

# Series de Fourier

①

Sea  $f(t)$  una función periódica de periodo  $T$ , entonces existe una representación

$$f(t) = \sum_{n=-\infty}^{+\infty} c_n e^{j \frac{2\pi n}{T} t} \quad (1)$$

donde

$$c_n = \frac{1}{T} \int_{-T/2}^{T/2} f(t) e^{-j \frac{2\pi n}{T} t} dt \quad (2)$$

→ IM604 - Serie Fourier Cn.m

Empleando  $e^{j\theta} = \cos\theta + j\sin\theta$  se obtiene  $\downarrow$  "optional"

$$f(t) = \frac{a_0}{2} + \sum_{n=1}^{+\infty} \left\{ a_n \cos\left(\frac{2\pi n t}{T}\right) + b_n \sin\left(\frac{2\pi n t}{T}\right) \right\}$$

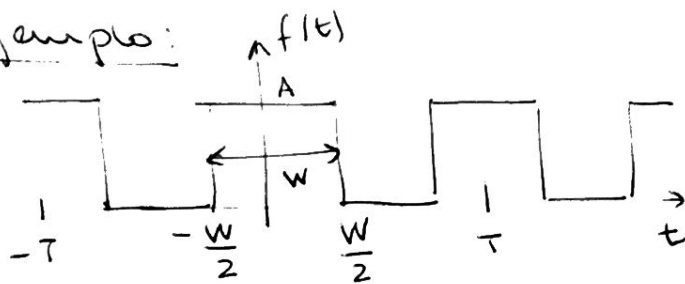
donde

$$a_0 = \frac{2}{T} \int_{-T/2}^{T/2} f(t) dt$$

$$a_n = \frac{2}{T} \int_{-T/2}^{T/2} f(t) \cos\left(\frac{2\pi n t}{T}\right) dt$$

$$b_n = \frac{2}{T} \int_{-T/2}^{T/2} f(t) \sin\left(\frac{2\pi n t}{T}\right) dt$$

Ejemplo:



← función periódica de periodo  $T$ .

$$a_0 = \frac{2}{T} \int_{-T/2}^{T/2} f(t) dt = \frac{2}{T} \int_{-W/2}^{W/2} A dt = \frac{2A}{T} t \Big|_{-W/2}^{W/2} = \frac{2AW}{T}$$

$$a_n = \frac{2}{T} \int_{-T/2}^{T/2} f(t) \cos\left(\frac{2\pi n t}{T}\right) dt = \frac{2A}{T} \int_{-W/2}^{W/2} \cos\left(\frac{2\pi n t}{T}\right) dt$$

$$= \frac{2A}{T} \cdot \frac{T}{2\pi n} \sin\left(\frac{2\pi n t}{T}\right) \Big|_{-W/2}^{W/2}$$

$$= \frac{2AT}{2\pi n T} \left( \sin\left(\frac{2\pi n W}{2T}\right) - \sin\left(-\frac{2\pi n W}{2T}\right) \right)$$

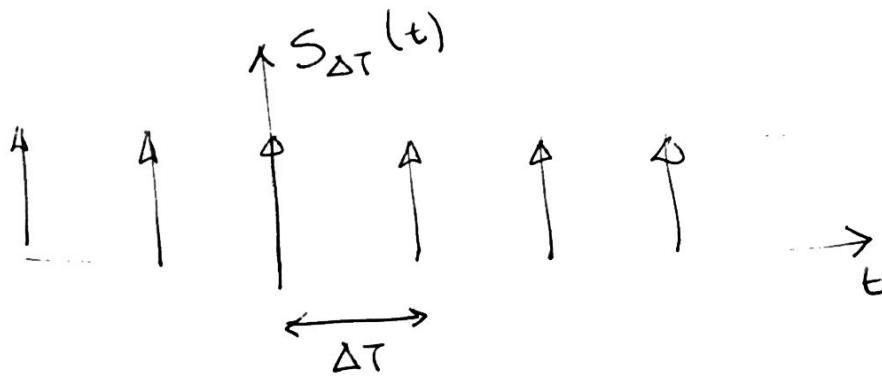
$$\boxed{a_n = \frac{2A}{\pi n} \sin\left(\frac{\pi n W}{T}\right)}$$

$$b_n = \frac{2}{T} \int_{-T/2}^{T/2} f(t) \sin\left(\frac{2\pi n t}{T}\right) dt = 0 \quad \text{porque } f(t) \sin\left(\frac{2\pi n t}{T}\right) \text{ es impar.}$$

$$f(t) = \frac{AW}{T} + \frac{2A}{\pi} \sum_{n=1}^{+\infty} \sin\left(\frac{\pi n W}{T}\right) \cos\left(\frac{2\pi n t}{T}\right)$$

→ IM604-Serie Fourier.m

# TREN DE IMPULSOS



$$S_{\Delta T}(t) = \sum_{n=-\infty}^{+\infty} \delta(t - n\Delta T) \quad (3)$$

Periodo =  $\Delta T$

$S_{\Delta T}(t)$  puede ser representada usando series de Fourier (1):

$$S_{\Delta T}(t) = \sum_{n=-\infty}^{+\infty} c_n e^{j \frac{2\pi n}{\Delta T} t}$$

$$c_n = \frac{1}{\Delta T} \int_{-\Delta T/2}^{\Delta T/2} S_{\Delta T}(t) e^{-j \frac{2\pi n t}{\Delta T}} dt$$

$$= \frac{1}{\Delta T} \int_{-\Delta T/2}^{\Delta T/2} \delta(t) e^{-j \frac{2\pi n t}{\Delta T}} dt$$

$$c_n = \frac{1}{\Delta T} e^0 = \frac{1}{\Delta T}$$

$$\Rightarrow S_{\Delta T}(t) = \frac{1}{\Delta T} \sum_{n=-\infty}^{+\infty} e^{j \frac{2\pi n t}{\Delta T}} \quad (4)$$

→ IM604\_TrenImpulsos.m

## TRANSFORMADA DE FOURIER DE FUNCIONES IMPULSO:

$$\delta(t) \xrightarrow{\bullet} \int_{-\infty}^{+\infty} \delta(t) e^{-j2\pi\omega t} dt = e^0 = 1 \quad (5)$$

$$\delta(t - t_0) \xrightarrow{\bullet} \int_{-\infty}^{+\infty} \delta(t - t_0) e^{-j2\pi\omega t} dt = e^{-j2\pi\omega t_0} \quad (6)$$

## TRANSFORMADA INVERSA:

$$\delta(\omega) \xrightarrow{\bullet} \int_{-\infty}^{+\infty} \delta(\omega) e^{j2\pi\omega t} d\omega = e^0 = 1 \quad (7)$$

$$\delta(\omega - \omega_0) \xrightarrow{\bullet} \int_{-\infty}^{+\infty} \delta(\omega - \omega_0) e^{j2\pi\omega t} d\omega = e^{j2\pi\omega_0 t} \quad (8)$$

EJEMPLO:

Calcular la Transformada de

$$f(t) = \frac{1}{2} [\delta(t - t_0) + \delta(t + t_0)]$$

$$\begin{aligned} f(t) \xrightarrow{\bullet} F(\omega) &= \frac{1}{2} [e^{-j2\pi\omega t_0} + e^{j2\pi\omega t_0}] \\ &= \cos(2\pi\omega t_0) \end{aligned} \quad (9)$$

de igual manera:

$$f(t) = \cos \omega_0 t \xrightarrow{\bullet} F(\omega) = \frac{1}{2} [\delta(\omega - \omega_0) + \delta(\omega + \omega_0)] \quad (10)$$

# TRANSFORMADA DE FOURIER DEL TIEN DE IMPULSOS

$$F \{ S_{\Delta T}(t) \} = \int_{-\infty}^{+\infty} \frac{1}{\Delta T} \sum_{n=-\infty}^{+\infty} e^{j \frac{2\pi n t}{\Delta T}} e^{-j 2\pi \omega t} dt$$

$$= \frac{1}{\Delta T} \sum_{n=-\infty}^{+\infty} \underbrace{\int_{-\infty}^{+\infty} e^{j \frac{2\pi n t}{\Delta T}} e^{-j 2\pi \omega t} dt}_{\text{Transformada de Fourier de } e^{j \frac{2\pi n t}{\Delta T}}}$$

Transformada de Fourier de  $e^{j \frac{2\pi n t}{\Delta T}}$  de (8)

$$= \frac{1}{\Delta T} \sum_{n=-\infty}^{+\infty} F \left\{ e^{j \frac{2\pi n t}{\Delta T}} \right\} = \frac{1}{\Delta T} \sum_{n=-\infty}^{+\infty} \delta \left( \omega - \frac{n}{\Delta T} \right)$$

Esto quiere decir que la transformada de un tren de impulsos es un tren de impulsos!

$$S_{\Delta T}(t) \longleftrightarrow \frac{1}{\Delta T} S_{1/\Delta T}(\omega) \quad (11)$$

$$\sum \delta(t - n\Delta T) \longleftrightarrow \frac{1}{\Delta T} \sum \delta\left(\omega - \frac{n}{\Delta T}\right) \quad (12)$$

# TRANSFORMADA DE LA CONVOLUCIÓN

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$$f(t) * g(t) = \int_{-\infty}^{+\infty} f(\tau) g(t-\tau) d\tau$$

$$\begin{aligned} F\{f(t) * g(t)\} &= \int_{-\infty}^{+\infty} \int_{-\infty}^{+\infty} f(\tau) g(t-\tau) d\tau e^{-j2\pi\omega t} dt \\ &= \int_{-\infty}^{+\infty} f(\tau) \underbrace{\int_{-\infty}^{+\infty} g(t-\tau) e^{-j2\pi\omega t} dt}_{A} d\tau \quad (*) \end{aligned}$$

con  $u = t - \tau$

$$A = \int_{-\infty}^{+\infty} g(u) e^{-j2\pi\omega(u+\tau)} du$$

$$= e^{-j2\pi\omega\tau} \underbrace{\int_{-\infty}^{+\infty} g(u) e^{-j2\pi\omega u} du}_{G(\omega)}$$

$G(\omega)$  : Transformada de Fourier de  $g$

$A = e^{-j2\pi\omega\tau} G(\omega) \dots$  reemplazando  $A$  en  $(*)$

$$\begin{aligned} F\{f(t) * g(t)\} &= \int_{-\infty}^{+\infty} f(\tau) e^{-j2\pi\omega\tau} G(\omega) d\tau \\ &= \left[ \int_{-\infty}^{+\infty} f(\tau) e^{-j2\pi\omega\tau} d\tau \right] G(\omega) \\ &= F(\omega) G(\omega) \end{aligned}$$

$f(t) * g(t) \longrightarrow F(\omega) G(\omega)$

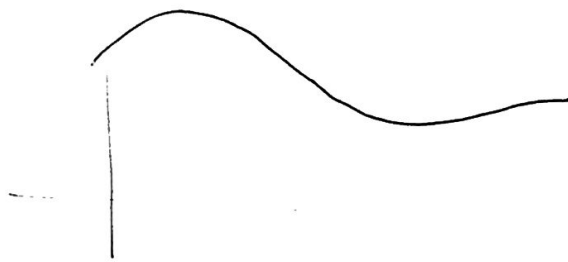
(13)

La transformada de la convolución de dos funciones es la multiplicación de sus transformadas!

# TEOREMA DEL MUESTREO

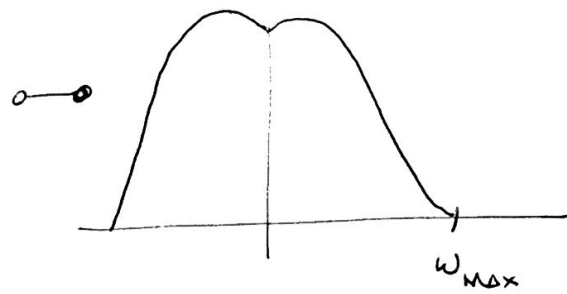
(7)

$\wedge f(t)$



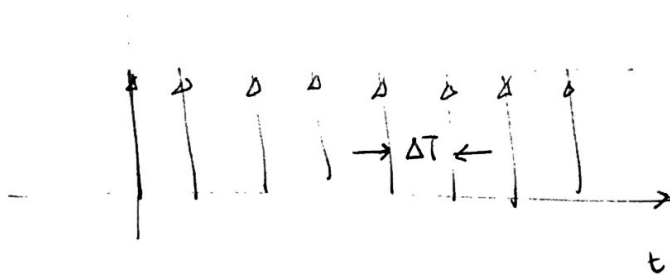
$t$

$\wedge F(\omega)$



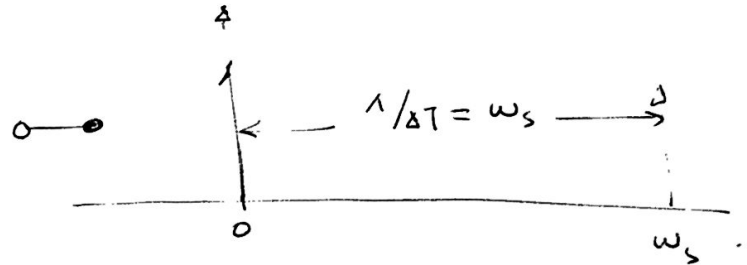
$\omega$

$\wedge s(t)$



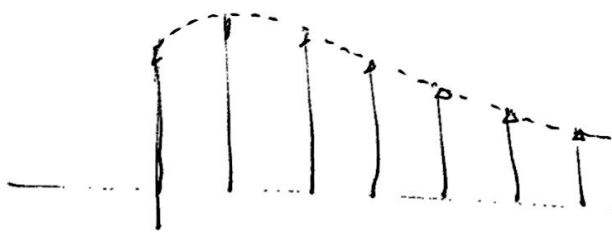
$t$

$S(\omega)$

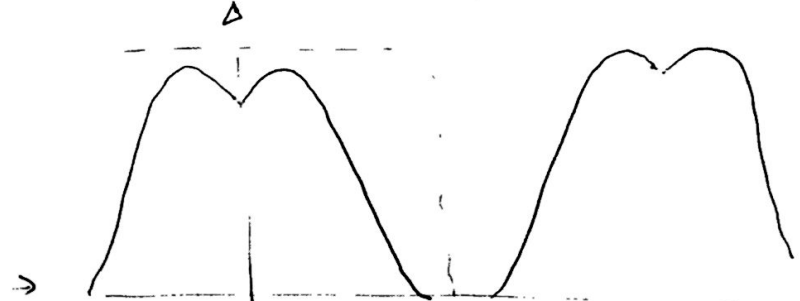


$\omega_s$

$\Delta s(t) \cdot f(t)$



$\Delta F(\omega) * S(\omega)$



Para recuperar  $f(t)$  a partir de un muestreo  $s(t)f(t)$  es necesario que  $\omega_s > 2\omega_{max}$  (Teorema de Nyquist). Se usa un filtro pasa bajos de ancho  $\omega_{max} + \epsilon$ .