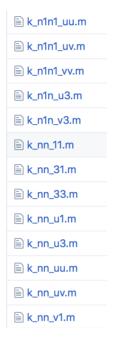
# **Understanding 2D heat equation code**

### files in .../Kernels



## files starting with k are covariance functions:

 $k.m = k_{u.u}^{n+1,n+1}$ 

 $k_n1nv3.m = k^{n+1,n}{v,3}$ 

...

There are corresponding to the functions in the Raissi's paper "Numerical Gaussian Processes for Time-dependent and Non-linear Partial Differential Equations"

$$\begin{array}{lll} k_{u,v}^{n+1,n+1} & = & \frac{d}{dx_2'} k_{u,u}^{n+1,n+1}, & (54) \\ k_{u,3}^{n+1,n} & = & k_{u,u}^{n+1,n+1} - \frac{1}{2} \Delta t \frac{d^2}{dx_1'^2} k_{u,u}^{n+1,n+1} - \frac{1}{2} \Delta t \frac{d^2}{dx_2'^2} k_{u,u}^{n+1,n+1}, \\ k_{v,v}^{n+1,n+1} & = & \frac{d}{dx_2} \frac{d}{dx_2'} k_{u,u}^{n+1,n+1}, \\ k_{v,3}^{n+1,n} & = & \frac{d}{dx_2} k_{u,u}^{n+1,n+1} - \frac{1}{2} \Delta t \frac{d}{dx_2} \frac{d^2}{dx_1'^2} k_{u,u}^{n+1,n+1} - \frac{1}{2} \Delta t \frac{d}{dx_2} \frac{d^2}{dx_2'^2} k_{u,u}^{n+1,n+1}, \\ k_{u,v}^{n,n} & = & \frac{d}{dx_2'} k_{u,u}^{n,n}, \\ k_{u,3}^{n,n} & = & -\frac{1}{2} \Delta t \frac{d^2}{dx_1'^2} k_{u,u}^{n,n} - \frac{1}{2} \Delta t \frac{d^2}{dx_2'^2} k_{u,u}^{n,n}, \\ k_{u,1}^{n,n} & = & k_{u,u}^{n,n}, \\ k_{v,v}^{n,n} & = & \frac{d}{dx_2} \frac{d}{dx_2'} k_{u,u}^{n,n}, \\ k_{v,u}^{n,n} & = & -\frac{1}{2} \Delta t \frac{d}{dx_2} \frac{d^2}{dx_1'^2} k_{u,u}^{n,n} - \frac{1}{2} \Delta t \frac{d}{dx_2} \frac{d^2}{dx_2'^2} k_{u,u}^{n,n}, \\ k_{v,1}^{n,n} & = & -\frac{1}{2} \Delta t \frac{d}{dx_2} \frac{d^2}{dx_1'^2} k_{u,u}^{n,n} - \frac{1}{2} \Delta t \frac{d}{dx_2} \frac{d^2}{dx_2'^2} k_{u,u}^{n,n}, \\ k_{v,1}^{n,n} & = & \frac{d}{dx_2} k_{u,u}^{n,n}, \\ k_{v,1}^{n,n} & = & \frac{d}{dx_2} k_{u,u}^{n,n}, \end{array}$$

#### files starting with D are differentiation operators:

D3kDx2Dy1s.m = 
$$\frac{d}{dx_2} \frac{d^2}{x_1^{\prime 2}}$$

where 3 denotes third derivative, k is the covariance functions which is inserted, x2 denotes  $x_2$ , y denotes x', and s means the square of the variable.

Another example:

DskDx2Dy2.m = 
$$\frac{d}{dx_2} \frac{d}{x_2'}$$

# files in the main repository

#### likelihood.m

It is the definition of the NLML function of 2D heat equation.

In the paper, it looks like

$$\left[\begin{array}{c} \boldsymbol{u}_D^{n+1} \\ \boldsymbol{v}_N^{n+1} \\ \boldsymbol{u}_D^n \\ \boldsymbol{v}_N^n \\ \boldsymbol{u}_3^n \\ \boldsymbol{u}_1^n \end{array}\right] \sim \mathcal{N}\left(0,\boldsymbol{K}\right),$$

The matrix K used in the distribution (44) is given by

where D denotes Dirichlet boundary conditions:

$$u(t, 0, x_2) = u(t, 1, x_2) = 0, \quad u(t, x_1, 0) = 0,$$

while N corresponds to a Neumann-type boundary condition:

$$u_{x_2}(t, x_1, 1) = 0.$$

In the code, each K is assigned corresponding covariance function to obtain the matrix K:

```
K_n1n1_DD = k_n1n1_uu(x_D, x_D, hyp(1:end-1), 0) + eye(N_D).*jitter;
K_n1n1_DN = k_n1n1_uv(x_D, x_N, hyp(1:end-1), 0);
K_n1n_D3 = k_n1n_u3(x_D, x_u, hyp(1:end-1), 0);
K_n1n1_NN = k_n1n1_vv(x_N, x_N, hyp(1:end-1), 0) + eye(N_N).*jitter;
K_n1n_N3 = k_n1n_v3(x_N, x_u, hyp(1:end-1), 0);
K_nn_DD = k_nn_uu(x_D, x_D, hyp(1:end-1), 0) + eye(N_D).*jitter;
K_nn_DN = k_nn_uv(x_D, x_N, hyp(1:end-1), 0);
K_nn_D3 = k_nn_u3(x_D, x_u, hyp(1:end-1), 0);
K_nn_D1 = k_nn_u1(x_D, x_u, hyp(1:end-1), 0);
K_nn_NN = k_nn_vv(x_N, x_N, hyp(1:end-1), 0) + eye(N_N).*jitter;
K_nn_N3 = k_nn_v3(x_N, x_u, hyp(1:end-1), 0);
K_nn_N1 = k_nn_v1(x_N, x_u, hyp(1:end-1), 0);
K_n_33 = k_n_33(x_u, x_u, hyp(1:end-1), 0) + eye(N_u).*sigma_u + eye(N_u).*jitter;
K_nn_31 = k_nn_31(x_u, x_u, hyp(1:end-1), 0);
K_n_11 = k_n_11(x_u, x_u, hyp(1:end-1), 0) + eye(N_u).*sigma_u + eye(N_u).*jitter;
K = [K_n1n1_DD \ K_n1n1_DN \ zeros(N_D,N_D) \ zeros(N_D,N_N) \ K_n1n_D3 \ zeros(N_D,N_u);
     K_n1n1_DN' K_n1n1_NN zeros(N_N,N_D) zeros(N_N,N_N) K_n1n_N3 zeros(N_N,N_u);
     zeros(N\_D,N\_D) \ zeros(N\_D,N\_N) \ K\_nn\_DD \ K\_nn\_DN \ K\_nn\_D3 \ K\_nn\_D1;
     zeros(N_N,N_D) zeros(N_N,N_N) K_nn_DN' K_nn_NN K_nn_N3 K_nn_N1;
     K_n1n_D3' K_n1n_N3' K_nn_D3' K_nn_N3' K_nn_33 K_nn_31;
     zeros(N_u,N_D) zeros(N_u,N_N) K_nn_D1' K_nn_N1' K_nn_31' K_nn_11];
```