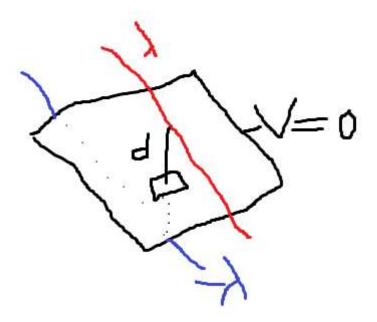
Phy 481 homework #7

1 The method of images

Griffiths Problem 3.10: A uniform line charge λ is placed on an infinite straight wire, a distance d above a grounded conducting plane. (Let's say the wire runs parallel to the x-axis and directly above it, and the conducting plane is the xy-plane.)

- 1. Find the potential in the region above the plane.
- 2. Find the charge density σ induced on the conducting plane.

Due to uniqueness, we can use the method of images to show how the grounded plane responds to the line charge density λ at height d above it. A line charge density of $-\lambda$ at -d would counteract the electric potential made by this.



1. In this situation, $\frac{dV}{dx}=-E_x=0$ since there the charge distribution is infinitely long in the x direction.

Using cylinder gaussian shells, we can find the electric field associated with any point due to a single charge distribution

$$abla \cdot E = rac{Q_{enc}}{\epsilon_0} \ 2\pi r l E = rac{\lambda l}{\epsilon_0} \ E(r) = rac{\lambda}{2\pi r \epsilon_0} \hat{r}$$

And the associated voltage is

$$\Delta V = -\int_{O}^{r} E(r) dr = -\frac{\lambda}{2\pi\epsilon_{0}} \int_{O}^{r} \frac{1}{r} dr = -\frac{\lambda}{2\pi\epsilon_{0}} \ln\left(\frac{r}{O}\right) = \frac{\lambda}{2\pi\epsilon_{0}} \ln\left(\frac{O}{r}\right)$$

The term r here is distance from the charge distribution, while O is a point where voltage is zero. By superposition, we can represent each charge distribution as this ΔV if we find the distance from each line in terms of y and z. r_1 is for λ while r_2 is for $-\lambda$

$$r_1 = \sqrt{y^2 + (z - d)^2}$$

$$r_2 = \sqrt{y^2 + (z + d)^2}$$

$$V(r_1, r_2) = \Delta V_1 + \Delta V_2 = \frac{\lambda}{2\pi\epsilon_0} \left[\ln\left(\frac{O}{r_1}\right) - \ln\left(\frac{O}{r_2}\right) \right]$$

$$= \frac{\lambda}{2\pi\epsilon_0} \ln\left(\frac{O/r_1}{O/r_2}\right) = \frac{\lambda}{2\pi\epsilon_0} \ln\left(\frac{r_2}{r_1}\right)$$

$$V(x, y) = \frac{\lambda}{2\pi\epsilon_0} \ln\left(\frac{\sqrt{y^2 + (z + d)^2}}{\sqrt{y^2 + (z - d)^2}}\right)$$

$$= \frac{\lambda}{4\pi\epsilon_0} \ln\left(\frac{y^2 + (z + d)^2}{y^2 + (z - d)^2}\right)$$

2. Using $\hat{n} = \hat{z}$,

$$\sigma = -\epsilon_0 \frac{\partial V}{\partial z}$$

Let's evaluate the derivative at z=0

$$\sigma = -\epsilon_0 rac{\lambda}{4\pi\epsilon_0} rac{\partial}{\partial z} \left[\ln \left(rac{y^2 + (z+d)^2}{y^2 + (z-d)^2}
ight)
ight]_{z=0}$$

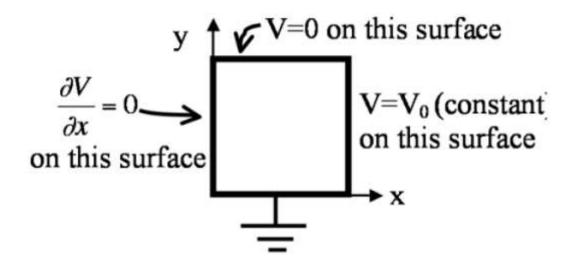
$$= rac{\lambda}{4\pi} \left[rac{2d}{y^2 + d^2} - rac{-2d}{y^2 + d^2}
ight]$$

$$= -rac{\lambda}{\pi} rac{d}{y^2 + d^2}$$

2 Rectangular Pipe: Separation of Variables-Cartesian-2D

A square rectangular pipe (sides of length a) runs parallel to the z-axis (from $-\infty$ to ∞) The 4 sides are maintained with boundary conditions given in the figure. (Each of the 4 sides is insulated from the others at the corners).

- 1. Find the potential V(x, y, z) at all points in this pipe.
- 2. Find the charge density $\sigma(x, y = 0, z)$ everywhere on the bottom conducting wall (y = 0). Check the units for your charge density (show us!).



1. Lets do it!

Can safely ignore z due to z invariance

$$abla^2 V = rac{\partial^2 V}{\partial x^2} + rac{\partial^2 V}{\partial y^2} = 0$$

Using
$$V(x,y) = X(x)Y(y)$$

$$= Y \frac{\partial^2 X}{\partial x^2} + X \frac{\partial^2 Y}{\partial Y^2}$$

$$\rightarrow \frac{1}{X} \frac{\partial^2 X}{\partial x^2} + \frac{1}{Y} \frac{\partial^2 YV}{\partial Y^2} = 0$$

$$\frac{1}{X} \frac{\partial^2 X}{\partial x^2} = -\frac{1}{Y} \frac{\partial^2 YV}{\partial Y^2}$$

Each term is constant as the other varies.

$$rac{1}{X}rac{\partial^2 X}{\partial x^2} = k^2$$

$$rac{1}{Y}rac{\partial^2 YV}{\partial Y^2} = -k^2$$

$$X(x) = Ae^{kx} + Be^{-kx}$$

$$Y(y) = C\cos(ky) + D\sin(ky)$$

These are standard second order differential equations. We can actually swap which is sine/cosine dependant and which is exponential. I chose these since they fit well with the the left wall boundary condition.

Apply boundary conditions

$$Y(0) = 0
ightarrow C = 0$$
 $Y(a) = 0
ightarrow ka = n\pi
ightarrow k = rac{n\pi}{a}$
 $Y(y) = D\sin(ky) = D\sin\left(rac{n\pi}{a}y
ight)$

Sine is zero at multiples of π . Next we'll apply the left wall boundary condition. Since it is dependant on x, Y will be constant.

$$\frac{\partial V}{\partial x}\Big|_{x=0} = 0 \to Yk \left[Ae^{k(0)} - Be^{-k(0)} \right] = 0 \to Yk[A - B] = 0 \to A = B$$

$$X(x) = Ae^{kx} + Ae^{-kx} = 2A\cosh(kx)$$

$$V(x,y) = \sum_{n} D_n \cosh\left(\frac{n\pi}{a}x\right) \sin\left(\frac{n\pi}{a}y\right)$$

Where we absorb 2A into D_n without loss of generality

To find the value of each D_n , we'll apply forier's trick using sine and integrating over the relevant length [0, a]. We'll also apply $V(x=a,y) = V_0$ to help solve for the coefficients.

$$V(x = a, y) = V_0 = \sum_n D_n \cosh(n\pi) \sin\left(\frac{n\pi}{a}y\right)$$

Multiply by $\sin(\frac{m\pi}{a}y)$ and integrate

$$\int_0^a V_0 \sin\left(\frac{m\pi}{a}y\right) dy = \sum_n D_n \cosh(n\pi) \int_0^a \sin\left(\frac{n\pi}{a}y\right) \sin\left(\frac{m\pi}{a}y\right) dy$$

LHS

$$\int_0^a V_0 \sin\left(\frac{m\pi}{a}y\right) dy = \frac{aV_0}{m\pi} - \frac{aV_0}{m\pi} \cos(m\pi)$$
$$= \begin{cases} 0 & m = 0 \pmod{2} \\ \frac{-2aV_0}{m\pi} & m = 1 \pmod{2} \end{cases}$$

$$\sum_n D_n \cosh(n\pi) \int_0^a \sin\Bigl(rac{n\pi}{a}y\Bigr) \sin\Bigl(rac{m\pi}{a}y\Bigr) dy = rac{a}{2} C_m \cosh(m\pi)$$

 $D_m = 0$ for even values of m_i , otherwise we have

$$rac{-2aV_0}{m\pi} = rac{a}{2}C_m\cosh(m\pi) \ C_m = rac{4V_0}{m\pi\cosh(m\pi)}$$

Now we have our voltage!

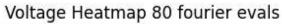
$$V(x,y) = \frac{4V_0}{\pi} \sum_{n=1,3,5,\dots} \frac{\cosh(\frac{n\pi}{a}x)}{\cosh(n\pi)} \frac{\sin(\frac{n\pi}{a}y)}{n}$$

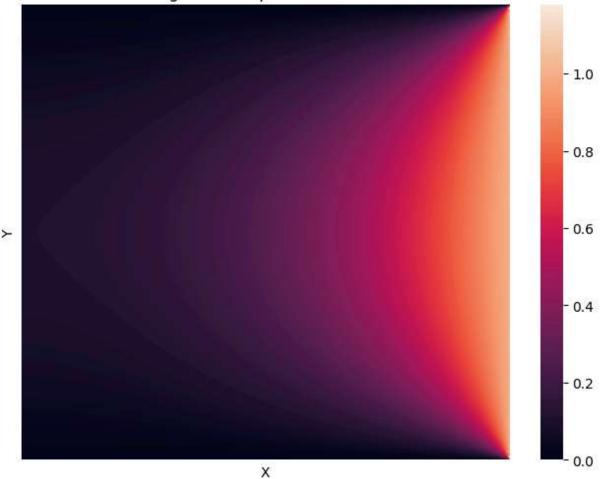
Here's what this looks like:

```
In [34]: import numpy as np
         from numpy import cosh, pi, sin
         import matplotlib.pyplot as plt
         import seaborn as sns
         from seaborn import cm
         fourier_size = 80 # watch out for overflow
         V0 = 1
         a = 1
         n2 = 2000
         x = np.linspace(0, a, n2)
         y = np.linspace(0, a, n2)
         X, Y = np.meshgrid(x, y)
         def phi(x, y, fourier_size, params):
             V0, a = params
             potential = 0
             for i in range(fourier_size):
                  n = 2*i + 1
                  potential += \cosh(n*pi*x/a) * \sin(n*pi*y/a) / (n*cosh(n*pi))
             return 4*V0/pi * potential
         params = (V0, a)
         meshout = phi(X, Y, fourier size, params)
         plt.figure(figsize=(8, 6))
         sns.heatmap(meshout, cmap=cm.rocket, xticklabels=False, yticklabels=False)
```

```
plt.title(f'Voltage Heatmap {fourier_size} fourier evals')
plt.xlabel('X')
plt.ylabel('Y')

plt.show()
```





2.
$$\hat{n} = \hat{x}$$

$$\sigma = -\epsilon_0 \frac{\partial V}{\partial x} \Big|_{y=0}$$

$$= -\epsilon_0 \frac{4V_0}{\pi} \sum_{n=1,3,5...} \frac{\sin(0)}{n \cosh(n\pi)} \frac{\partial}{\partial x} \cosh(\frac{n\pi}{a}x)$$

Since $\sin(0) = 0$, the result is zero.

Units:

$$[\sigma] = [C \cdot m^{-2}]$$
 $[\epsilon_0] = [C^2 \cdot N^{-1} \cdot m^{-2}]$ $[V] = [N \cdot m \cdot C^{-1}]$

For derivative, use $[m^{-1}]$

$$[\epsilon_0] \left[\frac{\partial}{\partial n} \right] [V] = [C^2 N^{-1} m^{-2} N m C^{-1} m^{-1}] = [C m^{-2}] = [\sigma]$$

3 Python: Method of Relaxation in 2D

ak_HW07-Method-of-relaxation-1D.ipynb