

# Magnetic interactions in AB-stacked kagome lattices: magnetic structure, symmetry, and duality

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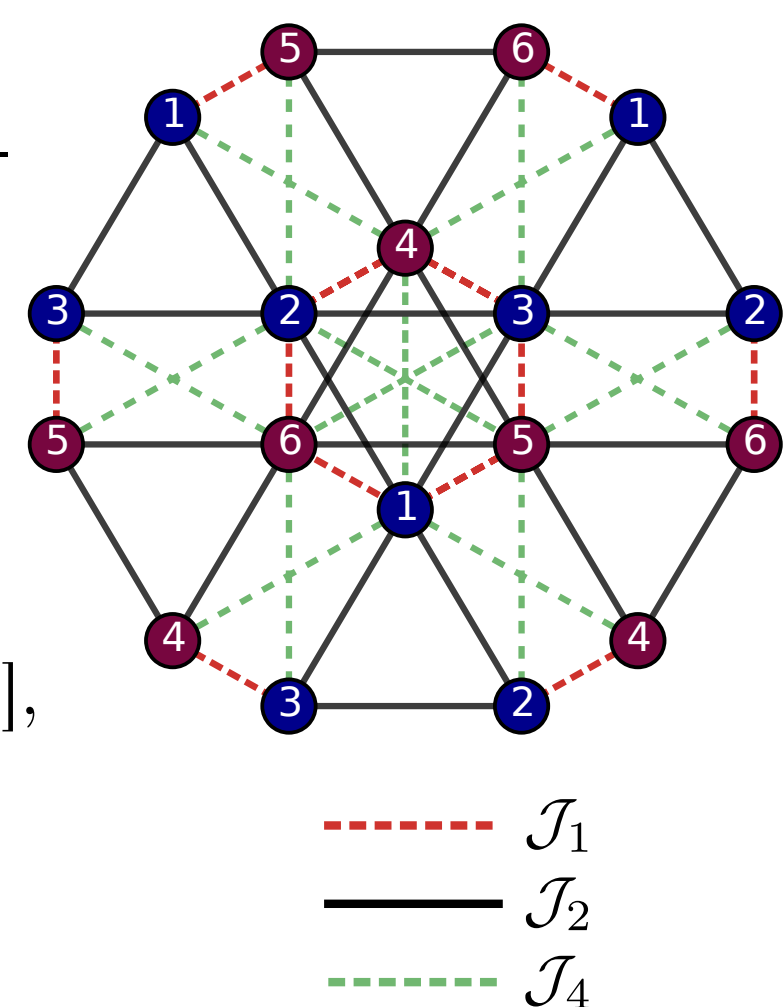
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## Background

Recently, large Anomalous Hall and Nernst effects (AHE) and (ANE) were experimentally observed in compounds with AB-stacked kagome layers and a general formula  $\text{Mn}_3\text{X}$ . These discoveries prompted theoretical and experimental studies of the magnetic properties in  $\text{Mn}_3\text{X}$  magnets. In our previous work [1], we derived a magnetic model for these magnetic compounds using general symmetry principles.



$$\mathcal{H} = \mathcal{H}_J + \mathcal{H}_D + \mathcal{H}_A + \mathcal{H}_K$$

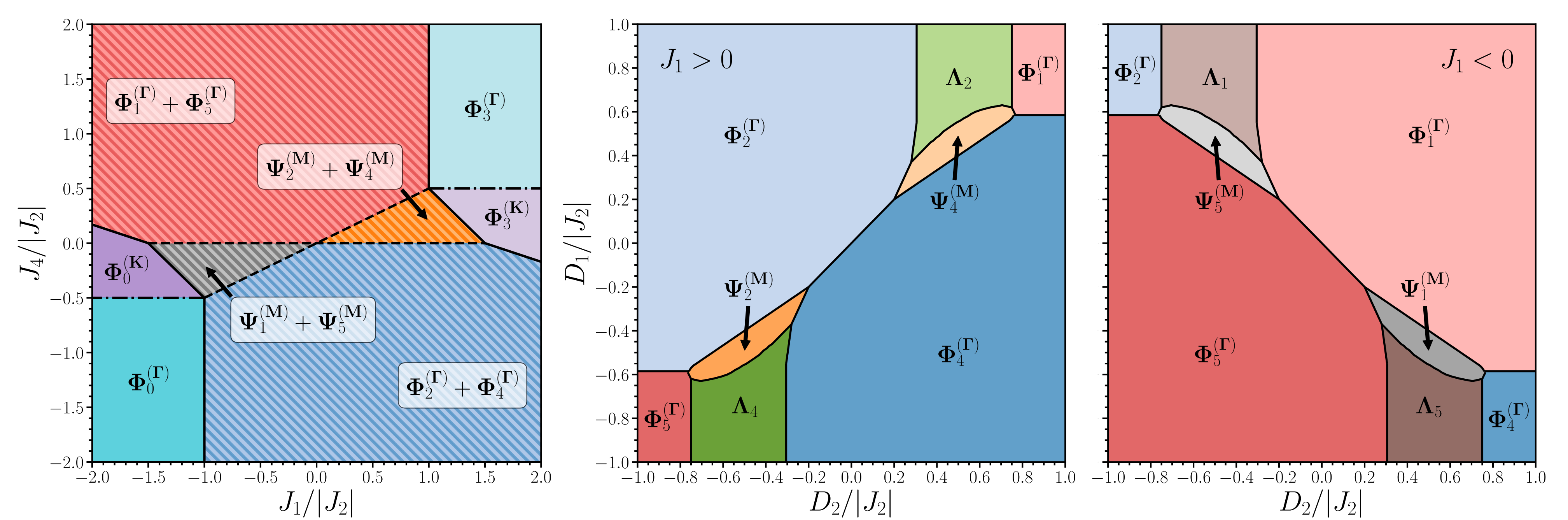
$$\mathcal{H}_J = \sum_{\mathbf{r}\mathbf{r}'} \sum_{ij} J_{ij}(\mathbf{r} - \mathbf{r}') \mathbf{S}_i(\mathbf{r}) \cdot \mathbf{S}_j(\mathbf{r}')$$

$$\mathcal{H}_D = \sum_{\mathbf{r}\mathbf{r}'} \sum_{ij} D_{ij}(\mathbf{r} - \mathbf{r}') \hat{\mathbf{z}} \cdot (\mathbf{S}_i(\mathbf{r}) \times \mathbf{S}_j(\mathbf{r}'))$$

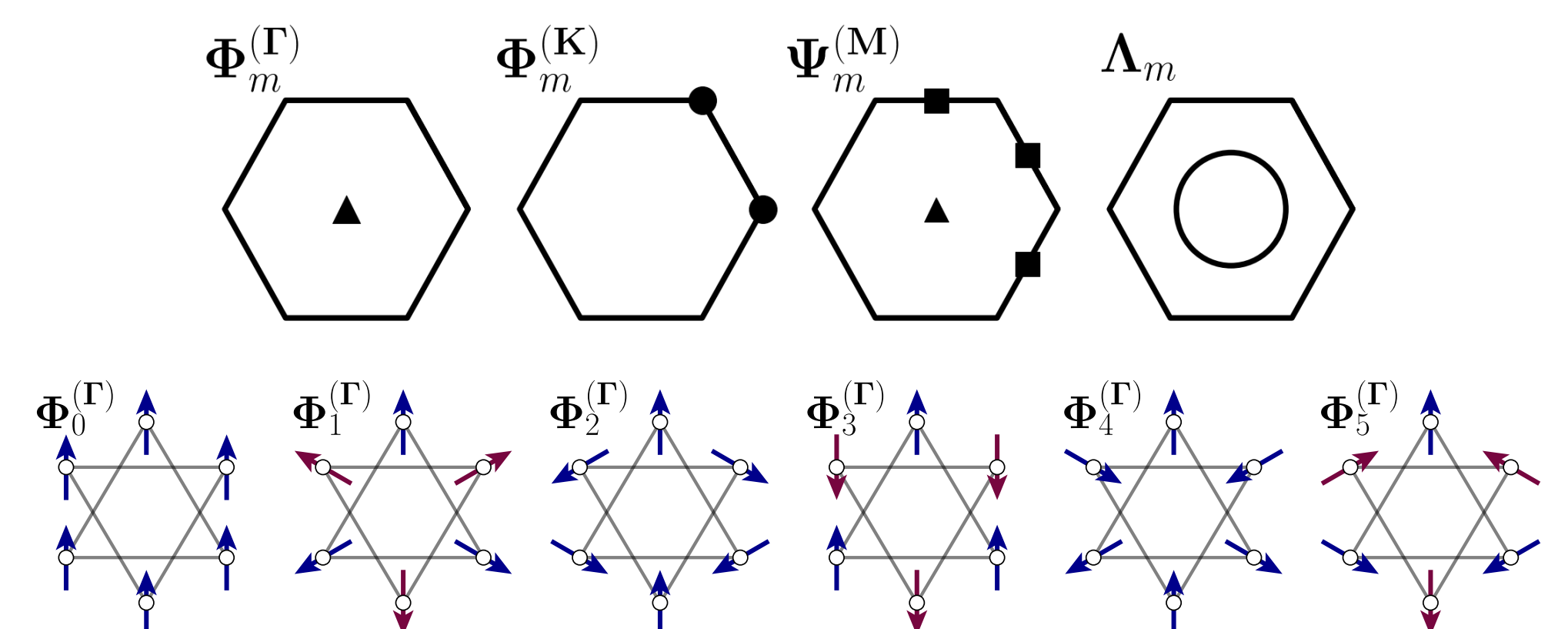
$$\mathcal{H}_A = \sum_{\mathbf{r}\mathbf{r}'} \sum_{ij} \sum_{\alpha} A_{ij\alpha}(\mathbf{r} - \mathbf{r}') (\hat{\mathbf{n}}_{i\alpha} \cdot \mathbf{S}_i(\mathbf{r})) (\hat{\mathbf{n}}_{j\alpha} \cdot \mathbf{S}_j(\mathbf{r}')),$$

$$\mathcal{H}_K = \sum_{\mathbf{r}} \sum_i \sum_{\alpha} K_{\alpha} (\hat{\mathbf{n}}_{i\alpha} \cdot \mathbf{S}_i(\mathbf{r}))^2.$$

## Magnetic phases

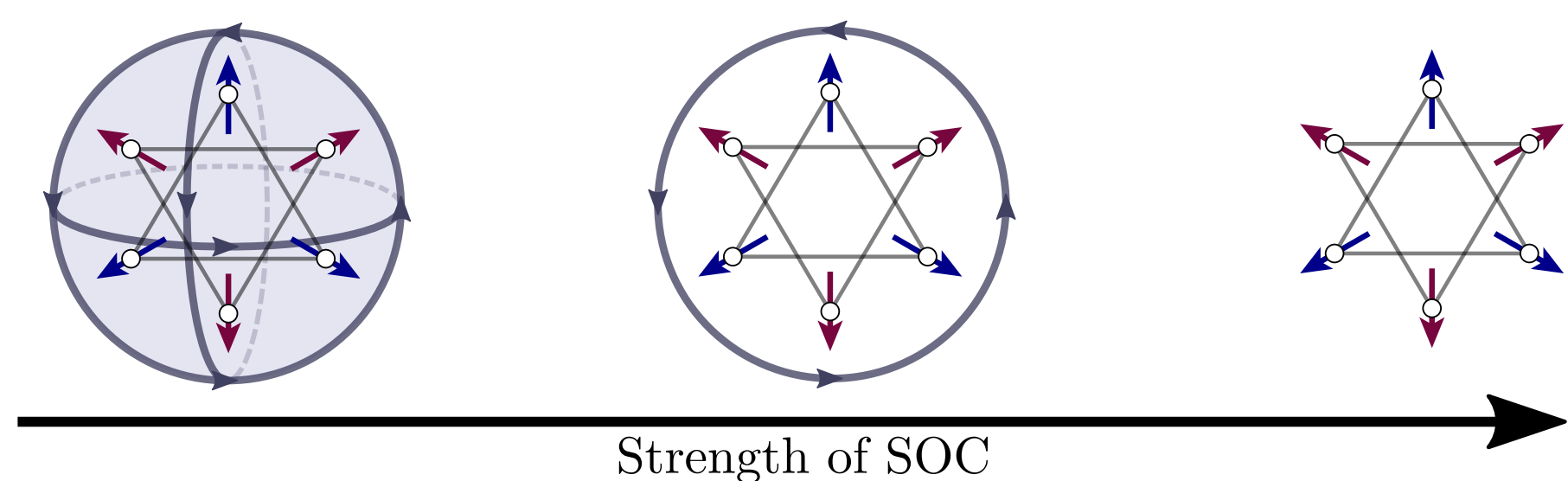


We determine the semi-classical ground state using a combination of analytical (Luttinger-Tisza) and numerical (simulated annealing) methods. The model gives rise to a wide range of magnetic phases, including configurations with single and multiple ordering wavevectors. For every type of magnetic order we find six states characterized by a "winding" number.



## Symmetry

Depending on the strength of the spin-orbit coupling (SOC), the effective symmetry group of the system can be much larger than the magnetic space group [2,3].



Magnetic space group + Full rotational symmetry  
 $\mathcal{H} \approx \mathcal{H}_J$

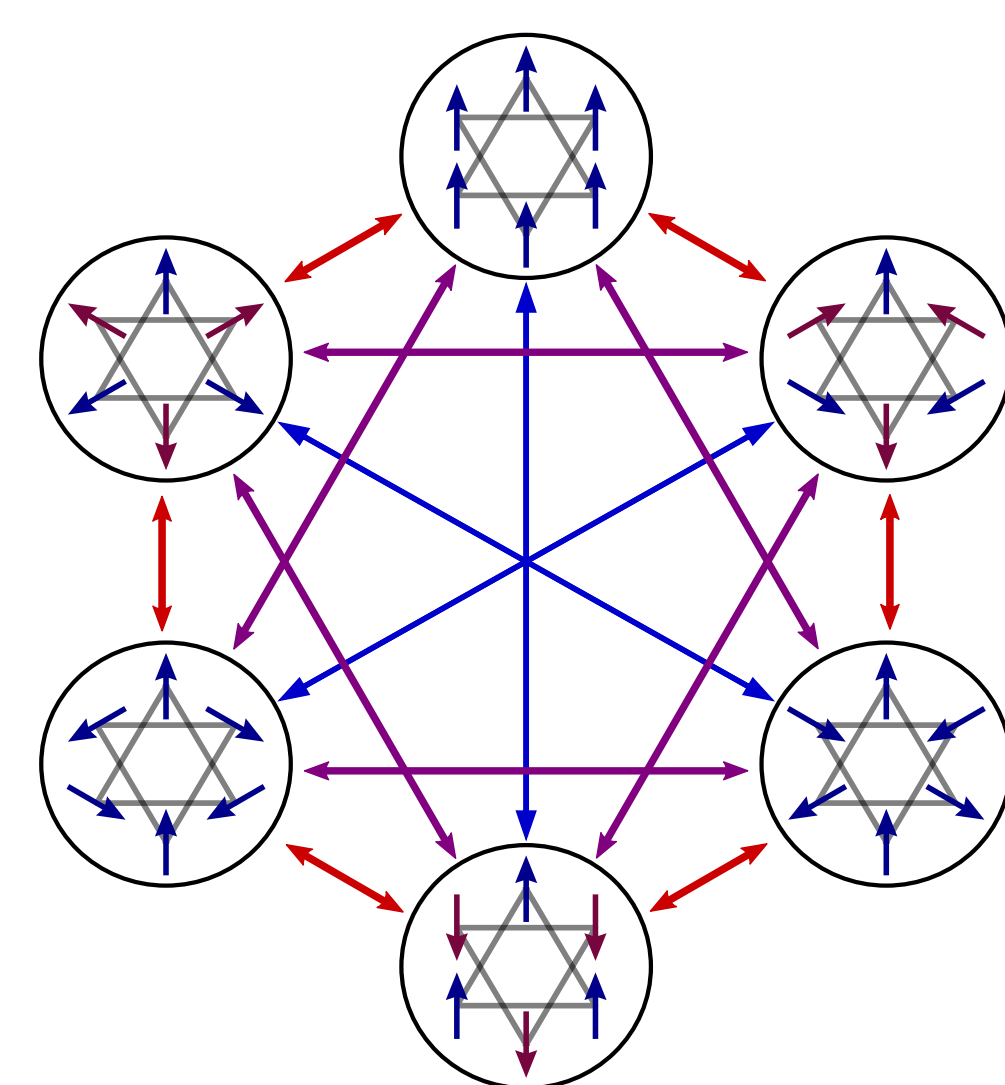
Magnetic space group + Axial rotational symmetry  
 $\mathcal{H} \approx \mathcal{H}_J + \mathcal{H}_D$

Magnetic space group

In the case of  $\text{Mn}_3\text{X}$ , there are three different cases corresponding to decoupled, weak, and intermediate SOC. Each case results in effective Hamiltonians with different spin symmetry properties.

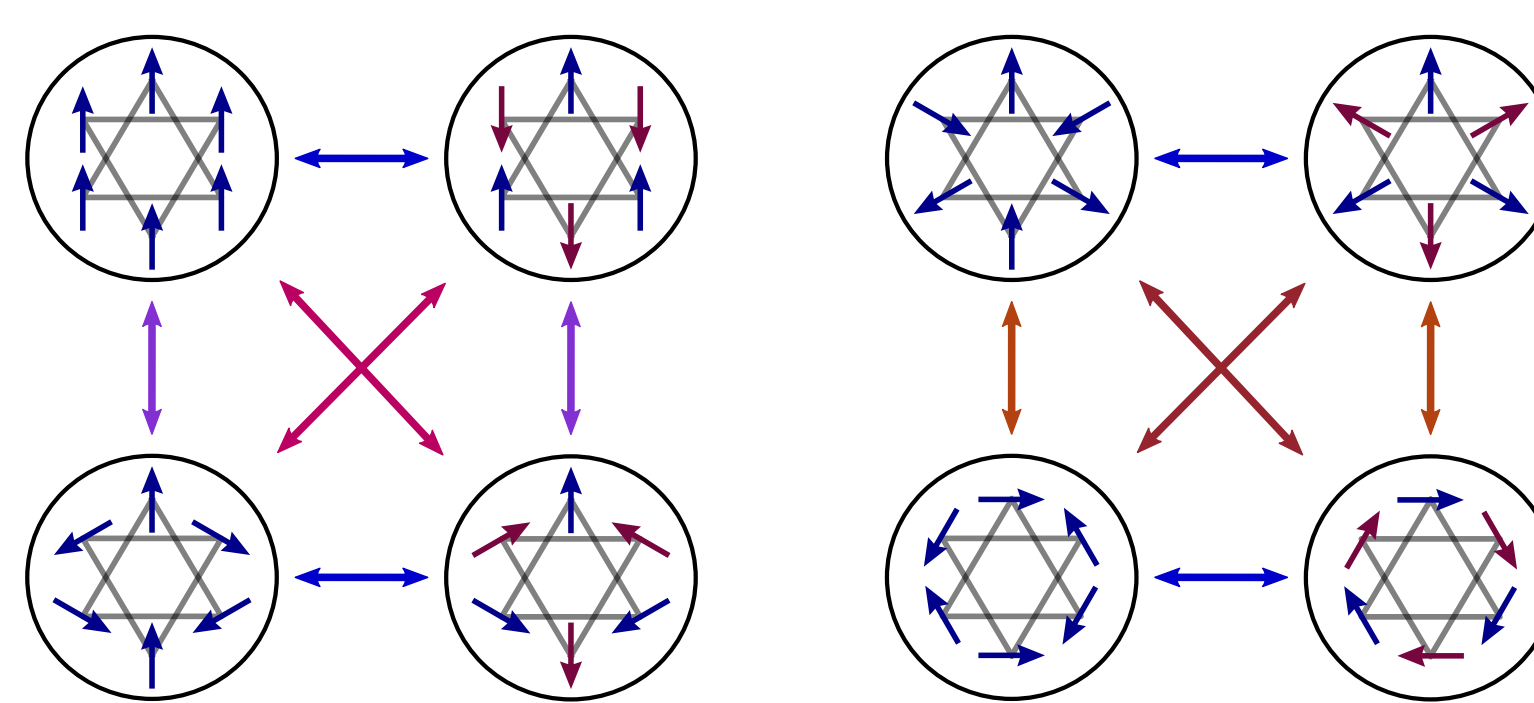
## Self-duality

Besides the symmetry operations, there is a set of local spin transformations that only change the values of the coupling constants. Since these transformations map the Hamiltonian onto itself, they are called **self-dualities** [4,5].



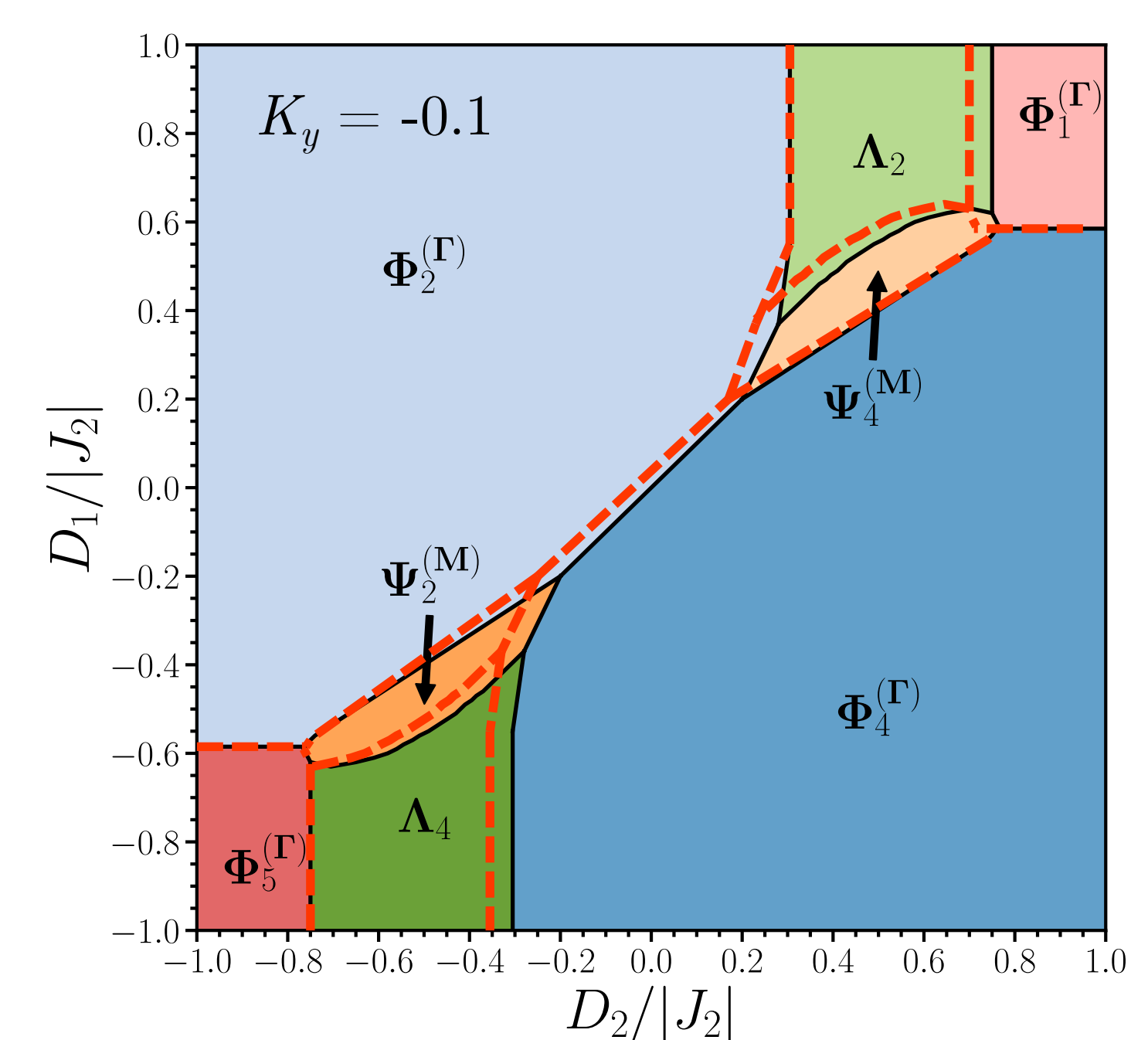
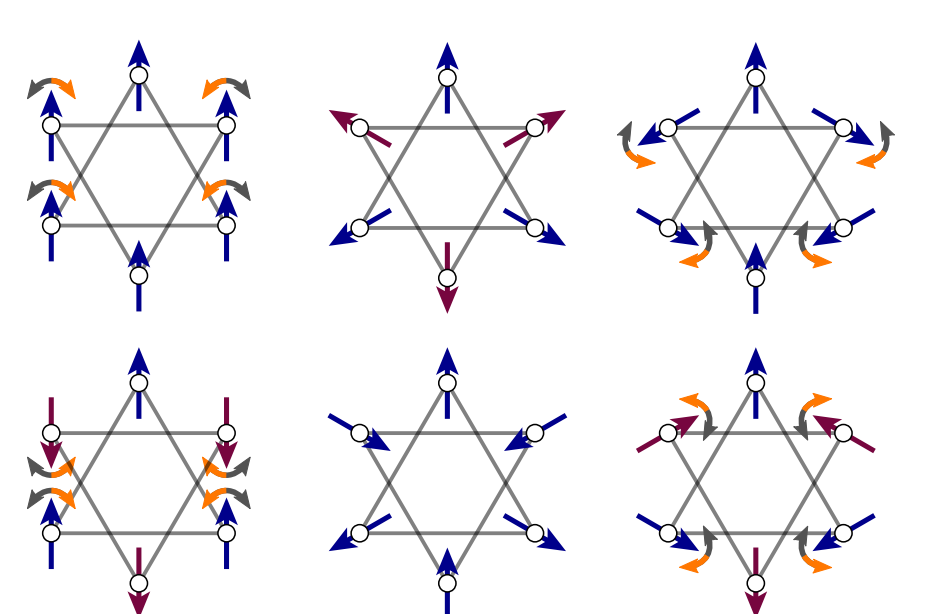
In the weak SOC limit, the self-dualities change the winding of the local axes.

In the intermediate SOC limit, there are still more than a dozen duality transformations.



## Anisotropy

Single-ion and bond-dependent anisotropy, which result from intermediate SOC, favour the states with spins collinear with the anisotropy axes, thus breaking the symmetry of the phase diagrams.



## Conclusions

The strength of the SOC has a direct effect on the **symmetry** of the magnetic Hamiltonian.

For  $\text{Mn}_3\text{X}$ , there are three distinct SOC limits, each characterized by different symmetry properties.

Competing magnetic interactions stabilize a wide range of magnetic phases.

The model exhibits a large number of self-dualities, which relate the properties of different phases.

The self-dualities remain even in the most general case of intermediate SOC.

## Acknowledgements

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## References

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