### Personal Statement

## Research Experience

As an undergraduate at the University of Central Florida (UCF) I worked for Joe Harrington for four years and earned BS degrees in Math and Physics. Joe's group studies exoplanets using transit, secondary eclipse, and radial velocity data.

My first year working with Joe, I used the group's orbits code (rv, Campo et al. 2011) which fits orbital parameters to data from stellar radial velocity measurements and transit and secondary eclipse timing using a Markov chain. I used this code to fit an orbit to the planet TrES-1b (Cubillos et al. 2014).

During my second and third years of working for Joe, I worked on a code called Bayesian Atmospheric Radiative Transfer (BART, Harrington et al. 2019). BART uses a Markov chain to constrain the atmospheres of hot-Jupiter exoplanets by fitting the radiative transfer model described by Rojo (2006) to transit and secondary eclipse spectroscopy data.

I optimized the radiative transfer model so it could run  $10^5$  to  $10^7$  times within the Markov chain. Patricio Cubillos and I went through all 15,000+ lines of code to determine what could be moved to an initialization function which would be run first, and what would have to be looped over for every iteration. After that, we used a code performance profiler to further optimize the runtime of the code.

During my third rnd fourth years of working for Joe, I used our Photometry for Orbits, Eclipses, and Transits (POET, Stevenson et al. (2012)) pipeline to fit a light curve to two secondary-eclipse observations of the planet HAT-P-30b. A companion star complicated this analysis, and I modified POET to be able to model and subtract out the flux from this companion star. I then used our orbit code to fit an orbit to HAT-P-30b. Finally, I used BART to fit an atmosphere to HAT-P-30b and published a first-author paper. (Foster et al. 2019).

## Outreach and Education Experience

At UCF, I ran the Astronomy Society. We held an open-to-the-public observing event every week (weather permitting) at the UCF Robinson Observatory 8-inch telescopes on the lawn and a tour inside of the observatory and its half-meter telescope. We would give additional private events for groups such as scout troops or local elementary schools in exchange for donations for our annual science project, a stratospheric weather balloon launch.

In grad school at Cornell, I am the outreach coordinator for the graduate students, organizing events at local libraries, the museum of the earth, youth groups, and local 4-H groups. I have also assisted with outreach events organized by Cornell's Spacecraft and Planetary Imaging Facility (SPIF).

I am in my sixth semester of being a graduate TA, including my second year running a class of my own design. I adapted the UCF Astronomy Society's stratospheric balloon project into an optional one credit lab class for our intro to astronomy students.

In lab 1 the students propose designs for a weather balloon payload that fits within a \$600 budget and 800g weight limit. The proposal includes constructing a prototype of the structure of the payload from old styrofoam containers and cardboard and writing a proposal for what electronics they want to use their budget on.

In lab 2 the lab groups negotiate these multiple proposals down to a single proposal within the timeframe of a single 2.5 hour lab session.

The third lab involves constructing and completing the payload, including soldering and programming the electronics, performing drop tests to determine the payload's terminal velocity, and learning to calculate predicted landing sites using a program that integrates high altitude wind profiles.

The fourth lab is launch and recovery of the weather balloon.

### Occultation Observations of Saturn with Cassini VIMS

### 1. Goals of the Investigation

During Cassini's 13 years in orbit around Saturn, the spacecraft's Visual and Infrared Mapping Spectrometer (VIMS) performed over 100 stellar occultation experiments on Saturn, covering a wide range of latitudes and seasons. Nicholson et al. (2018) describe the VIMS campaign of occultations, including the relevant instrumental setups for observing stellar and solar occultations, event predictions, geometric and photometric calibrations, and instrumental sensitivity.

Occultation observations at different observer geometries and wavelengths probe different levels of a planet's atmosphere. Spacecraft radio occultations are sensitive to differential refraction (Kliore et al. 2004), and probe the lower stratosphere and upper troposphere of Saturn at pressures of a few mbar to  $\sim 1$  bar (Lindal et al. 1985). Earth-based visible and near-infrared occultations of the outer planets are also sensitive to differential refraction and probe the upper stratosphere and mesosphere at pressures in the microbar range (?Elliot 1979; French et al. 1978). Spacecraft UV occultations ( $\lambda < 150$  nm) are dominated by molecular and atomic absorption and probe the thermosphere of Saturn at pressures in the nanobar to  $\sim 1 \,\mu \text{bar range}$  (Broadfoot et al. 1981; Sandel et al. 1982; Smith et al. 1983).

This proposal deals with occultations observed by the Cassini VIMS instrument, which took time series of spectra with 20-40 ms sample rates at near-infrared wavelengths (0.88–5.1  $\mu$ m in 256 spectral channels) as stars disappeared behind Saturn. These data cover an intermediate regime where some wavelengths are dominated by differential refraction and some by molecular absorption (methane and ethane). They span pressures from  $\sim 20~\mu$ bar to  $\sim 5~\text{mbar}$ , covering a large part of Saturn's stratosphere. See figure 1 for an example of what a VIMS occultation

looks like.

Modeling the effects of differential refraction reveals the temperature profile of the atmosphere. Modeling the effects of molecular absorption reveals the composition. This data is the first that allows us to retrieve both simultaneously from the same *Cassini* dataset and in the same region of the atmosphere.

The principle goals of observing Saturn stellar occultations with VIMS were to obtain light curves that could be inverted to obtain temperature-pressure profiles of the planet's stratosphere, and to obtain vertical profiles of CH<sub>4</sub> abundance in the upper stratosphere with which to test photochemical and eddy-diffusion models such as those of Moses et al. (2000) and Moses and Greathouse (2005). Some preliminary results have been presented by Nicholson et al. (2006) and Banfield et al. (2011), but so far no comprehensive study has been completed or published.

I propose to perform just such a comprehensive study that will attempt to answer the questions outlined in the previous paragraph and to address the NASA Planetary Science Research Program SMD 2014 Science Plan top-level goal to:

"Advance the understanding of how the chemical and physical processes in the Solar System operate, interact and evolve."

More specifically, we will address the topic:

"Investigate the origins, evolution, and properties of the atmospheres of planetary bodies (including satellites, small bodies, and exoplanets)"

as well as

"Enhance the scientific return of NASA Planetary Science Division missions through the analysis of data collected

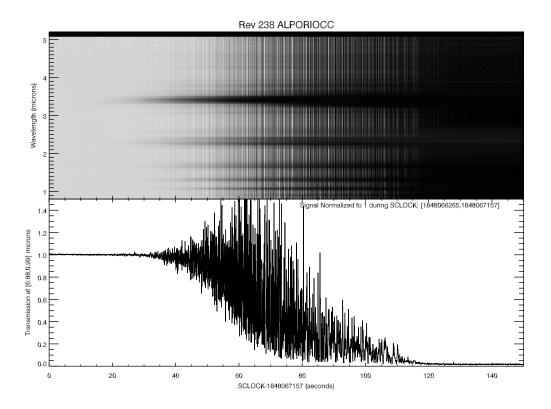


Fig. 1.— Occultation of  $\alpha$  Orionis during Cassini rev: 238 at planetocentric latitude  $-64.9^{\circ}$  from a range of 959,000 km with an integration time of 20 ms. **Upper panel** shows stellar signal with wavelength increasing vertically and time increasing horizontally. Stellar signal drops due to differential refraction independently of wavelength except for the molecular absorption in the narrow horizontal bands at 3.4, 2.3, 1.7, 1.4, 1.1 and 0.9  $\mu$ m. **The lower panel** shows the flux averaged between 0.9 and 1.0  $\mu$ mwhere molecular extinction is unimportant. In both panels, the stellar flux is normalized to unity in the period prior to the occultation.

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### 2. Scientific Significance

Pre-Cassini observations of Saturn's temperature show a peak in seasonal variation at around 3 mbar (Yanamandra-Fisher et al. 2005), the region that the VIMS stellar occultation data probe. CIRS observations by Flasar et al. (2004, 2005); Fletcher et al. (2007) confirm this. Our proposed analysis can provide far better altitude resolution than the CIRS results in this critical region. Furthermore, the CIRS results are derived from thermal emission in the methane band

at  $8\mu$ mand cannot easily disentangle variations in temperature from variations in methane mixing ratio. We expect methane abundances also to depend on latitude and season as described by Moses and Greathouse (2005) and Fouchet et al. (2009). The proposed analysis will allow us to recover temperature profiles and methane abundances independently of each other at various latitudes and at seasons spanning almost half of a Saturnian year. (Science Project 1: Verify CIRS Results)

The chemical abundance profiles calcu-

lated in Moses and Greathouse (2005) match and Eddy-Diffusion) the "methane cycle" described in D.F. Strobel (1969). Methane diffuses up through the stratosphere to the homopause, slowly dropping in abundance through the region of the stratosphere to which VIMS occultations are sensitive. Far above this region of the stratosphere, methane undergoes photodissociation and is converted to higher-order hydrocarbons which diffuse back down to the troposphere where they are converted back to methane. The gases diffuse from their source to their sink via turbulence- or wind-driven eddy diffusion. The rates of diffusion, production, and destruction of these gases determine the shape of the abundance profiles through the stratosphere seen in figure 2.

As noted by Fouchet et al. (2009), important byproducts of these multiple profiles will be a better determination of the altitude of the CH<sub>4</sub> homopause on Saturn (which is known to be guite variable with latitude from UV occultations Koskinen et al. 2013, 2015). This is a function of the rate of photochemistry, which destroys CH4 in the upper atmosphere (Fouchet et al. 2009), and the speed of zonal winds in the stratosphere which distort the shape of the planet through the centrifugal force (Merritt and Nicholson 2019). (Science Project 2: Measure the Homopause)

Photodissociated methane is partially converted into ethane on its way to becoming aerosols composed of higher order hydrocarbons. Chemical abundances in the stratosphere impact the thermal profile of the entire atmosphere since trace gasses are the major absorbers of radiation from the sun and Saturn's blackbody emissions from lower levels. The proposed analysis will constrain both the methane and ethane mixing ratios, giving us insight into seasonal variations in photochemistry and eddy-diffusion rates that we can compare to models such as those described in Moses and Greathouse (2005). (Science Project 3: Test Models of Chemistry

Photochemistry also leads to aerosol pro-Although not very spectrally active duction. in our observations, aerosols are thought to be chemically important as a factor in cloud formation because they sediment downwards to serve as cloud-condensation nuclei for condensible volatiles in the troposphere (Fletcher et al. 2018). We will constrain photochemistry and eddy-diffusion rates through a direct measurement of the relative mixing ratios of methane to ethane, which will help future investigators understand aerosol production in the thermosphere and cloud formation processes in the convective regions of the troposphere. We will measure these abundances with high radial resolution over the course of nearly half of a Saturnian year at a multitude of latitudes thanks to Cassini's long-lasting mission in the Saturn system.

Finally, uncertainties in the production and optical properties of photochemical haze are a major stumbling blocks in current hot-Jupiter exoplanet atmospheric models and retrievals, as they are thought to be the cause of the flattening of the transit spectra of gaseous exoplanets (Fraine et al. 2013). Through providing a ground-truth understanding of the atmosphere of one of our local giant planets whose atmospheric chemistry is also photochemically driven, our analysis will help future studies constrain the atmospheres of these exciting bodies that we will soon be able to probe in greater detail with the James Webb Space Telescope, ELT-class groundbased telescopes, and other future telescopes.

# Technical Background Differential Refraction

Consider a light ray from the occulted star incident on a spherical planet of radius R, at an impact parameter  $\rho$ . As it penetrates more deeply into the atmosphere, where the refractive index of the gas is larger, the ray gradually refracts towards the planet's center until it reaches a minimum radius, after which it traces a mirror-

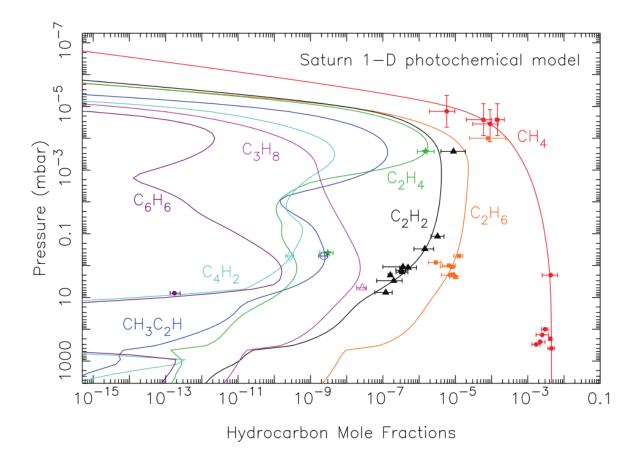


Fig. 2.— Figure taken from Fouchet et al. (2009). Hydrocarbon mole fractions as a function of pressure in Saturn's upper atmosphere as derived from the "Model C" 1-D steady-state photochemical model of Moses and Greathouse (2005). The solid curves represent the model profiles for the individual hydrocarbons (as labeled), and the symbols with associated error bars represent current infrared and ultraviolet observations. Methane is created in the troposphere and diffuses upwards to the homopause and the thermosphere where it is photochemically destroyed. Ethane and acetylene are produced at around  $\times 10^{-4}$  mbar and diffuse down to the troposphere where they are converted back to methane. This is the methane cycle described in D.F. Strobel (1969)

image path back out of the atmosphere.

The formal expression for the overall angular deflection of the ray once it emerges from the atmosphere is:

$$\theta(\rho) = -2\rho \int_{r_0}^{\infty} \frac{dn/dr}{n\sqrt{n^2r^2 - \rho^2}} dr, \qquad (1)$$

where n(r) is the refractive index profile of the atmosphere, and the ray's minimum radius  $r_0$  is specified by the condition  $n(r_0)r_0 = \rho$ .

Adjacent light rays are deflected by progressively larger amounts as  $\rho$  decreases and the rays

penetrate more deeply into the atmosphere. As a result of this differential refraction, the flux of light from the source is reduced by a factor of  $\Phi$ :

$$\Phi(\rho) = \left(1 - D\frac{d\theta}{d\rho}\right)^{-1} \left(1 - \frac{D\theta}{\rho}\right)^{-1}, \quad (2)$$

where D is the distance from the planet to the observer. The second factor represents the focusing effect that occurs near the center of the geometric shadow and can be neglected for our data where  $D\theta \ll \rho$ .

As described in Elliot et al. (1977)

and French et al. (1978), we can use these two expressions to perform an inverse Abel transform and invert our occultation lightcurves in regions of the spectra dominated by differential refraction to acquire multiple radial profiles of the temperature and pressure in the planet's stratosphere.

This technique is standard for inversions of differential-refraction dominated stellar occultations. Due to Saturn's highly oblate shape, the above expressions are adjusted in practice to use a local spherical fit to the atmosphere.

We are confident that the attenuation of the starlight at wavelengths where hydrocarbon molecular absorption is negligible is dominated by the above-described impact of differential refraction and not aerosol-driven scattering extinction. Although aerosols are an important chemical component of Saturn's atmosphere, these particles are likely sub-micron in size and their scattering efficiency should decrease rapidly with decreasing wavelength, perhaps scaling as  $\lambda^{-q}$  where q is between 1 and 4. The aerosols on Titan follow such a power law with  $q \simeq 1.8 \pm 0.5$  (Bellucci et al. 2009). The extinction in the VIMS Saturn occultations at wavelengths outside of the strong hydrocarbon bands is observed to be essentially independent of wavelength, strongly suggesting that aerosol extinction is negligible at the few mbar level and above and therefore we may neglect them in our models.

# 3.2. Morphology of an Occultation

Let's examine equation 2, neglect the second term, and substitute the appropriate isothermal expression for  $\theta(\rho)$  to get an approximate expression of  $\Phi(\rho)$  that can give us some intuition for how flux is lost over the course of an occultation:

$$\Phi(\rho) \simeq \left(1 + \frac{D\theta}{H}\right)^{-1}.$$
(3)

where the atmosphere's scale height  $H = \mathcal{R}T/\mu g$ , T is the temperature,  $\mu$  is the mean molecular weight, g is the local gravitational ac-

celeration, and  $\mathcal{R}$  is the universal gas constant.

Note that  $\Phi(\rho) \simeq \frac{1}{2}$  when  $D\theta = H$ , indicating that the impact parameter for which the flux from the source has halved is dependent on the distance to the observer D. This is the reason that Earth-based stellar occultations probe much higher atmospheric levels than do those observed by orbiting spacecraft at similar wavelengths. For typical VIMS Saturn occultations, the refractive half-light level is at  $p \sim 3$  mbar, assuming that  $\bar{T}_{stratosphere} \sim 140 \mathrm{K}$ .

The bending angle  $\theta$  scales with the refractivity at the point of deepest penetration of the ray which increases exponentially with depth, so the attenuation of the rays increases rapidly as they drop below the half-light level. However, the observed brightness of the occulted star falls off much more gradually because the light rays are also deflected further into the planet's geometric shadow, prolonging the star's visibility. The result is that such a purely-refractive lightcurve has a long tail that extends well beyond the half-light time. In our observations this tail is truncated because the image of the star is refracted out of the single  $0.25 \times 0.5$  mrad VIMS pixel, leading to a "floor" beyond which we cannot probe. This occurs at  $\Phi \simeq \frac{1}{3.5}$ , for  $\theta \simeq 0.25$  mrad,  $D \simeq 5 \times 10^5$  km, and  $H \simeq 50$  km. This is empirically confirmed by a small number of VIMS occultations observed in the instrument's imaging-mode, which has a larger field-of-view at the expense of less-frequent timesampling.

### 3.3. Molecular Absorption

Let's consider an "onion-skin" model of an atmosphere for which each thin radial layer is of uniform composition and density. As a ray of star light passes through each of these layers, it is attenuated according to the equation:

$$\Phi(\lambda) = \Phi_0 \ e^{-\delta\tau(\lambda)},\tag{4}$$

where  $\Phi_0$  is the attenuation experienced before entering that layer and  $\delta \tau(\lambda)$  is the optical depth of the layer as a function of wavelength, described by the equation:

$$\delta \tau(\lambda) = \kappa(\lambda) L d \tag{5}$$

Here,  $\kappa$  is the absorption coefficient, L is the slant pathlength, and d is the density of the layer which is known from the inversion of the lightcurve at the refraction-dominated continuum wavelengths.

The data for a given occultation can be thought of as a time-series of spectra, each of which cuts a deeper path through the atmosphere than its predecessor. Each time-step can be used to constrain the opacity of an additional deeper layer, and then be used as an input to the next timestep since that layer will also be probed during the subsequent observations. Datapoints can be binned in time to increase signal to noise at the expense of the radial resolution of the resulting abundance profile. We call this technique "onion-peeling" and a preliminary proof-of-concept using a single wavelength bin was presented by Banfield et al. (2011).

We will fit a model to a full-spectrum  $\kappa(\lambda)$  calculated in each layer from our data. A model spectrum will be calculated with a line-by-line approach using Voigt profiles calculated using the pressures and temperatures from our refraction inversions described above. We will use line list databases such as the HITEMP database from HITRAN <sup>1</sup> (Gordon et al. 2017) for each gas relevant to Saturn's atmosphere. The contributions of each gas (chiefly CH<sub>4</sub> and C<sub>2</sub>H<sub>6</sub>, but possibly also C<sub>2</sub>H<sub>2</sub>) to the final opacity are independent and proportional to their mixing ratios, which will be fit to the data.

In practice, hydrogen and helium are essentially transparent in the near-infrared for any reasonable path length through the stratosphere<sup>2</sup>, so the opacity is dominated by trace gases, methane, ethane, and acetylene (Moses and Greathouse 2005).

### 4. Project Summary

Our project is to write a code that will calculate an onion-skin model of Saturn's stratosphere from the VIMS occultation data.

End of 2019 to Early 2020: We will first carry out an atmospheric inversion on the refraction-dominated parts of the occultation spectra to acquire the temperature, pressure, and density of each layer of the onion, at a radius accurately known from the occultation geometry.

Rest of 2020: We will then "peel" that onion to calculate molecular abundances in each layer as the occultation descends through Saturn's atmosphere. We will accomplish this by fitting spectra calculated to those observed over the course of the occultation.

This procedure will provide local atmospheric profiles on Saturn at each of the locations and times probed by an occultation.

**2021:** We will look for trends in latitude and season within this collection of profiles with three major publishable goals in mind:

- 1. Verify CIRS results
- 2. Measure CH<sub>4</sub> homopause
- 3. Constrain chemistry models and eddydiffusion rates

**2022:** Publish papers and defend my thesis by mid-summer.

<sup>&</sup>lt;sup>1</sup>https://hitran.org/hitemp

 $<sup>^{2}</sup>$ H<sub>2</sub> does have a fundamental vibrational transition at 2.1  $\mu$ m, which is detectable in the spectra of the Jovian planets, but this is a collision-induced absorption that scales as  $p^{2}$ , and so is very weak in the stratosphere, even for very long path lengths.

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