

# Effect of Delay on the Stability of a Coupled Reactor-Flash System Sustaining an Elementary Non-isothermal Reaction

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# Introduction

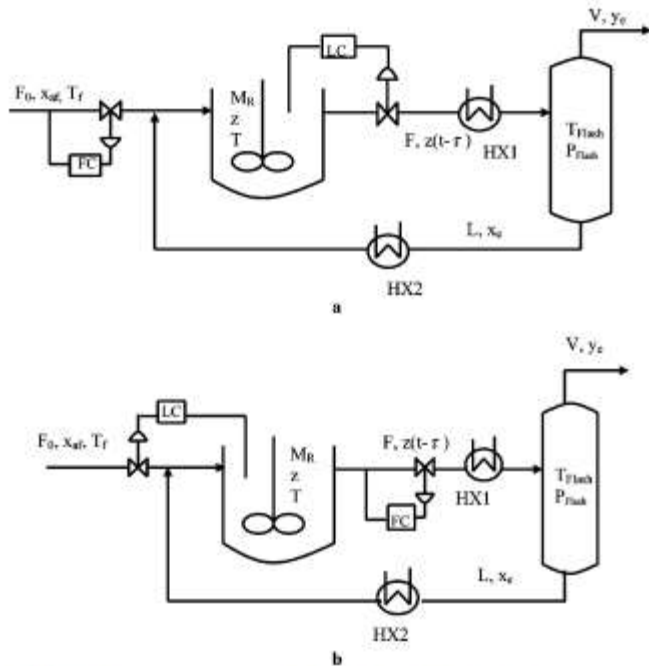
- In a typical chemical industry, a reactor is coupled with a separator downstream so that a recycle stream rich in the unreacted reactant can be fed back to the reactor upstream.
- Hence , the reactor and separator unit are coupled by the recycle stream and changes in any one of them affects the other.
- Therefore, it is important to study the stability and dynamics of such processes as it is undesirable to operate the system using process variables which may lead to unstable steady states.
- Along with this, we will also consider the effect of transportation delay from the reactor to the separator - The fluid leaving the reactor with a given concentration does not instantaneously reach the separator.

# Objective of the work

- We will be taking the example of a non isothermal CSTR coupled with an isothermal and isobaric flash.
- The reactor and separator are coupled both by mass as well as energy but here, we will consider only the effect of mass coupling.
- The stability and dynamic analysis without the effect of delay will be first considered.
- The main objective is to study changes occurring in the stability of the steady states when the effect of transportation delay from the CSTR to the flash column is considered.
- We will consider the effect of delay in the transportation from the reactor effluent to the flash column while neglecting the delay in the recycle stream.

# Description of the system

- The system consists of a non isothermal CSTR, coupled with an isothermal and isobaric flash downstream along with two heat exchangers as shown below with an elementary first order reaction.



**Figure 1.** Schematic diagram for a non-isothermal CSTR coupled to an isothermal, isobaric flash when (a)  $F_0$  is flow controlled and (b)  $F$  is flow controlled.

- The heat exchanger HX1 is used to bring the temperature of the reactor effluent stream to the flash temperature.
- Similarly the heat exchanger HX2 is used to change the temperature of the recycle stream from the flash temperature to the reactor feed temperature.
- In this way, the temperature of the reactor and separator are independent and the two units are hence energetically decoupled.
- This will allow us to study only the effect of mass coupling.

- Two cases of this system will be considered –
- Molar holdup ( $M_R$ ) and fresh feed flow rate ( $F_0$ ) is fixed while the reactor effluent flow rate ( $F$ ) is controlled.
  - Molar holdup ( $M_R$ ) and reactor effluent flow rate ( $F$ ) is fixed while the fresh feed flow rate ( $F_0$ ) is controlled.

# Model equations

- The species mole balance and energy balance for the CSTR is given as:

$$\frac{dz}{dt} = \left(\frac{F_0}{M_R}\right)x_{af} + \left(\frac{L}{M_R}\right)x_e - \left(\frac{F}{M_R}\right)z + r_A \quad (1)$$

$$\frac{dT}{dt} = \left(\left(\frac{F_0}{M_R}\right) + \left(\frac{L}{M_R}\right)\right)T_f - \left(\frac{F}{M_R}\right)T + \frac{(-\Delta H)r_A}{C_p} - \left(\frac{UA_h}{C_p M_R}\right)(T - T_c) \quad (2)$$

- The total and species mole balance for the flash column is given as:

$$F = V + L \quad (3)$$

$$Fz(t - \tau) = Vy_e + Lx_e \quad (4)$$

- As there is no accumulation of moles in the entire system, we get:

$$F_0 = V \quad (5)$$

# Model equations

The derivation of these equations is given as follows:

## 1. Species Mole balance in CSTR

Derivation of governing equations -

For the CSTR :

① Species mole balance on A

$$\frac{dN_A}{dt} = \dot{Q}_{in} C_{Ain} - \dot{Q}_{out} C_{Aout} + r_A V M_R$$

For this system,

$F, z$  - Molar flow rate and mole fraction of A in CSTR effluent stream

$F_0, x_{AF}$  - Molar flow rate and mole fraction in fresh feed

$L, x_e$  - Molar flow rate and mole fraction in recycle stream

$M_R$  - Molar holdup (No. of moles in the CSTR)

$$N_A = z M_R$$

$$\dot{Q}_{in} C_{Ain} = F_0 x_{AF} + L x_e$$

$$\dot{Q}_{out} C_{Aout} = Fz$$

Considering constant molar holdup  $M_R$ , we can take it out from the  $\frac{d}{dt}$  on the LHS and dividing on both sides, we get

$$\frac{dz}{dt} = \left(\frac{F_0}{M_R}\right) x_{AF} + \left(\frac{L}{M_R}\right) x_e - \left(\frac{F}{M_R}\right) z + r_A$$

# Model equations

## 2. Energy balance for non isothermal CSTR

② Energy balance for non isothermal CSTR

$$M_R C_p \frac{dT}{dt} = (F_0 + L) T_F - F T - \Delta H M_R r_A + \dot{Q}_h$$

$$\dot{Q}_h = UA(T_c - T), \quad T_c \text{ is temperature of coolant.}$$

$$\therefore M_R C_p \frac{dT}{dt} = (F_0 + L) T_F - F T - \Delta H M_R r_A + UA(T_c - T)$$

$$\therefore \frac{dT}{dt} = \left[ \left( \frac{F_0}{M_R} \right) + \left( \frac{L}{M_R} \right) \right] T_F - \left( \frac{F}{M_R} \right) T + \frac{(-\Delta H) r_A}{C_p} - \frac{UA(T - T_c)}{C_p M_R}$$

Constitutive relation for  $r_A$

For the elementary reaction  $A \rightarrow B$ ,

$$r_A = -k_0 z e^{-E/RT}$$

$r_A$  is defined wrt. mole fraction  $z$ , hence rate of generation term in the governing equation is  $r_A M_R$ .

# Model equations

- Frank Kamenetzskii Approximation -

For the temperature dependence of the rate, we will linearize the Arrhenius equation with respect to a reference temperature as below.

Frank Kamenetzskii approximation:

The temperature dependence of  $\tau_A$  i.e.  $e^{-E/RT}$  is linearized around reference temperature  $T_R$  as

$$\frac{1}{T} = \frac{1}{T_R} + \frac{(T - T_R)}{T_R^2}$$

$$\therefore \tau_A = k_0 z \exp\left(-\frac{E}{RT}\right)$$

$$= k_0 z \exp\left[-\frac{E}{R} \left(\frac{1}{T_R} + \frac{T - T_R}{T_R^2}\right)\right]$$

$$\therefore \tau_A = k_0 z \exp\left(-\frac{E}{RT_R}\right) \exp\left[-\frac{E(T - T_R)}{RT_R^2}\right]$$



# Model equations

## 3. Species and total mole balances on the flash column

For the flash column :

① Total mass balance

$$F = L + V$$

② Species mass balance

$$Fz(t-\tau) = Lx_e + Vy_e$$

Here,  $V, y_e$  = Flow rate and mole fraction in the distillate  
 $z(t-\tau)$  introduces the delay. The composition of reactor at a time instant earlier by  $\tau$  units reaches the flash at time  $t$ .

As we assume an isothermal, isobaric flash the compositions  $x_e$  and  $y_e$  are fixed depending on  $T, P$  of the column because they are in equilibrium.

As there is no accumulation in the system,

$$F_0 = V$$

# Degree of Freedom Analysis

- We have 5 equations and one constitutive relation for the temperature dependence of the reaction rate.
- $X_{af}$ ,  $T_f$ ,  $T_{flash}$ ,  $P_{flash}$ ,  $M_R$  are fixed. As the streams exiting the flash column are assumed to be in equilibrium, for fixed  $T_{flash}$  and  $P_{flash}$ ,  $x_e$  and  $y_e$  are also fixed.
- Hence there are 6 independent variables -  $z$ ,  $T$ ,  $F$ ,  $F_0$ ,  $L$ ,  $V$ .
- Number of model equations = 5.
- Therefore,  $DOF = 1$  and hence we can specify one of the 6 variables which leads to the two cases which we will be analyzing further -
  1. Fixed  $M_R$  and  $F_0$
  2. Fixed  $M_R$  and  $F$
- The governing equations in these cases can be simplified and converted to a non dimensional form as follows

# Dimensionless governing equations

## 1. Case 1: Fixed $M_R$ and $F_0$

- In this case, the variables  $F$ ,  $V$ ,  $L$  can be eliminated using equations 3-5 and  $F$  in terms of  $F_0$  can be substituted in equations 1 and 2.
- The equations obtained can be non dimensionalized as follows using the dimensionless variables given.

$$\frac{dz}{dt_1^*} = (x_{af} - z) + \left( \frac{z(t_1^* - \tau^*) - y_e}{x_e - z(t_1^* - \tau^*)} \right) (x_e - z) - Da_1 z e^\theta \quad (6)$$

$$\frac{d\theta}{dt_1^*} = -\theta \left( 1 + \beta_1 + \left( \frac{z(t_1^* - \tau^*) - y_e}{x_e - z(t_1^* - \tau^*)} \right) \right) + B_h Da_1 z e^\theta \quad (7)$$

where

$$Da_1 = \frac{M_R z k_0 e^{-(E/RT_f)}}{F_0}; \quad t_1^* = \frac{t F_0}{M_R}; \quad \theta = \frac{E}{RT_f^2} (T - T_f);$$

$$B_h = \frac{(-\Delta H)}{C_p} \left( \frac{E}{RT_f} \right); \quad \beta_1 = \frac{UA_h}{F_0 C_p M_R}$$

## 2. Case 2: Fixed $M_R$ and $F$

- In this case, the variables  $F_0$ ,  $V$ ,  $L$  can be eliminated using equations 3-5 and  $F_0$  in terms of  $F$  can be substituted in equations 1 and 2.
- The equations obtained can be non dimensionalized as follows using the dimensionless variables given.

$$\frac{dz}{dt_2^*} = (x_{af} - z) + \left( \frac{z(t_2^* - \tau^*) - y_e}{x_e - y_e} \right) (x_e - x_{af}) - Da_2 z e^\theta \quad (8)$$

$$\frac{d\theta}{dt_2^*} = -\theta (1 + \beta_2) + B_h Da_2 z e^\theta \quad (9)$$

where

$$Da_2 = \frac{M_R z k_0 e^{-(E/RT_f)}}{F}; \quad t_2^* = \frac{t F}{M_R}; \quad \beta_2 = \frac{UA_h}{F C_p M_R}$$

# Details of the simulation

## 1. Case 1: Fixed $M_R$ , $F_0$

The parameters used for the simulation are:

- $X_{af} = 1$
- $Y_e = 0.1$
- $X_e = 0.9$
- $B_h = 8.8$
- $\beta_1 = 4.0$

## 2. Case 2: Fixed $M_R$ , $F$

The parameters used for the simulation are:

- $X_{af} = 0.9$
  - $Y_e = 0.2$
  - $X_e = 0.8$
  - $B_h = 14.0$
  - $\beta_2 = 4.2$
- 
- In case 1,  $Da_1$  which depends on  $F_0$  and in case 2,  $Da_2$  which depends on  $F$  are varied for the bifurcation analysis.

# Details of the simulation

For the linear stability analysis of the delay differential equations, we use the following method:

- The characteristic equation to find the eigen values is given as  $\det(sI - J - J_\tau e^{s\tau}) = 0$
- If we consider that the real part of  $s = 0$ , we can get the value of  $\tau$  critical for which the system shows a transition in stability.
- Here,  $J$  is the Jacobian w.r.t the normal variables and  $J_\tau$  is the Jacobian w.r.t to the delayed variables evaluated at the steady state.
- Sample calculations to derive the Jacobians and the characteristic equation are shown for case 1 which can be repeated for case 2 as well.

For case 1,

$$\frac{dz}{dt^*} = (\alpha_0 - z) + \left[ \frac{z(t^* - \tau^*) - y_e}{\alpha_e - z(t^* - \tau^*)} \right] (\alpha_e - z) - D_{A_1} z e^{\theta}$$

$$\frac{d\theta}{dt^*} = -\theta \left\{ 1 + \beta_1 + \left[ \frac{z(t^* - \tau^*) - y_e}{\alpha_e - z(t^* - \tau^*)} \right] \right\} + B_h D_{A_1} z e^{\theta}$$

$$J = \begin{bmatrix} \frac{\partial f_1}{\partial z} & \frac{\partial f_1}{\partial \theta} \\ \frac{\partial f_2}{\partial z} & \frac{\partial f_2}{\partial \theta} \end{bmatrix} = \begin{bmatrix} -1 - \left( \frac{z - y_e}{\alpha_e - z} \right) - D_{A_1} e^{\theta} & -D_{A_1} z e^{\theta} \\ B_h D_{A_1} e^{\theta} & -1 - \beta - \left( \frac{z - y_e}{\alpha_e - z} \right) + B_h D_{A_1} z e^{\theta} \end{bmatrix}$$

$$J_\tau = \begin{bmatrix} \frac{\partial f_1}{\partial z(t-\tau)} & \frac{\partial f_1}{\partial \theta(t-\tau)} \\ \frac{\partial f_2}{\partial z(t-\tau)} & \frac{\partial f_2}{\partial \theta(t-\tau)} \end{bmatrix} = \begin{bmatrix} \frac{\alpha_e - y_e}{\alpha_e - z} & 0 \\ -\frac{\theta(\alpha_e - y_e)}{(\alpha_e - z)^2} & 0 \end{bmatrix}$$

Here,  $\theta(t-\tau)$  is not in any equation as we have not considered delay in  $\theta$ .

When  $s$  is purely imaginary,  $s = i\omega$ .

$$\therefore \det[sI - J - J_\tau e^{s\tau}] = \det[i\omega I - J - J_\tau e^{i\omega\tau}]$$

$$\therefore \det[i\omega I - J - J_\tau (\cos \omega\tau + i\sin \omega\tau)] = 0$$

This is converted to 2 equations by getting real and imaginary part = 0. Similarly these equations can be derived for case 2.

# Details of the simulation

- These two equations along with the steady state equations can be solved to find the critical  $\tau$  for a given value of Da for both the cases.
- The equations for the linear stability analysis derived by this method for are as follows:

$$0 = x_{af} - y_e - Da_1 z e^\theta \quad (13)$$

$$0 = -\theta \left( 1 + \beta_1 + \left( \frac{z - y_e}{x_e - z} \right) \right) + B_h Da_1 e^\theta \quad (14)$$

$$0 = -\omega^2 + \omega \left( \frac{y_e - x_e}{x_e - z} \right) \sin \omega \tau + \left( \left( \frac{y_e - x_e}{x_e - z} \right) \cos \omega \tau - \left( \frac{y_e - x_e}{x_e - z} \right) + Da_1 e^\theta \right) \left( \beta_1 + \left( \frac{x_e - y_e}{x_e - z} \right) - B_h Da_1 z e^\theta \right) - Da_1 e^\theta z \left( \left( \frac{\theta(x_e - y_e)}{(x_e - z)^2} \right) \cos \omega \tau - B_h Da_1 e^\theta \right) \quad (15)$$

$$0 = \omega \left( \left( \beta_1 + \left( \frac{x_e - y_e}{x_e - z} \right) - B_h Da_1 z e^\theta \right) + \left( \frac{y_e - x_e}{x_e - z} \right) \times (\cos \omega \tau - 1) + Da_1 e^\theta \right) - \left( \frac{y_e - x_e}{x_e - z} \right) \sin \omega \tau \left( \beta_1 + \left( \frac{x_e - y_e}{x_e - z} \right) - B_h Da_1 z e^\theta \right) + Da_1 e^\theta z \left( \frac{\theta(x_e - y_e)}{(x_e - z)^2} \right) \sin \omega \tau \quad (16)$$

$$0 = \frac{(x_e - z)(x_{af} - y_e)}{(x_e - y_e)} - Da_2 z e^\theta \quad (17)$$

$$0 = -(1 + \beta_2)\theta + B_h Da_2 z e^\theta \quad (18)$$

$$0 = -\omega^2 - \omega \left( \frac{x_e - x_{af}}{x_e - y_e} \right) \sin \omega \tau + \left( 1 + Da_2 e^\theta - \left( \frac{x_e - x_{af}}{x_e - y_e} \right) \cos \omega \tau \right) (1 + \beta_2 - B_h Da_2 z e^\theta) + B_h Da_2^2 z e^{2\theta} \quad (19)$$

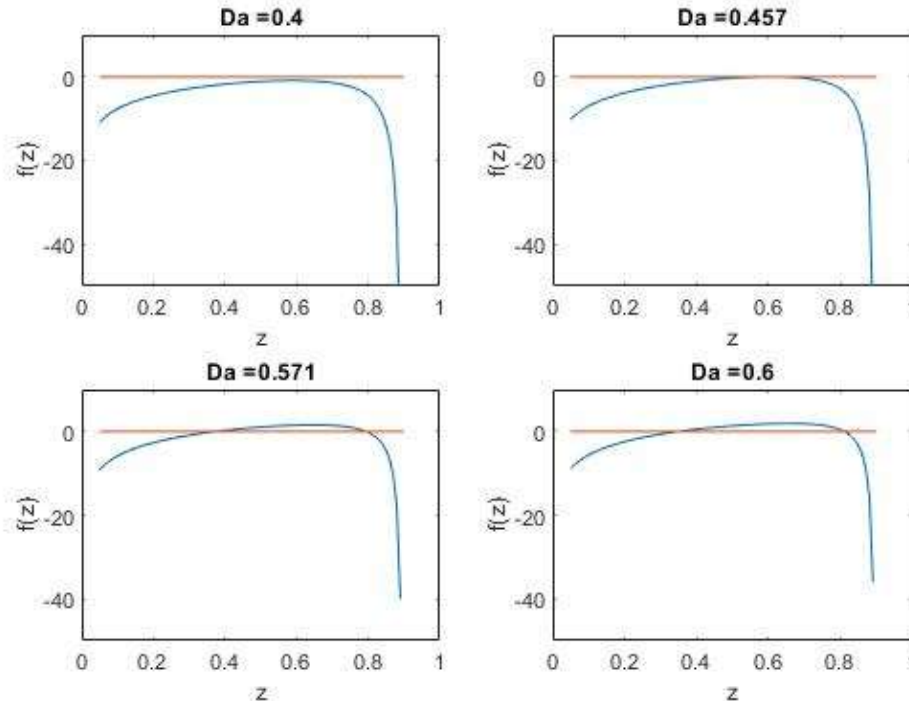
$$0 = \omega \left( 2 + \beta_2 + Da_2 e^\theta - B_h Da_2 z e^\theta - \left( \frac{x_e - x_{af}}{x_e - y_e} \right) \cos \omega \tau \right) + \left( \left( \frac{x_e - x_{af}}{x_e - y_e} \right) \sin \omega \tau \right) (1 + \beta_2 - B_h Da_2 z e^\theta) \quad (20)$$

- To simulate the dynamic response of the system with delay, the governing equations are a system of delay differential equations which are solved using the dde23 function in MATLAB.

# Results

1. Case 1: Fixed  $M_R, F_0$ 
  - Linear Stability analysis without delay:

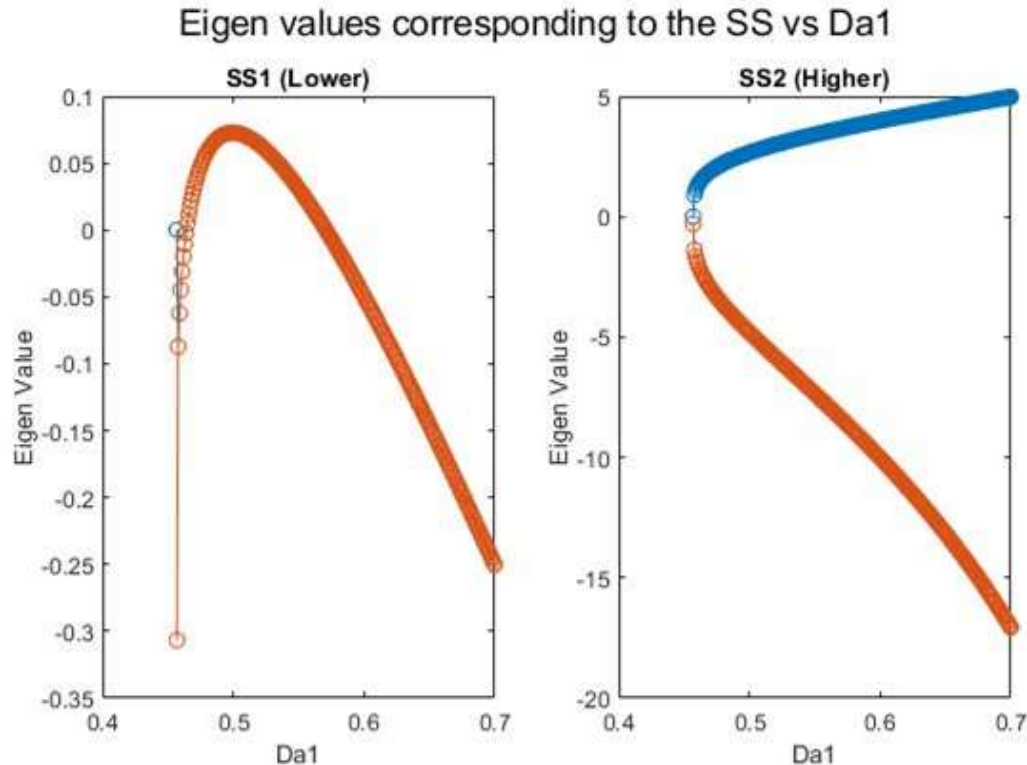
Graphical Method to find number of Steady States



- The variable  $\theta$  was eliminated from the two model equations at SS to get a single equation in terms of only  $z$  which has been plotted as  $f(z)$  vs  $z$  to find the number of steady states which occur as  $f(z) = 0$ .
- There is no steady state for  $Da_1 < 0.457$
- There is only one steady state at  $Da_1 = 0.457$  beyond which there are 2 SS for every  $Da_1$ .

# Results

1. Case 1: Fixed  $M_R$ ,  $F_0$ 
  - Linear Stability analysis without delay:

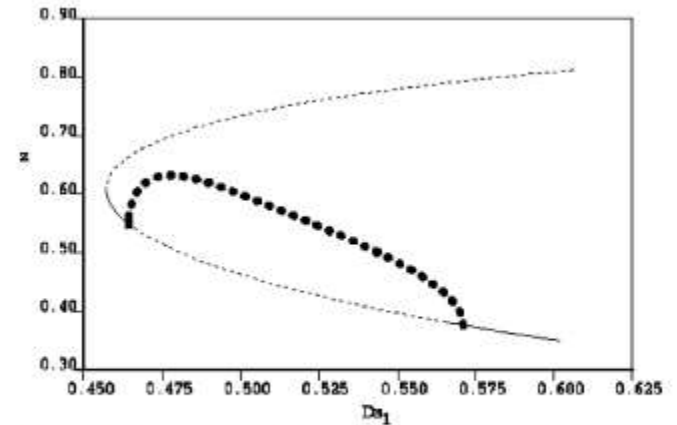
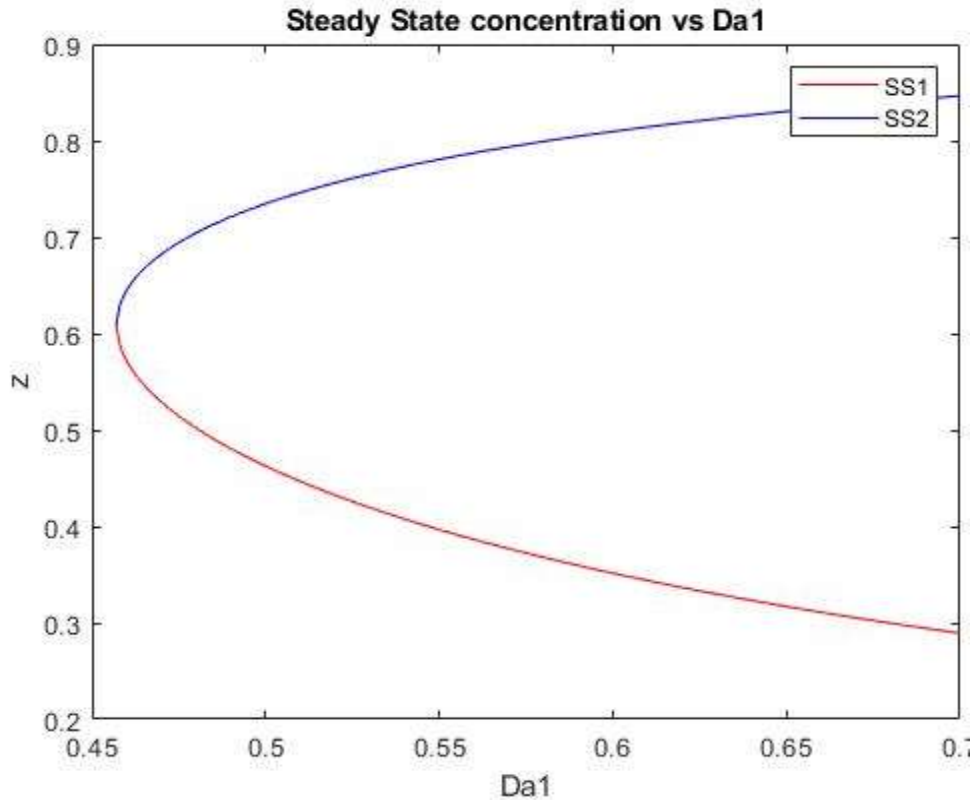


- For SS2, one eigen value is positive and the other is negative for all Da1, so the higher steady state is always unstable.
- For SS1, except for at  $Da1 = 0.457$ , the eigen values become complex with the same real part for both eigen values, hence we get only one curve.
- When  $0.464 < Da1 < 0.571$ , both eigen values are positive and hence the SS1 is also unstable, otherwise it is stable.



# Results

1. Case 1: Fixed  $M_R, F_0$ 
  - Bifurcation diagram without delay:

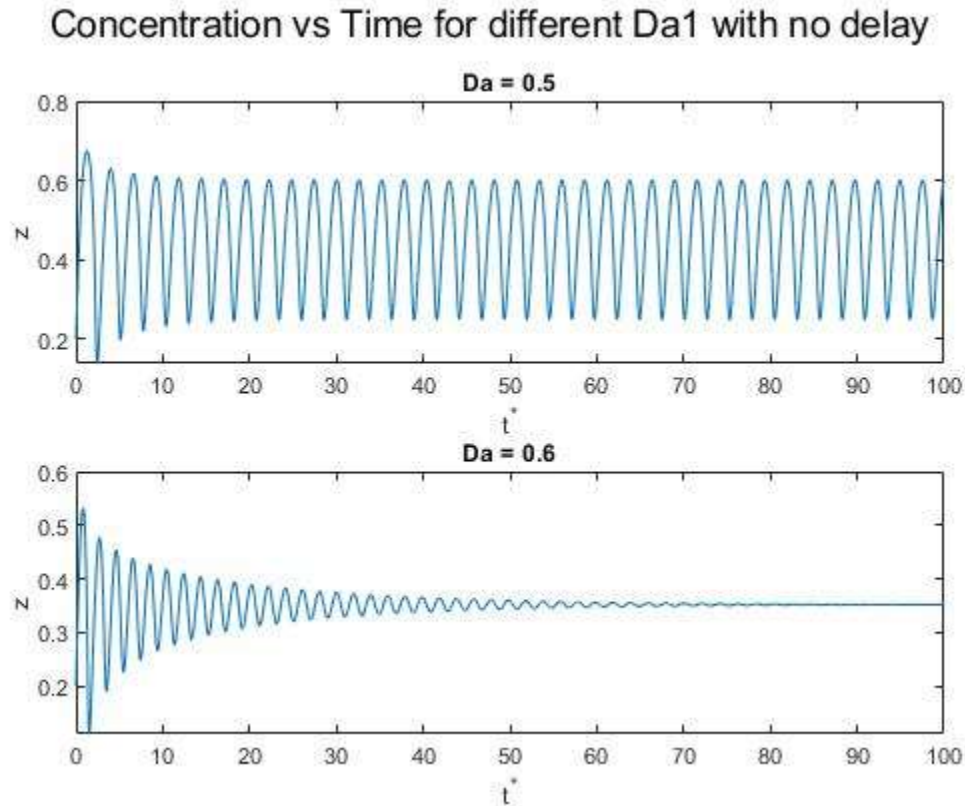


**Figure 2.** Bifurcation diagram of  $z$  versus  $Da_1$  for the fixed  $F_0, M_R$  case with  $x_{air} = 1.0, y_o = 0.1, x_o = 0.9, B_h = 8.8$ , and  $\beta_1 = 4.0$  (in the absence of delay).

- From the analysis of eigen values, we can see that for  $0.464 < Da_1 < 0.571$  both the steady states are unstable.
- For other values of  $Da_1$ , SS1 is stable while SS2 is unstable.
- This is also seen in the plot to the right where the dotted lines show the unstable SS and solid lines show stable SS.

# Results

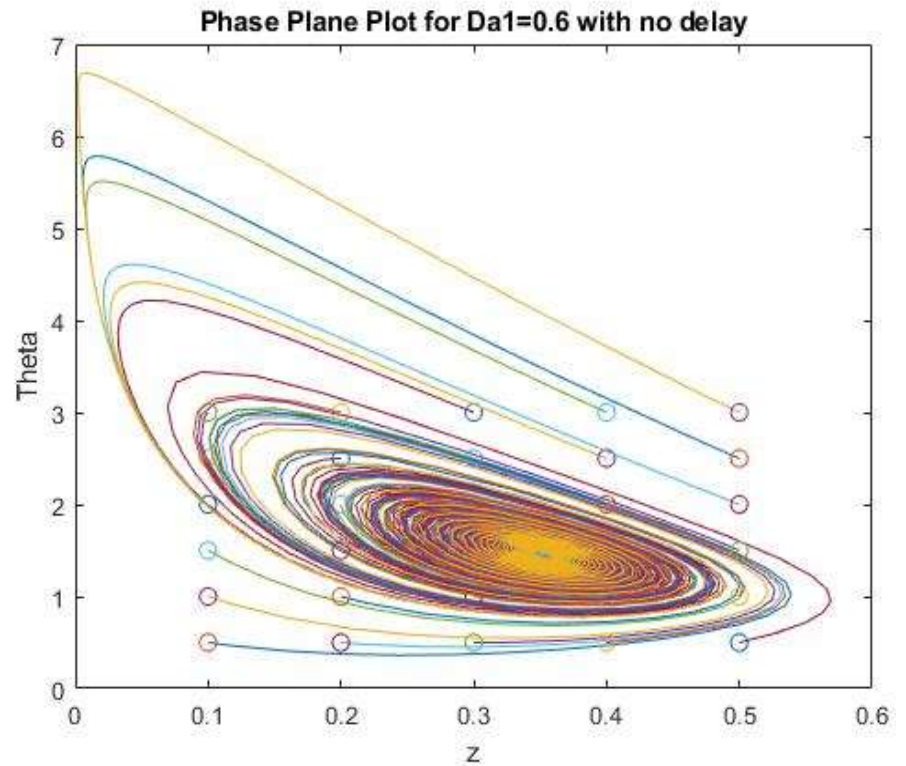
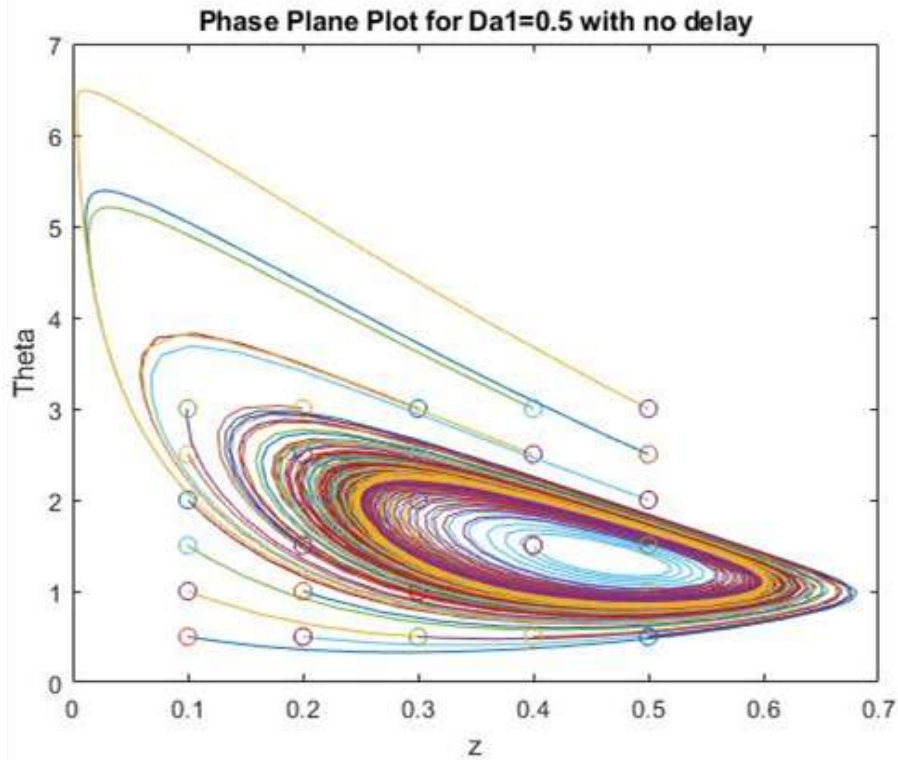
1. Case 1: Fixed  $M_R, F_0$ 
  - Dynamic analysis without delay :



- For  $Da1 = 0.5$ , we know that there is no stable SS and hence we get the oscillations.
- For  $Da1 = 0.6$ , only SS1 is stable and the system reaches SS1 as shown.

# Results

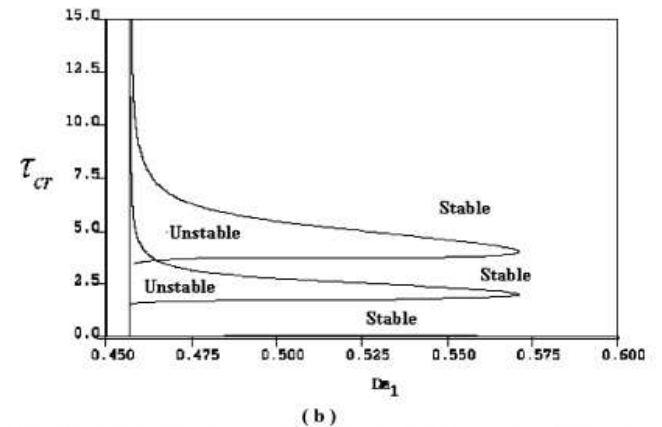
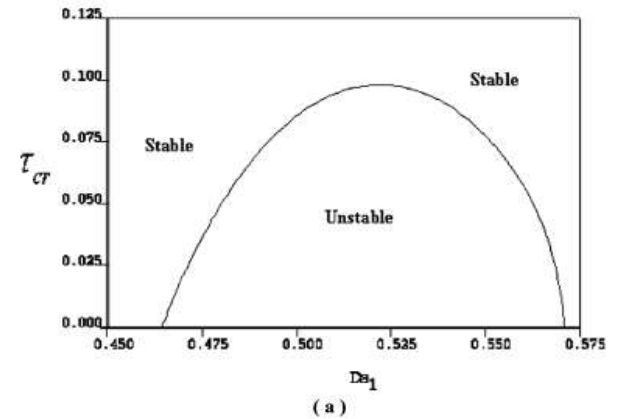
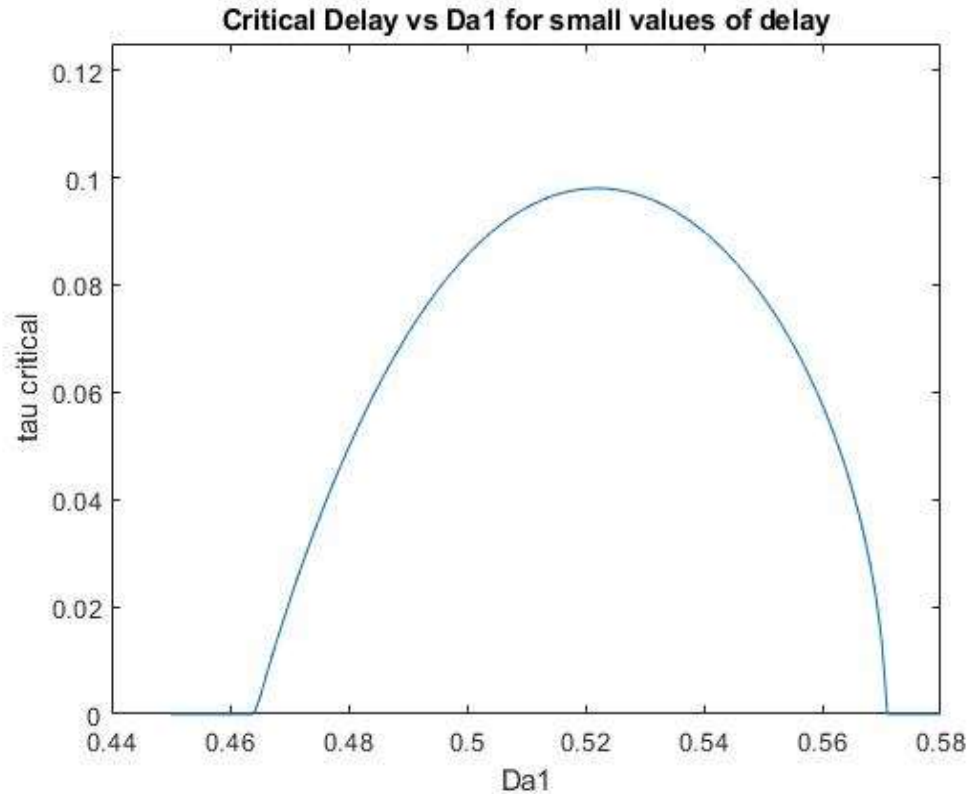
1. Case 1: Fixed  $M_R, F_0$ 
  - Phase Plane plots without delay:



- This again shows the previous result that a SS is attained for  $Da1 = 0.6$ , not for  $Da1 = 0.5$ .

# Results

1. Case 1: Fixed  $M_R$ ,  $F_0$ 
  - Bifurcation diagram with delay:



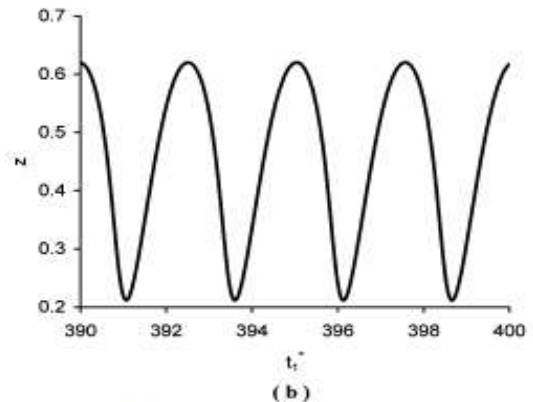
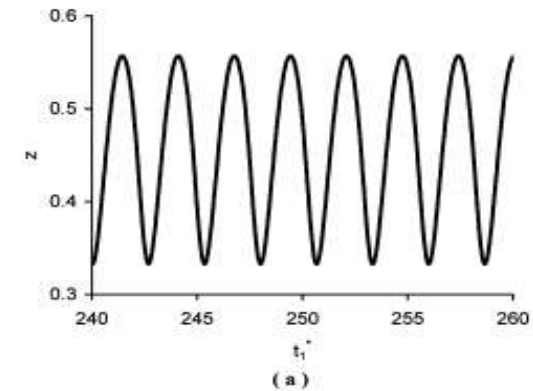
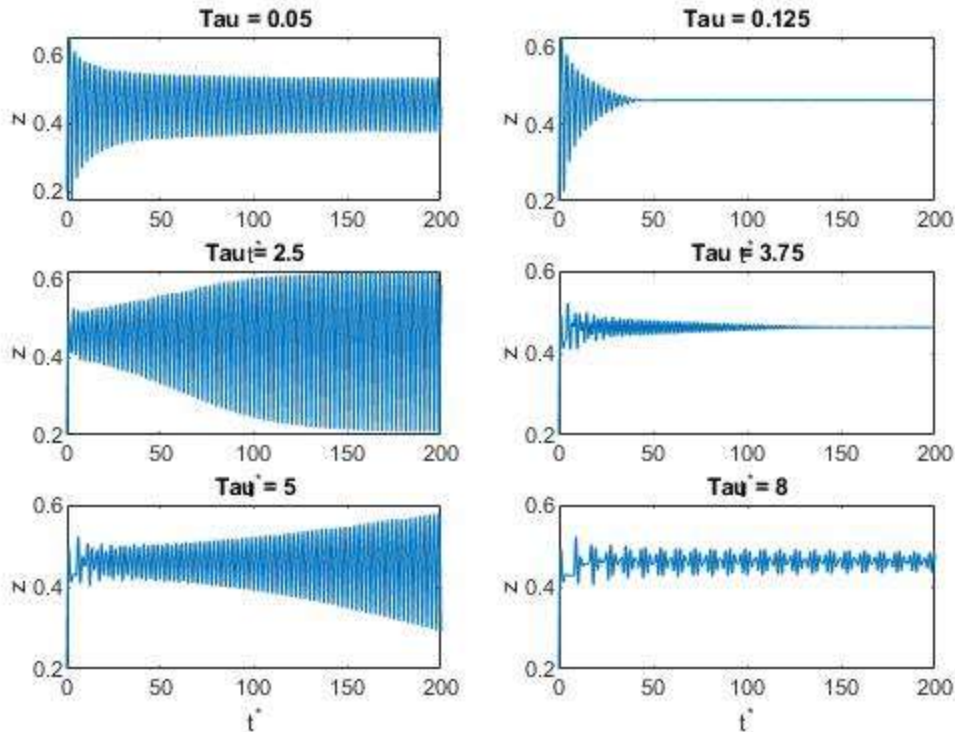
**Figure 3.** Dependency of dimensionless critical delay on  $Da_1$  for the fixed  $F_0$ ,  $M_R$  case with  $x_{af} = 1.0$ ,  $y_e = 0.1$ ,  $x_e = 0.9$ ,  $B_h = 8.8$ , and  $\beta_1 = 4.0$  for (a) small delay and (b) large delay.

- Depending on the value of the delay, we have different regions of stability as shown in the plot on the right. Below the critical delay, the system is unstable and beyond the critical delay, it is stable.
- As the delay increases, for small values the width of the region of dynamic instability decreases.

# Results

1. Case 1: Fixed  $M_R, F_0$ 
  - Dynamic analysis with delay:

Concentration vs Time for different  $Da_1$  with delay



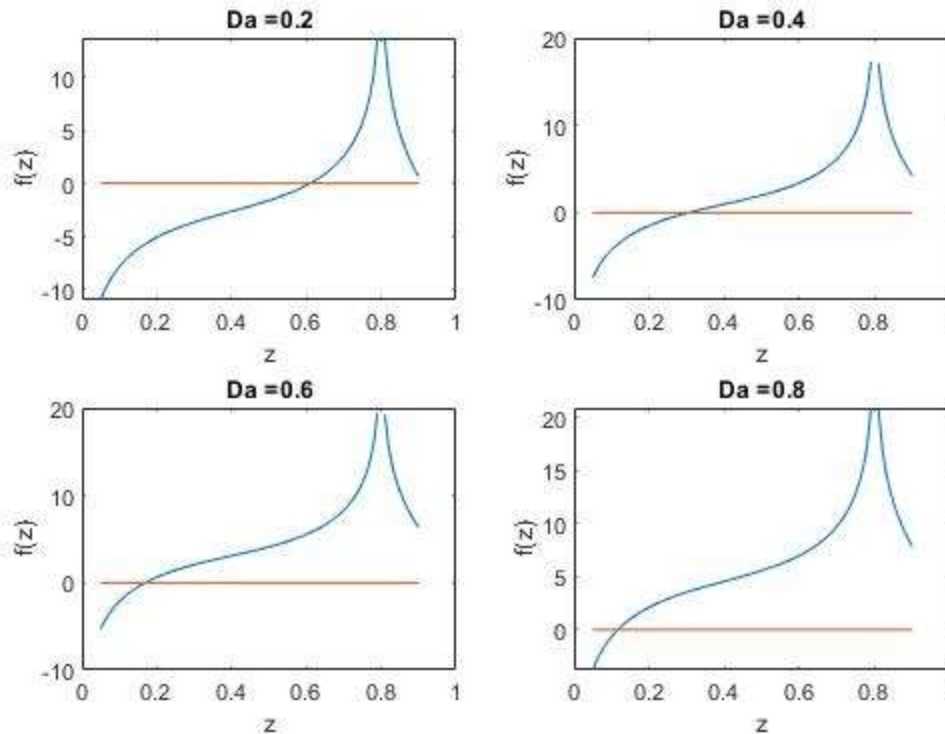
**Figure 4.** Mole fraction of **A** versus dimensionless time plot for the fixed  $F_0, M_R$  case showing sustained oscillations with  $x_{af} = 1.0, y_0 = 0.1, x_0 = 0.9, B_h = 8.8, \beta_1 = 4.0$ , and  $Da_1 = 0.5$  and with (a)  $\tau = 0.05$  and (b)  $\tau = 2.5$ .

- As shown, the stability is now dependent on the delay.
- The plots are shown for  $Da_1 = 0.5$  which lead to 2 unstable SS in the case without delay and the stability changes with delay which is also seen from the bifurcation diagram.
- For this case if we consider a isothermal CSTR, the stability becomes independent of the delay.

# Results

2. Case 2: Fixed  $M_R$ ,  $F$ 
  - Linear Stability analysis without delay:

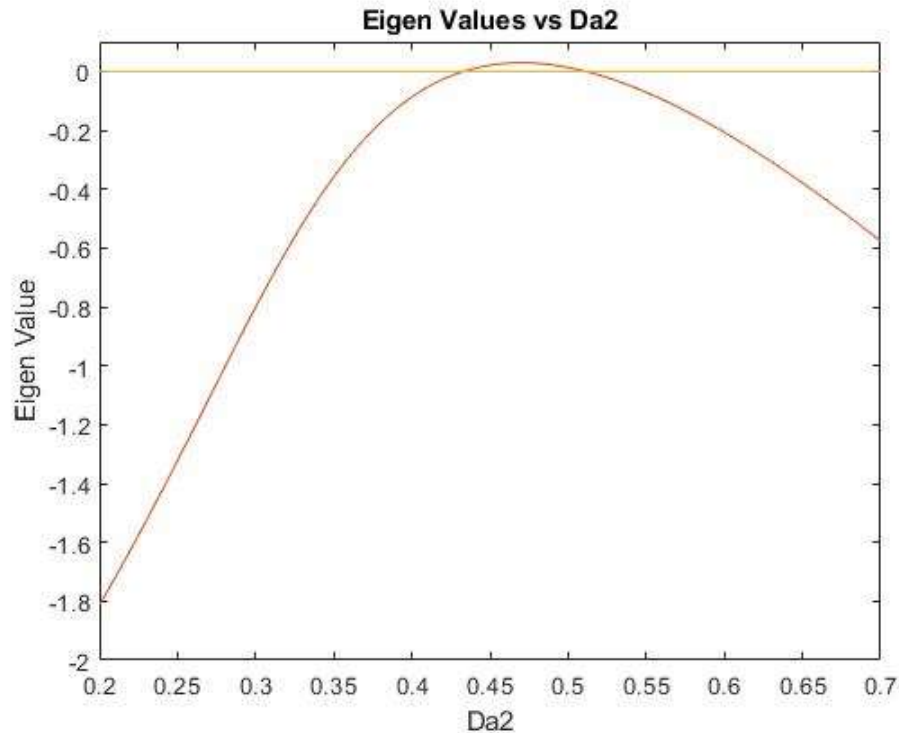
Graphical Method to find number of Steady States



- The variable  $\theta$  was eliminated from the two model equations at SS to get a single equation in terms of only  $z$  which has been plotted as  $f(z)$  vs  $z$  to find the number of steady states which occur as  $f(z) = 0$ .
- In this case, there is only one steady state for all values of  $Da$ .

# Results

2. Case 2: Fixed  $M_R$ ,  $F$
- Linear Stability analysis without delay:



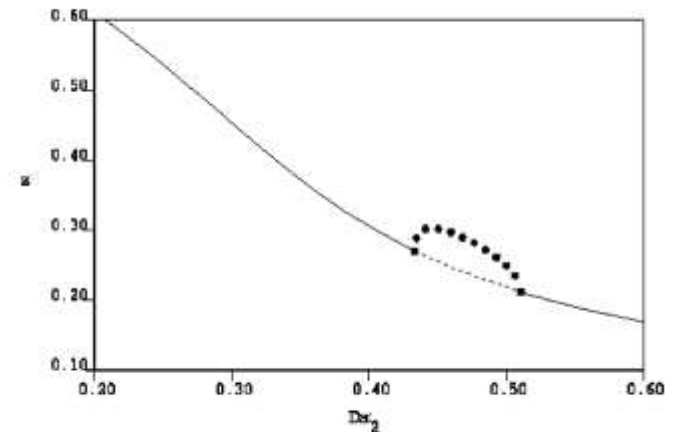
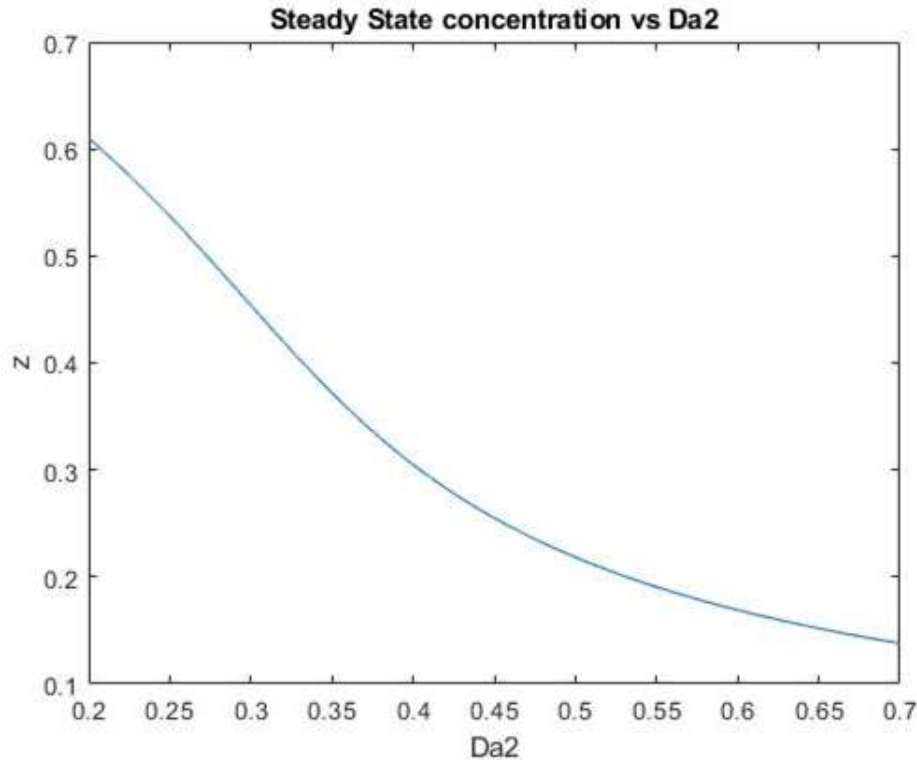
- Here, both the eigen values for the SS are imaginary with equal real parts, hence we get only one curve.
- It can be seen that for  $0.433 < Da_2 < 0.511$ , both eigen values are positive and so the SS is unstable.
- For other values of  $Da_2$ , both the eigen values are negative and we get a stable SS.



# Results

## 2. Case 2: Fixed $M_R$ , $F$

- Bifurcation diagram without delay:



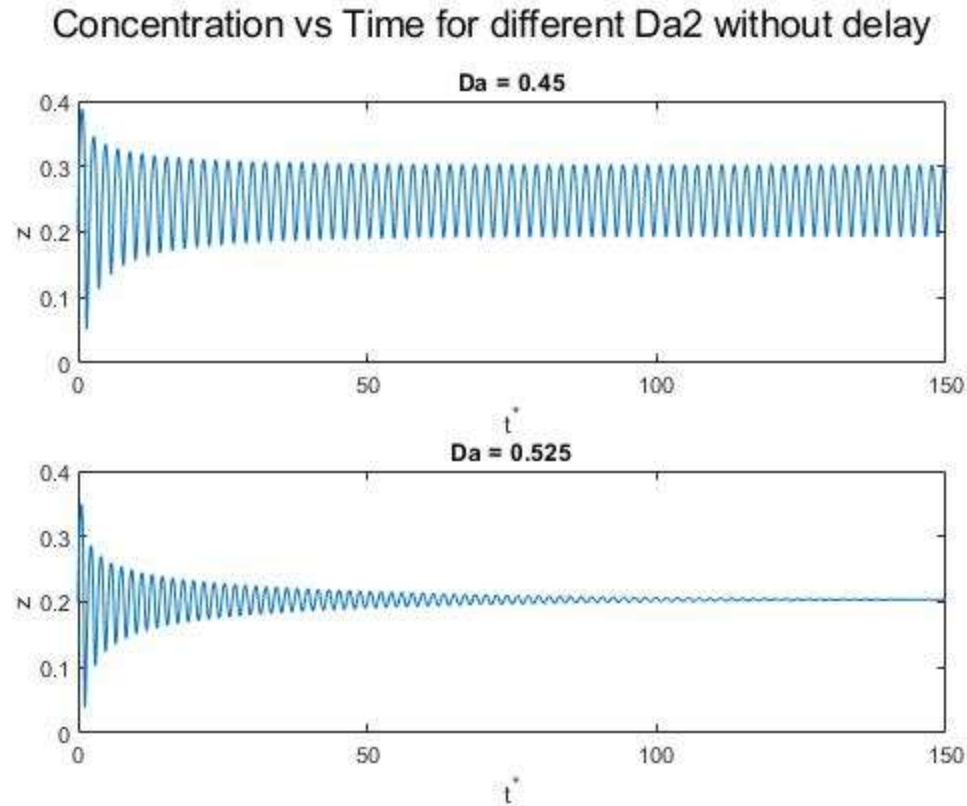
**Figure 5.** Bifurcation diagram of  $z$  versus  $Da_2$  for the fixed  $F$ ,  $M_R$  case with  $x_{af} = 0.9$ ,  $y_0 = 0.2$ ,  $x_0 = 0.8$ ,  $B_h = 14.0$ , and  $\beta_2 = 4.2$  (in the absence of delay).

- From the analysis of eigen values, we can see that for  $0.433 < Da_2 < 0.511$  the steady state is unstable and for other values, it is stable.
- This is also seen in the plot to the right where the dotted lines show the unstable SS and solid lines show stable SS.



# Results

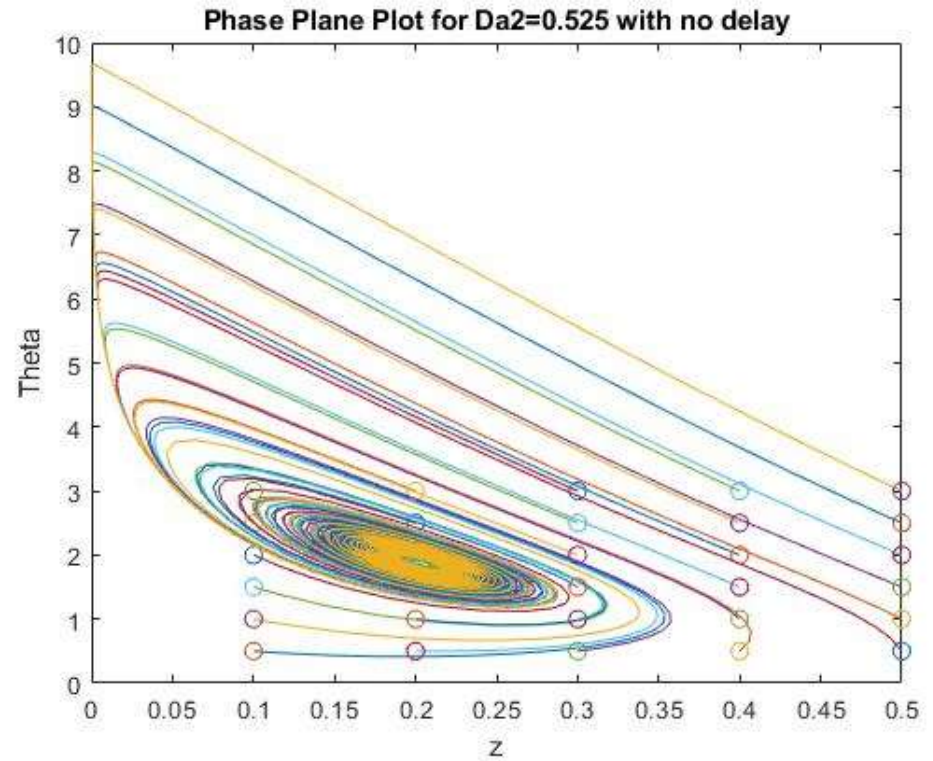
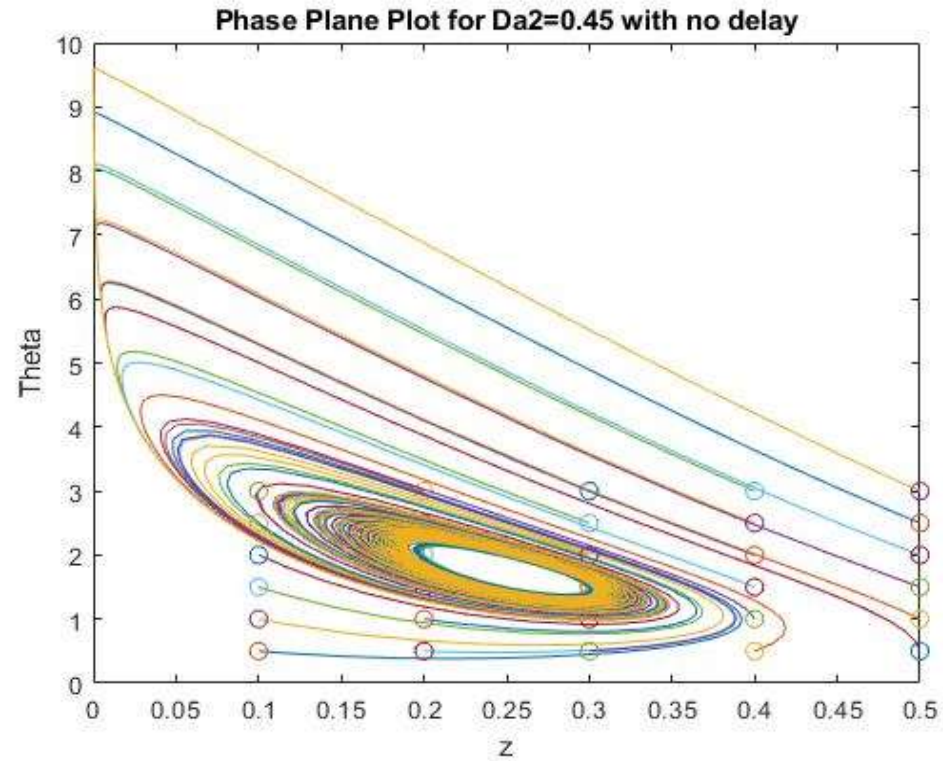
2. Case 2: Fixed  $M_R$ ,  $F$
- Dynamic analysis without delay :



- For  $Da_2 = 0.45$ , we know that there is no stable SS and hence we get the oscillations.
- For  $Da_2 = 0.525$ , there is a stable steady state which the system attains as shown.

# Results

2. Case 2: Fixed  $M_R$ ,  $F$ 
  - Phase Plane plots without delay:

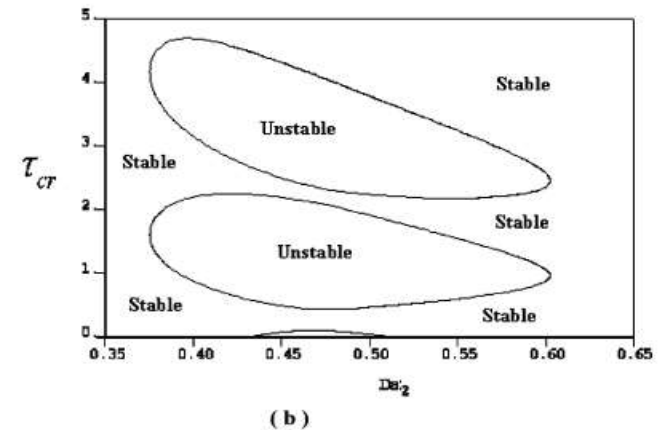
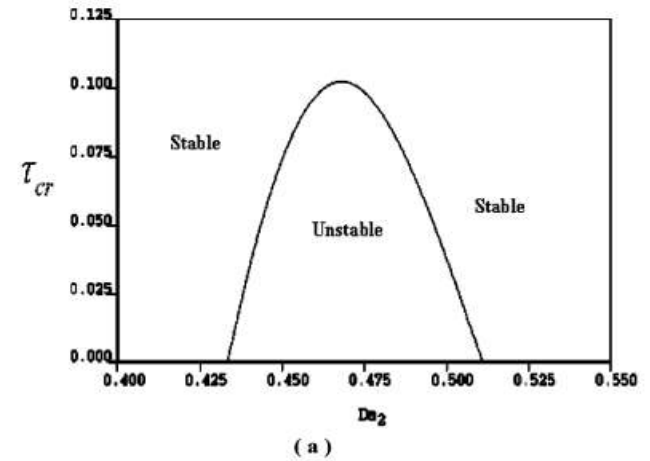
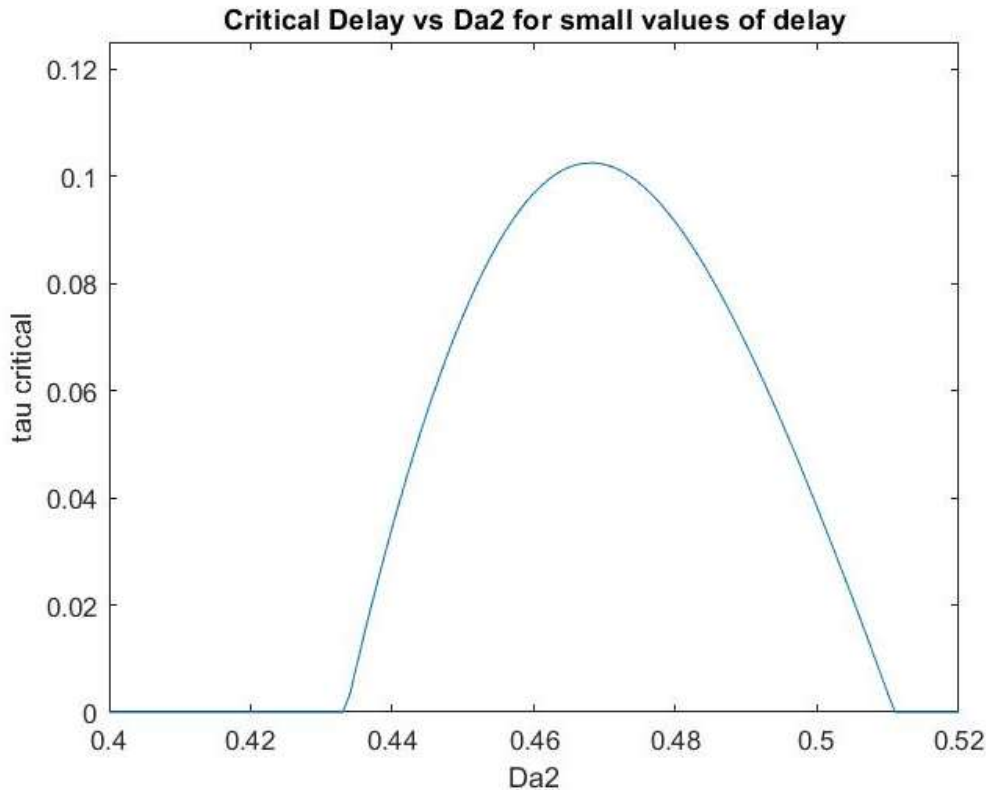


- This again shows the previous result that a SS is attained for  $Da_2= 0.525$ , not for  $Da_2 = 0.45$ .

# Results

## 2. Case 2: Fixed $M_R$ , $F$

- Bifurcation diagram with delay:



**Figure 6.** Dependency of dimensionless critical delay on  $Da_2$  for the fixed  $F$ ,  $M_R$  case with  $x_{af} = 0.9$ ,  $y_e = 0.2$ ,  $x_e = 0.8$ ,  $B_h = 14.0$ , and  $\beta_2 = 4.2$  for (a) small delay and (b) large delay.

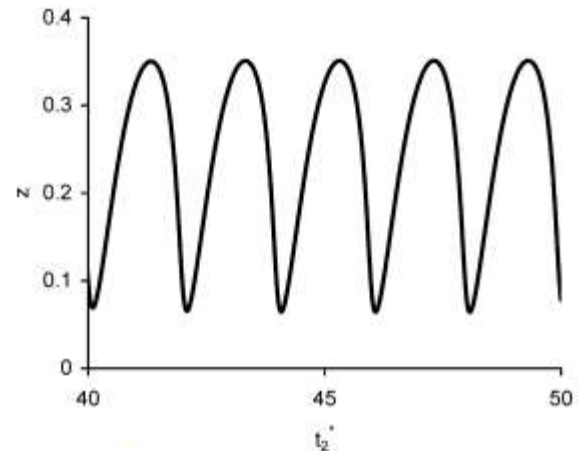
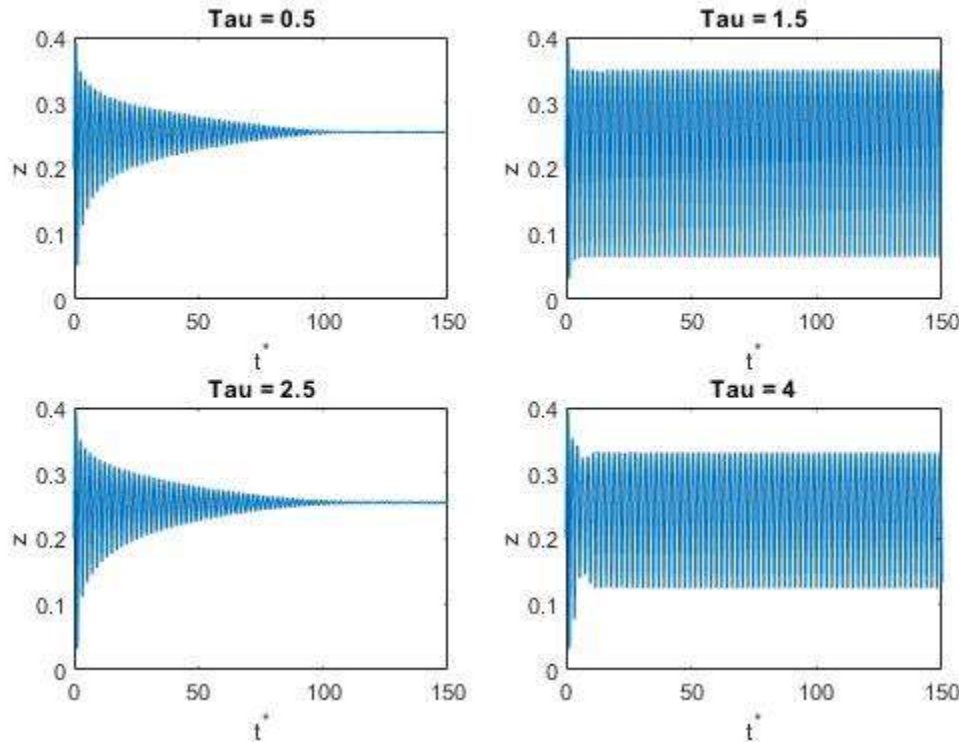
- Depending on the value of the delay, we have different regions of stability as shown in the plot on the right. Below the critical delay, the system is unstable and beyond the critical delay, it is stable.
- As the delay increases, for small values the width of the region of dynamic instability decreases.

# Results

## 2. Case 2: Fixed $M_R$ , $F$

- Dynamic analysis with delay:

Concentration vs Time for different  $Da_2$  with delay



**Figure 7.** Mole fraction of A versus dimensionless time plot for the fixed  $F$ ,  $M_R$  case showing sustained oscillations with  $x_{in} = 0.9$ ,  $y_e = 0.2$ ,  $x_e = 0.8$ ,  $B_0 = 14.0$ ,  $\beta_2 = 4.2$ ,  $Da_2 = 0.45$ , and  $\tau = 1.5$ .

- As shown, the stability is now dependent on the delay.
- The plots are shown for  $Da_2 = 0.45$  which lead to an unstable SS in the case without delay and the stability changes with delay which is also seen from the bifurcation diagram.
- For this case if we consider a isothermal CSTR, delay cannot introduce a new instability.

# Conclusion

- For a non isothermal reactor, the stability of a given steady state of a system changes with the transportation delay in the reactor effluent stream.
- For small values of the delay, the width of the region of unstable steady state decreases.
- In the case without delay, the instability is an effect of the non linearities in the Arrhenius temperature dependency of the rate.
- But in the case with delay, new zones of instability are introduced solely due to the effect of this lag .
- Hence, the delay has an significant effect on the stability and dynamics of this non isothermal reactor separator system which needs to be considered while designing process plants.

## Extension of the work

- A similar analysis can be done considering the effect of temperature coupling between the reactor and separator which we have avoided in this system with the use of two heat exchangers.
- The lag in the temperature has also been ignored which can be considered in the model.
- The lag in the transport of the recycle stream has been ignored here as the bottoms of the flash column gives an liquid assumed to be incompressible, through which changes in the flow rate are transmitted instantaneously. A new model extending this study can consider this effect as well.
- During the oscillations in the region of instabilities, it is observed that  $z$  falls below  $y_e$  which is not possible. This is because at  $z = y_e$ , the flow rate of the recycle stream is zero and for  $z < y_e$  the flow rate becomes negative which is not possible. For this a new model can be developed which imposes the condition that  $L = 0$  when  $z < y_e$ .