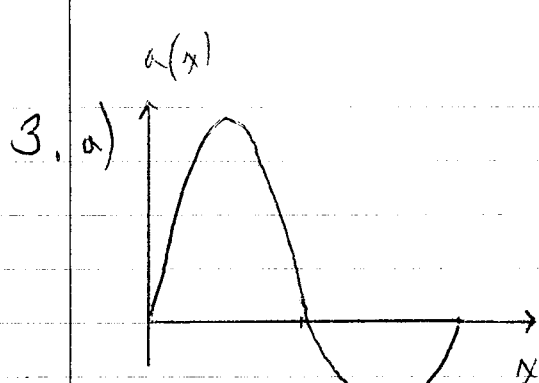
$$P\left[\frac{L}{4}, \frac{3L}{4}\right] = \int_{L/4}^{3L/4} \frac{2}{L} \sin^2 \frac{\pi x}{L} dx = \frac{2}{L} \cdot \frac{1}{2} \int_{L/4}^{3L/4} (1 - \cos \frac{2\pi x}{L}) dx$$

$$= \frac{1}{L} \cdot x \Big|_{L/4}^{3L/4} - \frac{1}{L} \cdot \frac{L}{2\pi} \sin \frac{2\pi x}{L} \Big|_{L/4}^{3L/4}$$

$$= \frac{1}{L} \cdot \frac{L}{2} - \frac{1}{2\pi} \left(\sin \frac{3\pi}{2} - \sin \frac{\pi}{2} \right) = \frac{1}{2} + \frac{1}{\pi} = \underline{\underline{0.818}}$$



In the ground state the probability of finding the particle in the middle half of the well is large.



The inner products are the integrals of these functions

i) $\langle a|b \rangle > 0$ ii) $\langle a|c \rangle < 0$ iii) $\langle b|c \rangle < 0$

b) $|\langle a|c \rangle| > |\langle a|b \rangle| = |\langle b|c \rangle|$

$\langle a|c \rangle$ is largest in magnitude because $a(x)c(x)$ has the largest bump

$\langle a|b \rangle = -\langle b|c \rangle$ because $c(x) = -a(x)$

so $|\langle a|b \rangle| = |\langle b|c \rangle|$