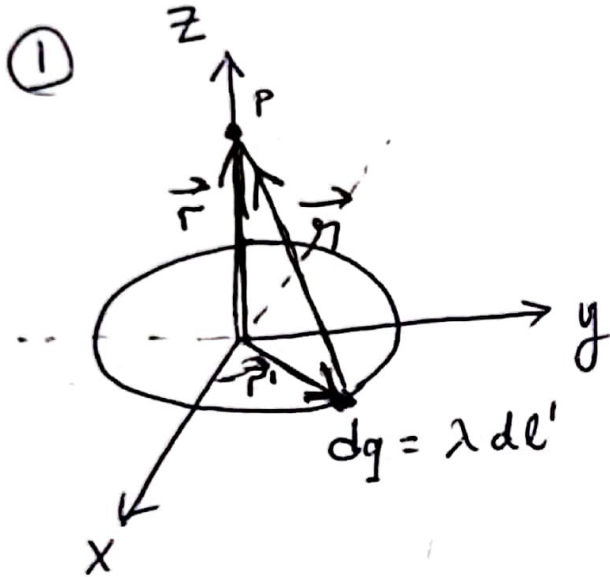


Electric Field of a ring of charge



②

$$\vec{r} = z \hat{z}$$
$$\vec{r}' = R \hat{s}$$
$$\vec{r} = \vec{r} - \vec{r}' = z \hat{z} - R \hat{s}$$
$$\hat{r} = \frac{\vec{r}}{|\vec{r}|} = \frac{z \hat{z} - R \hat{s}}{\sqrt{z^2 + R^2}}$$

$$dl' = R d\phi'$$

③

$$\hat{r} = \frac{z \hat{z} - R (\cos \phi \hat{x} + \sin \phi \hat{y})}{\sqrt{z^2 + R^2}}$$

$$= -\frac{R \cos \phi}{\sqrt{z^2 + R^2}} \hat{x} - \frac{R \sin \phi}{\sqrt{z^2 + R^2}} \hat{y} + \frac{z}{\sqrt{z^2 + R^2}} \hat{z}$$

$$\begin{aligned}
 \textcircled{4} \quad \vec{E} &= \frac{1}{4\pi\epsilon_0} \int \frac{\lambda}{r^2} \hat{r} \, dl' \\
 &= \frac{\lambda}{4\pi\epsilon_0} \int \frac{-R\cos\phi \hat{x} - R\sin\phi \hat{y} + z \hat{z}}{(\sqrt{z^2+R^2})^3} R \, d\phi
 \end{aligned}$$

$$\textcircled{5} \quad E_x = - \frac{1}{4\pi\epsilon_0} \frac{\lambda R^2}{(z^2+R^2)^{3/2}} \underbrace{\int_0^{2\pi} \cos\phi \, d\phi}_0 = 0$$

$$E_y = - \frac{1}{4\pi\epsilon_0} \frac{\lambda R^2}{(z^2+R^2)^{3/2}} \underbrace{\int_0^{2\pi} \sin\phi \, d\phi}_0 = 0$$

$$E_z = \frac{1}{4\pi\epsilon_0} \frac{\lambda z R}{(z^2+R^2)^{3/2}} \underbrace{\int_0^{2\pi} d\phi}_{2\pi} = \frac{1}{4\pi\epsilon_0} \frac{\lambda(2\pi R) z}{(z^2+R^2)^{3/2}}$$

There is a symmetry in x and y . So it makes sense that $E_x=0$ and $E_y=0$. Now we have

$$\vec{E} = \frac{Q}{4\pi\epsilon_0} \frac{z}{(z^2+R^2)^{3/2}} \hat{z}$$

Checking your result

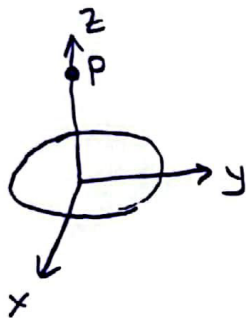
① $\vec{E} = \frac{1}{4\pi\epsilon_0} \frac{\lambda 2\pi R z}{(z^2 + R^2)^{3/2}} \hat{z}$

$$\left[\frac{1}{4\pi\epsilon_0} \right] = \left[\frac{Nm^2}{C^2} \right], \quad [\lambda] = \left[\frac{C}{m} \right], \quad [R \cdot z] = [m^2]$$

$$[(z^2 + R^2)^{3/2}] = [m^3]$$

$$\frac{\left[\frac{Nm^2}{C^2} \right] \left[\frac{C}{m} \right] [m^2]}{[m^3]} = \left[\frac{N}{C} \right] \quad \checkmark$$

②



Check 1: At $z=0$, we expect $E=0$ by symmetry.

Check 2: At $z \gg R$, we expect the charge on the ring looks like a point charge.

③ • At $z=0$: $E(z=0) = \frac{1}{4\pi\epsilon_0} \frac{\lambda 2\pi R z}{(z^2 + R^2)^{3/2}} \overset{0}{=} 0 \quad \checkmark$

• At $z \gg R$: $(z^2 + R^2)^{-3/2} = z^{-3} \left(1 + \frac{R^2}{z^2} \right)^{-3/2} \approx z^{-3} \left(1 - \frac{3}{2} \frac{R^2}{z^2} \right)$

Using $(1 \pm \epsilon) \approx 1 \pm n\epsilon$ if $\epsilon \ll 1$

$$\text{So } E \approx \frac{1}{4\pi\epsilon_0} \lambda 2\pi R z \cdot z^{-3} \left(1 - \frac{3}{2} \frac{R^2}{z^2} \right) = \frac{1}{4\pi\epsilon_0} \frac{Q}{z^2} \left(1 - \frac{3}{2} \frac{R^2}{z^2} \right)$$

point charge small correction

So: units \checkmark

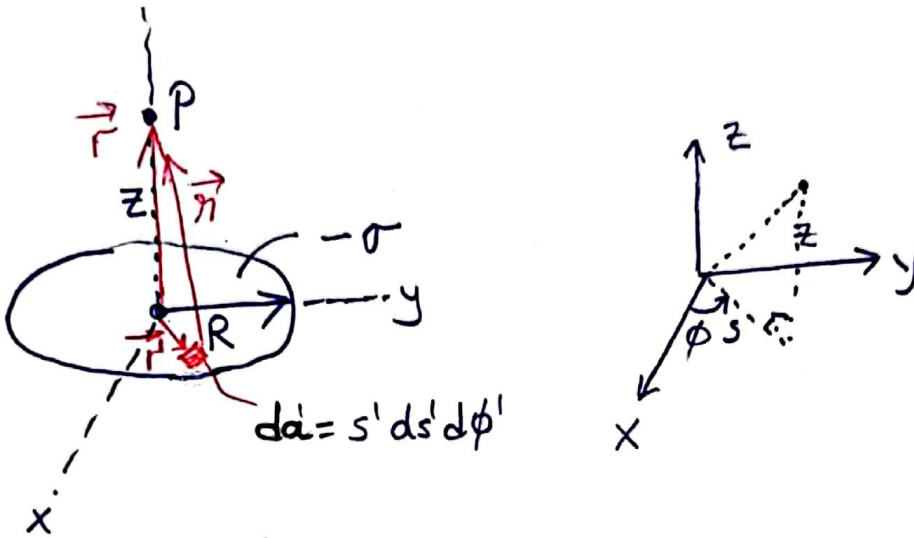
$E \rightarrow 0$ as $z \rightarrow 0 \quad \checkmark$

$E \rightarrow$ field due to point charge as $z \gg R \quad \checkmark$

Based on those checks, I believe this result is correct.

Electric field of a disk of charge

①



$$\vec{r} = z \hat{z}$$

$$\vec{r}' = s' \hat{s}$$

$$\vec{r} = \vec{r} - \vec{r}' = z \hat{z} - s' \hat{s} = z \hat{z} - s' \cos \phi' \hat{x} - s' \sin \phi' \hat{y}$$

$$|\vec{r}| = \sqrt{z^2 + s'^2}$$

$$dq = -\sigma da' = -\sigma s' ds' d\phi'$$

$$\vec{E} = \frac{1}{4\pi\epsilon_0} \int_S \left(-\sigma da' \cdot \frac{\vec{r}}{|\vec{r}|^3} \right) = \frac{-\sigma}{4\pi\epsilon_0} \int \frac{-s' \cos \phi' \hat{x} - s' \sin \phi' \hat{y} + z \hat{z}}{(z^2 + s'^2)^{3/2}} s' ds' d\phi'$$

$$E_x = \frac{-\sigma}{4\pi\epsilon_0} \int_0^R \frac{-s'^2 ds'}{(z^2 + s'^2)^{3/2}} \underbrace{\int_0^{2\pi} \cos \phi' d\phi'}_0 = 0$$

Similarly $E_y = 0$ because $\int_0^{2\pi} \sin \phi' d\phi' = 0$

$$E_z = \frac{-\sigma}{4\pi\epsilon_0} \int_0^R \underbrace{\frac{z s' ds'}{(z^2 + s'^2)^{3/2}}}_{-\frac{z}{\sqrt{z^2 + s'^2}} \Big|_{s'=0}^{s'=R}} \underbrace{\int_0^{2\pi} d\phi'}_{2\pi} = \frac{-\sigma \cdot 2\pi}{4\pi\epsilon_0} \left(\frac{-z}{\sqrt{R^2 + z^2}} + \frac{z}{|z|} \right)$$

$$\text{Hence } \vec{E} = \frac{-\sigma}{2\epsilon_0} \left(\frac{z}{|z|} - \frac{z}{\sqrt{R^2 + z^2}} \right) \hat{z}$$

② If $z \gg R$, then $\frac{R}{z} \ll 1$ (Assume $z > 0$)

$$\frac{z}{\sqrt{R^2 + z^2}} = \frac{z}{z \sqrt{\frac{R^2}{z^2} + 1}} = \left(1 + \frac{R^2}{z^2} \right)^{-1/2} \approx 1 - \frac{R^2}{2z^2}$$

using $(1 \pm \epsilon)^n \approx 1 \pm n\epsilon$
where $\epsilon \ll 1$.

$$\text{So } \frac{z}{|z|} - \frac{z}{\sqrt{R^2 + z^2}} \approx 1 - \left(1 - \frac{R^2}{2z^2} \right) = \frac{R^2}{2z^2} \quad (\text{for } z > 0)$$

$$|\vec{E}|_{\text{far}} \approx -\frac{\sigma}{2\epsilon_0} \frac{R^2}{2z^2} = \frac{(-\sigma \pi R^2)}{4\pi\epsilon_0 z^2} = \frac{-|Q|}{4\pi\epsilon_0 z^2} \quad (\text{for } z > 0)$$

As expected, $|\vec{E}|$ behaves like E-field of a point charge

③ If $z \ll R$, then $\frac{z}{R} \ll 1$ (Assume $z > 0$)

$$\frac{z}{\sqrt{R^2 + z^2}} = \frac{z}{R \sqrt{1 + \frac{z^2}{R^2}}} = \frac{z}{R} \left(1 + \frac{z^2}{R^2}\right)^{-1/2} \approx \frac{z}{R} \left(1 - \frac{z^2}{2R^2}\right)$$

$$= \frac{z}{R} - \frac{z^3}{2R^3} \approx \frac{z}{R}$$

↪ very small

$$\text{So, } \frac{z}{|z|} - \frac{z}{\sqrt{R^2 + z^2}} \approx 1 - \frac{z}{R} \quad (\text{for } z > 0)$$

$$\text{Therefore, } |\vec{E}|_{\text{near}} \approx -\frac{\sigma}{2\epsilon_0} \left[1 - \frac{z}{R} \right]$$

field from infinite sheet of charge + small correction

$$\textcircled{4} \quad \vec{E} = -\frac{\sigma}{2\epsilon_0} \left[\frac{z}{|z|} - \frac{z}{\sqrt{R^2 + z^2}} \right] \hat{z}$$

At $z=0 \rightarrow$ discontinuity because of $\frac{z}{|z|}$

For $z > 0 \rightarrow \frac{z}{|z|} = 1$

For $z < 0 \rightarrow \frac{z}{|z|} = -1$

