

# Computational Design of Materials Lab 2

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# 1 Problem 1: Cutoff Convergence of Absolute Energies

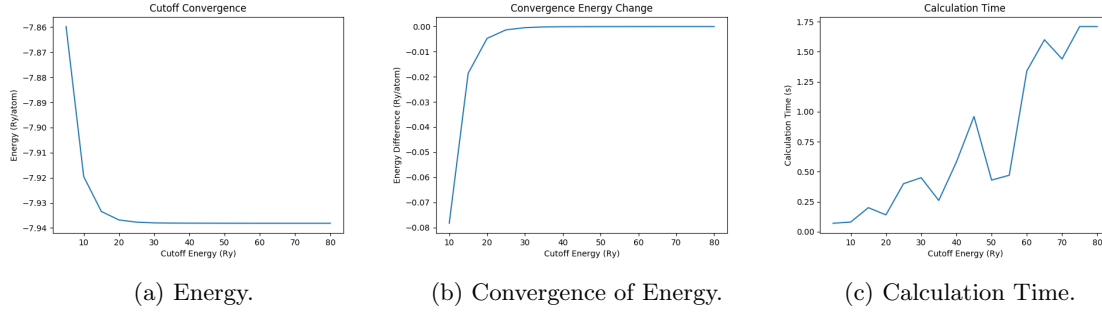


Figure 1: Absolute Energy vs. Kinetic Energy Cutoff.

## 1.1 A

In this section, we examine the convergence behavior of the Quantum Espresso DFT code for a germanium lattice with diamond cubic structure with respect to the kinetic energy cutoff, which sets the maximum kinetic energy (and by extension, the largest multiple of the reciprocal lattice vector) included in the approximation. We use a fixed k-point grid of  $4 \times 4 \times 4$  by setting the  $nk$  variable equal to 4, a fixed lattice parameter of  $5.0 \text{ \AA}$ , and vary the energy between 5 and 80 Ry in increments of 5 Ry. This leads to the convergence behavior in Figure 1. The energy converges to a value of  $-7.94 \text{ Ry/atom}$ , and the second plot shows that the incremental change in energy decreases below  $5 \text{ meV/atom}$  (or  $0.0004 \text{ Ry/atom}$ ) at a value of 35 Ry. Note that we define convergence as the difference between the final energy value and the energy value at each energy cutoff point.

## 1.2 B

The third plot contains the calculation time as a function of the cutoff energy, showing a clear positive relationship between the two. This is expected, as higher kinetic energy cutoffs means more values are included for each k-point, thereby increasing the total computational cost.

## 1.3 C

One advantage of using the primitive cell over the unit cell is obvious- the primitive cell contains less atoms (in this case, 2 vs. 8 for diamond cubic structure) and a smaller overall volume. In many cases, however, the symmetries present in the unit cell are not necessarily reflected in the primitive cell. Therefore, while a primitive cell may allow for fewer calculations due to a smaller number of atoms, a unit cell may alternatively facilitate fewer calculations through symmetries that may not otherwise be apparent. Symmetries can in turn reduce the number of k-points required to effectively sample the Brillouin zone.

# 2 Problem 2: K-Point Convergence of Absolute Energies

## 2.1 A

As before, we use a diamond cubic Ge lattice with a lattice parameter of  $5.0 \text{ \AA}$ , and fix our kinetic energy cutoff to a value of 30 Ry, which should be nearly sufficient in light of the trend seen in the previous problem. The resultant energy is recorded for each input number of k-points, specified as the number in each direction (so an input value of 2 corresponds to a  $2 \times 2 \times 2$  grid), along with the computation time and the number of unique k-points, as some are removed due to the symmetries of the crystal. Figures 2 and 3 display these results as a function of the input grid size and the number of unique k-points, respectively.

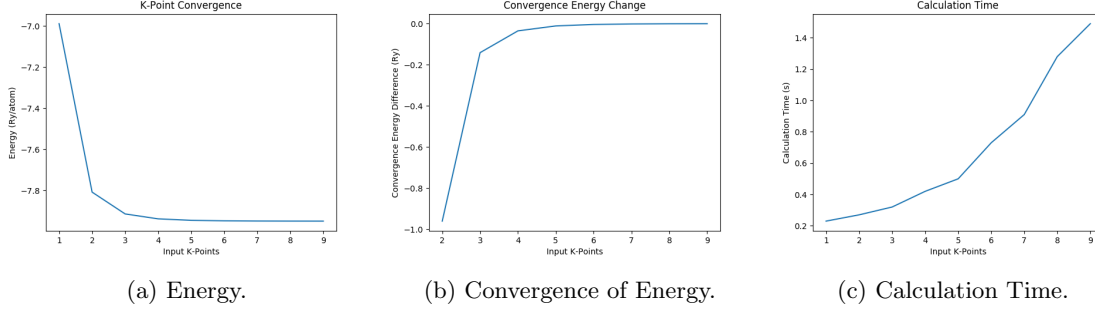


Figure 2: Absolute Energy vs. Input K-Point Grid Size.

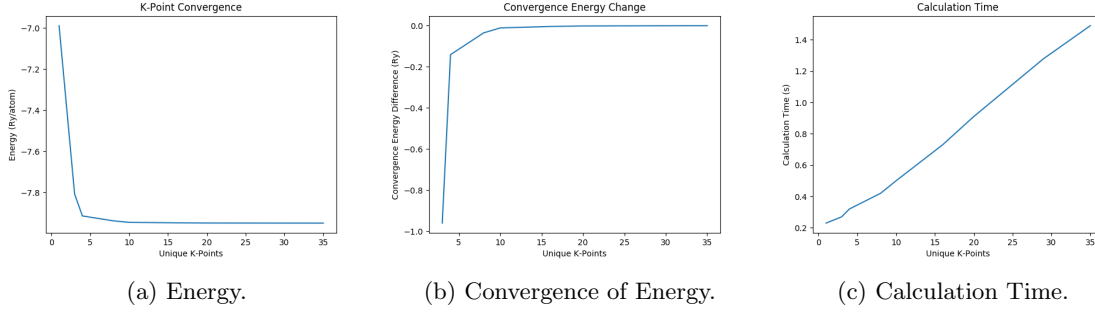


Figure 3: Absolute Energy vs. Unique Number of K-Points.

## 2.2 B

A trend is evident for both the calculated energies and the computation times. For the former, the energy converges to a constant value of -7.95 Ry/atom around 35 unique k-points, or an input grid of 9 x 9. Beyond this value, increasing the number of k-points does not appear to significantly change the approximation, as all of the major energy terms (e.g. Hartree, exchange-correlation, and kinetic energy) have been accounted for, though the computation time continues to increase since more k-points are still being tested, even though the result is not much better.

## 3 Problem 3: Cutoff Convergence of Forces

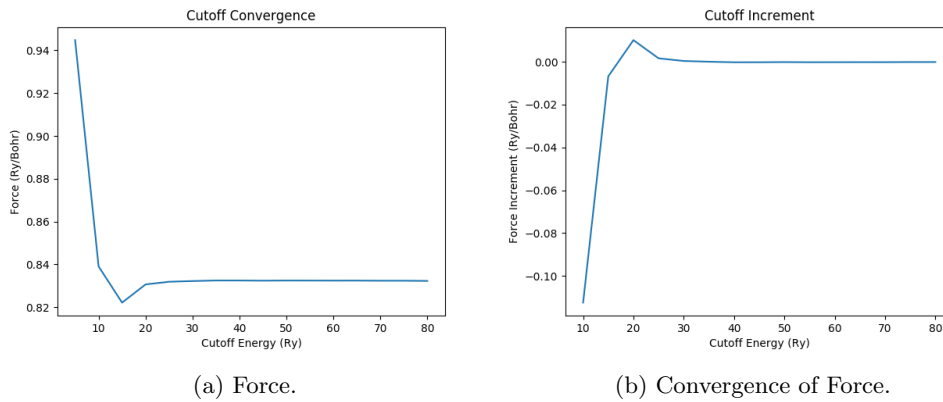


Figure 4: Force vs. Cutoff Energy.

We repeat the same process as in Problem 1 using the same lattice parameter ( $5 \text{ \AA}$ ), crystal structure, and k-point grid size ( $4 \times 4 \times 4$ ), as well as the same range of kinetic energies. In this case, however, we shift the position of the first Ge atom in the lattice by adding 0.05 to the third coordinate of its position through the `.positions` attribute, corresponding to a fractional shift of 0.05 times the lattice parameter in the  $z$  direction. Figure 4 displays the results from these calculations. The force converges to a value of  $0.83 \text{ Ry}/(\text{Bohr-atom})$  at a value of  $30 \text{ Ry}$ , in which convergence is defined as above but with a threshold of  $10 \text{ meV}/(\text{\AA-atom})$ , or  $0.0004 \text{ Ry}/(\text{Bohr-atom})$ .

## 4 Problem 4: K-Point Convergence of Forces

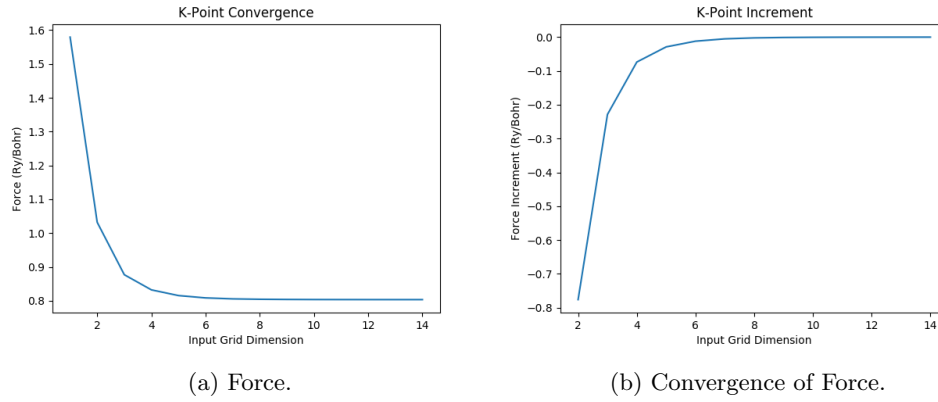


Figure 5: Force vs. Input K-Point Grid Size.

Using the same setup as in Problem 2 (lattice parameter of  $5 \text{ \AA}$  and cutoff energy of  $30 \text{ Ry}$ ) with the same displacement scenario as in the previous problem, we obtain the results shown in Figure 5. Here, the force converges to a value of  $0.80 \text{ Ry}/(\text{Bohr-atom})$  at a grid size of  $11 \times 11 \times 11$ , or 216 unique k-points.

## 5 Problem 5: Cutoff Convergence of Energy Differences

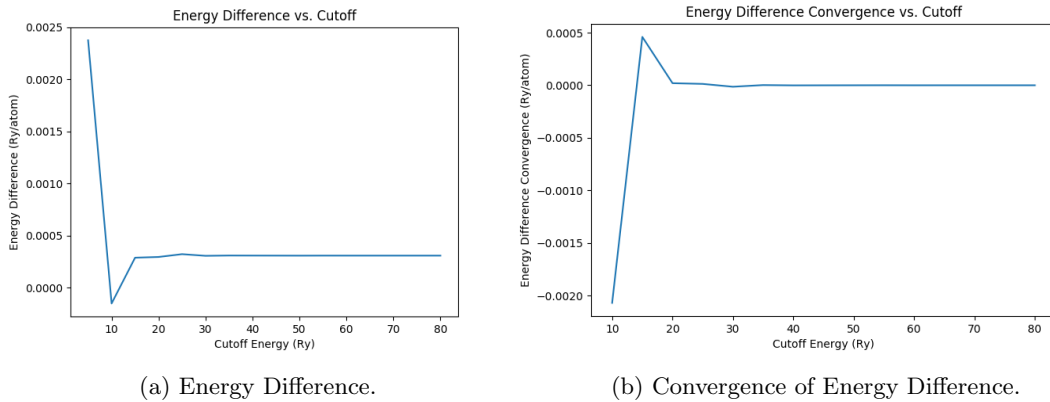


Figure 6: Force vs. Cutoff Energy.

Using a k-point grid of  $4 \times 4 \times 4$  and the same lattice parameter as in all previous problems, we perform cutoff convergence as in Problem 1 for two different lattice parameters:  $10.70 \text{ Bohr}$  (or  $\text{\AA}$ ) for one, and  $10.75 \text{ Bohr}$  (or  $\text{\AA}$ ) for the other. Figure 6 shows the calculated energy difference as a function of kinetic energy

cutoff between these two cells (the one with the smaller lattice parameter is subtracted from the one with the larger lattice parameter). We see convergence to a final value of 0.0004 Ry/atom at a comparably low kinetic energy cutoff of 20 Ry.

## 6 Problem 6: Discussion

Based on the above results, we can observe opposite trends in convergence with respect to the kinetic energy cutoff and k-point grid size for absolute energies and forces. For absolute energies, we see a slightly larger kinetic energy cutoff required for convergence (35 Ry) compared to that necessary for force convergence (30 Ry). Conversely, an input grid dimension of 9 is sufficient to allow for convergence of absolute energies with respect to k-points, while a larger input dimension of 11 is necessary for forces. This latter behavior cannot be further contrasted with respect to energy differences, since k-point convergence was not examined. However, we have found that energy differences appear to converge more quickly with respect to the kinetic energy cutoff than either of the other two properties, requiring a cutoff of only 20 Ry. The larger k-point grid requirement with respect to forces seems to make sense, as the force was introduced by displacing the atom from its equilibrium position. As a consequence, some of the symmetries of the crystal may be broken, which could in turn require more k-points to be sampled before convergence. It is possible that the larger kinetic energy cutoff for absolute energies reflects a larger importance of core electrons in computing absolute energies over forces, though the difference between these two cutoffs is fairly small. The significantly lower cutoff energy for the energy differences may be explained by considering potential errors or fluctuations that arise during computation. When taking a difference between two structures, it is possible that

## 7 Problem 7: Equilibrium Lattice Constant and Bulk Modulus

As we have seen that forces and energy differences appear to show higher cutoff criteria than absolute energies, we increase the size of our calculation to a k-point grid size of 11 and an energy cutoff of 35 Ry.

### 7.1 A: Lattice Parameter

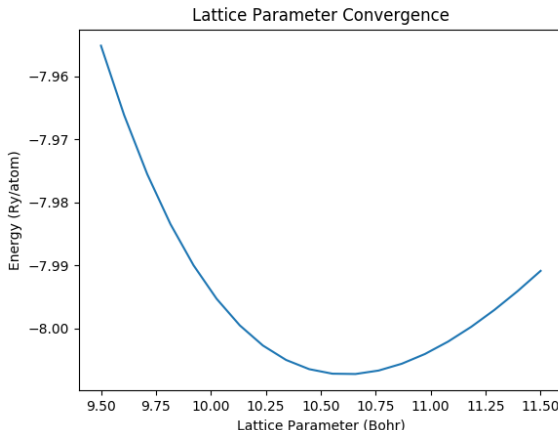


Figure 7: Lattice Parameter Optimization.

We vary the lattice parameter of Ge in the diamond cubic crystal structure between 9.5 and 12.5 Bohr in increments of 0.1, and obtain the results in Figure 7. The k-point grid is set to 11 and the kinetic energy cutoff to 35 to ensure convergence with respect to the previously tested properties. The minimum value of -8.007 Ry/atom is found at 10.61 Bohr, nearly consistent with the experimental results, but off by almost 1%, consistent with the fact that DFT often overbinds systems.

## 7.2 B: Bulk Modulus

To compute the bulk modulus, we begin with the given relationship:

$$B = -V_0 \frac{\partial P}{\partial V}$$

Here, the proportionality constant is the equilibrium volume of the unit cell,  $B$  is the bulk modulus, and  $P$  and  $V$  are pressure and volume, respectively. Commonly, bulk modulus calculations are solved by invoking a second relationship relating energy and pressure,

$$P = -\frac{\partial E}{\partial V}$$

Combining these results yields the relationship between the bulk modulus and the curvature of the energy-volume relationship:

$$B = V_0 \frac{\partial}{\partial V} \frac{\partial E}{\partial V} = V_0 \frac{\partial^2 E}{\partial V^2}$$

We then need a relation between the volume of the primitive cell and the lattice parameter. A simple empirical relationship can be derived using the output files, or from the geometry of the primitive cell. The volume can be found to be related to the lattice parameter as  $V = \frac{a^3}{4}$ . With these formulae, we can construct a simple method for obtaining the bulk modulus. We begin by taking the lattice parameter at the minimum energy found above, 10.62 Bohr, as *alat\_min*, and find its corresponding volume,  $V_{min}$ . In order to facilitate unit conversions and input into the lattice scan function, we convert both of these parameters to Å instead. We approximate the second derivative using a fourth-order expansion of the central difference theorem, such that:

$$f''(x_0) = \frac{-f(x_0 + 2h) + 16f(x_0 + h) - 30f(x_0) + 16f(x_0 - h) - f(x_0 - 2h)}{12h^2}$$

We use a value of  $h = 1$  cubic Angstrom, so that the denominator simply becomes 12, and for each perturbation from the central value  $x_0$ , we find the lattice parameter corresponding to the desired volume by inverting the above formula, so  $a = (4V)^{1/3}$ , and pass this as an input into the DFT code. We obtain the resultant energy from the code in units of electron volts, and obtain the approximation for the second derivative using the above formula. We then multiply by the  $V_{min}$  parameter to yield a final result with units of electron volt per cubic Angstrom. We then convert to units of  $J/m^3 = Pa$  for the final result. Using this methodology, we obtain a bulk modulus of 70 GPa, fairly close to the experimental value of 76 GPa. Some of the error in this result may not be due to the simulation itself, but rather a consequence of the finite difference method used to approximate the second derivative. Options to improve this accuracy include using higher-order expansions in deriving the formula or fitting the data (with more points) near the minimum to an equation of state and extracting the value from this equation. With respect to the former option, grids were tested from 0.01 to 1 and from 0.5 to 4 cubic Å to examine the effect of the volume increment parameter on the results, and Figure 8 displays these results. The former grid shows that the values have converged by the initial test of 1, and the latter shows that the method appears to get less accurate for larger values, which is expected as these will be less accurate in approximating the true form of the second derivative, though these fluctuations in the latter case are quite small.

## 7.3 C: Automatic Optimization

In order to turn on the automatic structure optimization in Quantum Espresso, we change the calculation type in the input file to 'vc-relax' instead of 'scf', which allows the structure to vary during the calculations. We also alter the 'ion dynamics' and 'cell dynamics' parameters in the IONS and CELL subfields of the input file to 'bfgs' to specify the algorithm used for nonlinear optimization. Using an input lattice parameter at the lower bound of the range we tested before (9.5 Bohr), a k-point input dimension of 11, and an energy cutoff of 35 Ry, the output file shows a final lattice parameter of 10.61 Bohr, identical to the value found from the automatic optimization. The energy per atom is equal out to three decimal places at -8.007 Ry/atom as well.

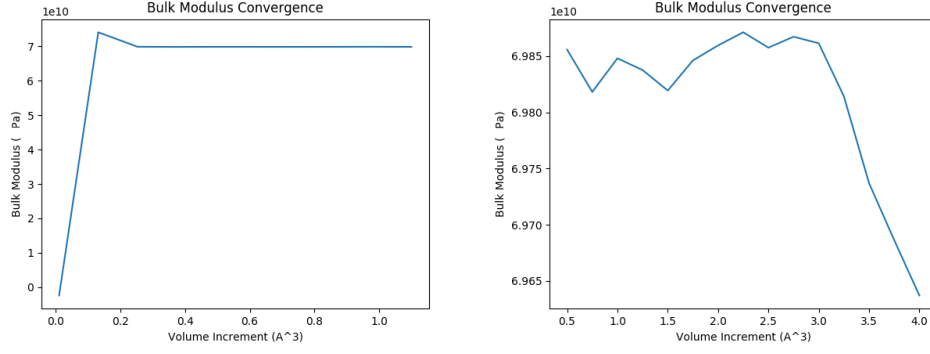


Figure 8: Bulk Modulus Convergence.

## 8 Bonus 1: Elastic Constants

A cubic crystal can be described by three elastic constants, denoted  $C_{11}$ ,  $C_{12}$ , and  $C_{44}$ , which represent normal stress-strain relationship in the same direction, normal strain in response to a stress in a different direction, and shear relationship, respectively. Using a complete unit cell of Ge with diamond cubic structure, we follow the procedure of Mehl et al. (*Intermetallic Compounds: Principles and Practice*, 1994) in "First principles calculations of elastic properties of metals," who perform volume-conserving transformations on a tetragonal cell to uncover the same elastic constants. However, these strains are performed with respect to the FCC basis vectors, which are defined as:

$$\begin{aligned}\vec{a}_1 &= \frac{a_0}{2}\vec{j} + \frac{a_0}{2}\vec{k} \\ \vec{a}_2 &= \frac{a_0}{2}\vec{i} + \frac{a_0}{2}\vec{k} \\ \vec{a}_3 &= \frac{a_0}{2}\vec{i} + \frac{a_0}{2}\vec{j}\end{aligned}$$

Here,  $a_0$  is the typical cubic lattice parameter. The function used in Python from ASE to create the structure (crystal) does not appear to create the correct structure if these basis vectors are passed in using the "cell" parameter when the spacegroup is also defined, instead creating a cell with the correct number of atoms but with a volume that is four times too small. Instead, we use the following relations to obtain the three typical orthorhombic vectors ( $\vec{a}$ ,  $\vec{b}$ , and  $\vec{c}$ ), from the primitive basis vectors:

$$\begin{aligned}\vec{a} &= \vec{a}_2 + \vec{a}_3 - \vec{a}_1 \\ \vec{b} &= \vec{a}_1 + \vec{a}_3 - \vec{a}_2 \\ \vec{c} &= \vec{a}_1 + \vec{a}_2 - \vec{a}_3\end{aligned}$$

With the bulk modulus already determined, we need two more constants, which are typically taken as the shear moduli  $C_{44}$  and  $C_{11} - C_{12}$ . The relation  $B = \frac{C_{11} + 2C_{12}}{3}$  can be used to complete the set and obtain all three elastic constants. To obtain  $C_{11} - C_{12}$ , we utilize a volume-conserving orthorhombic strain of the form:

$$\begin{aligned}e_1 &= -e_2 = x \\ e_3 &= x^2/(1 - x^2) \\ e_4 &= e_5 = e_6 = 0\end{aligned}$$

The strain tensor,  $\vec{\epsilon}$  is then constructed as:

$$\begin{bmatrix} e_1 & e_6/2 & e_5/2 \\ e_6/2 & e_2 & e_4/2 \\ e_5/2 & e_4/2 & e_3 \end{bmatrix}$$

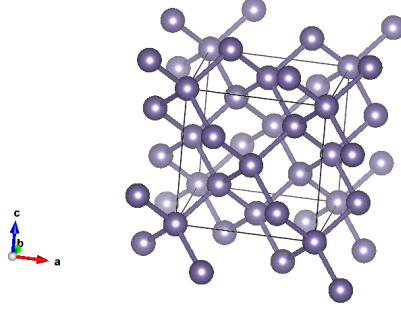


Figure 9: Germanium Unit Cell.

To obtain the new primitive vectors, we perform the matrix multiplication (where  $\vec{I}$  is the identity matrix):

$$\begin{bmatrix} \vec{a}'_1 \\ \vec{a}'_2 \\ \vec{a}'_3 \end{bmatrix} = \begin{bmatrix} \vec{a}_1 \\ \vec{a}_2 \\ \vec{a}_3 \end{bmatrix} \cdot [\vec{I} + \vec{\epsilon}]$$

Using this strain field, the volume of the cell remains constant, and the elastic constants can be found through:

$$\Delta E(x) = V(C_{11} - C_{12})x^2$$

We use a similar methodology for obtaining the shear modulus  $C_{44}$ , but our strains are instead defined as:

$$\begin{aligned} e_6 &= x \\ e_3 &= x^2/(4 - x^2) \\ e_1 &= e_2 = e_4 = e_5 = 0 \end{aligned}$$

This yields an energy difference of:

$$\Delta E(x) = \frac{Vx^2}{2}C_{44}$$

We implement these strain fields in two different runs of the Ge script 2 python code. In the first of the two runs, the strain field for  $C_{44}$  is used. Two energies are calculated- one corresponding to the normal lattice, and one to the strained lattice. The normal lattice is displayed above in Figure 9. Note that, for all calculations in this problem, the "nosym" parameter in the &SYSTEM input field is set to True, so no symmetries are assumed which could disrupt computation of the strained cell. The energy difference obtained through this calculation is found to be 0.00468 eV. Inserting this into the formula above, with a strain of 0.01 and a cell volume of 176.55 cubic Å, we obtain  $C_{44} = 84GPa$ . A similar procedure is used for the other strain field, for which we obtain an energy difference of 0.00898 eV. Using the same cell volume and strain, this yields  $C_{11} - C_{12} = 81GPa$ . Invoking our previously obtained bulk modulus of 70 GPa and the above formula for bulk modulus as a function of the two elastic constants, and solving the two equations simultaneously, we obtain  $C_{12} = 43GPa$  and  $C_{11} = 124GPa$ . Comparing our results with the literature, we consider the work of Bond et al. (*Physical Review*, 1950) "The Elastic Constants of Germanium Single Crystals." In this work, the elastic constants are obtained as  $C_{11} = 130GPa$ ,  $C_{12} = 49GPa$ , and  $C_{44} = 67GPa$ . Of our results, the only one that differs substantially is  $C_{44}$ . Given that we only test one strain value (due to the high computational cost of running a calculation with eight atoms without any symmetry, a cutoff of 35 Ry, and an 11 x 11 x 11 k-point grid), it is possible that this value had not converged with respect to the strain, and this may be one source of discrepancy. It is also worth noting that use of the experimental bulk modulus value, 76 GPa, yields values for the other two elastic constants that exactly match the above cited experimental results.



## 9 Bonus 2: Band Structure

To uncover the band structure of Ge, we change the calculation parameter to *bands* and specify a list of k-points along high symmetry directions. These directions are chosen as those specified in the ASE documentation, which are in turn drawn from Setyawana and Curtarolo (*Computational Materials Science*, 2010). In the Python script, these can be obtained using the `get special points` function from `ase.dft.kpoints`. The k-point path is chosen to align with that used in the literature ( $W-\Gamma-X-W-L-\Gamma$ ). We also specify the `nbnd` parameter to be 10 to account for the valence and first four conduction bands in the `&SYSTEM` input field. The parameter `disk io` is changed to 'low' instead of 'none,' which causes creation of wavefunction and energy density files along with the typical input and output files. Since multiple calculations from `pw.x` are necessary, the command line interface is used in lieu of a Python script. Prior to changing the calculation type, the typical 'scf' calculation is run first in the same folder, and then a 'bands' calculation is run using the same directory. On the command line, `pw.x < pwsf.in > pwsf.out` runs the first of these calculations, and `pw.x < pwsfband.in >` runs the second. The `bands.x` function is then used to sort the bands using the input file `bands.in`, and finally `plotband.x` is used to plot the results. The Fermi level was found to be -3.3 eV by shifting the value around until it corresponded with the approximate location of the band gap. The input files are included in the appendix below.

Figure 10 displays the result of the DFT band structure calculation as well as results from a thesis, *Development of Non-Local Density Functional Methods* by Dominik Bogdan Jochym at the University of Durham. Qualitatively, the results are somewhat similar, with the QE results showing more fluctuations in the individual paths. Importantly, the QE results do not show a band gap, while a small one is present in the thesis results. Therefore, using this method would incorrectly predict that germanium is metallic rather than a semiconductor.

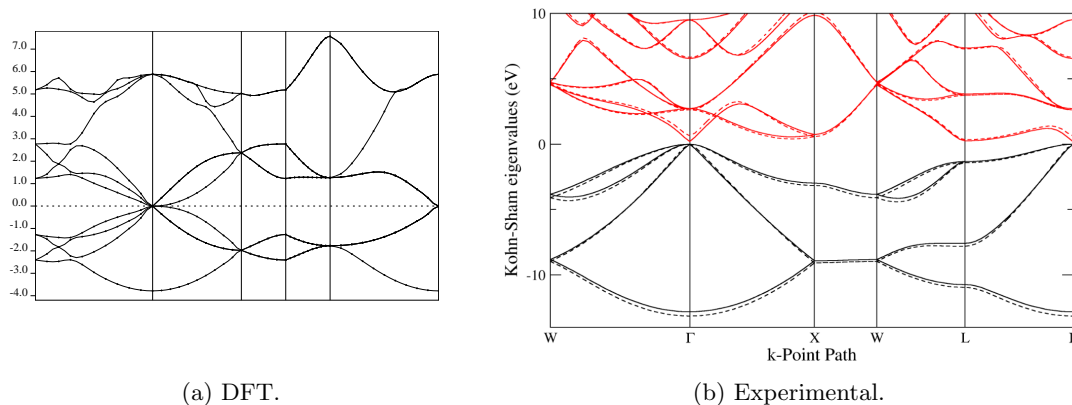


Figure 10: Germanium Band Structure.

## 10 Appendix

Attached below is the source code for the modified Ge script file. The major modifications are in the main function, where the Q parameter specifies which sub-region to enter. These values correspond to the question to be answered (and Q = 8 is simply used to run any value for the automatic optimization in part 7C). The displacement value in the make struc function is currently set to 0.00, and can be adjusted to 0.05 for problems 3 and 4, to add forces. The runpath in compute energy can also be changed to save to different folders depending on the question. An updated version of pwscf.py was also constructed that extracts the walltime in addition to energy, force, stress, and pressure to obtain the calculation times. The added code in lattice scan parses the obtained string to extract only the numbers and decimal point.

```
1  from my_labutil.src.plugins.pwscf import *
2  from ase.spacegroup import crystal
3  from ase.build import *
4  from ase.io import write
5  import matplotlib.pyplot as plt
6  import numpy as np
7
8
9  def make_struc(alat):
10     """
11     Creates the crystal structure using ASE.
12     :param alat: Lattice parameter in angstrom
13     :return: structure object converted from ase
14     """
15     # set primitive_cell=False if you want to create a simple cubic unit cell with 8 atoms
16     gecell = crystal('Ge', [(0, 0, 0)], spacegroup=227, cellpar=[alat, alat, alat, 90, 90,
17     90], primitive_cell=True)
18     gecell.positions[0][2] = gecell.positions[0][2] + 0.00
19     # change above to +0.05 for numbers 3 and 4
20     # check how your cell looks like
21     write('s.cif', gecell)
22     structure = Struc(ase2struc(gecell))
23     return structure
24
25  def compute_energy(alat, nk, ecut):
26     """
27     Make an input template and select potential and structure, and the path where to run
28     """
29     potname = 'Ge.pz-bhs.UPF'
30     pseudopath = os.environ['ESPRESSO_PSEUDO']
31     potpath = os.path.join(pseudopath, potname)
32     pseudopots = {'Ge': PseudoPotential(name=potname, path=potpath, ptype='uspp', element='
33     Ge', functional='LDA')}
34     struc = make_struc(alat=alat)
35     kpts = Kpoints(gridsize=[nk, nk, nk], option='automatic', offset=False)
36     runpath = Dir(path=os.path.join(os.environ['WORKDIR'], "Lab2/Problem7", str(ecut)+'.'+
37     str(nk)+'.'+str(alat)))
38     input_params = PWscf.inparam({
39         'CONTROL': {
40             'calculation': 'scf',
41             'pseudo_dir': pseudopath,
42             'outdir': runpath.path,
43             'tstress': True,
44             'tprnfor': True,
45             'disk_io': 'none',
46         },
47         'SYSTEM': {
48             'ecutwfc': ecut,
49         },
50         'ELECTRONS': {
51             'diagonalization': 'david',
52             'mixing_beta': 0.5,
53             'conv_thr': 1e-7,
```

```

53         'IONS': {
54
55         },
56         'CELL': {
57
58         },
59
60     })
61
62     output_file = run_qe_pwscf(runpath=runpath, struc=struc, pseudopots=pseudopots,
63                               params=input_params, kpoints=kpts)
64     output = parse_qe_pwscf_output(outfile=output_file)
65     return output
66
67 def lattice_scan(nk, ecut, alat):
68     #nk = 3
69     #ecut = 30
70     #alat = 5.0
71     output = compute_energy(alat=alat, ecut=ecut, nk=nk)
72     energy = output['energy']
73     kpts = output['kpoints']
74     forces = output['force']
75     walltime = output['walltime']
76     wallt = []
77     for t in walltime:
78         if str.isdigit(t) or '.' in t:
79             wallt.append(str(t))
80     #print(energy)
81     #print(wallt)
82     return energy, float(''.join(wallt)), kpts, forces
83
84
85 if __name__ == '__main__':
86     # put here the function that you actually want to run
87     Q = 7
88
89     if Q == 1:
90         nk = 4
91         energy = []
92         ediff = []
93         caltime = []
94         forces = []
95         alat = 5.0
96         for ecut in range(5, 85, 5):
97             energy.append(lattice_scan(nk=nk, ecut=ecut, alat=alat)[0]/(2*13.605698066))
98             caltime.append(lattice_scan(nk=nk, ecut=ecut, alat=alat)[1])
99             forces.append(lattice_scan(nk=nk, ecut=ecut, alat=alat)[3])
100         for en in range(0, len(energy)-1):
101             ediff.append(energy[len(energy)-1]-energy[en])
102
103         plt.plot(range(5, 85, 5), energy)
104         plt.xlabel('Cutoff Energy (Ry)')
105         plt.ylabel('Energy (Ry/atom)')
106         axes = plt.gca()
107         #axes.set_ylim([0, 0.2])
108         plt.title("Cutoff Convergence")
109         plt.show()
110
111         plt.plot(range(10, 85, 5), ediff)
112         plt.xlabel('Cutoff Energy (Ry)')
113         plt.ylabel('Energy Difference (Ry/atom)')
114         axes = plt.gca()
115         # axes.set_ylim([0, 0.2])
116         plt.title("Convergence Energy Change")
117         plt.show()
118
119         plt.plot(range(5, 85, 5), caltime)
120         plt.xlabel('Cutoff Energy (Ry)')

```

```

121     plt.ylabel('Calculation Time (s)')
122     axes = plt.gca()
123     # axes.set_ylim([0, 0.2])
124     plt.title("Calculation Time")
125     plt.show()
126
127 elif Q == 2:
128     ecut = 30
129     energy = []
130     ediff = []
131     caltime = []
132     kpts=[]
133     alat=5.0
134     for nk in range(1, 13):
135         energy.append(lattice_scan(nk=nk, ecut=ecut, alat=alat)[0] / (2 * 13.605698066))
136         caltime.append(lattice_scan(nk=nk, ecut=ecut, alat=alat)[1])
137         kpts.append(lattice_scan(nk=nk, ecut=ecut, alat=alat)[2])
138     for en in range(0, len(energy) - 1):
139         ediff.append(energy[len(energy)-1] - energy[en])
140     plt.plot(kpts, energy)
141     plt.xlabel('Unique K-Points')
142     plt.ylabel('Energy (Ry/atom)')
143     axes = plt.gca()
144     # axes.set_ylim([0, 0.2])
145     plt.title("K-Point Convergence")
146     plt.show()
147
148     plt.plot(kpts[1:len(kpts)], ediff)
149     plt.xlabel('Unique K-Points')
150     plt.ylabel('Convergence Energy Difference (Ry)')
151     axes = plt.gca()
152     # axes.set_ylim([0, 0.2])
153     plt.title("Convergence Energy Change")
154     plt.show()
155
156     plt.plot(kpts, caltime)
157     plt.xlabel('Unique K-Points')
158     plt.ylabel('Calculation Time (s)')
159     axes = plt.gca()
160     # axes.set_ylim([0, 0.2])
161     plt.title("Calculation Time")
162     plt.show()
163
164     plt.plot(range(1,13), energy)
165     plt.xlabel('Input K-Points')
166     plt.ylabel('Energy (Ry/atom)')
167     axes = plt.gca()
168     # axes.set_ylim([0, 0.2])
169     plt.title("K-Point Convergence")
170     plt.show()
171
172     plt.plot(range(2,13), ediff)
173     plt.xlabel('Input K-Points')
174     plt.ylabel('Convergence Energy Difference (Ry)')
175     axes = plt.gca()
176     # axes.set_ylim([0, 0.2])
177     plt.title("Convergence Energy Change")
178     plt.show()
179
180     plt.plot(range(1,13), caltime)
181     plt.xlabel('Input K-Points')
182     plt.ylabel('Calculation Time (s)')
183     axes = plt.gca()
184     # axes.set_ylim([0, 0.2])
185     plt.title("Calculation Time")
186     plt.show()
187
188     plt.plot(range(1,13), kpts)

```

```

189     plt.xlabel('Input K-Points')
190     plt.ylabel('Unique K-Points')
191     axes = plt.gca()
192     # axes.set_ylim([0, 0.2])
193     plt.title("K-Points")
194     plt.show()
195
196 elif Q == 3:
197     forces=[]
198     fdiff=[]
199     alat=5.0
200     for ecut in range(5, 85, 5):
201         forces.append(lattice_scan(nk=4, ecut=ecut, alat=alat)[3]/2)
202         print(forces)
203
204     for en in range(0, len(forces)-1):
205         fdiff.append(forces[len(forces)-1]-forces[en])
206
207     plt.plot(range(5, 85, 5), forces)
208     plt.xlabel('Cutoff Energy (Ry)')
209     plt.ylabel('Force (Ry/Bohr)')
210     axes = plt.gca()
211     # axes.set_ylim([0, 0.2])
212     plt.title("Cutoff Convergence")
213     plt.show()
214
215     plt.plot(range(10, 85, 5), fdiff)
216     plt.xlabel('Cutoff Energy (Ry)')
217     plt.ylabel('Force Increment (Ry/Bohr)')
218     axes = plt.gca()
219     # axes.set_ylim([0, 0.2])
220     plt.title("Cutoff Increment")
221     plt.show()
222
223 elif Q == 4:
224     forces = []
225     fdiff = []
226     alat=5.0
227     for nk in range(1,15):
228         forces.append(lattice_scan(nk=nk, ecut=30, alat=alat)[3]/2)
229         print(forces)
230
231     for en in range(0, len(forces) - 1):
232         fdiff.append(forces[len(forces)-1] - forces[en])
233
234     plt.plot(range(1,15), forces)
235     plt.xlabel('Input Grid Dimension')
236     plt.ylabel('Force (Ry/Bohr)')
237     axes = plt.gca()
238     # axes.set_ylim([0, 0.2])
239     plt.title("K-Point Convergence")
240     plt.show()
241
242     plt.plot(range(2,15), fdiff)
243     plt.xlabel('Input Grid Dimension')
244     plt.ylabel('Force Increment (Ry/Bohr)')
245     axes = plt.gca()
246     # axes.set_ylim([0, 0.2])
247     plt.title("K-Point Increment")
248     plt.show()
249
250 elif Q == 5:
251     e1=[]
252     e2=[]
253     alat1=(10.70*0.529177249)
254     alat2=(10.75*0.529177249)
255     ediff=[]
256     ediff2=[]

```

```

257 nk=4
258 for ecut in range(5,85,5):
259     eltemp=lattice_scan(nk=nk,ecut=ecut,alat=alat1)[0]/(2 * 13.605698066)
260     e1.append(eltemp)
261     e2temp=lattice_scan(nk=nk,ecut=ecut,alat=alat2)[0]/(2 * 13.605698066)
262     e2.append(e2temp)
263     ediff.append(e2temp-e1temp)
264
265 for en in range(0, len(ediff) - 1):
266     ediff2.append(ediff[len(ediff)-1] - ediff[en])
267
268 plt.plot(range(5,85,5),ediff)
269 plt.xlabel('Cutoff Energy (Ry)')
270 plt.ylabel('Energy Difference (Ry/atom)')
271 axes = plt.gca()
272 plt.title('Energy Difference vs. Cutoff')
273 plt.show()
274
275 plt.plot(range(10, 85, 5), ediff2)
276 plt.xlabel('Cutoff Energy (Ry)')
277 plt.ylabel('Energy Difference Convergence (Ry/atom)')
278 axes = plt.gca()
279 plt.title('Energy Difference Convergence vs. Cutoff')
280 plt.show()
281
282 elif Q == 7:
283     nk=11
284     ecut=35
285     energy=[]
286     for alat in np.linspace(10,11.1,10):
287         alat=alat*0.529177249
288         energy.append(lattice_scan(nk=nk,ecut=ecut,alat=alat)[0]/(2 * 13.605698066))
289
290 plt.plot(np.linspace(10,11.1,10),energy)
291 plt.xlabel('Lattice Parameter (Bohr)')
292 plt.ylabel('Energy (Ry/atom)')
293 plt.title('Lattice Parameter Convergence')
294 plt.show()
295
296 B=[]
297 for v in np.linspace(0.01,1.1,10):
298     ind_min = np.argmin(energy)
299     alat_min=0.529177249*np.linspace(10,11.1,10)[ind_min]
300     V_min=alat_min**3/(4)
301     alatf2=(4*(V_min+2*v))**(1/3)
302     alatf1=(4*(V_min+v))**(1/3)
303     alatfn1=(4*(V_min-v))**(1/3)
304     alatfn2=(4*(V_min-2*v))**(1/3)
305     f2=lattice_scan(nk=nk,ecut=ecut,alat=alatf2)[0]
306     f1=lattice_scan(nk=nk,ecut=ecut,alat=alatf1)[0]
307     f0=(2 * 13.605698066)*energy[ind_min]
308     fn1=lattice_scan(nk=nk,ecut=ecut,alat=alatfn1)[0]
309     fn2=lattice_scan(nk=nk,ecut=ecut,alat=alatfn2)[0]
310
311     d2EdV2=(-f2+16*f1-30*f0+16*fn1-fn2)/(12*v**2)
312     B.append(V_min*((10**(10))**3*(1.6*10**(-19)))*d2EdV2)
313     print(alat_min)
314     print(V_min)
315     print(f2)
316     print(f1)
317     print(f0)
318     print(fn1)
319     print(fn2)
320     print(B)
321
322 plt.plot(np.linspace(0.01,1.1,10), B)
323 plt.xlabel('Volume Increment (A^3)')
324 plt.ylabel('Bulk Modulus (GPa)')

```

```

325     plt.title('Bulk Modulus Convergence')
326     plt.show()
327
328     elif Q == 8:
329         lattice_scan(11,35,5.0)
330

```

Included below is a modified version of the above script to be used with the bonus questions. Here, a number of notable changes are made for each question. For the first bonus, the input to the crystal function is changed to a 3 x 3 matrix instead of the [alat,alat,alat,90,90,90] form used previously, in order to specify the strained lattice vectors properly. Each function from lattice scan to compute energy, to make struc now accepts the matrix a, which contains these vectors, in addition to the lattice parameter alat. For this question, the primitive argument is also set to False to obtain the complete unit cell, and the nosym parameter in &SYSTEM is set to True so no symmetries are assumed. Both strain fields are present, with the current state being that used to obtain the quantity  $C_{11} - C_{12}$ . The commented version can be pasted in for  $C_{44}$ . The matrix a contains the unstrained lattice vectors, and the lines using at are simply there to ensure that the transformation from primitive basis vectors to unit cell vectors reconstructs the correct matrix. The an matrix performs the same transformation to obtain the strained unit cell vectors from the strained basis vectors.

```

1  from my_labutil.src.plugins.pwscf import *
2  from ase.spacegroup import crystal
3  from ase.build import *
4  from ase.io import write
5  import matplotlib.pyplot as plt
6  import numpy as np
7  import ase.dft.kpoints as k
8
9
10 def make_struc(alat, a):
11     """
12     Creates the crystal structure using ASE.
13     :param alat: Lattice parameter in angstrom
14     :return: structure object converted from ase
15     """
16     # set primitive_cell=False if you want to create a simple cubic unit cell with 8 atoms
17     gecell = crystal('Ge', [(0, 0, 0)], spacegroup=227, cell=a, primitive_cell=False)
18     #check how your cell looks like
19     gecell.positions[0][2] = gecell.positions[0][2] + 0.00
20     ge = bulk('Ge', 'diamond', a=5.61)
21     points = k.get_special_points('fcc', ge.cell)
22     print(points)
23     write('s2.cif', gecell)
24     structure = Struc(ase2struc(gecell))
25     return structure
26
27
28 def compute_energy(alat, nk, ecut, a):
29     """
30     Make an input template and select potential and structure, and the path where to run
31     """
32     potname = 'Ge.pz-bhs.UPF'
33     pseudopath = os.environ['ESPRESSO_PSEUDO']
34     potpath = os.path.join(pseudopath, potname)
35     pseudopots = {'Ge': PseudoPotential(name=potname, path=potpath, ptype='uspp', element='
Ge', functional='LDA')}
36     struc = make_struc(alat=alat, a=a)
37     kpts = Kpoints
38     (6, [[0, 0, 0, 1], [0.375, 0.375, 0.75, 1], [0.5, 0.5, 0.5, 1], [0.625, 0.625, 0.625, 1], [0.5, 0.25, 0.75, 1], [0.5, 0, 0.5,
option='crystal_b', offset=False])
39     print(kpts)
40     runpath = Dir(path=os.path.join(os.environ['WORKDIR'], "Lab2/Problem9/", str(ecut)+'.'+
str(nk)+'.'+str(alat)))
41     input_params = PWscf.inparam({
42         'CONTROL': {

```

```

43         'pseudo_dir': pseudopath,
44         'outdir': runpath.path,
45         'tstress': True,
46         'tprnfor': True,
47         'disk_io': 'low',
48     },
49     'SYSTEM': {
50         'ecutwfc': ecut,
51         'nosym': False,
52         'nbnd': 5
53     },
54     'ELECTRONS': {
55         'diagonalization': 'david',
56         'mixing_beta': 0.5,
57         'conv_thr': 1e-7,
58     },
59     'IONS': {
60
61     },
62     'CELL': {
63
64     },
65     'K.POINTS': {
66         '{crystal_b}'
67         '6'
68         '0 0 0 1'
69         '0.375 0.375 0.75 1'
70         '0.5 0.5 0.5 1'
71         '0.625 0.25 0.625 1'
72         '0.5 0.25 0.75 1'
73         '0.5 0 0.5 1'
74     }
75 })
76
77
78 output_file = run_qe_pwscf(runpath=runpath, struc=struc, pseudopots=pseudopots,
79                             params=input_params, kpoints=kpts)
80 output = parse_qe_pwscf_output(outfile=output_file)
81 return output
82
83
84
85 def lattice_scan(nk, ecut, alat, a):
86     #nk = 3
87     #ecut = 30
88     #alat = 5.0
89     output = compute_energy(alat=alat, ecut=ecut, nk=nk, a=a)
90     energy = output['energy']
91     kpts = output['kpoints']
92     forces = output['force']
93     walltime = output['walltime']
94     wallt = []
95     if isinstance(walltime, int):
96         wallout=walltime
97     else:
98         for t in walltime:
99             if str.isdigit(t) or '.' in t:
100                 wallt.append(str(t))
101             wallout=float(''.join(wallt))
102     #print(energy)
103     #print(wallt)
104     return energy, wallout, kpts, forces
105
106
107 if __name__ == '__main__':
108     # put here the function that you actually want to run
109     Q = 9
110

```



```

111     if Q == 8:
112         x=0.01
113         alat=5.61
114         e1=x
115         e2=-x
116         e3=x**2/(1-x**2)
117         e4=0
118         e5=0
119         e6=0
120         e=np.array ([[ e1 , e6 /2 , e5 /2 ] , [ e6 /2 , e2 , e4 /2 ] , [ e5 /2 , e4 /2 , e3 ]])
121         a1=[0, alat , alat ]
122         a2=[alat , 0 , alat ]
123         a3=[alat , alat , 0]
124         a=np.array ([ a1 , a2 , a3 ])
125         a=a/2
126         print (a)
127         a2=a
128         at1=a[1]+a[2]-a[0]
129         at2=a[2]+a[0]-a[1]
130         at3=a[0]+a[1]-a[2]
131         a=np.array ([ at1 , at2 , at3 ])
132         e=e+np.identity (3)
133         print (e)
134         anew=np.dot (a2 , e)
135         print (anew)
136         an1=anew[1]+anew[2]-anew[0]
137         an2=anew[2]+anew[0]-anew[1]
138         an3=anew[0]+anew[1]-anew[2]
139         an=np.array ([ an1 , an2 , an3 ])
140
141         print (a)
142         print (an)
143         E1=lattice_scan (nk=11,ecut=35,alat=5.61, a=a) [0]
144         E2=lattice_scan (nk=11,ecut=35,alat=5.61, a=an) [0]
145         print (E1)
146         print (E2)
147
148         ', '
149         For C44 instead ,
150
151         e1=0
152         e2=0
153         e4=0
154         e5=0
155         e6=x
156         e3=(x)**2/(4-x**2)
157         ', '
158
159     elif Q == 9:
160         a=1
161         lattice_scan (nk=11,ecut=50,alat=5.61, a=a)
162

```

For the final question, not all of the code is included below. Instead, the input files for pw.x and bands.x are provided. Initially, an scf calculation is run, using only the primitive lattice in the same setup as the previous (non-bonus) questions with no displacement. Once this is complete, the calculation is changed to 'bands', the 'nbnd' parameter in &SYSTEM is set to 5, the 'disk\_io' parameter in &CONTROL is set to 'low' so the necessary wavefunction and pwscf.save files are created, and the K\_POINTS input is added below CELL to specify the special k-point directions sampled in the band structure calculation. The first section is used for the typical scf calculation with a cutoff of 35 Ry and an 11 x 11 x 11 grid.

```

1 &CONTROL
2     disk_io = 'low'
3     outdir = '/home/bond/WORK/Lab2/Problem9/35.11.5.61'
4     tprnfor = .true.
5     tstress = .true.

```

```

6      calculation = 'scf'
7      pseudo_dir = '/home/bond/Software/qe-6.0/pseudo'
8  /
9  &SYSTEM
10     ecutwfc = 35
11     nat = 8
12     nosym = .false.
13     ntyp = 1
14     ibrav = 0
15  /
16  &ELECTRONS
17     conv_thr = 1e-07
18     diagonalization = 'david'
19     mixing_beta = 0.5
20  /
21  &IONS
22  /
23  &CELL
24  /
25  K_POINTS {automatic}
26  11 11 11 0 0 0
27  ATOMIC_SPECIES
28  Ge 72.61 Ge.pz-bhs.UPF
29  CELL_PARAMETERS {angstrom}
30  5.61 0.0 0.0
31  3.435134271608326e-16 5.61 0.0
32  3.435134271608326e-16 3.435134271608326e-16 5.61
33  ATOMIC_POSITIONS {angstrom}
34  Ge 0.00000 0.00000 0.00000
35  Ge 0.00000 2.80500 2.80500
36  Ge 2.80500 2.80500 0.00000
37  Ge 2.80500 0.00000 2.80500
38  Ge 4.20750 1.40250 4.20750
39  Ge 1.40250 1.40250 1.40250
40  Ge 1.40250 4.20750 4.20750
41  Ge 4.20750 4.20750 1.40250

```

The second input file is used for the band structure calculation with 10 bands and the previously specified k-point path in reciprocal space. The weights next to each k-point were chosen empirically so that the number of points along each section of the path were similar.

```

1  &CONTROL
2      disk_io = 'low'
3      outdir = '/home/bond/WORK/Lab2/Problem9/35.11.5.61'
4      tprnfor = .true.
5      tstress = .true.
6      calculation = 'bands'
7      pseudo_dir = '/home/bond/Software/qe-6.0/pseudo'
8  /
9  &SYSTEM
10     ecutwfc = 35
11     nat = 8
12     nosym = .false.
13     ntyp = 1
14     nbnd = 10
15     ibrav = 0
16  /
17  &ELECTRONS
18     conv_thr = 1e-07
19     diagonalization = 'david'
20     mixing_beta = 0.5
21  /
22  &IONS
23  /
24  &CELL
25  /
26  K_POINTS crystal_b
27  6

```

```

28      0.5      0.25      0.75      10.0
29      0.0000000000      0.0000000000      0.0000000000      10.0
30      0.5      0.00      0.5      5.0
31      0.5      0.25      0.75      10.0
32      0.5      0.5      0.5      10.0
33      0.0      0.0      0.0      10.0
34 ATOMIC_SPECIES
35   Ge 72.61 Ge.pz-bhs.UPF
36 CELL_PARAMETERS {angstrom}
37   5.61 0.0 0.0
38   3.435134271608326e-16 5.61 0.0
39   3.435134271608326e-16 3.435134271608326e-16 5.61
40 ATOMIC_POSITIONS {angstrom}
41   Ge 0.00000 0.00000 0.00000
42   Ge 0.00000 2.80500 2.80500
43   Ge 2.80500 2.80500 0.00000
44   Ge 2.80500 0.00000 2.80500
45   Ge 4.20750 1.40250 4.20750
46   Ge 1.40250 1.40250 1.40250
47   Ge 1.40250 4.20750 4.20750
48   Ge 4.20750 4.20750 1.40250

```

Finally, the bands.x input file specifies the directory, filename, output filename, and a parameter which prevents crossovers during plotting.

```

1 &BANDS
2   outdir = '/home/bond/WORK/Lab2/Problem9/35.11.5.61/ ',
3   prefix = 'pwsf',
4   filband = 'Ge.bands.dat',
5   no_overlap= .true.
6 /

```