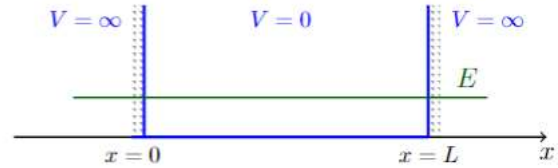


Deep Dive #10 - Quantum Mechanics

Its a good day to be a physics major!

$$-\frac{\hbar^2}{2m} \psi''(x) = E \psi(x),$$

$$\psi(0) = 0, \quad \psi(L) = 0.$$



Question 1:

(1a) (20 points) Find all the eigenvalues and eigenfunctions in the problem above.

$$\Psi''(x) + \frac{2mE}{\hbar^2} \Psi(x) = 0$$

$$\lambda = \frac{2mE}{\hbar^2}$$

Can energy be negative? In this situation that would yield an exponential Ψ function which cannot be zero at two points. For $E = 0$ this is trivial, so we'll consider $E > 0$

The characteristic polynomial is

$$r^2 + \frac{2mE}{\hbar^2} = 0$$

$$r = \pm i \frac{\sqrt{2mE}}{\hbar}$$

This yields

$$\Psi(x) = c_+ \cos\left(\frac{\sqrt{2mE}}{\hbar} x\right) + c_- \sin\left(\frac{\sqrt{2mE}}{\hbar} x\right)$$

$$\Psi(0) = 0 \implies 0 = c_+ \cos(0) \implies c_+ = 0$$

$$\Psi(L) = 0 \implies 0 = c_- \sin\left(\frac{\sqrt{2mE}}{\hbar} L\right)$$

With $c \neq 0$, we have $\sin(n\pi) = 0$ this gives us

$$\frac{\sqrt{2mE}}{\hbar} L = n\pi$$

$$\frac{\sqrt{2mE}}{\hbar} = \frac{n\pi}{L}$$

$$\frac{2mE}{\hbar^2} = \left(\frac{n\pi}{L}\right)^2$$

Thus all eigenvalues are $\left(\frac{n\pi}{L}\right)^2$, and the allowed energies are

$$E = \frac{n^2 \pi^2 \hbar^2}{2mL^2}$$

and the eigenfunctions are

$$y_n(x) = \sin(\sqrt{\lambda_n}x) = \sin\left(\frac{n\pi}{L}x\right)$$

for $c_- = 1$ (eigenfunctions can be scaled)

(1b)

$$L = a = 3$$

$$\alpha = 1, 2, 3, \dots, n \implies \lambda_n = \left(\frac{\pi}{3}\right)^2, \left(\frac{2\pi}{3}\right)^2, \left(\frac{3\pi}{3}\right)^2, \dots, \left(\frac{n\pi}{L}\right)^2$$

(2a) (20 points) Split the solution function ψ as follows,

$$\psi(x) = \begin{cases} \psi_1(x) & \text{for } 0 \leq x \leq a, \\ \psi_2(x) & \text{for } x \geq a, \end{cases}$$

Show that if you split the Schrödinger differential equation in the same way you can write it as

$$\begin{aligned} \psi_1'' &= -k^2 \psi_1, & 0 \leq x \leq a, \\ \psi_2'' &= \kappa^2 \psi_2, & x \geq a, \end{aligned}$$

where k and κ are appropriate positive constants given in terms of m , \hbar , V_0 , and $|E|$. Find the formulas for $k > 0$ and $\kappa > 0$ in terms of the mass m , Planck constant \hbar , potential $V_0 > 0$, and energy $|E| > 0$. Also find one the boundary condition for ψ_1 and one boundary condition for ψ_2 , obtained from the boundary conditions on ψ at $x = 0$ and at $x \rightarrow \infty$.

For $x < a$, $E - V > 0$ and we are in the well. For $x > a$, $E - V < 0$ we are outside the well, classically impossible.

Schrodigers equation can be split in the same way as $V(x)$

$$\psi''(x) = \begin{cases} \frac{2m}{\hbar^2}(-V_0 - E) \psi(x) & \text{for } 0 \leq x \leq a \\ -\frac{2m}{\hbar^2} E \psi(x) & \text{for } x \geq a \end{cases}$$

Remember that $E < 0$ and can be represented correctly as $E = -|E|$

$$\psi''(x) = \begin{cases} -\frac{2m}{\hbar^2}(V_0 - |E|) \psi(x) & \text{for } 0 \leq x \leq a \\ \frac{2m}{\hbar^2} |E| \psi(x) & \text{for } x \geq a \end{cases}$$

This gets us our k -values

$$-k^2 = -\frac{2m}{\hbar^2}(V_0 - |E|)$$

$$\kappa^2 = \frac{2m}{\hbar^2}|E|$$

Also $\psi_1(0) = \psi_2(\infty) = 0$

2b

To be precise, there is one condition $E = -|E|$

The characteristic polynomial for ψ_1 yields

$$r_1^2 = -k^2$$

$$r_1 = \pm ik$$

$$\psi_1(x) = c_+ \cos(kx) + c_- \sin(kx)$$

And for ψ_2 yield

$$r_2^2 = \kappa^2$$

$$r_2 = \pm \kappa$$

$$\psi_2(x) = c_0 e^{\kappa x} + c_1 e^{-\kappa x}$$

From applying boundary conditions we see

$$\psi_1(0) = 0 \implies c_+ = 0$$

$$\psi_2(\infty) = 0 \implies c_0 = 0$$

There's another condition at $x = a$, but we get to that later

proportional to an exponential; let's call the proportionality factor d . The third part of the problem is to match the functions ψ_1 and ψ_2 found in (2b) at $x = 0$. Impose the matching conditions

$$\psi_1(a) = \psi_2(a), \quad \psi_1'(a) = \psi_2'(a).$$

From these equations find a relation between k and κ and between the scaling factor for ψ_1 , which we called it c , and the scaling factor for ψ_2 , which we called it d , of the form

$$\frac{k}{\kappa} = -f(ka), \quad d = c g(k, \kappa) e^{\kappa a},$$

Find the functions $f(ka)$ and $g(k, \kappa)$. These functions do not depend on c or d .

Lets do it.

$$\begin{aligned}
\psi_1(a) &= c \sin(k * a) \\
\psi_2(a) &= d e^{-\kappa a} \\
\psi'_1(a) &= c k \cos(k * a) \\
\psi'_2(a) &= -d \kappa e^{-\kappa * a} \\
\implies c k \cos(k * a) + d \kappa e^{-\kappa * a} &= 0 \\
\implies c \sin(k * a) - d e^{-\kappa a} &= 0 \\
\kappa(c \sin(k * a) - d e^{-\kappa a}) &
\end{aligned}$$

We can add the zeroes together

$$\begin{aligned}
c k \cos(k * a) + d \kappa e^{-\kappa * a} + c \kappa \sin(k * a) - \kappa d e^{-\kappa a} &= 0 \\
c k \cos(k * a) &= -c \kappa \sin(k * a) \\
\frac{k}{\kappa} &= -\frac{\sin(ka)}{\cos(ka)} - f(ka)
\end{aligned}$$

g is simple

$$\begin{aligned}
d e^{-\kappa a} &= c \sin(k * a) \\
d &= \frac{c \sin(k * a)}{e^{-\kappa a}} = c \sin(k * a) e^{\kappa a} \\
g(k) &= \sin(k * a)
\end{aligned}$$

Wait, is it simple?

$$\begin{aligned}
\implies c k \cos(k * a) + d \kappa e^{-\kappa * a} &= 0 \\
d \kappa e^{-\kappa * a} &= -c k \cos(k * a) \\
d &= c \left(-\frac{k}{\kappa} \cos(k * a) \right) e^{\kappa a} \\
g(k, \kappa) &= -\frac{k}{\kappa} \cos(k * a)
\end{aligned}$$

I'm not quite sure if the first formula is correct, but I'll leave it in anyways since I know the second one is what we were going for.

(3) there are two -9.2, -3.2, also trivially 0?