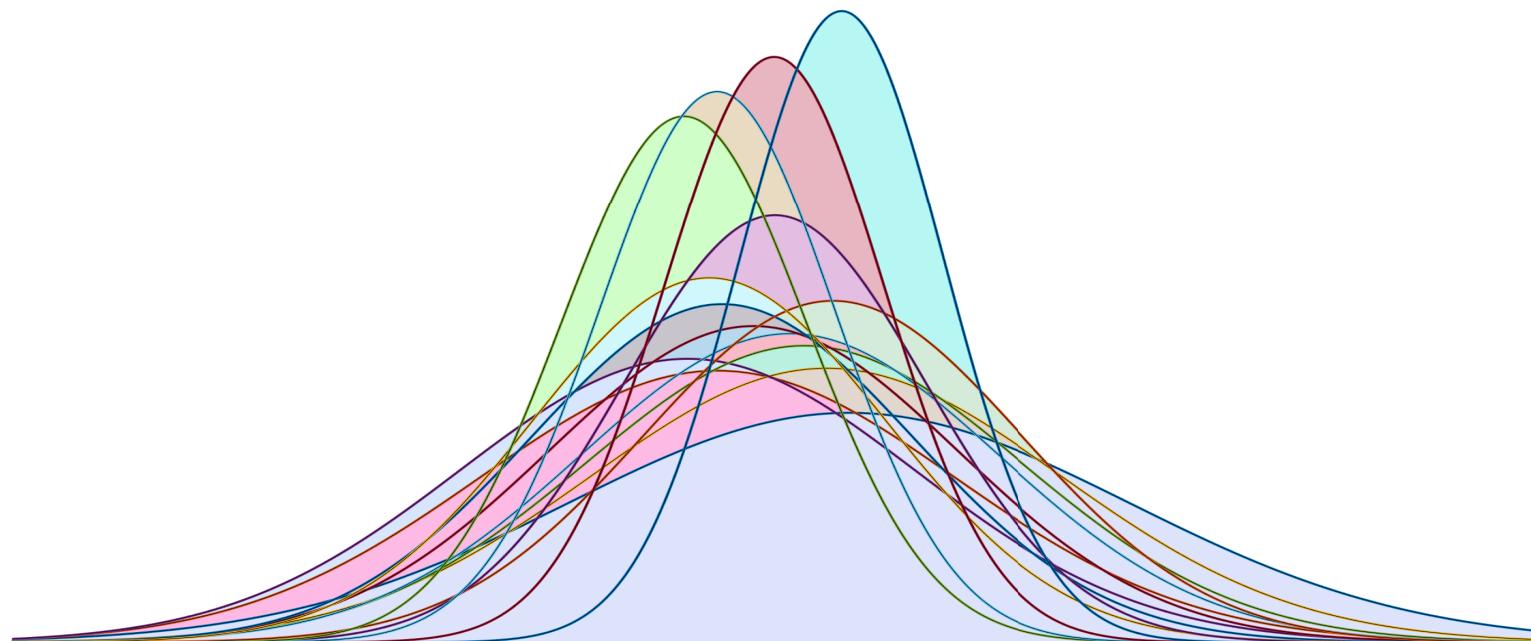
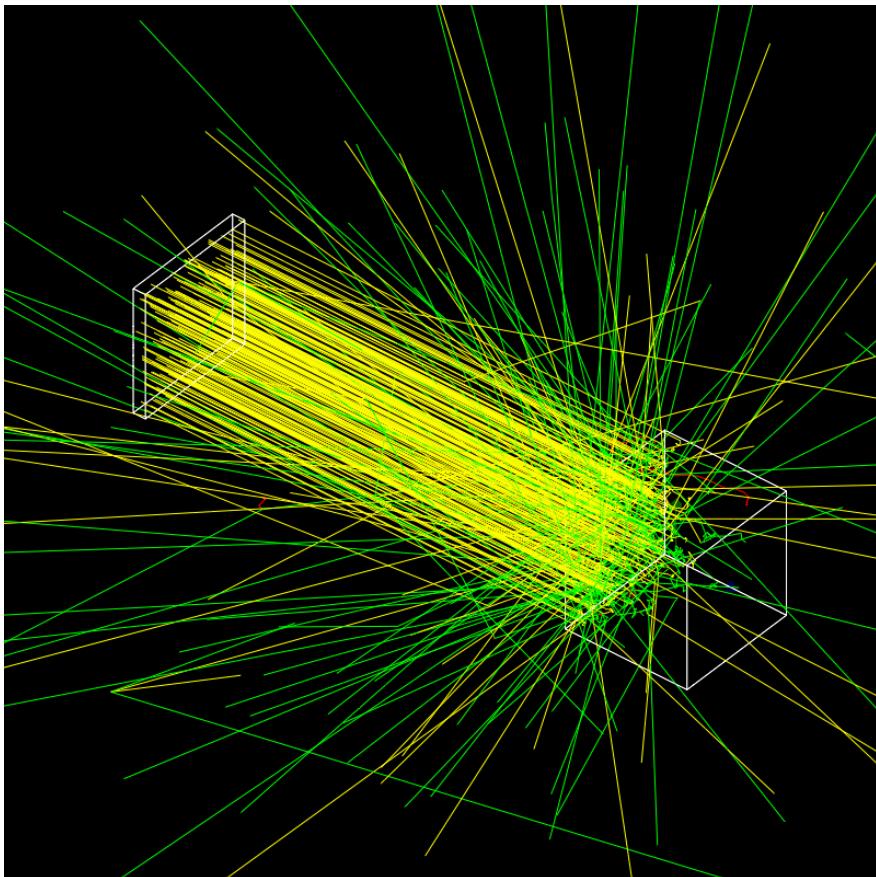


Towards an interval particle transport monte carlo method



Ander Gray
Institute for Risk and Uncertainty
University of Liverpool

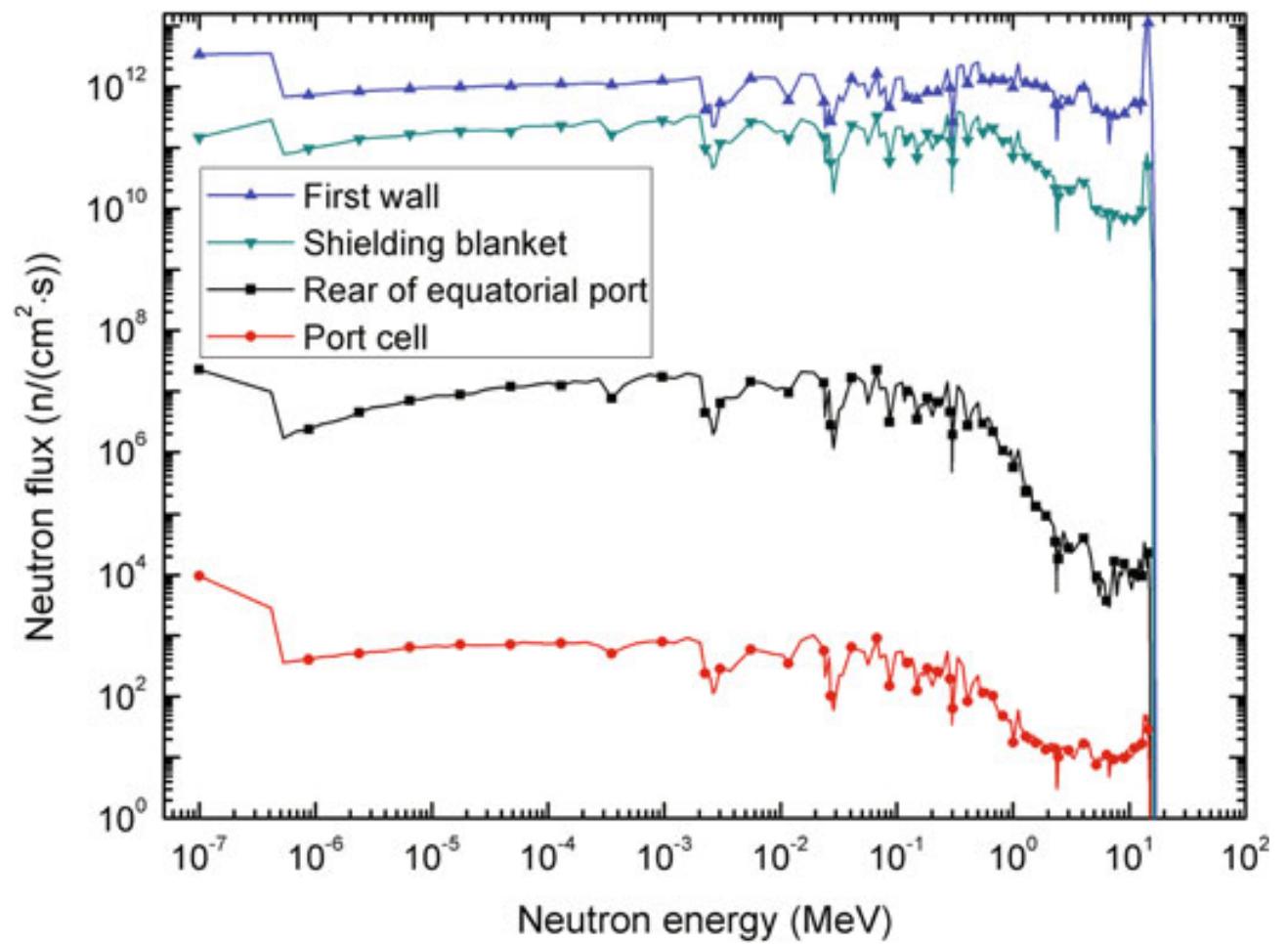
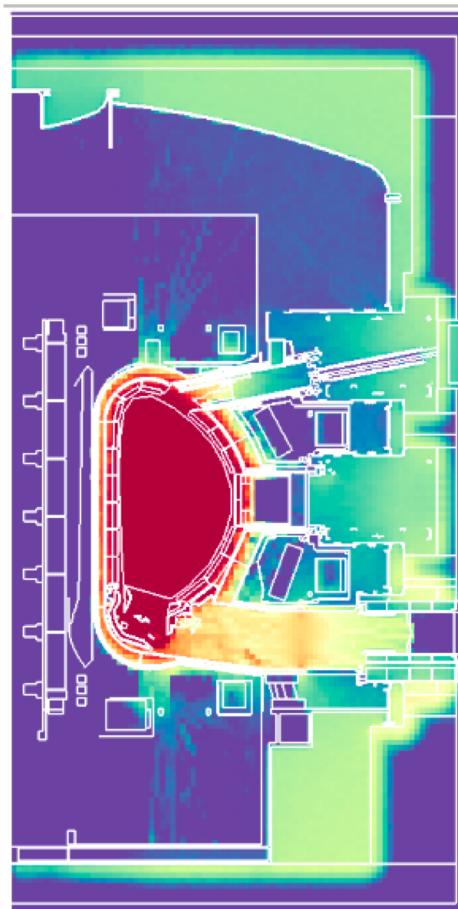
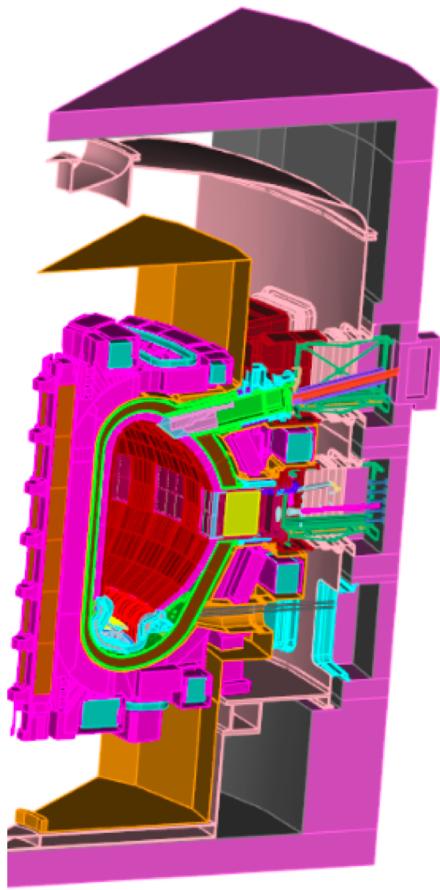
Neutron transport monte carlo



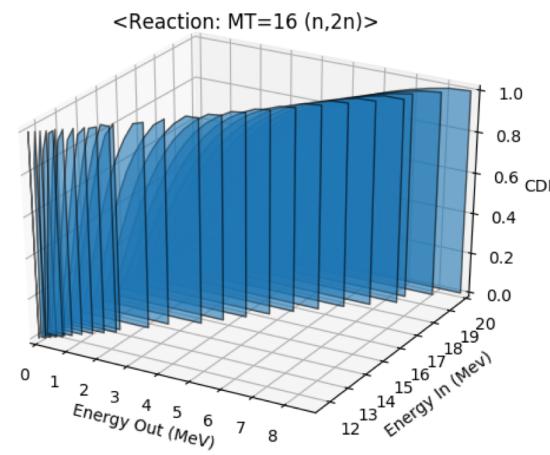
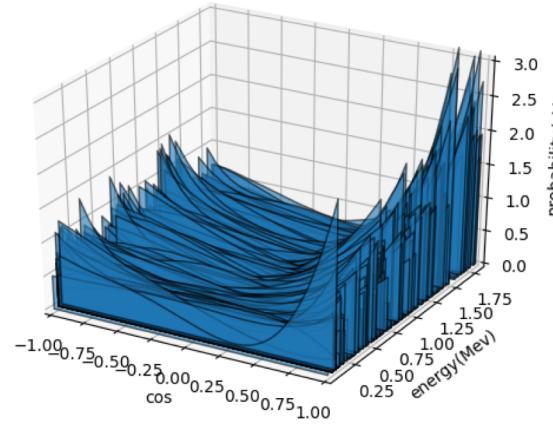
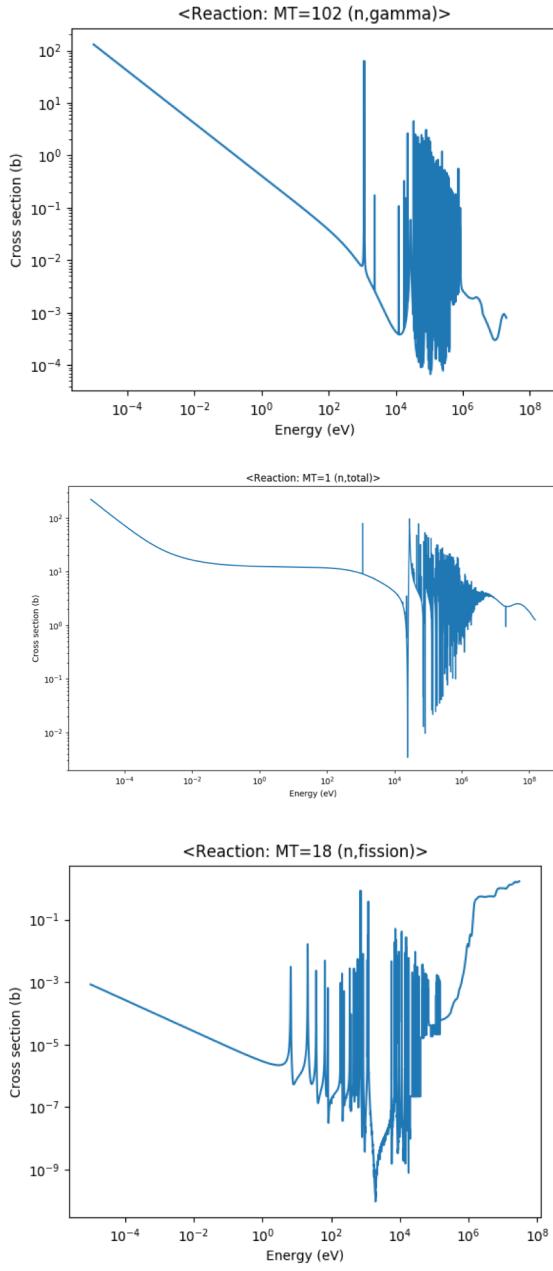
Inputs:

- Particle source
- Geometry
- A definition of particle/matter interaction rules (Nuclear Data)

Neutron transport monte carlo



Nuclear Data



- For Tokamaks you would need to consider 100s of different nuclides
- Each nuclide has ~ 30 different reaction channels (events)
- All of which are uncertain

~ 10^{12} histories needed for reasonable monte carlo error

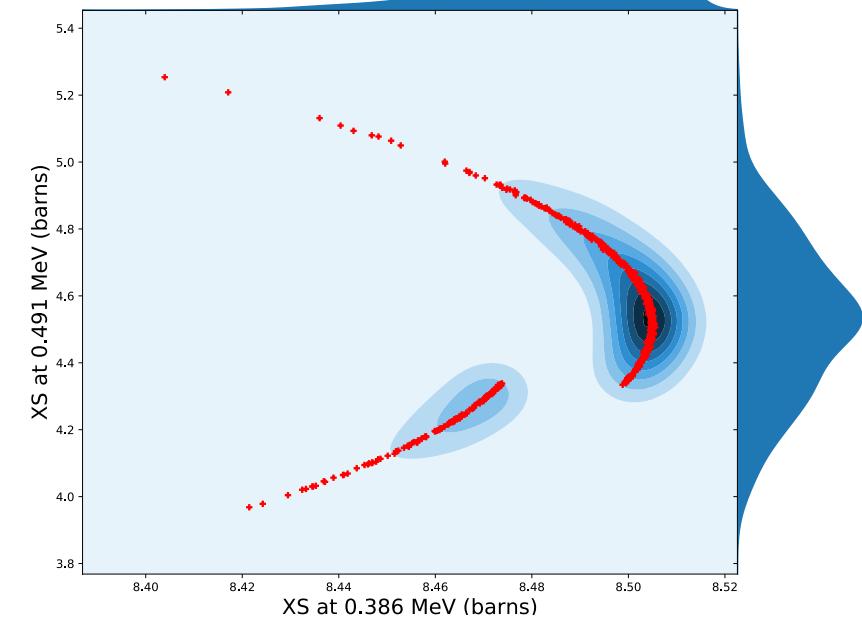
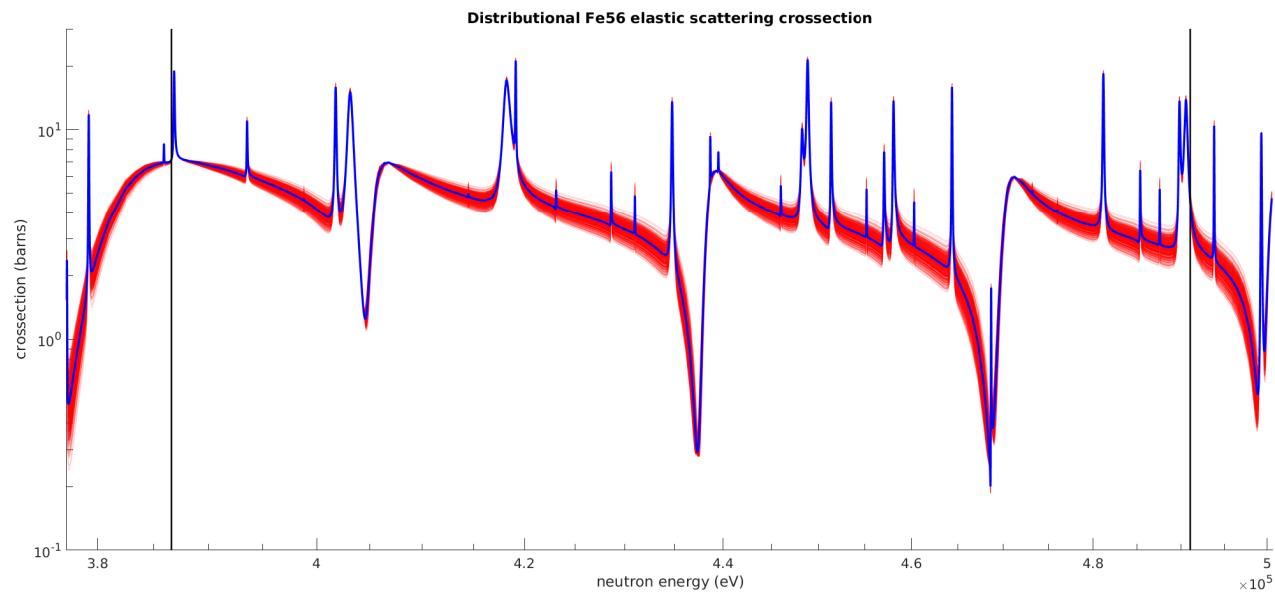
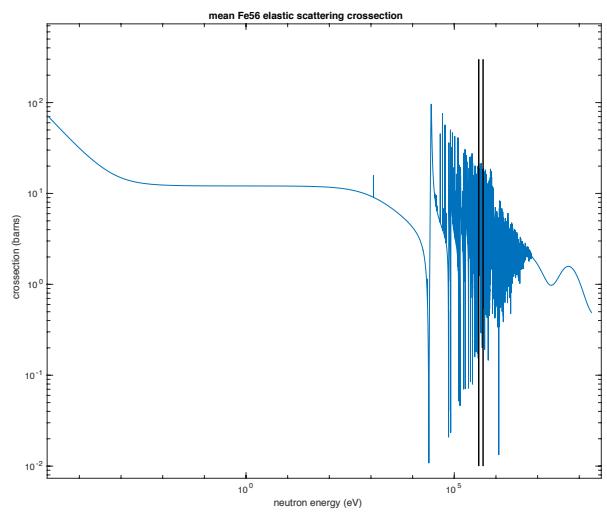
We can simulate ~1000 histories per s

On 1000 cpus

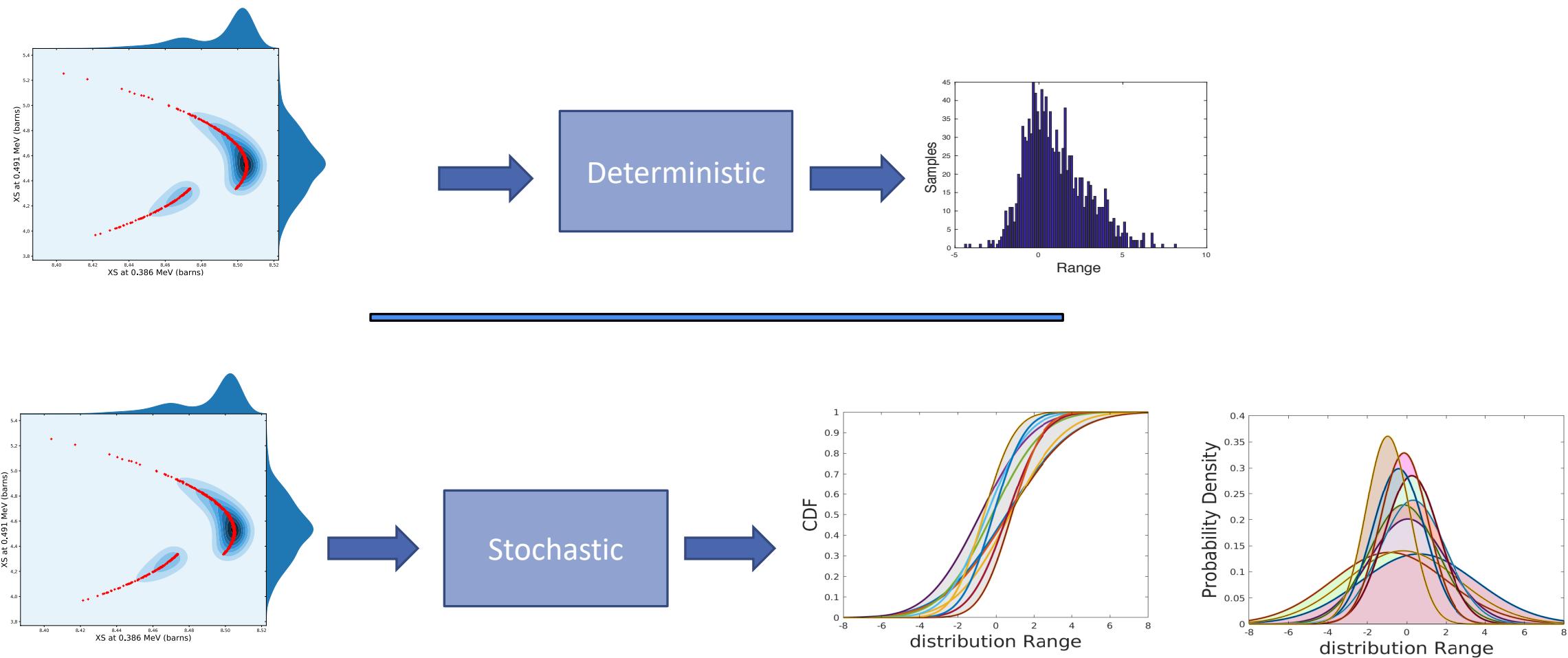


~1.5 weeks for one simulation!

Uncertainty in cross sections



Monte Carlo uncertainty propagation

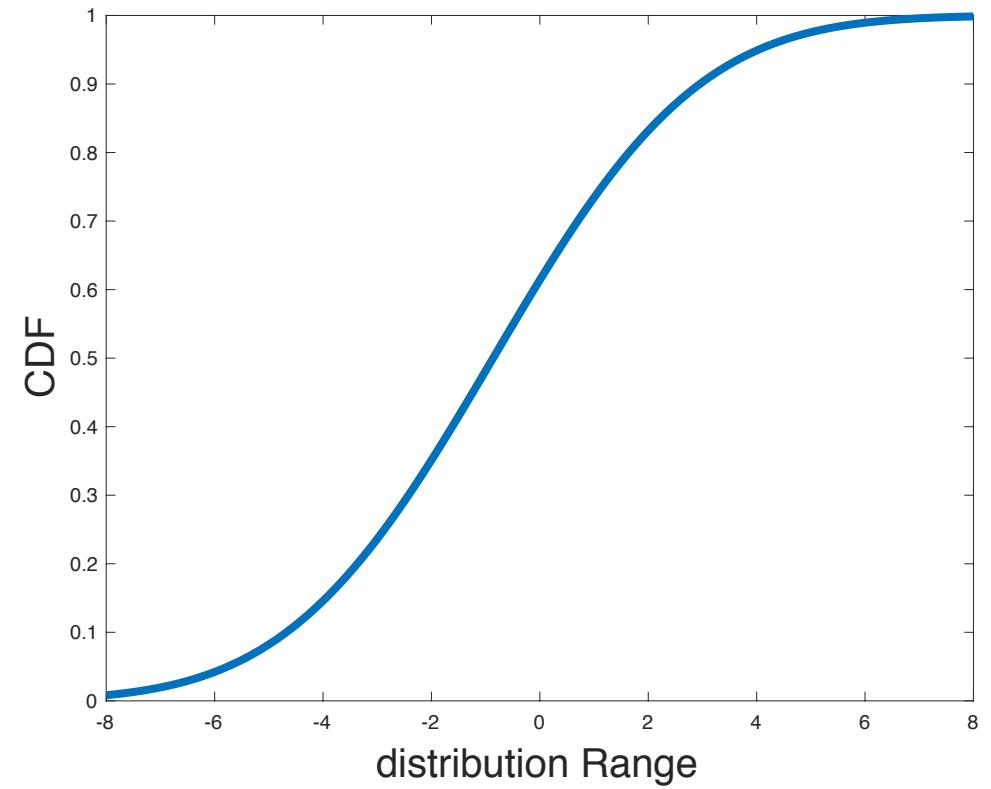


2nd order distributions

Consider:

$$X \sim N(\mu, \sigma)$$

μ, σ are scalar

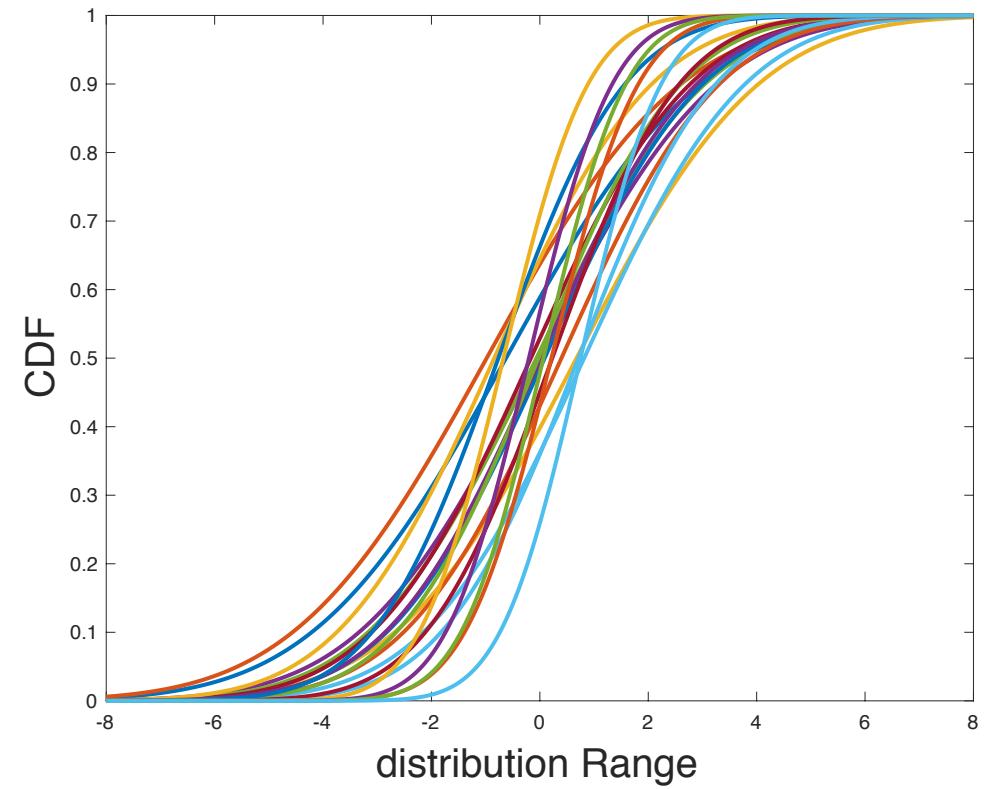


2nd order distributions

Consider:

$$X \sim N(\mu, \sigma)$$

μ, σ have a joint distribution

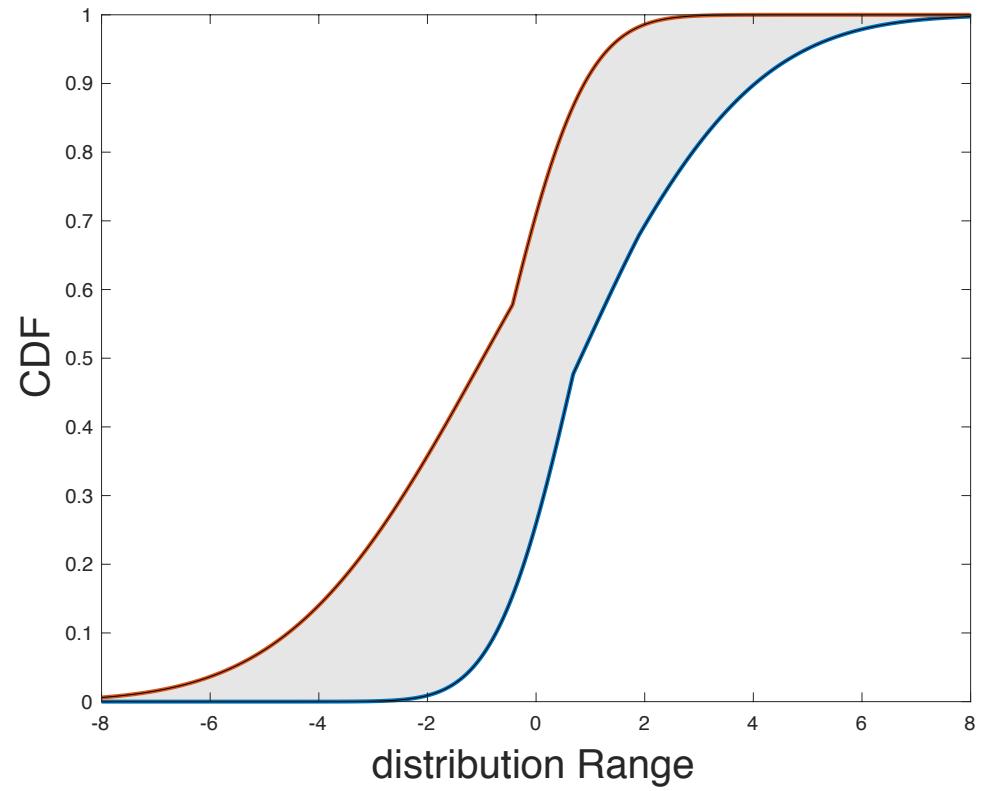


2nd order distributions

Consider:

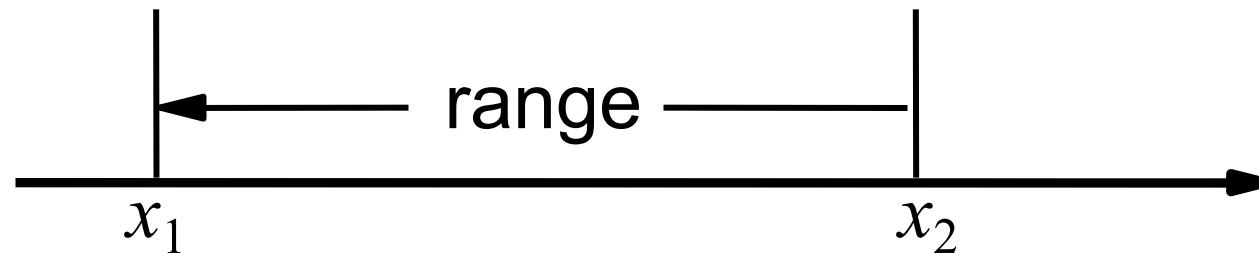
$$X \sim N(\mu, \sigma)$$

μ, σ are intervals



The interval

$$x = [x_1, x_2] = [\underline{x}, \bar{x}]$$



- Used when upper and lower bounds are known reliably
- Sometimes absolutely sure, sometimes just a judgment

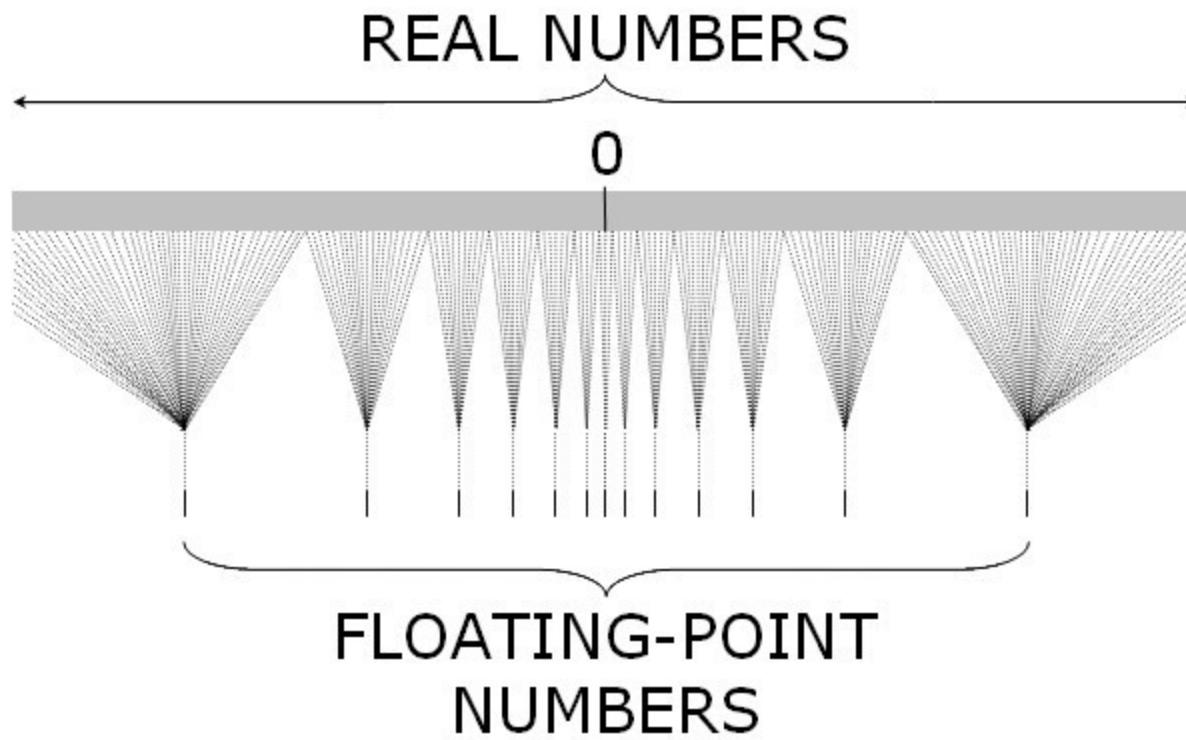
Why intervals are a good idea

- Intervals are easy to understand
- Calculations with them are easy
- They work with even the crappiest data
- Often sufficient for a decision
- Quicker than Monte Carlo

Intervals are rigorous

- The computations are guaranteed to enclose the true results (so long as the inputs do)
- You can still be wrong, but the *method* won't be the reason if you are
- Outward directed rounding

Intervals are rigorous



- Outward directed rounding

Interval arithmetic

$$x + y = [x_1 + y_1, x_2 + y_2]$$

$$x - y = [x_1 - y_2, x_2 - y_1]$$

$$x \times y = [\min(x_1y_1, x_1y_2, x_2y_1, x_2y_2), \max(x_1y_1, x_1y_2, x_2y_1, x_2y_2)]$$

$$x \div y = x \times [1/y_2, 1/y_1], \text{ if } y \text{ doesn't include zero}$$

Interval arithmetic

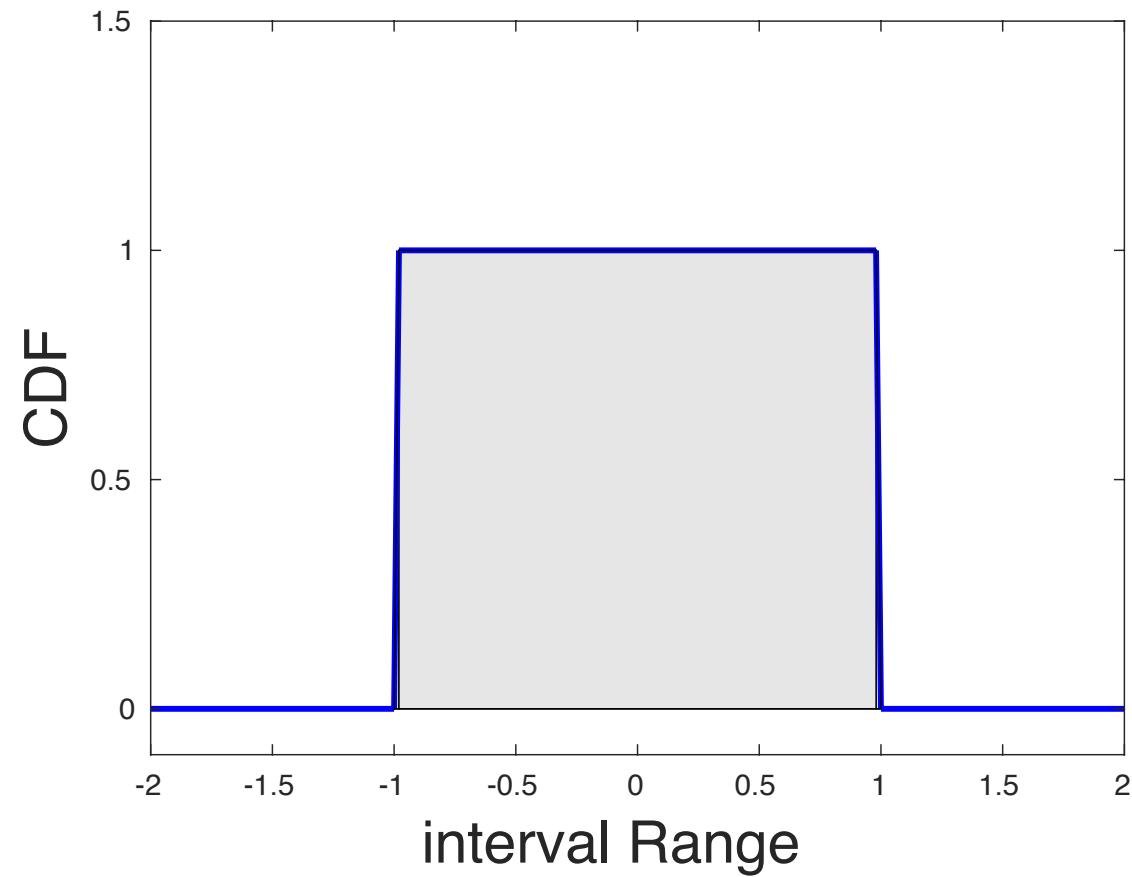
$$\min(x, y) = [\min(x_1, y_1), \min(x_2, y_2)]$$

$$\max(x, y) = [\max(x_1, y_1), \max(x_2, y_2)]$$

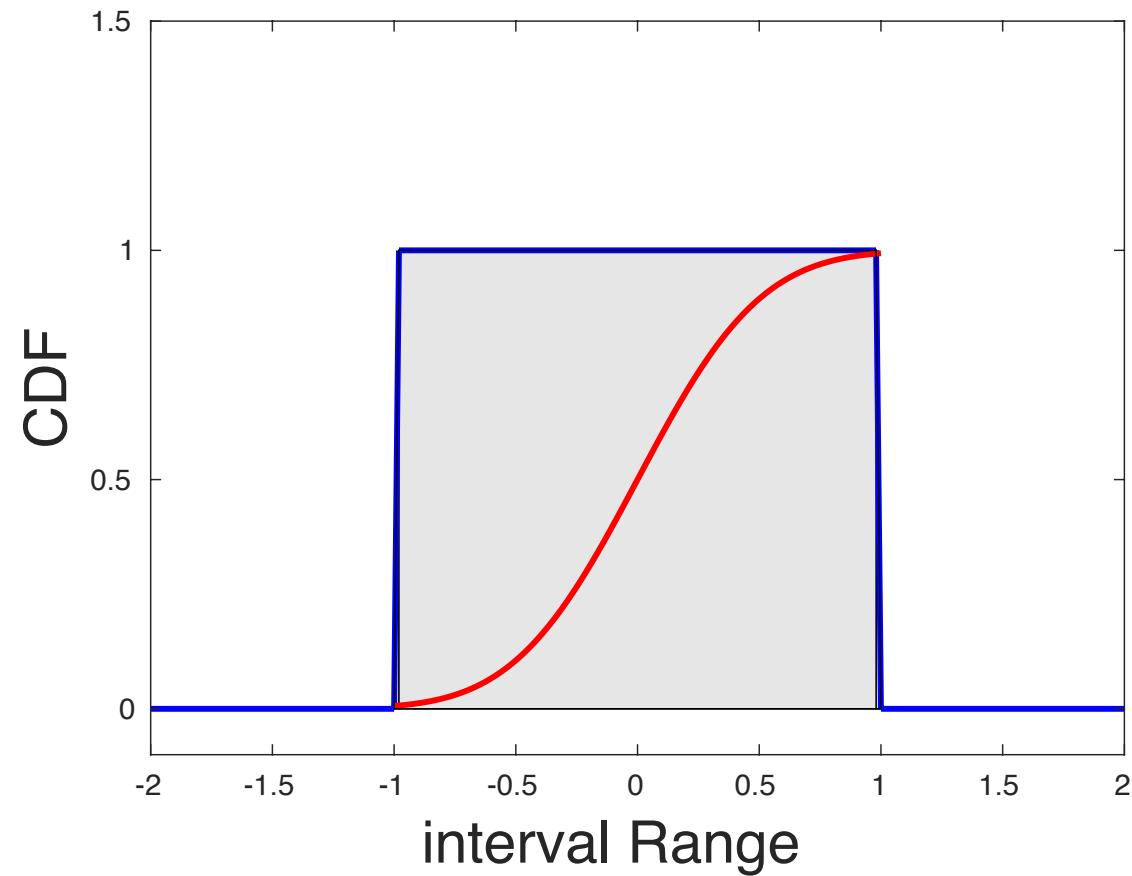
$$\text{env}(x, y) = [\min(x_1, y_1), \max(x_2, y_2)]$$

$$\text{intersect}(x, y) = [\max(x_1, y_1), \min(x_2, y_2)], \text{ if } x \text{ and } y \text{ overlap}$$

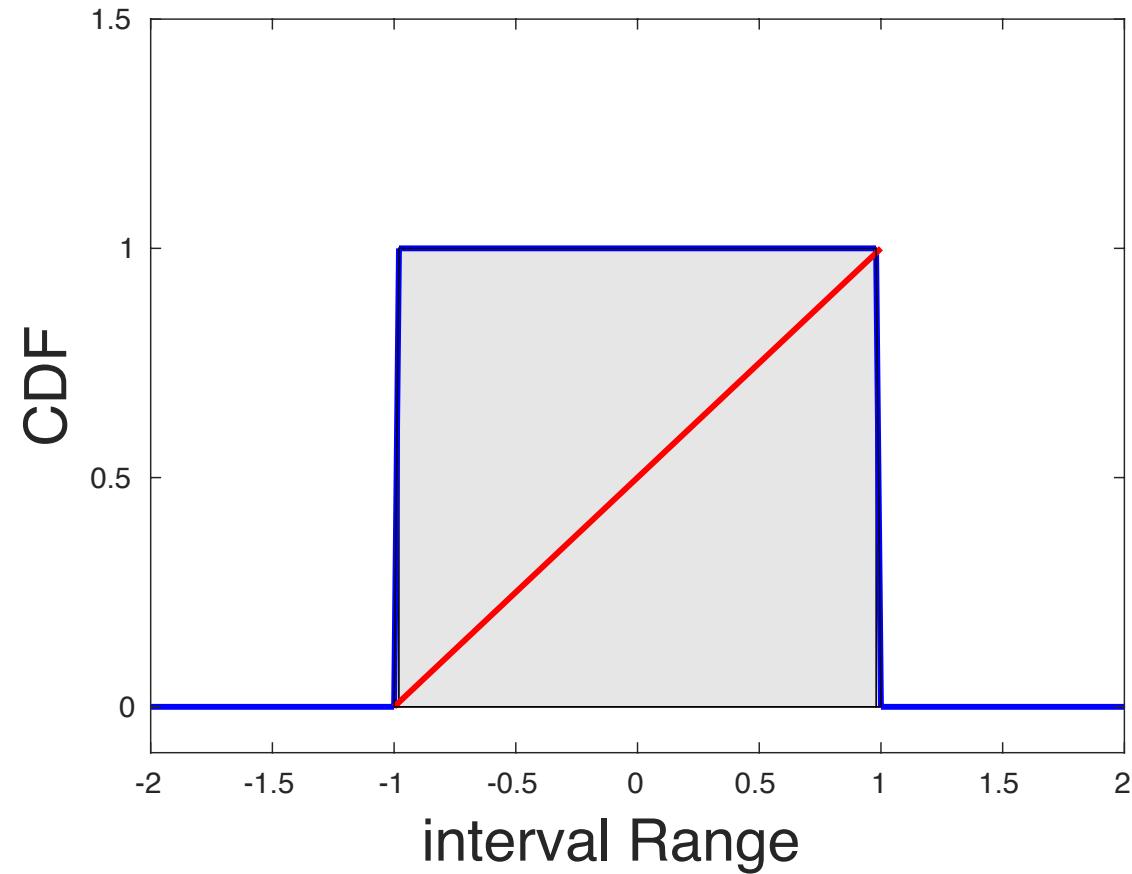
Intervals as a model for uncertainty



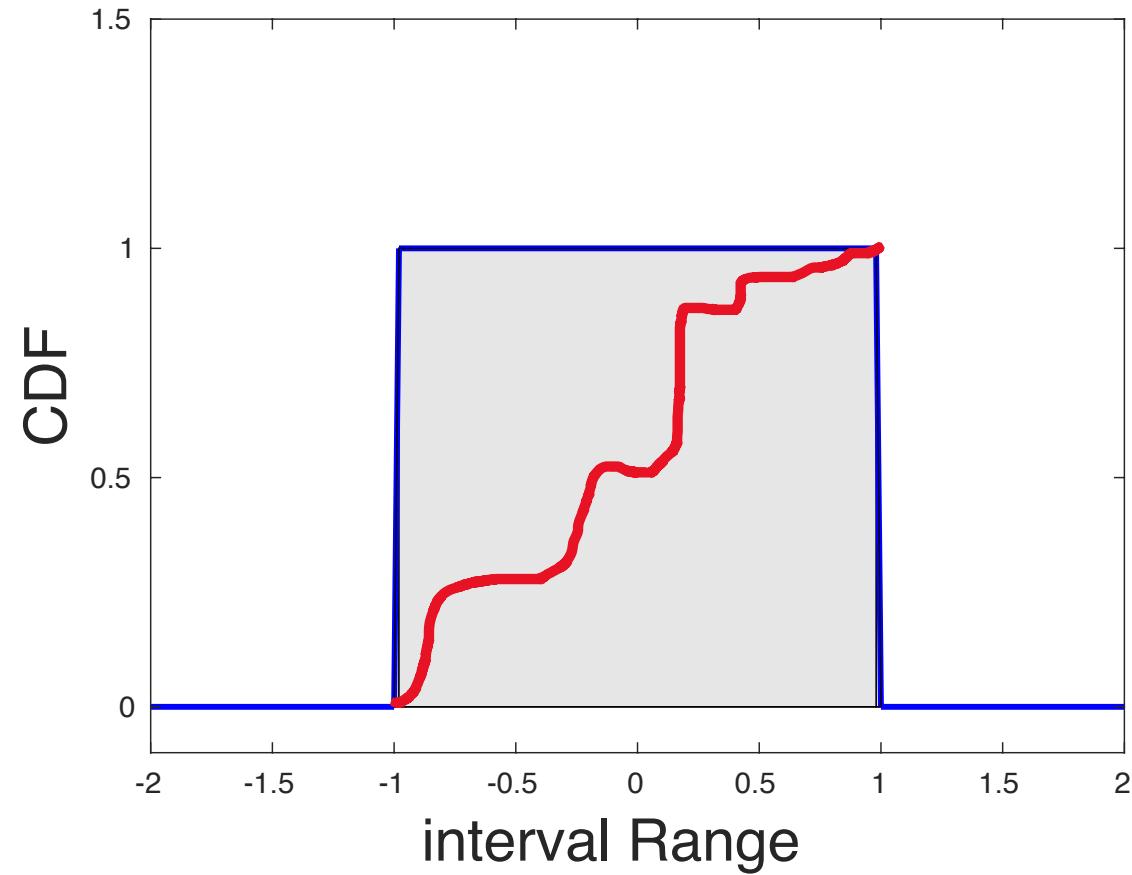
Intervals as a model for uncertainty



Intervals as a model for uncertainty

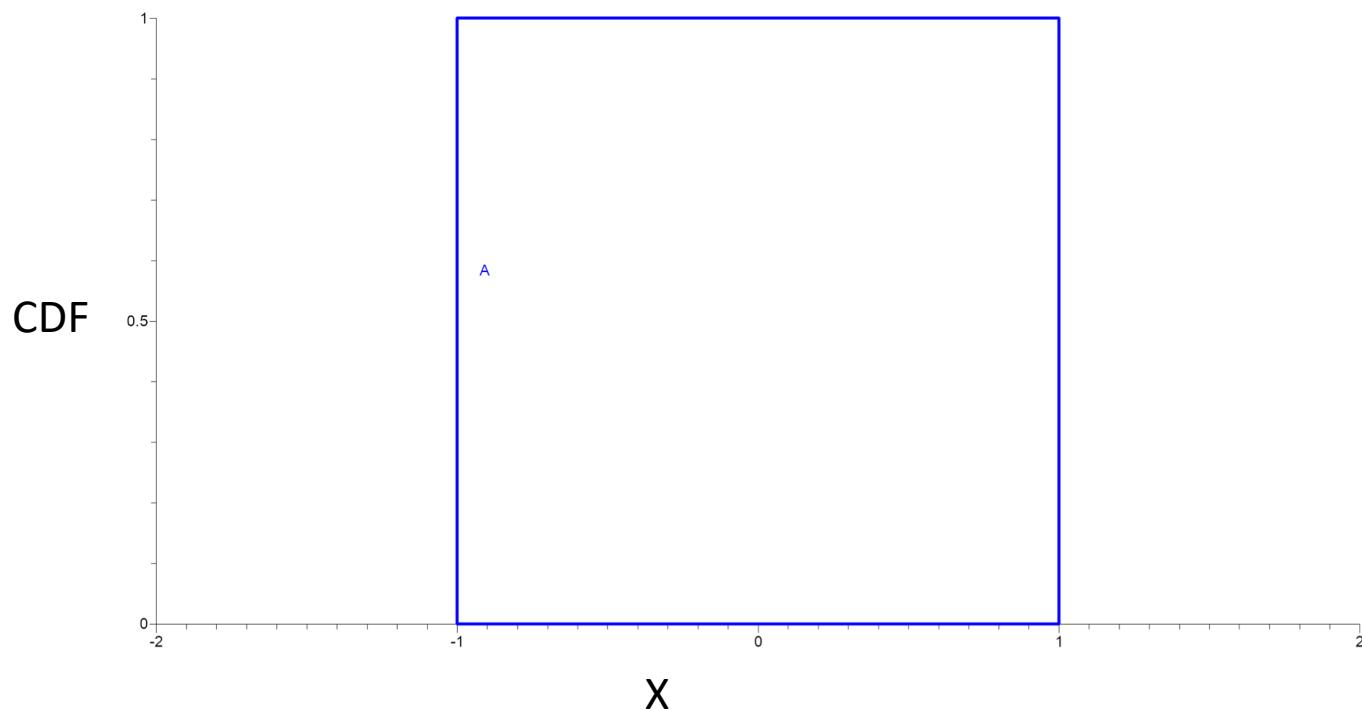


Intervals as a model for uncertainty



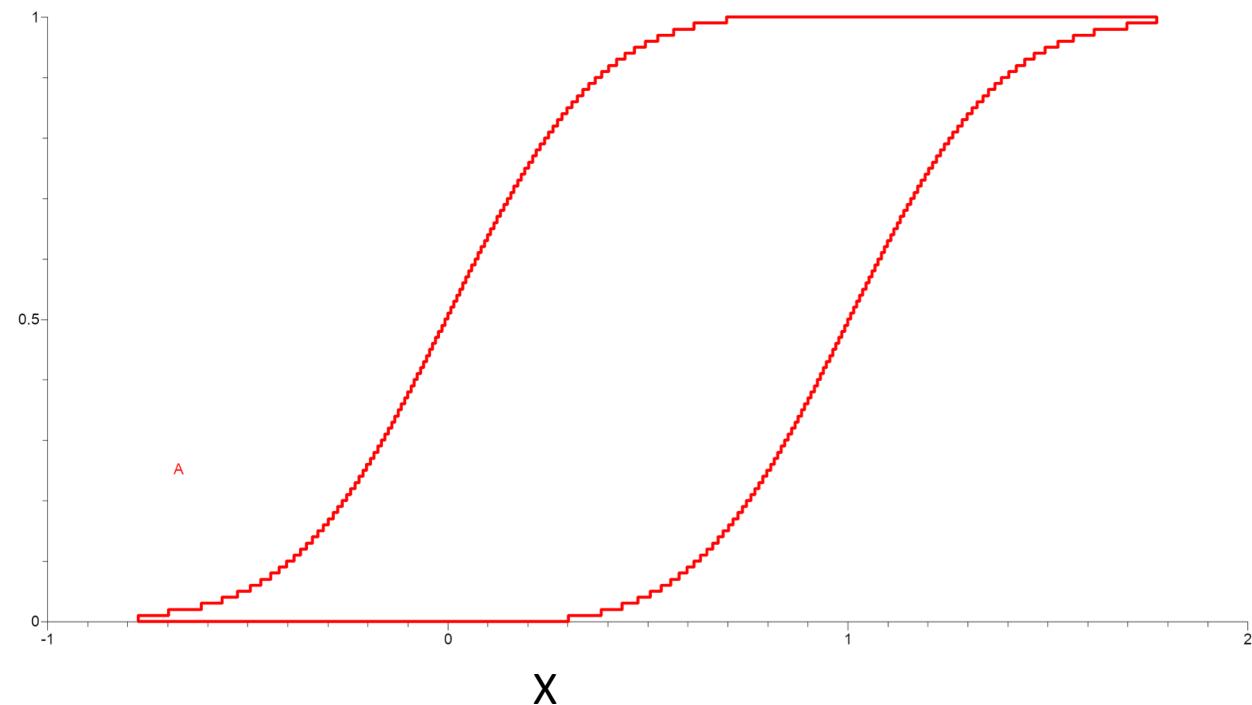
Intervals and probability distributions

$$\underline{x} < x < \bar{x}$$



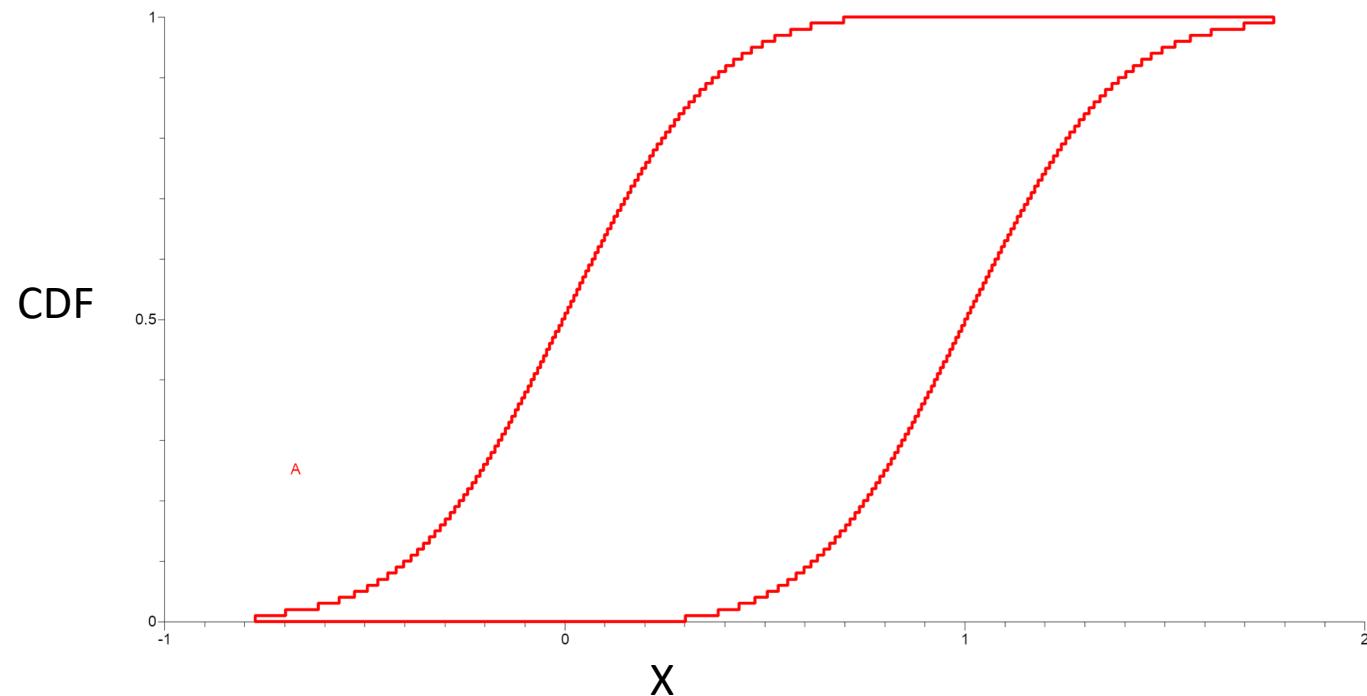
Intervals and probability distributions

$$\underline{f(x)} < f(x) < \overline{f(x)}$$



Examples of uncertain distributions

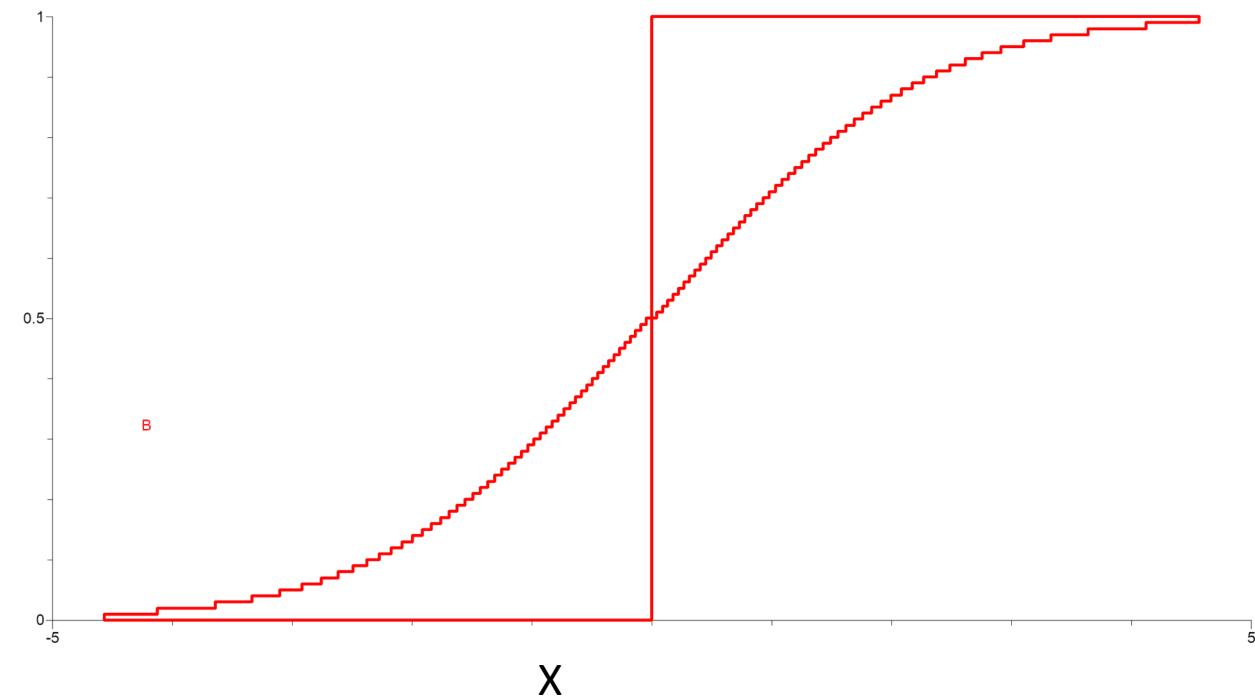
$$X = N([-1,1], 0.3)$$



Examples of uncertain distributions

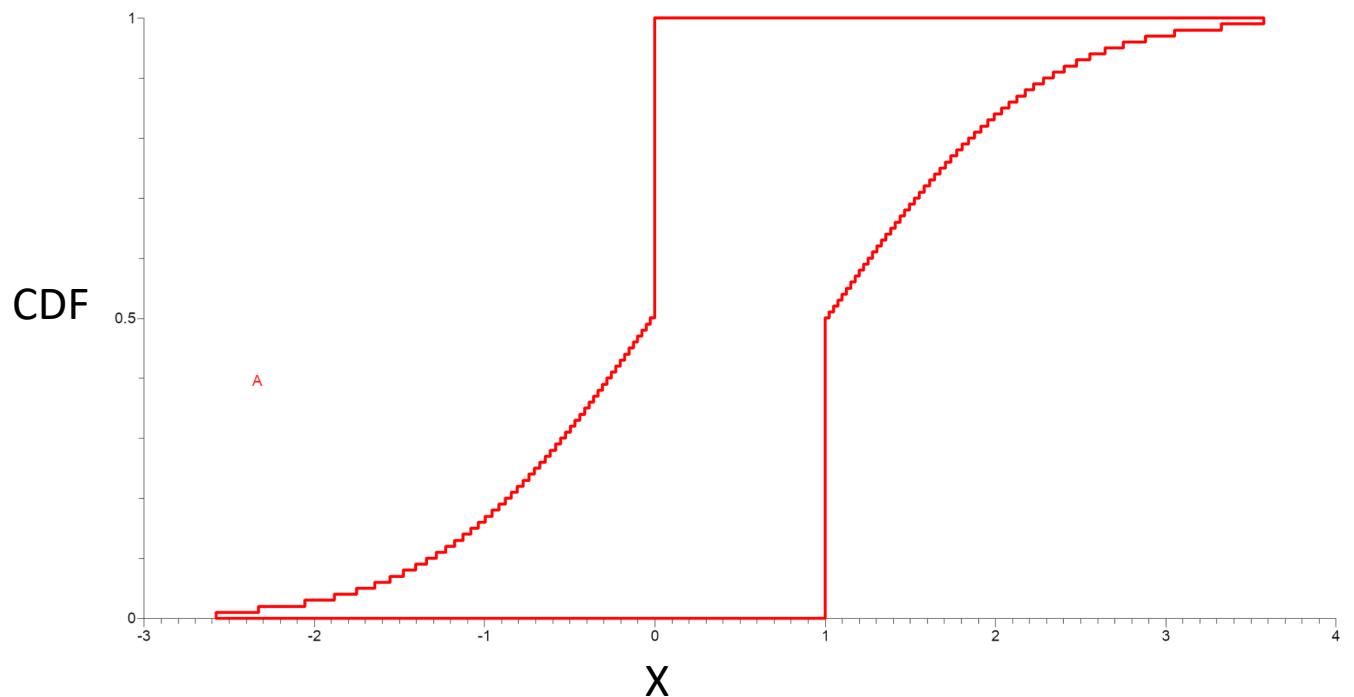
$$A = N([-1,1], 0.3)$$

$$X = N(0, A)$$



Examples of uncertain distributions

$$X = N([0,1], [0,1])$$

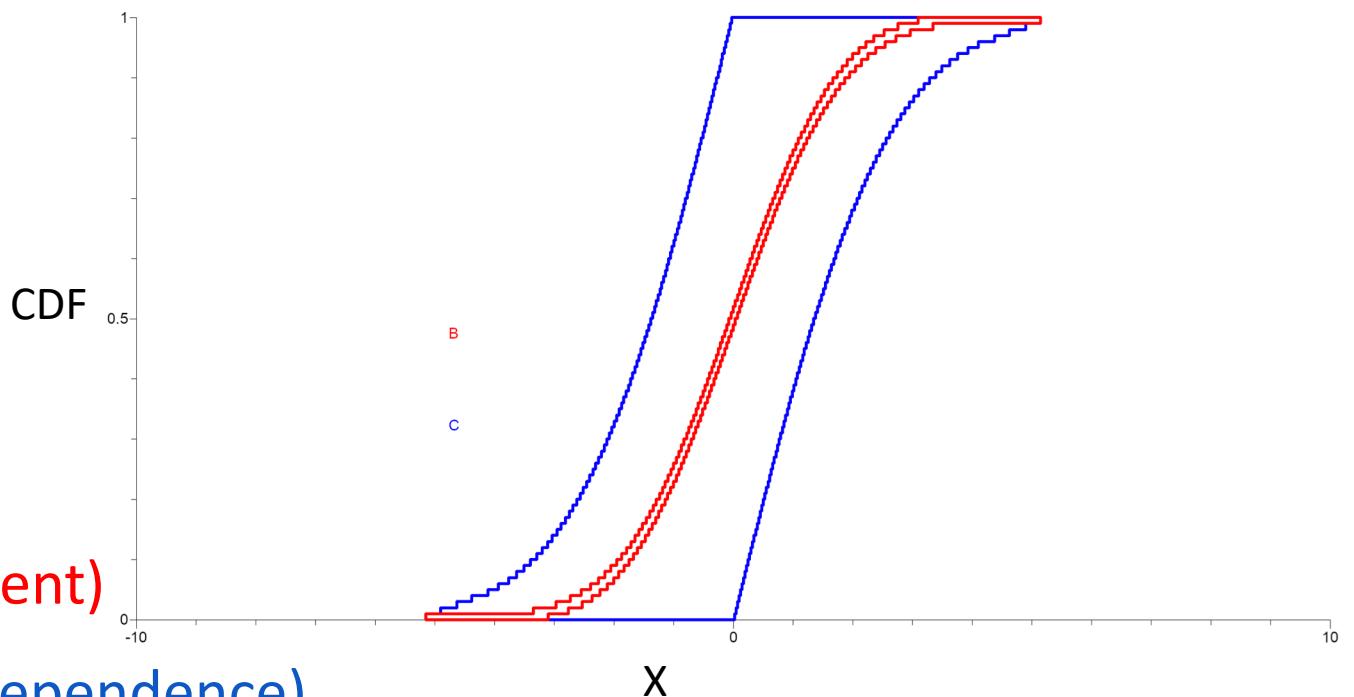


Examples of uncertain distributions

$$A = N(0,1)$$
$$B = N(0,1)$$

$$X = A \mid+ \mid B \quad (\text{Independent})$$

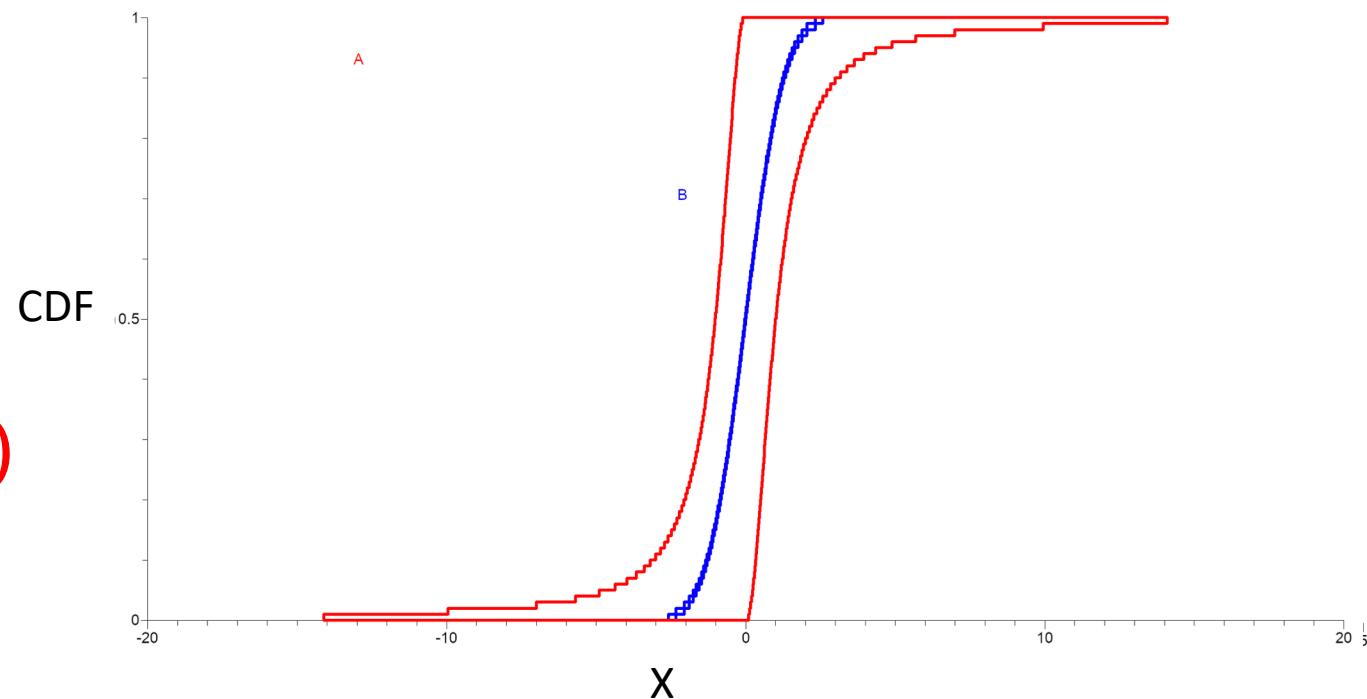
$$X = A + B \quad (\text{Unknown dependence})$$



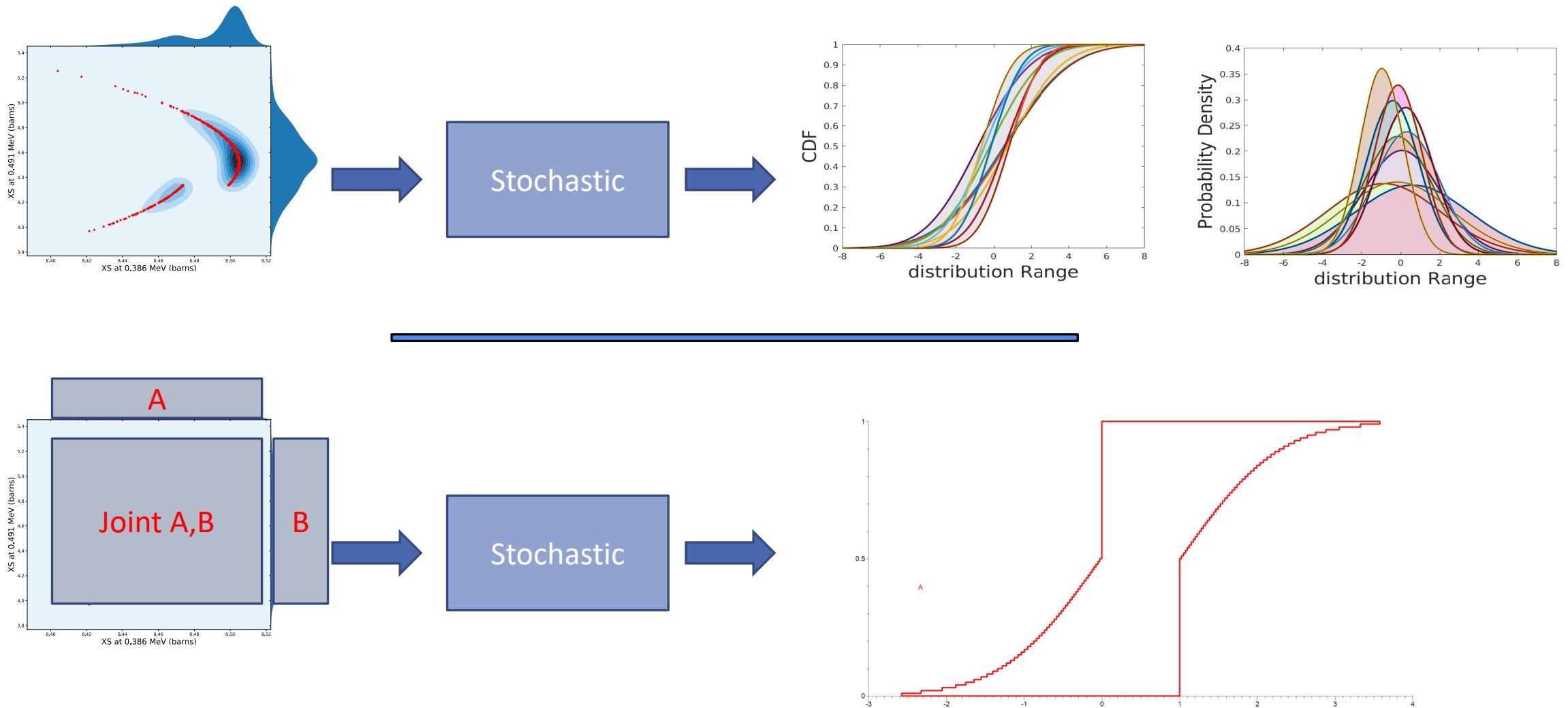
Examples of uncertain distributions

$X = N(0,1)$

$X = Chebyshev(0,1)$



=> Bounds on all distributions (parametric and non-parametric) with a known mean and variance



Interval particle transport

- 1) Specify particle attributes as intervals
- 2) Characterize samplers as imprecise distributions
- 3) Propagate imprecise events concurrently

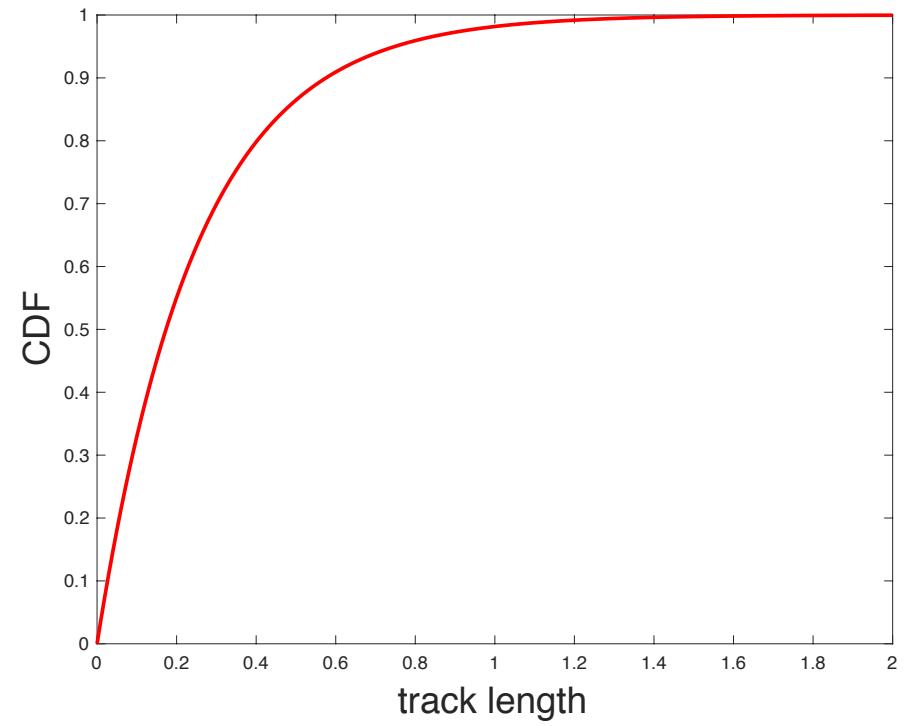
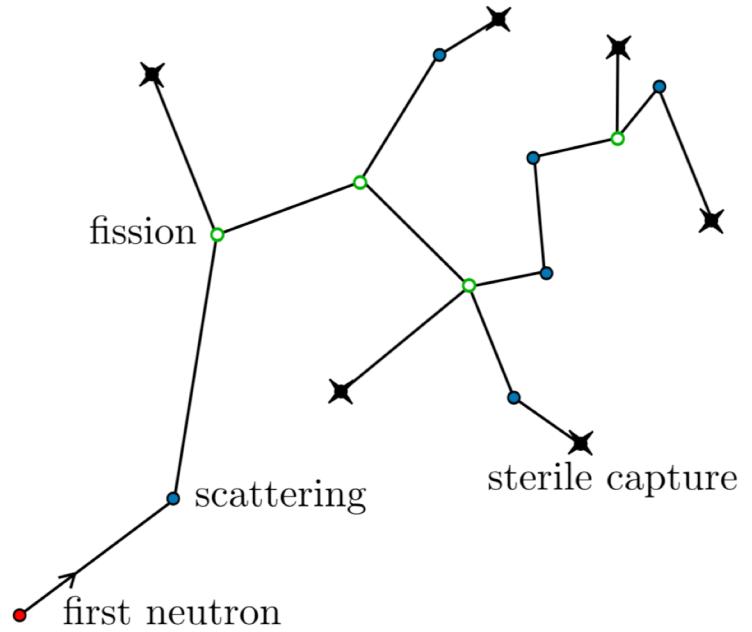
Particle attributes

- Energy (scalar)
- Position (3 vector)
- Direction (3 vector)
- Alive state (Boolean)
- Statistical weight (scalar)



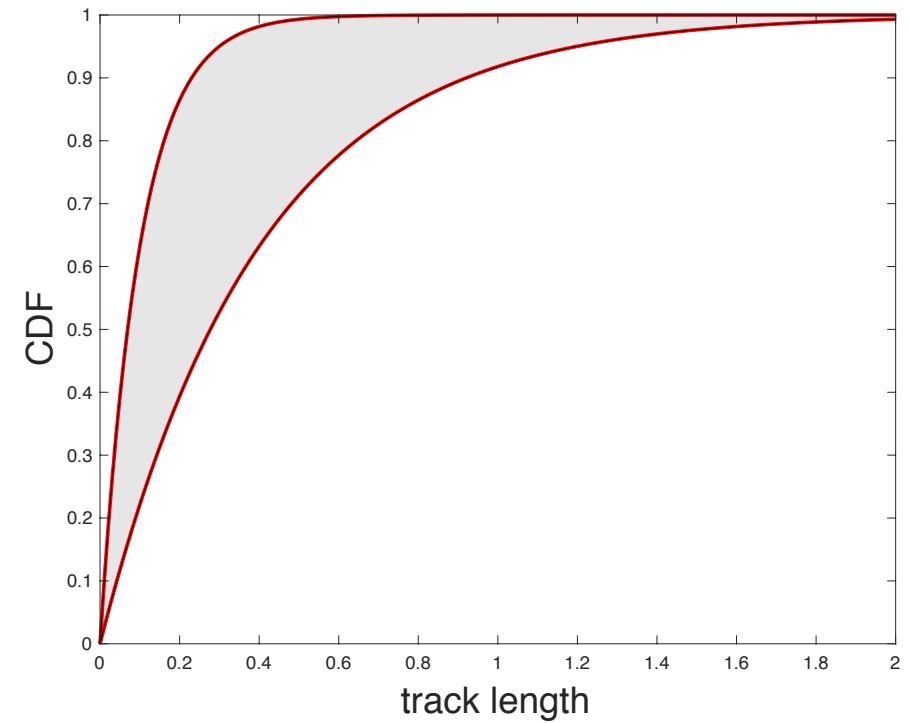
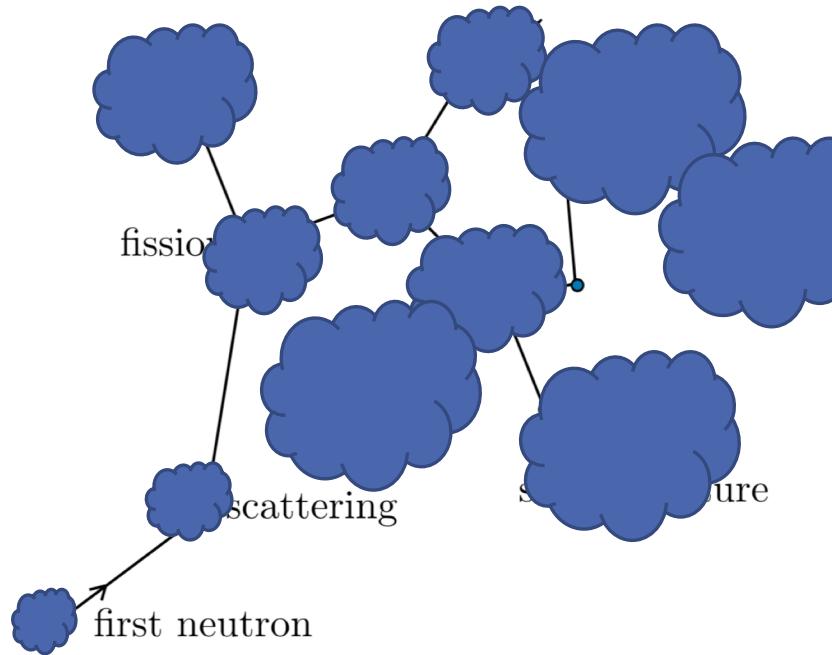
Random initial interval states can be created by specifying the particle source as an imprecise distribution

Imprecise exponential random walks



$$d = -\frac{\ln(\varepsilon)}{\sum_T}$$

Imprecise exponential random walks



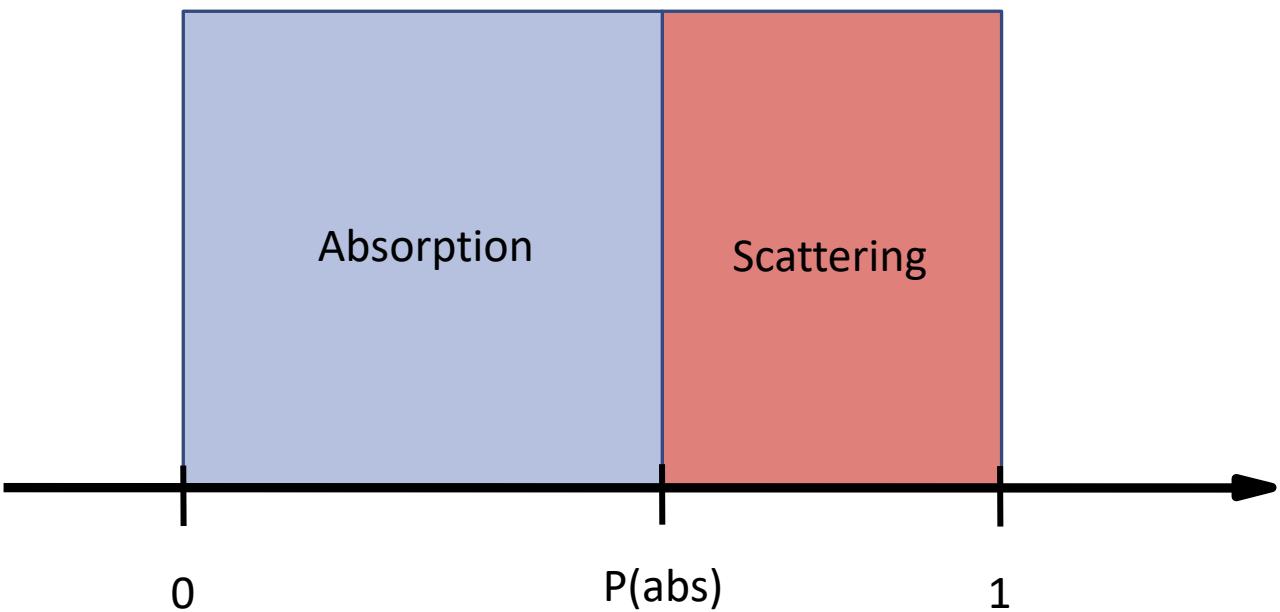
$$[\underline{d}, \bar{d}] = \left[-\frac{\ln(\varepsilon)}{\bar{\Sigma}_T}, -\frac{\ln(\varepsilon)}{\underline{\Sigma}_T} \right]$$

Imprecise events

Consider 2 events: Absorption and scattering

$$P(abs) = \frac{\sigma_{abs}}{\sigma_{scat} + \sigma_{abs}}$$

$$P(scat) = 1 - P(abs)$$

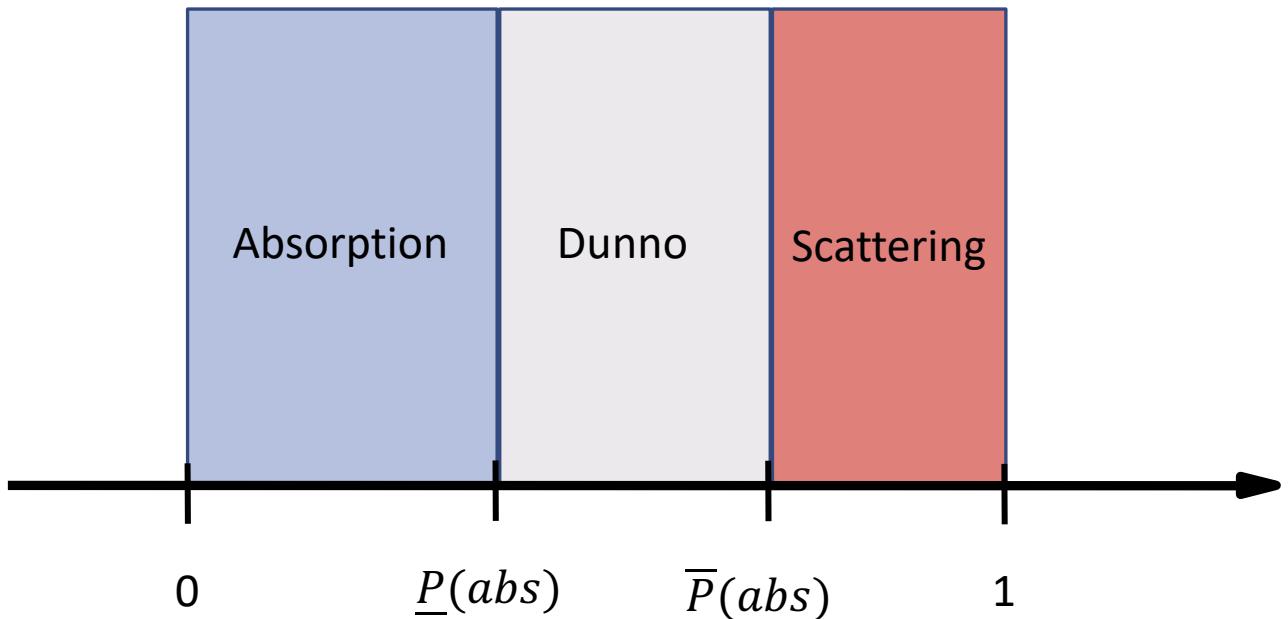


Imprecise events

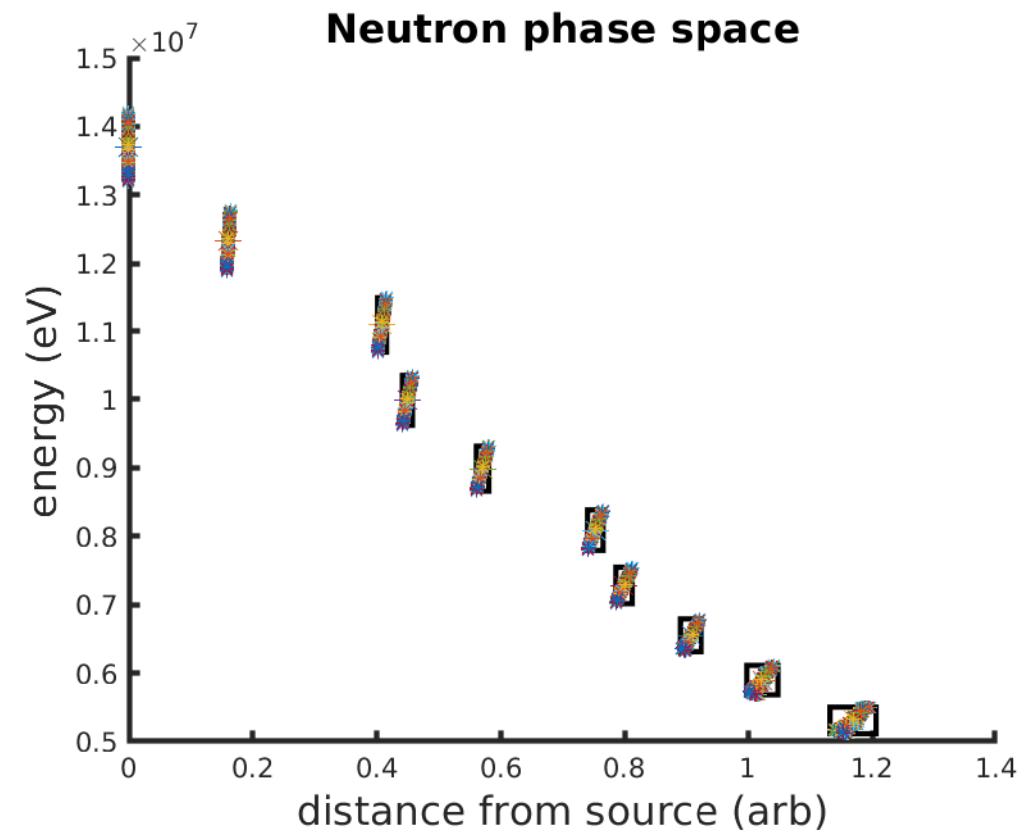
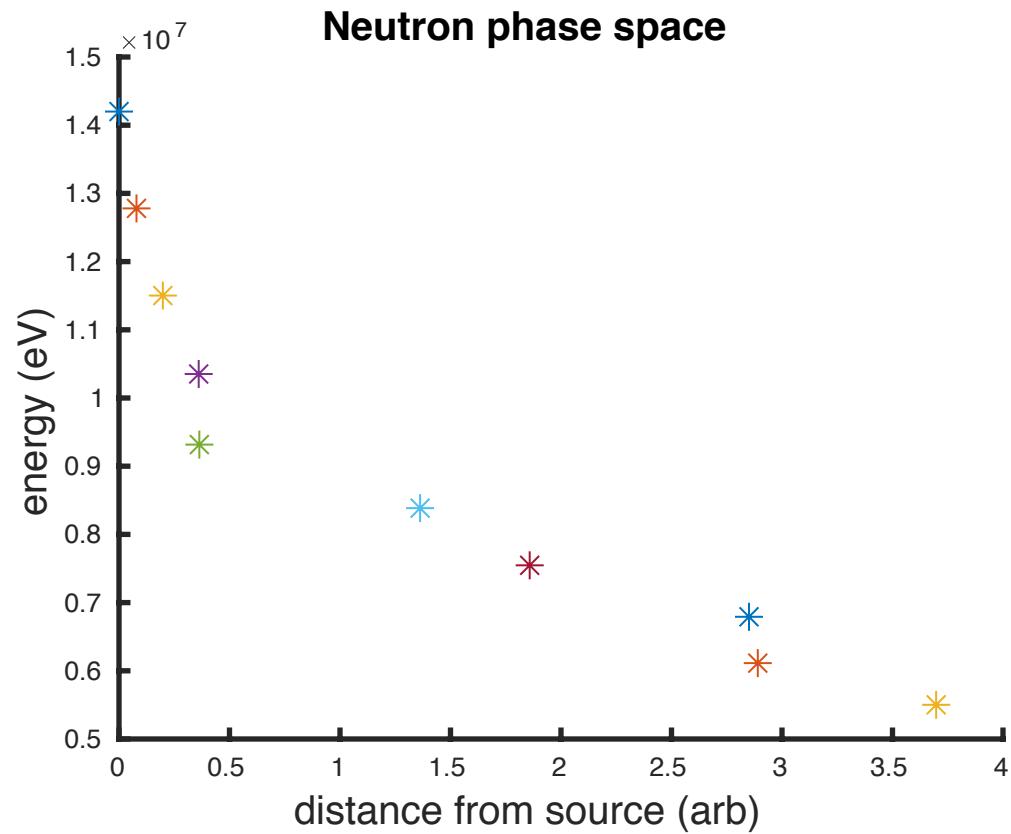
Consider 2 events: Absorption and scattering

$$[\underline{P}(abs), \bar{P}(abs)] = \left[\frac{1}{1 + \frac{\sigma_{scat}}{\sigma_{abs}}}, \frac{1}{1 + \frac{\underline{\sigma_{scat}}}{\underline{\sigma_{abs}}}} \right]$$

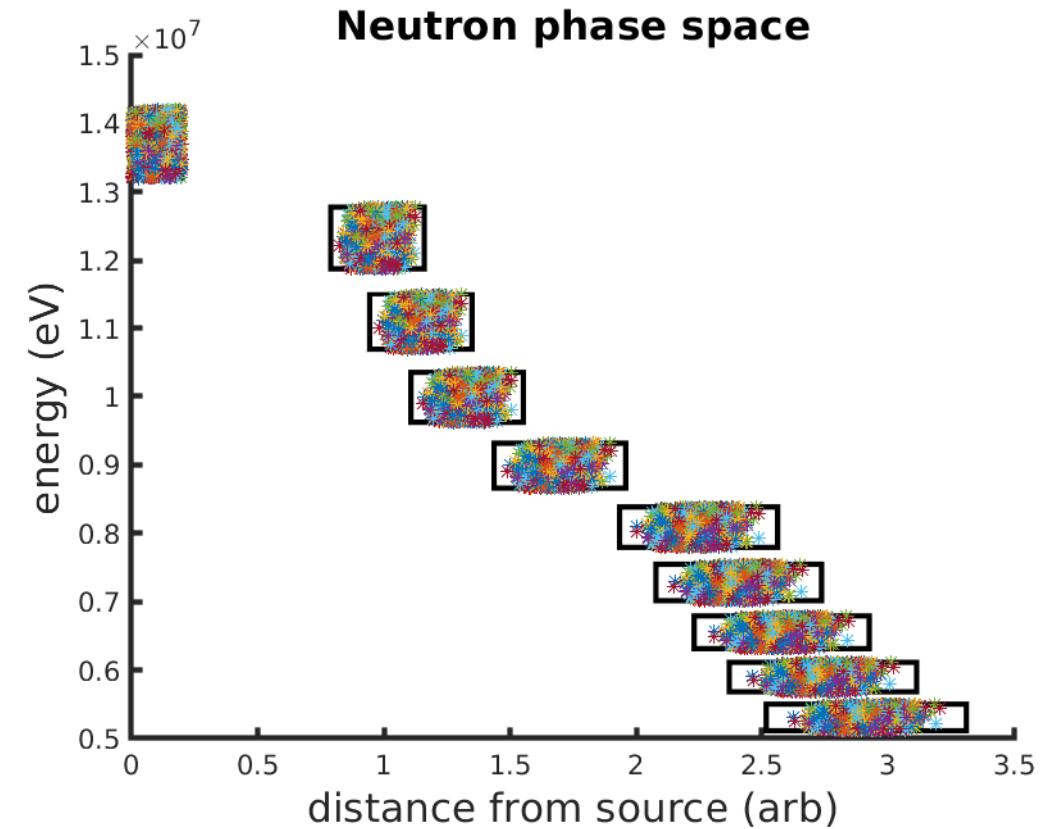
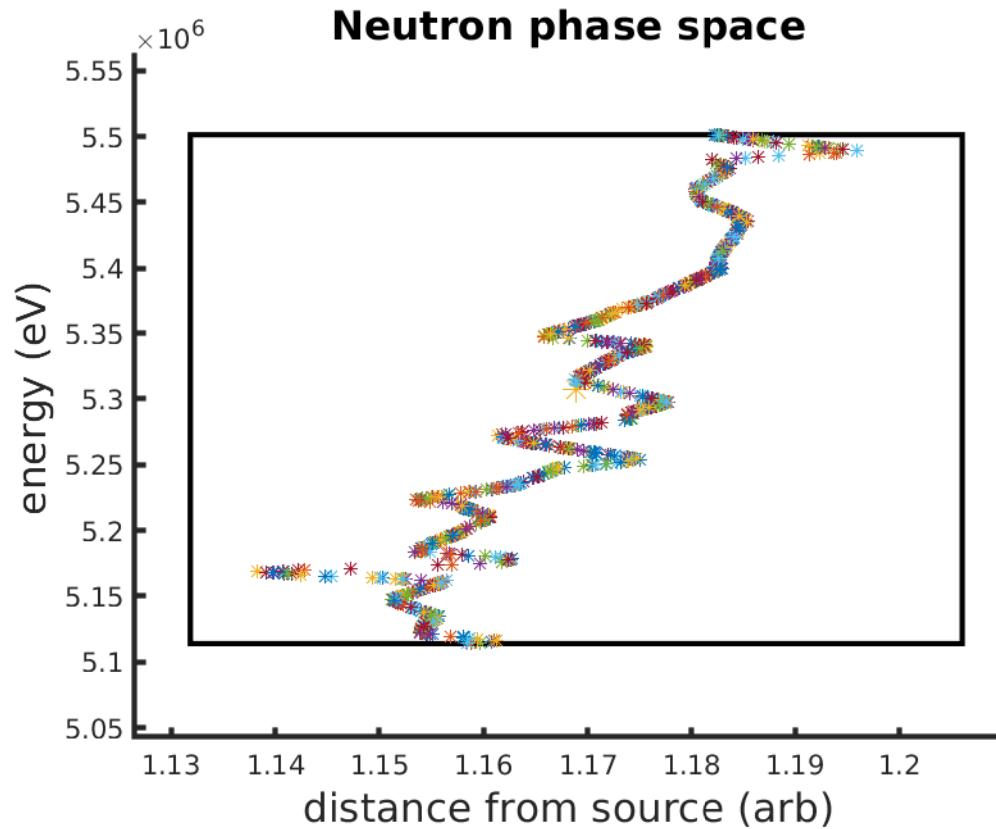
$$P(scat) = 1 - P(abs)$$



Imprecise histories



Imprecise histories



Tallies

The monte carlo estimator for the particle flux is:

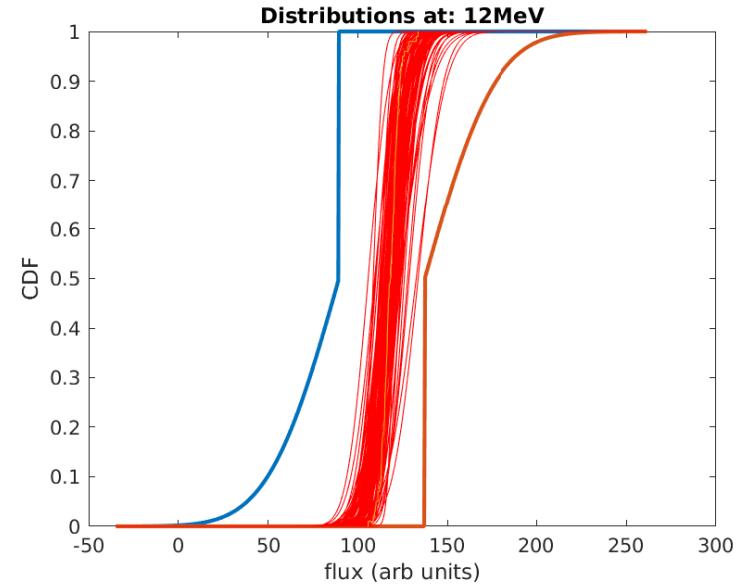
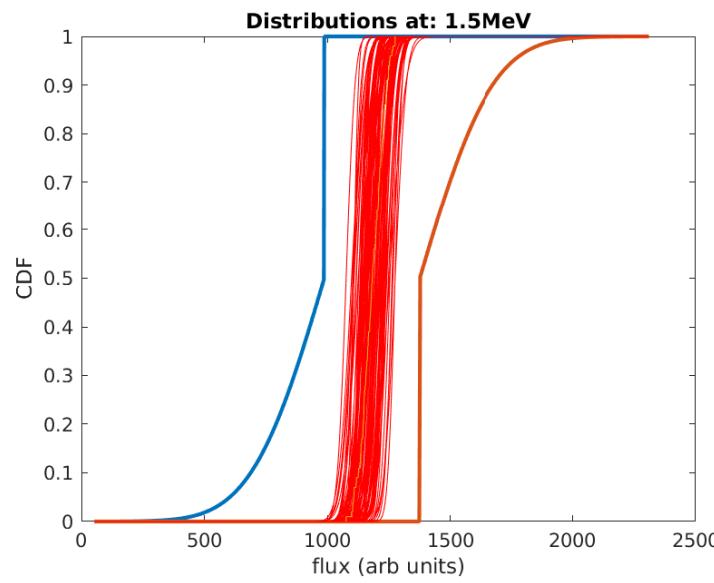
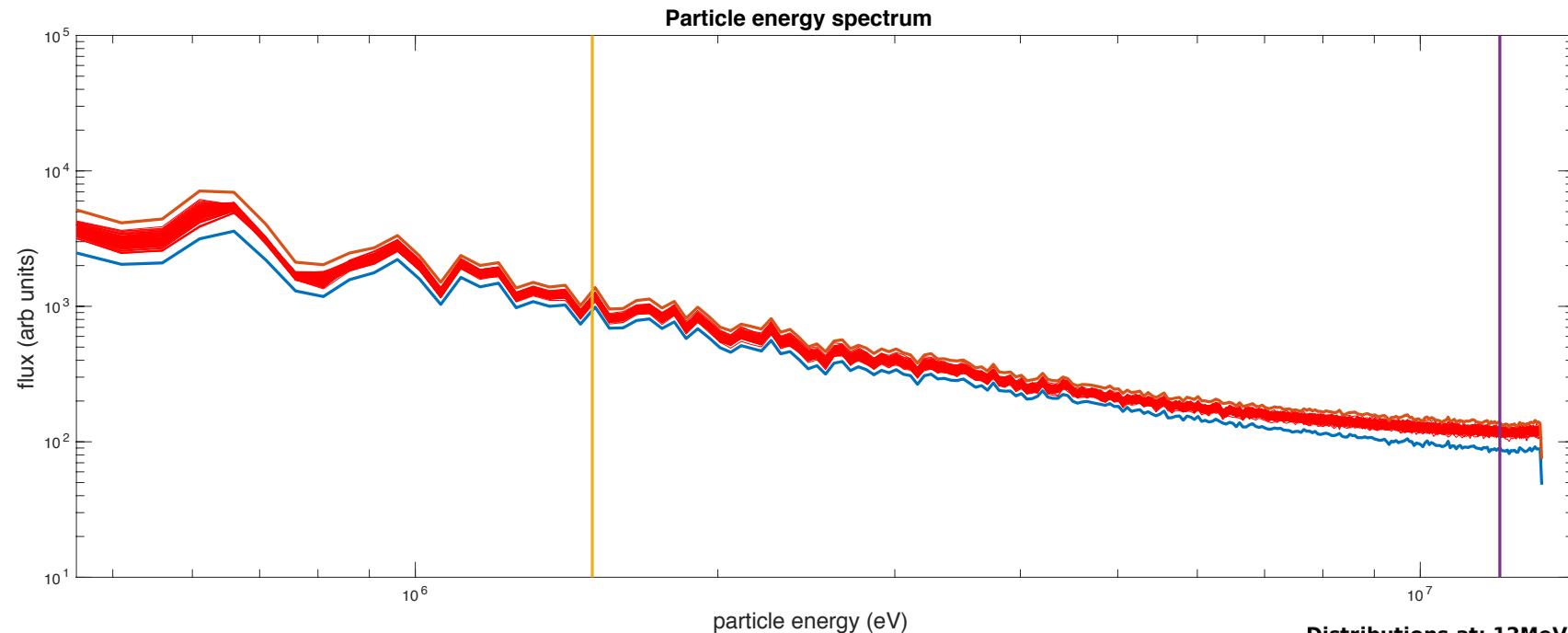
$$\phi(E) = \frac{1}{N} \sum_n d_i$$

Tallies

The monte carlo estimator for the particle flux is:

$$\bar{\phi}(E) = \frac{1}{N} \sum_n \bar{d}_i$$

$$\underline{\phi}(E) = \frac{1}{N} \sum_n \underline{d}_i$$



Conclusions

- We are working towards the (near) automatic propagation of uncertainty in particle transport monte carlo
- We have shown that this works under very simple simulation conditions
- For the method to be generally applicable:
 - The action to be taken when the particle intersects a geometric barrier
 - The simulation of a material with multiple nuclides
 - The concurrent propagation of interval hulls which have been split from imprecise events