

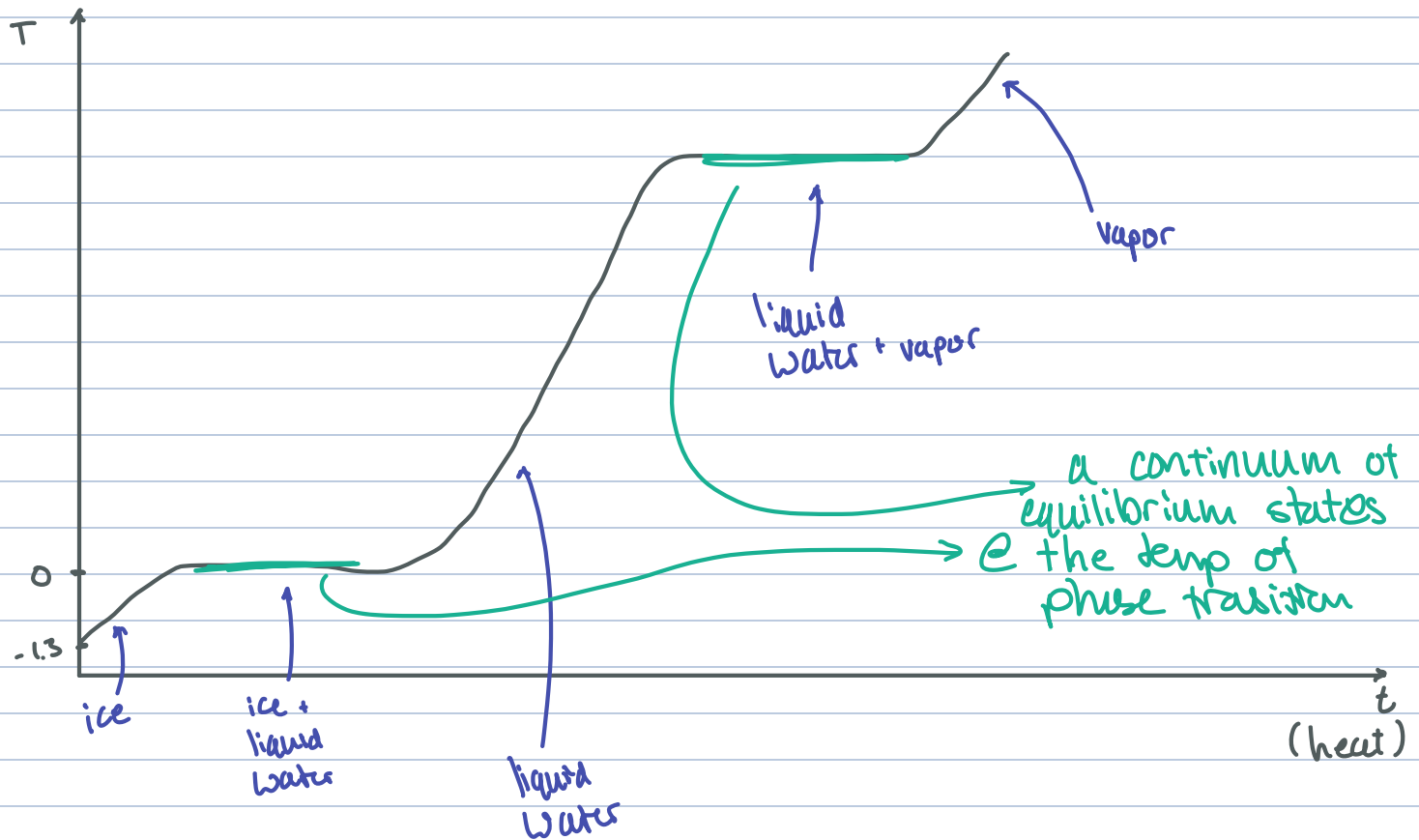
Moodle PW: kasteleyn.

Advertisement: lecture M.Sc. Physics 3 ECTS,

Order in disorder: stochastic geometry in physics.

Friday, 2-4 pm: Office hours (complements on class; background).

Monthly review session.



Motivating question: what is the state of water as a function of time/heat/temperature?

State.

Quantum \leadsto density operator ρ . Hermitian.

$$\begin{pmatrix} p_1 & & \\ & p_2 & \\ & & \ddots \end{pmatrix} \leadsto \sum p_i = \text{tr } \rho = 1, \text{ to accurately repr. probabilities}$$

Assume thermal equilibrium:

H = Hamiltonian of water

$\beta = \frac{1}{kT}$ inverse temperature

$$\leadsto p = \frac{e^{-\beta H}}{\text{tr } e^{-\beta H}} \quad \left. \vphantom{\frac{e^{-\beta H}}{\text{tr } e^{-\beta H}}} \right\} \text{only one state per temperature, contradictory flat parts of graph above.}$$

Why should this make us feel uneasy? 

Soln: $e^{-\beta H} / \text{tr } e^{-\beta H}$ is divergent / requires regularization:

1. $\text{trace } \sum_i e^{-\beta E_i} = \infty$

not unique, matters.

2. H infinite in thermodynamic limit (some extensive term, eg binding energy)

Phase transition is obtained from details of regularization. Control is required to capture physics of the transition.

The motivation is: how to abstractly treat systems w/ poly many dots. At one end, Ising model ~ at the other, QFT. In CM/SP case, the ∞ s arise because of their infinite extent.

Books. Classical: Sabine ~  SM of Lattice..., slides.

Quantum: Robert ~ • Brattelli, Robinson;
2 vols (we mainly use volume 2).

• Thirring: vol 4

• Robert lecture notes.

Video lectures: Imucast.

Classical Stat Mech.

Overview math part:

1. Ising model: existence of TL of pressure/free energy.

FV : Peierls argument
ch 3 : Phase transition

2. Gibbs measures in infinite volume.

DLR conditions: Dobrushin, Lanford, Ruelle.
FV Ch 6.

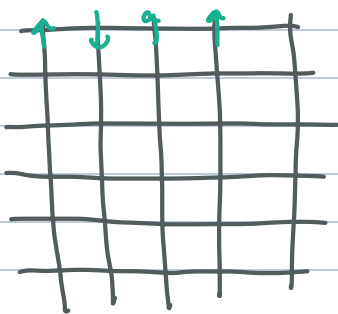
3. Mermin-Wagner thm, absence of cts sym. breaking, $d=1, 2$ FV Ch. 9.

(4.) Reflection positivity/ex. of sym. breaking $d=3$
FV Ch 10.

1. The Ising model.

Reference: Hugo Duminil-Copin: 100 years of (critical) Ising model (arXiv 2022).

1.1. The model.



$\uparrow = +1$
 $\downarrow = -1$

Notn:

$\Lambda \in \mathbb{Z}^d$: $\Lambda \subset \mathbb{Z}^d$
finite, non-empty

Often: $\Lambda = \{1, \dots, L\}^d$, $L \in \mathbb{N} \leadsto$ cubes.

Configuration space. $\Omega_\Lambda := \{-1, +1\}^\Lambda$.

$\omega \in \Omega_\Lambda$, $\omega = (\omega_i)_{i \in \Lambda}$, $\omega_i \in \{\pm 1\}$.

$h \in \mathbb{R}$ external magnetic field,

$$\mathcal{H}_{\Lambda;h} : \Omega_{\Lambda} \rightarrow \mathbb{R},$$

$$\omega = (\omega_i)_{i \in \Lambda} \mapsto \mathcal{H}_{\Lambda;h}(\omega) = - \sum_{\substack{\{i,j\} \subset \Lambda \\ i,j \text{ n.n.}}} \omega_i \omega_j - h \sum_{i \in \Lambda} \omega_i$$

↑
nearest nbs in \mathbb{Z}^d

Minimizer of \mathcal{H} ?

$$h > 0 : \omega_i = 1 \forall i$$

$$h < 0 : \omega_i = -1 \forall i$$

$$h = 0 : 2 \text{ minimizers, all } +1 \text{ or all } -1$$

Partition function (Zustandssumme)

$$Z_{\Lambda}(\beta, h) = \sum_{\omega \in \Omega_{\Lambda}} e^{-\beta \mathcal{H}_{\Lambda;h}(\omega)}$$

Gibbs measure: prob measure on Ω_{Λ}

$$\mu_{\Lambda; \beta, h}(\{\omega\}) = \frac{1}{Z_{\Lambda}(\beta, h)} e^{-\beta \mathcal{H}_{\Lambda;h}(\omega)}$$

$\beta = 0$: uniform distribution

$\beta \rightarrow \infty$: measure concentrates on minimizer(s).

Pressure / free energy:

$$\Psi_{\Lambda}(\beta, h) := \frac{1}{\beta |\Lambda|} \log Z_{\Lambda}(\beta, h)$$

Total magnetization:

$$M_{\Lambda} : \Omega_{\Lambda} \rightarrow \mathbb{R},$$

$$\omega \mapsto M_{\Lambda}(\omega) = \sum_{i \in \Lambda} \omega_i$$

Obs: $\frac{\partial}{\partial h} \Psi_{\Lambda}(\beta, h) = \frac{1}{|\Lambda|} \sum_{\omega \in \Omega_{\Lambda}} M_{\Lambda}(\omega) e^{-\beta \mathcal{H}_{\Lambda;h}(\omega)} \frac{1}{Z_{\Lambda}(\beta, h)}$

$$= \frac{1}{|\Lambda|} \langle M_{\Lambda} \rangle_{\Lambda, \beta, h} =: m_{\Lambda}(\beta, h)$$

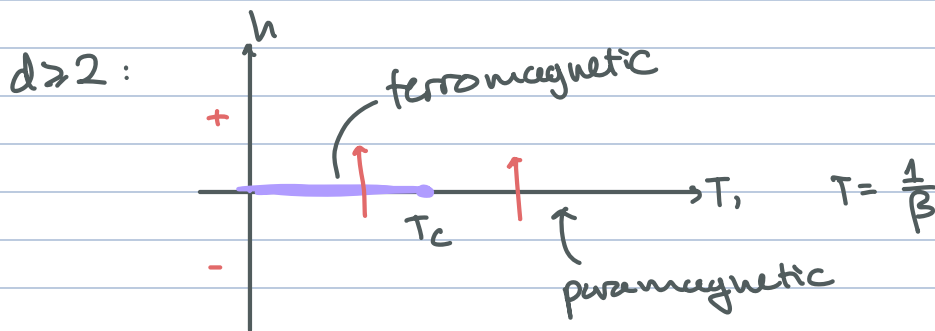
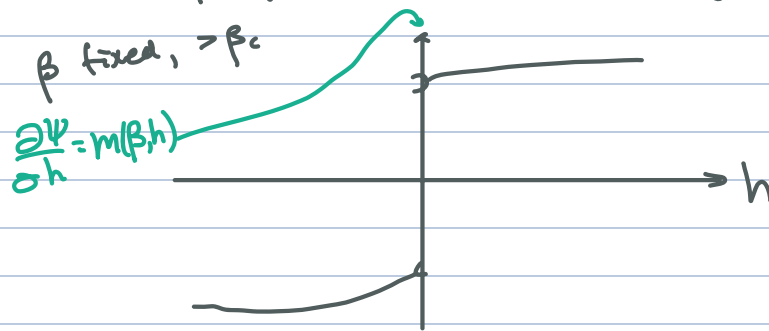
↙ average magnetization per unit volume

Preview: ① $\Psi(\beta, h) = \lim_{\Lambda \uparrow \mathbb{Z}^d} \Psi_\Lambda(\beta, h)$

$$(\log Z_\Lambda(\beta, h) \approx \beta |\Lambda| \Psi(\beta, h) + \text{correction terms}) \ll |\Lambda|$$

② $\lim_{n \rightarrow \infty} m_{\Lambda_n}(\beta, h) \stackrel{?}{=} \frac{\partial \Psi}{\partial h}(\beta, h)$
 ↑ does it exist?

③ $\exists \beta_c < \infty$ s.t. $\beta > \beta_c$, $\frac{\partial \Psi}{\partial h}$ has jump discontinuity at $h=0$.



Will there be a jump, step change in something as h is changed from $- \rightarrow 0^+$, below and above T_c ?