

# Blueprint for a Finite-Difference Eigensolver for the 2D Schrödinger Equation

## 1 Introduction

The finite-difference method approximates derivatives using differences between values at discrete points. Think of a smooth curve: instead of computing the exact slope at one point, you measure the change between two close points and divide by the distance between them. For example, for a function  $f(x)$ ,

$$f'(x_i) \approx \frac{f(x_{i+1}) - f(x_{i-1}))}{2\Delta x}.$$

For the second derivative, which appears in the Laplacian,

$$f''(x_i) \approx \frac{f(x_{i+1}) - 2f(x_i) + f(x_{i-1}))}{\Delta x^2}.$$

In our problem we solve the 2D time-independent Schrödinger equation

$$-\frac{1}{2}\nabla^2\psi(x, y) + V(x, y)\psi(x, y) = E\psi(x, y),$$

where  $\nabla^2$  is the Laplacian operator. We discretize the domain (i.e., replace the continuous region by a grid of points) and approximate the derivatives using finite differences.

## 2 Problem Setup and Discretization

We consider a rectangular domain of size  $L_x \times L_y$  centered at the origin. The wavefunction  $\psi$  is set to zero on the boundary (Dirichlet conditions). For simplicity, let us use:

$$L_x = 2, \quad L_y = 2, \quad N_x = 3, \quad N_y = 3,$$

which means we have 3 *interior* points in each direction. The grid spacing is

$$\Delta x = \frac{L_x}{N_x + 1} = \frac{2}{4} = 0.5, \quad \Delta y = 0.5.$$

The interior points for  $x$  are:

$$x = [-0.5, 0, 0.5],$$

and similarly for  $y$ .

## 3 Step-by-Step Code with Matrix Visualizations

Below, we explain each code block and show how the matrices look before and after any updates.

## 1. Grid Setup

Listing 1: Grid Setup

```
dx = Lx / (Nx + 1) # 2/4 = 0.5
dy = Ly / (Ny + 1) # 2/4 = 0.5

x = np.linspace(-Lx/2 + dx, Lx/2 - dx, Nx) # x = [-0.5, 0, 0.5]
y = np.linspace(-Ly/2 + dy, Ly/2 - dy, Ny) # y = [-0.5, 0, 0.5]

X, Y = np.meshgrid(x, y, indexing='ij')
```

**Visualization:** The 2D grid (meshgrid) is

$$X = \begin{pmatrix} -0.5 & -0.5 & -0.5 \\ 0 & 0 & 0 \\ 0.5 & 0.5 & 0.5 \end{pmatrix}, \quad Y = \begin{pmatrix} -0.5 & 0 & 0.5 \\ -0.5 & 0 & 0.5 \\ -0.5 & 0 & 0.5 \end{pmatrix}.$$

Here, each pair  $(X_{ij}, Y_{ij})$  represents a point in the domain where the wavefunction  $\psi$  will be approximated.

## 2. Potential Evaluation

Assume the potential is zero everywhere:

Listing 2: Potential Evaluation

```
V = Vfun(X, Y, *args) # For our case, V(x,y)=0.
```

**Before:** The function  $Vfun$  is applied to each grid point. **After:**

$$V = \begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix}.$$

This matrix represents the potential at each interior grid point.

## 3. Constructing the Finite-Difference Laplacian

The Laplacian is approximated using the 5-point stencil:

$$\nabla^2 \psi \approx \frac{\psi_{i+1,j} + \psi_{i-1,j} - 2\psi_{i,j}}{\Delta x^2} + \frac{\psi_{i,j+1} + \psi_{i,j-1} - 2\psi_{i,j}}{\Delta y^2}.$$

Since there are  $N_{\text{total}} = N_x \times N_y = 9$  unknowns, we will construct arrays that represent the diagonals of a  $9 \times 9$  matrix.

### a) Main Diagonal Initialization

Listing 3: Main Diagonal of Laplacian

```
N_total = Nx * Ny # 9 points
lap_main = (-2.0/dx**2 - 2.0/dy**2) * np.ones(N_total)
```

**Explanation:** With  $\Delta x = \Delta y = 0.5$ , we have:

$$\frac{1}{dx^2} = \frac{1}{0.25} = 4.$$

So,

$$-2 \left( \frac{1}{dx^2} \right) - 2 \left( \frac{1}{dy^2} \right) = -2(4) - 2(4) = -8 - 8 = -16.$$

**Before:** The call `np.ones(9)` creates:

$$[1, 1, 1, 1, 1, 1, 1, 1, 1].$$

**After:** Multiplying by  $-16$  gives:

$$\text{lap\_main} = [-16, -16, -16, -16, -16, -16, -16, -16, -16].$$

Each element represents the contribution from the center point in the finite-difference formula.

## b) Off-Diagonals in the $y$ -Direction

Listing 4: Off-Diagonals in  $y$ -Direction

```
lap_off_y = 1.0/dy**2 * np.ones(N_total - 1)
```

**Before Update:** `np.ones(8)` creates:

$$[1, 1, 1, 1, 1, 1, 1, 1].$$

Multiplying by  $1/dy^2 = 4$  gives initially:

$$\text{lap\_off\_y} = [4, 4, 4, 4, 4, 4, 4, 4].$$

These values represent the connection between a grid point and its immediate neighbor in the  $y$ -direction.

**Removing Wrap-Around Connections:** We must not connect the last point in one row to the first point in the next row. This is done by:

Listing 5: Removing Wrap-Around in  $y$ -Direction

```
for i in range(1, Nx):
    idx = i*Ny - 1
    lap_off_y[idx] = 0.0
```

**Updates:**

- For  $i = 1$ :  $\text{idx} = 1 \times 3 - 1 = 2$ . Set `lap_off_y[2] = 0` so that index 2 is not connected to index 3.
- For  $i = 2$ :  $\text{idx} = 2 \times 3 - 1 = 5$ . Set `lap_off_y[5] = 0`.

**After Update:**

$$\text{lap\_off\_y} = [4, 4, 0, 4, 4, 0, 4, 4].$$

Each element here connects adjacent points in the  $y$ -direction, except at the row boundaries where we set the connection to zero.

### c) Off-Diagonals in the $x$ -Direction

Listing 6: Off-Diagonals in  $x$ -Direction

```
lap_off_x = 1.0/dx**2 * np.ones(N_total - Ny)
```

**Before and After:** `np.ones(6)` creates:

`[1, 1, 1, 1, 1, 1]`.

Multiplying by 4 gives:

`lap_off_x = [4, 4, 4, 4, 4, 4]`.

These represent the connection between points in the  $x$ -direction (neighbors above and below in the grid).

### d) Assembling the Laplacian Matrix

We combine the diagonals into a sparse matrix:

Listing 7: Assembling the Laplacian

```
diagonals = [lap_main, lap_off_y, lap_off_y, lap_off_x, lap_off_x]
offsets    = [0, 1, -1, Ny, -Ny]
L = sp.diags(diagonals, offsets, shape=(N_total, N_total), format='csr')
```

**Interpretation:**

- The **main diagonal** (offset 0) has the values from `lap_main` (all  $-16$ ).
- The two **off-diagonals** with offsets  $+1$  and  $-1$  use `lap_off_y` to connect neighboring points in the  $y$ -direction.
- The off-diagonals with offsets  $+Ny$  and  $-Ny$  (i.e.,  $+3$  and  $-3$ ) use `lap_off_x` to connect points in the  $x$ -direction.

**Dense Form of  $L$ :** For clarity, the full  $9 \times 9$  matrix (if converted to a dense matrix) would look like

$$L = \begin{pmatrix} -16 & 4 & 0 & 4 & 0 & 0 & 0 & 0 & 0 \\ 4 & -16 & 4 & 0 & 4 & 0 & 0 & 0 & 0 \\ 0 & 4 & -16 & 0 & 0 & 4 & 0 & 0 & 0 \\ 4 & 0 & 0 & -16 & 4 & 0 & 4 & 0 & 0 \\ 0 & 4 & 0 & 4 & -16 & 4 & 0 & 4 & 0 \\ 0 & 0 & 4 & 0 & 4 & -16 & 0 & 0 & 4 \\ 0 & 0 & 0 & 4 & 0 & 0 & -16 & 4 & 0 \\ 0 & 0 & 0 & 0 & 4 & 0 & 4 & -16 & 4 \\ 0 & 0 & 0 & 0 & 0 & 4 & 0 & 4 & -16 \end{pmatrix}.$$

Notice that the off-diagonal zeros correspond to the blocked connections (wrap-around issues).

## 4. Hamiltonian Assembly

The Hamiltonian is given by

$$H = -\frac{1}{2}L + V,$$

with  $V$  being added as a diagonal matrix.

Listing 8: Hamiltonian Assembly

```
H = -0.5 * L + sp.diags(V.ravel(order='C'), 0, format='csr')
```

**Before:**  $L$  is as shown above, and  $V$  is a  $3 \times 3$  matrix of zeros (flattened to a 9-element vector). **After:** Every element of  $L$  is multiplied by  $-0.5$ . For example:

$$-0.5 \times (-16) = 8 \quad (\text{main diagonal}), \quad -0.5 \times 4 = -2 \quad (\text{off-diagonals}).$$

Thus, the dense form of  $H$  becomes

$$H = \begin{pmatrix} 8 & -2 & 0 & -2 & 0 & 0 & 0 & 0 & 0 \\ -2 & 8 & -2 & 0 & -2 & 0 & 0 & 0 & 0 \\ 0 & -2 & 8 & 0 & 0 & -2 & 0 & 0 & 0 \\ -2 & 0 & 0 & 8 & -2 & 0 & -2 & 0 & 0 \\ 0 & -2 & 0 & -2 & 8 & -2 & 0 & -2 & 0 \\ 0 & 0 & -2 & 0 & -2 & 8 & 0 & 0 & -2 \\ 0 & 0 & 0 & -2 & 0 & 0 & 8 & -2 & 0 \\ 0 & 0 & 0 & 0 & -2 & 0 & -2 & 8 & -2 \\ 0 & 0 & 0 & 0 & 0 & -2 & 0 & -2 & 8 \end{pmatrix}.$$

Each element now represents the combined effect of the kinetic energy operator (from the Laplacian) and the potential (which is zero in this example).

## 5. Solving the Eigenvalue Problem and Normalization

The eigensolver finds solutions to

$$H\psi = E\psi.$$

Listing 9: Solving the Eigenvalue Problem

```
E, psi_flat = spla.eigsh(H, k=neigs, sigma=E0, which='LM')
```

**Explanation:** The vector  $\psi_{\text{flat}}$  contains eigenfunctions in a flattened (one-dimensional) form. Each eigenvector has 9 components corresponding to our 9 grid points.

Then, each eigenvector is reshaped into a  $3 \times 3$  matrix and normalized:

Listing 10: Reshaping and Normalizing

```
psi = np.empty((neigs, Nx, Ny), dtype=complex)
for n in range(neigs):
    psi_n = psi_flat[:, n].reshape((Nx, Ny), order='C')
    norm = np.sqrt(np.sum(np.abs(psi_n)**2) * dx * dy)
    psi[n, :, :] = psi_n / norm
```

**Before Normalization:** An eigenvector might look like

$$\psi_{\text{flat}}^{(n)} = \begin{pmatrix} a_0 \\ a_1 \\ \vdots \\ a_8 \end{pmatrix}.$$

**After Reshaping:** It becomes a  $3 \times 3$  matrix:

$$\psi^{(n)} = \begin{pmatrix} a_0 & a_1 & a_2 \\ a_3 & a_4 & a_5 \\ a_6 & a_7 & a_8 \end{pmatrix}.$$

**After Normalization:** The entries are scaled so that

$$\sum_{i,j} |\psi_{ij}^{(n)}|^2 \Delta x \Delta y = 1.$$

## 4 Explanation of `lap_off_x` and `lap_off_y` Arrays

Below we provide a detailed, visual explanation of the off-diagonals for the  $x$  and  $y$  directions.

### Visualizing the Grid

Consider the 3x3 grid with flattened indices:

$$\begin{array}{lll} 0 : (-0.5, -0.5) & 1 : (-0.5, 0) & 2 : (-0.5, 0.5) \\ 3 : (0, -0.5) & 4 : (0, 0) & 5 : (0, 0.5) \\ 6 : (0.5, -0.5) & 7 : (0.5, 0) & 8 : (0.5, 0.5) \end{array}$$

### Horizontal Connections (`lap_off_y`)

- **Purpose:** Connect each grid point to its left and right neighbors.
- **Coefficient:** Each connection gets a value  $1/\Delta y^2 = 4$  (for  $\Delta y = 0.5$ ).

**Connections:**

$$\begin{array}{ll} 0 \leftrightarrow 1, & 1 \leftrightarrow 2, & \text{(Row 0)} \\ 3 \leftrightarrow 4, & 4 \leftrightarrow 5, & \text{(Row 1)} \\ 6 \leftrightarrow 7, & 7 \leftrightarrow 8, & \text{(Row 2)} \end{array}$$

**The For Loop:** A loop sets the connection between the last index of a row and the first index of the next row to zero:

- For row 0: Index 2 (end) should not connect to index 3 (start of row 1).
- For row 1: Index 5 (end) should not connect to index 6 (start of row 2).

Thus, the `lap_off_y` array becomes:

$$[4, 4, 0, 4, 4, 0, 4, 4].$$

### Vertical Connections (`lap_off_x`)

- **Purpose:** Connect each grid point to its above and below neighbors.
- **Coefficient:** Each connection gets a value  $1/\Delta x^2 = 4$  (for  $\Delta x = 0.5$ ).

**Connections:**

$$\begin{array}{lll} 0 \leftrightarrow 3, & 1 \leftrightarrow 4, & 2 \leftrightarrow 5, \\ 3 \leftrightarrow 6, & 4 \leftrightarrow 7, & 5 \leftrightarrow 8. \end{array}$$

In the flattened array, a jump of  $N_y$  (here, 3) corresponds to moving vertically in the grid.

## 5 Conclusion and Final Words

This document provides a comprehensive, step-by-step explanation of the finite-difference method for solving the 2D Schrödinger equation. We illustrated the grid, showed how the Laplacian is built using finite differences, and explained the off-diagonals (`lap_off_y` and `lap_off_x`) with visual aids. The use of a loop in setting up `lap_off_y` ensures that the physical structure of the grid is respected by preventing connections that cross row boundaries.