

CW 1 - Intro to Class. Mech. ①

Outline: What is Classical Mechanics?

How do we formulate it?

What are the essential physics models for single particles?

What mathematics do we need to get started?

What's Classical Physics ?

- the study of slow, large things

slow? no relativity; no QFT

large? no quantum; no stat-mech

What about Classical Mechanics?

- We now add "mechanical" to our conditions and so we exclude electro magnetic systems.

⇒ Not always. We can describe (2)
the force on a charged
particle using a classical model,

$$\vec{F}_{\text{Lorentz}} = q(\vec{E} + \vec{v} \times \vec{B})$$

How do we formulate Class. Mech.?

We first consider how have seen
classical mechanics in the past.

$$\vec{F}_{\text{net}} = m\vec{a}$$

Newton's 2nd
law

Notice this formulation is vector
based. That is, the relationship
between pushes and accelerations
are vectorial. Namely,

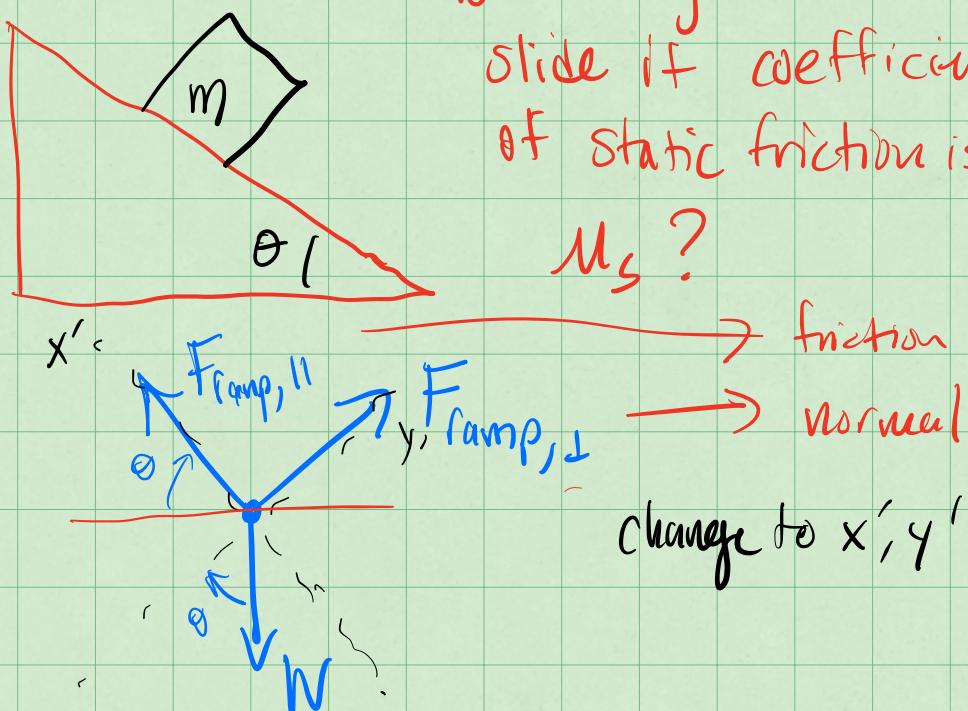
$$F_x = m a_x \quad F_y = m a_y \quad F_z = m a_z$$

Each push in a Cartesian direction ③

results in a proportional response—an acceleration in the same direction as the net push.

Ex.: Box on a plane with friction.

What angle does it slide if coefficient of static friction is μ_s ?



$$\vec{F}_{\text{net}} = m\vec{a} = 0 \quad \text{static}$$

$$\text{Max friction force} = F_{\parallel} = \mu_s F_{\perp}$$

(4)

$$\sum F_{xi} = F_{\text{ramp}, \parallel} - W \sin \theta = 0$$

$$\sum F_{yj} = F_{\text{ramp}, \perp} - W \cos \theta = 0$$

$$W = Mg \quad \text{so that ,}$$

$$F_{\text{ramp}, \parallel} = Mg \sin \theta \quad F_{\text{ramp}, \perp} = Mg \cos \theta$$

$$\text{But } F_{\text{ramp}, \parallel, \max} = \mu_s Mg \cos \theta$$

so,

$$Mg \sin \theta = \mu_s Mg \cos \theta @ \max!$$

$$\tan \theta = \mu_s$$

$$\theta = \tan^{-1} (\mu_s)$$

$$\mu_{\text{steel}} = 0.16$$

$$\sim 90^\circ$$

$$\mu_{\text{rubber}} = 0.8$$

$$\sim 39^\circ$$

Notes: - this was a static problem

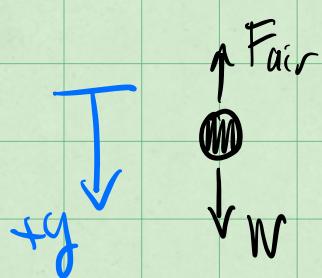
(5)

$$\vec{F}_{\text{net}} = 0$$

- We rotated the coordinate system
to match our ramp

- We still used Cartesian coords.

Ex: Falling Ball in 1D Predict its motion



models for F_{air} ?

$$\text{let } F_{\text{air}} = F(v)$$

just some
function of v

In 1D,

$$F_{\text{net},y} = m a_y = +mg - F(v)$$

Assume low v , why?

\Rightarrow Classical Mechanics!

Taylor Expand $F(v)$ & keep low terms (b)

$$f(x) \approx \sum_{n=0}^N \frac{1}{n!} \left. \frac{d^n f}{dx^n} \right|_{x=a} (x-a)^n$$

formula for Taylor Expansion around
 $x=a$.

$$f(x) \approx f(a) + f'(a)(x-a) + \frac{1}{2} f''(a)(x-a)^2 + \dots$$

Corl. Let's do that for $F(v)$ around
 $v=0$ when $F_{\text{drag}} = 0$

$$F(v) = \underbrace{F(0)}_0 + \underbrace{F'(0)v}_b + \underbrace{\frac{F''(0)}{2} v^2}_c + \dots$$

these are just #'s



quadratic drag

$$F(v) \approx b v + c v^2$$

linear drag

Back to Newton 2,

(7)

$$F_y = ma_y = mg - bv - cv^2$$

and thus,

$$a_y = g - \frac{b}{m}v - \frac{c}{m}v^2 \quad \text{OOF.}$$

How do we solve this?

$$a_y = g - \frac{b}{m}v - \frac{c}{m}v^2$$

$$\frac{dy}{dt^2} = g - \frac{b}{m}\left(\frac{dy}{dt}\right) - \frac{c}{m}\left(\frac{dy}{dt}\right)^2$$

$$\ddot{y} = g - \frac{b}{m}\dot{y} - \frac{c}{m}\dot{y}^2$$

$$\dot{v} = g - \frac{b}{m}v - \frac{c}{m}v^2$$

We will come back to it.

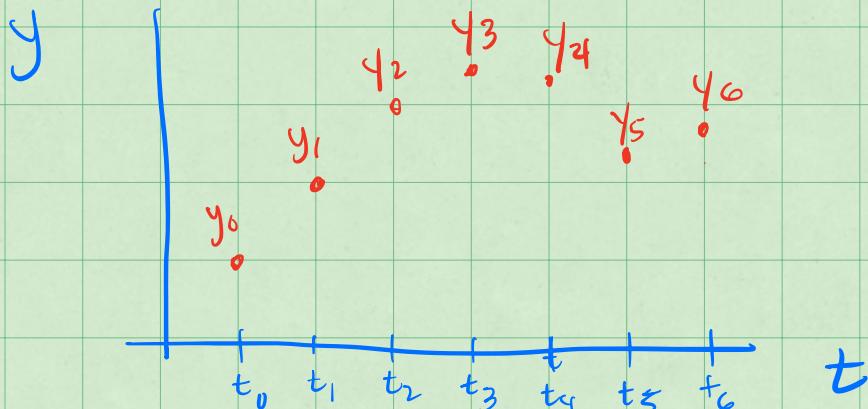
- Note : - this is a dynamic 1D problem (8)
- this is a nonlinear problem
 - we are stuck @ the moment

Enter Discretization ← another formulation

We posit discrete time, like snapshots of the motion where a given measure of time, t_i exists in a discrete set, from $t_0 \rightarrow t_f$
(initial → final)

$$t \in [t_0, t_f]$$

thus we conceive of a plot of motion as discrete,



| t | y |
|----------|----------|
| t_0 | y_0 |
| t_1 | y_1 |
| t_2 | y_2 |
| \vdots | \vdots |

if these are equally spaced (9)

then,

or

$$\Delta t = t_{i+1} - t_i = \frac{t_f - t_0}{n}$$

$$\text{Thus, } t_i = t_0 + i \Delta t$$

$$y(t_i) = y_i$$

Great, but what can we do with this?

let's define an average velocity over

a time step, Δt , like this,

$$v(t) = \frac{y(t + \Delta t) - y(t)}{\Delta t}$$

Avg velocity

$$v(t_i) = v_i \quad \leftarrow \text{discrete } v \text{ also.}$$

$$v_i = \frac{y_{i+1} - y_i}{\Delta t}$$

Avg velocity
(Discrete)

Note if we take the limit of $\Delta t \rightarrow 0$
we have the instantaneous velocity

$$\lim_{\Delta t \rightarrow 0} \frac{y(t + \Delta t) - y(t)}{\Delta t} = \frac{dy}{dt} = \dot{y}$$

fundamental
theorem of
calculus

Ok what about the acceleration?

We can also define an average accel over
an interval Δt ,

$$a(t) = \frac{v(t + \Delta t) - v(t)}{\Delta t}$$

average
acceleration

$$a(t_i) = a_i$$

discrete a ,

$$a_i = \frac{v_{i+1} - v_i}{\Delta t}$$

average acceleration
(discrete)

Again we can take the limit as $\Delta t \rightarrow 0$
to show the instantaneous acceleration

$$\lim_{\Delta t \rightarrow 0} \frac{v(t + \Delta t) - v(t)}{\Delta t} = \frac{dv}{dt} = \ddot{v} = \ddot{y}$$

again
FTC

Discrete Formulation of Mechanics

Let there be a 1D net force, $F(x)$

Here the force changes with location, x ,
a position dependent force.

$F(x_i) = F_i \rightarrow$ discretize force.

$$a_i = F_i/m \rightarrow \text{Newton 2}$$

$$a_i = \frac{v_{i+1} - v_i}{\Delta t} \Rightarrow v_{i+1} = v_i + a_i \Delta t$$

$$v_{i+1} = v_i + \frac{F_i}{m} \Delta t$$

predict the new
velocity just a
bit later.

Nice! Now we can predict the new velocity, v_{i+1} , a little time later. (12)

We will pause here and derive these methods for numerical integration later.

The discrete formulation is quite powerful and will help us solve our equations of motion like,

$$a_y = g - \frac{c}{m} v - \frac{d}{m} v^2$$

What mathematical ideas are we going to need? Obviously, algebra & geometry ← lots

Coordinate Sys & transforms ← lots

Differential & Integral Calculus ← lots

Vectors and Vector operations ← lots

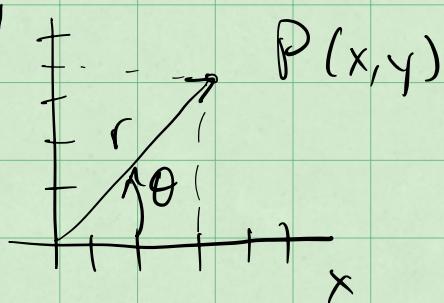
Discrete Calculus ← same

Complex Analysis ← a little

Vectors

(13)

in 2D



$$x^2 + y^2 = r^2$$

$$x = r \cos \theta$$

$$y = r \sin \theta$$

Plane Polar Coordinates (r, θ)

$$\vec{r} = x\hat{x} + y\hat{y} = x\hat{e}_x + y\hat{e}_y = x\hat{i} + y\hat{j}$$

↳ ↳ ↳ ↳ ↳ ↳ ↳ ↳

Unit vectors - for Cartesian, fixed in space/time

Claim

$$\vec{r} = |\vec{r}| \hat{r}$$

no $\hat{\theta}$

$$\vec{r} = x\hat{x} + y\hat{y}$$

$$|\vec{r}| = \sqrt{x^2 + y^2}$$

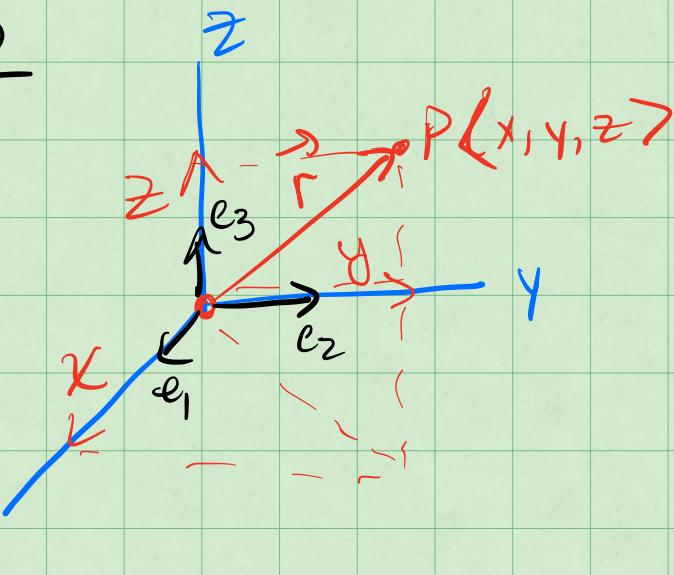
$$\hat{r} = \frac{\vec{r}}{|\vec{r}|} = \frac{x\hat{x} + y\hat{y}}{\sqrt{x^2 + y^2}}$$

$$\vec{r} = |\vec{r}| \hat{r} = x\hat{x} + y\hat{y}$$



In 3D

(14)



$$\vec{r} = \hat{x}\hat{x} + \hat{y}\hat{y} + \hat{z}\hat{z} = \hat{x}\hat{e}_1 + \hat{y}\hat{e}_2 + \hat{z}\hat{e}_3 \text{ etc.}$$

Unit Vectors

Cartesian unit vectors are fixed in space/time in inertial frames.

Magnitude = 1 $|\hat{x}|=1$ $|\hat{e}_1|=1$ etc.

The are orthogonal \Rightarrow their dot product vanishes b/c they are \perp

$$\hat{x} \cdot \hat{y} = 0$$

$$\hat{z} \cdot \hat{z} = 1$$

$$\hat{e}_1 \cdot \hat{e}_3 = 0$$

$$\hat{e}_2 \cdot \hat{e}_2 = 1$$

etc.

etc...

Dot Products (Inner Products)

(15)

$$\vec{a} \cdot \vec{b} = \langle a_x, a_y, a_z \rangle \cdot \langle b_x, b_y, b_z \rangle$$

$$= a_x b_x + a_y b_y + a_z b_z = |\vec{a}| |\vec{b}| \cos \theta_{AB}$$



The dot product is distributive,

$$\vec{a} \cdot (\vec{b} + \vec{c}) = \vec{a} \cdot \vec{b} + \vec{a} \cdot \vec{c}$$

PROOF:

$$\vec{a} \cdot [\vec{b} + \vec{c}] = \langle a_x, a_y, a_z \rangle \cdot \langle b_x + c_x, b_y + c_y, b_z + c_z \rangle$$

$$= a_x(b_x + c_x) + a_y(b_y + c_y) + a_z(b_z + c_z)$$

$$= (a_x b_x + a_x c_x) + (a_y b_y + a_y c_y) + (a_z b_z + a_z c_z)$$

(16)

$$= (a_x b_x + a_y b_y + a_z b_z) + (a_x c_x + a_y c_y + a_z c_z)$$

$$= \vec{a} \cdot \vec{b} + \vec{a} \cdot \vec{c} = \vec{a} \cdot (\vec{b} + \vec{c}) \quad \checkmark$$