

DC 2232 : Assignment 2

1) Given :- $V(u) = -\frac{\hbar^2 \alpha}{m} \operatorname{sech}^2(\alpha u)$

i) Show that $\psi(u) = A \operatorname{sech}(\alpha u)$ is a bound state

1D. SE.

$$-\frac{\hbar^2}{2m} \frac{\partial^2}{\partial u^2} \psi(u) - \frac{\hbar^2 \alpha^2}{m} \operatorname{sech}^2(\alpha u) \psi(u) = E \psi(u).$$

∴ consider LHS,

$$\frac{\partial}{\partial u} A \operatorname{sech}(\alpha u) = \frac{d}{du} A \frac{1}{\operatorname{coth}(\alpha u)} = \left[\frac{\sinh(\alpha u)}{\operatorname{coth}(\alpha u)} \right] [-\alpha A] = -\alpha A \tanh(\alpha u) \operatorname{sech}(\alpha u).$$

$$\therefore \frac{d^2}{du^2} A \operatorname{sech}(\alpha u) = \frac{d}{du} \left[-\alpha A \tanh(\alpha u) \operatorname{sech}(\alpha u) \right]$$

$$= -\alpha A \left[\frac{d}{du} (\tanh(\alpha u)) \operatorname{sech}(\alpha u) + \tanh(\alpha u) \frac{d}{du} \frac{1}{\cosh(\alpha u)} \right]$$

$$\text{But } \frac{d}{du} [\tanh(\alpha u)] = \frac{d}{du} \left[\frac{\sinh(\alpha u)}{\cosh(\alpha u)} \right]$$

$$= \frac{\alpha - \alpha \sinh^2(\alpha u)}{\cosh^2(\alpha u)}$$

$$= \alpha - \alpha \tanh^2(\alpha u)$$

$$= \alpha \operatorname{sech}^2(\alpha u).$$

$$\therefore \frac{d^2}{du^2} A \operatorname{sech}(\alpha u) = -\alpha^2 A \left[\operatorname{sech}^3(\alpha u) - \tanh^2(\alpha u) \operatorname{sech}(\alpha u) \right]$$

$$= -\alpha^2 A \operatorname{sech}^3(\alpha u) [1 - \sinh^2(\alpha u)].$$

$$\therefore -\frac{\hbar^2}{2m} \frac{\partial^2}{\partial u^2} \psi(u) - \frac{\hbar^2 \alpha^2}{m} A \operatorname{sech}^3(\alpha u) = \frac{\hbar^2 \alpha^2}{2m} A \operatorname{sech}^3(\alpha u) [1 - \sinh^2(\alpha u) - 2].$$

$$= -\frac{\hbar^2 \alpha^2}{2m} A \operatorname{sech}^3(\alpha u) [1 + \sinh^2(\alpha u)]$$

$$= -\frac{\hbar^2 \alpha^2}{2m} A \operatorname{sech}^3(\alpha u) \cosh^2(\alpha u)$$

$$= -\frac{\hbar^2 \alpha^2}{2m} A \operatorname{sech}(\alpha u)$$

$$= -\frac{\hbar^2 \alpha^2}{2m} \psi_1(u).$$

$$\therefore E = -\frac{\hbar^2 \alpha^2}{2m} \quad \& \Rightarrow \text{proven}$$

Given them: $\tanh(u) = \frac{\sinh(u)}{\cosh(u)}$ $\operatorname{coth}(u) = \frac{1}{\operatorname{cosech}(u)}$

$$\frac{d}{du} \sinh(u) = \cosh(u) \quad | \quad \operatorname{cosech}^2(u) = 1 + \sinh^2(u).$$

$$\frac{d}{du} \cosh(u) = \sinh(u)$$

ii) show that $\psi_{\kappa}(x) = B \left(\frac{ik - \alpha \tanh(\alpha x)}{ik + \alpha} \right) e^{ikx}$ is a free particle solution.

To simplify the maths, consider the terms separately

$$\frac{d}{dx} [ik e^{ikx}] = -k^2 e^{ikx} \Rightarrow \frac{d^2}{dx^2} [ik e^{ikx}] = -k^3 e^{ikx}.$$

$$\frac{d}{dx} [\alpha \tanh(\alpha x) e^{ikx}] = \alpha^2 \operatorname{sech}^2(\alpha x) e^{ikx} + ik \alpha \tanh(\alpha x) e^{ikx}.$$

$$\begin{aligned} \frac{d^2}{dx^2} [\alpha \tanh(\alpha x) e^{ikx}] &= -\alpha^3 \operatorname{sech}^2(\alpha x) e^{ikx} + k^2 \alpha \tanh(\alpha x) e^{ikx} \\ &\quad + 2ik^3 \operatorname{sech}^2(\alpha x) \tanh(\alpha x) e^{ikx} + ik\alpha^3 \operatorname{sech}^2(\alpha x) e^{ikx}. \end{aligned}$$

$$\therefore -\frac{\hbar^2}{2m} \frac{d^2}{dx^2} \psi_{\kappa}(x) = \frac{\hbar^2}{2m} \frac{B}{ik + \alpha} \left[-ik^3 - 2ik\alpha^2 \operatorname{sech}^2(\alpha x) + k^2 \alpha \tanh(\alpha x) + 2\alpha^3 \operatorname{sech}^2(\alpha x) \tanh(\alpha x) \right] e^{ikx}$$

\therefore the other term of the LHS :-

$$V(x) \psi_{\kappa}(x) = -\frac{\hbar^2}{2m} \frac{B}{ik + \alpha} \left[2\alpha^2 \operatorname{sech}^2(\alpha x) [ik - \alpha \tanh(\alpha x)] \right] e^{ikx}$$

$$\begin{aligned} \therefore \text{LHS reduce to: } & \frac{-\hbar^2}{2m} \frac{B}{ik + \alpha} [-ik^3 + k^2 \alpha \tanh(\alpha x)] \\ &= \frac{\hbar^2 k^2}{2m} \frac{B}{ik + \alpha} [ik - \alpha \tanh(\alpha x)] \\ &= \frac{\hbar^2 k^2}{2m} \psi_{\kappa}(x). \end{aligned}$$

iii) Given that: $\tanh x = \frac{e^{-2x}}{1 - e^{-2x}}$.

Show that for a free particle incident on this potential, there is no reflected wave component.

$$\psi_{\kappa}(x) = B \left(\frac{ik + \alpha \tanh(\alpha x)}{ik + \alpha} \right) e^{ikx} \quad \therefore \text{when } x \rightarrow \infty,$$

Note that $x \rightarrow \infty$ does not really affect e^{ikx}

\therefore throughout all of x , if the incident particle is travelling right, it will continue travelling right

\therefore there is no reflection.

$$\lim_{x \rightarrow \infty} \tanh(\alpha x) = \frac{1 + e^{-2\infty}}{1 - e^{-2\infty}} = 1.$$

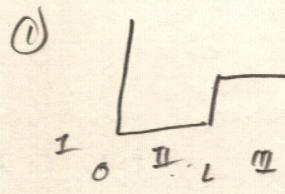
$$\therefore \psi_{\kappa}(x \rightarrow \infty) = \frac{\alpha(e^{ik} - 1)}{ik + \alpha} e^{ikx}$$

when $x \rightarrow -\infty$,

$$\lim_{x \rightarrow -\infty} \tanh(\alpha x) = \frac{e^{2\infty} + 1}{e^{2\infty} - 1} = -1$$

$$\therefore \psi_{\kappa}(x \rightarrow -\infty) = \frac{B(ik + 1)}{ik + \alpha} e^{ikx}$$

$$b) U(r) = \begin{cases} \infty & r < 0 \\ -\frac{32\hbar^2 n^2}{mL^2} & 0 \leq r \leq a \\ 0 & r > a. \end{cases}$$



i) How many bound states are there?

\Rightarrow Bound states $\rightarrow E < 0$.

\because the minimum E possible is $-\frac{32\hbar^2 n^2}{mL^2}$ (no E can exist below this potential),

\therefore The highest E for bound state will have a magnitude of $\frac{32\hbar^2 n^2}{mL^2}$

Recall for ∞ -sq. well, $E = \frac{n^2 h^2}{8mL^2}$.

\therefore if this were an infinite sq. well, $E' = \frac{32\hbar^2 n^2}{mL^2} = 32 \cdot \frac{\hbar^2}{4mL^2} = 64 \frac{\hbar^2}{8mL^2}$
 $\therefore n=4$.

$\because E$ of finite sq. well $< E$ of ∞ sq. well, (but only by a little),
highest bound state here would be $n=8$

ii) Solving for this potential for both cases.

Region I: 0

Region II: $t_{\text{out}} = A \sin kx + B \cos kx \rightarrow$ Bound. cond. $\Rightarrow 0$

Region III: $t_{\text{out}} = C e^{ikx} + D e^{-ikx}$.

Because it does not blow up

$$\therefore A \sin kL = D e^{-ikL} \quad \text{and } kL = \alpha C e^{ikL}.$$

Region I: 0

Region II: $t_{\text{out}} = A \sin kx + B \cos kx$

Region III: $t_{\text{out}} = F e^{ikx} + G e^{-ikx}$

$$A \sin kL = F e^{ikL}$$

$$B \cos kL = i k F e^{ikL}$$

$E < 0$.

$\left. \begin{array}{l} \text{we cannot throw} \\ \text{away} G e^{-ikx} \\ \text{term} \because \text{since} \\ \text{region I has } C \rightarrow 0, \\ \text{particle must go} \\ \text{incident from region} \\ \text{III}. \end{array} \right\} E < 0$

iii) Briefly discuss how to find energies of the bound states.

• Bound states only occur for $E < 0$.

• W/ the boundary cond. eq. obtain α & it relates k to α .

\therefore when those are related to E we can replace α in terms of E .

• Plot a graph of LHS wrt E w.r.t. E . Find intersections.

$$2a) f(E) = \frac{N_n}{d_n} = \frac{1}{e^{\frac{E-E_F}{k_B T}} + 1}$$

i) when $T=0$, $E_n < E_F$

$$f(E) = \frac{1}{e^{-\theta} + 1} = 1$$

$T=0$, $E_n > E_F$

$$f(E) = \frac{1}{e^{+\infty} + 1} = 0$$

ii) Fermi E is the energy value that, at 0 K, the highest E particle has.

Better phrased version: At 0 K, the highest E particle has the energy value : E_F .

* Fermi-Dirac distribution is basically the probability that a certain energy level E_n is filled at a certain temperature T .

b) Band gap = 0.67 eV.

i) $T_1 = 250\text{K}$, $T_2 = 300\text{K}$ and $T_3 = 350\text{K}$

$$0.67 \left[\frac{E_n}{e^{-\frac{E_F - E_n}{k_B T}} + 1} \right] \quad E_F = E_n - \frac{0.67}{2} \text{ eV}$$

Probability to state at the bottom of conduction band is occupied?

$$\therefore f(E_n) = \frac{1}{e^{-\frac{0.67}{2} + 1}}$$