

PC2232: Tutorial 6 solutions

Question 1: Harmonic oscillator ground state

Given that ψ satisfies

$$-\frac{\hbar^2}{2m} \frac{d^2\psi}{dx^2} + \frac{1}{2}m\omega^2 x^2 \psi = E\psi, \quad (1)$$

we can calculate some equations that will be useful later:

$$\frac{d^2\psi}{dx^2} = -\frac{2m}{\hbar^2} E\psi + \frac{m^2\omega^2}{\hbar^2} x^2 \psi, \quad (2)$$

$$\frac{d^3\psi}{dx^3} = -\frac{2m}{\hbar^2} E \frac{d\psi}{dx} + \frac{2m^2\omega^2}{\hbar^2} x\psi + \frac{m^2\omega^2}{\hbar^2} x^2 \frac{d\psi}{dx} \quad (3)$$

(a) A new wavefunction ϕ is defined by

$$\phi = \hbar \frac{d\psi}{dx} + m\omega x \psi. \quad (4)$$

Differentiating both sides of the equation gives

$$\frac{d\phi}{dx} = \hbar \frac{d^2\psi}{dx^2} + m\omega \psi + m\omega x \frac{d\psi}{dx}. \quad (5)$$

Differentiating again gives

$$\frac{d^2\phi}{dx^2} = \hbar \frac{d^3\psi}{dx^3} + m\omega \frac{d\psi}{dx} + m\omega \frac{d\psi}{dx} + m\omega x \frac{d^2\psi}{dx^2}. \quad (6)$$

Substituting (2) and (3) into (6), then rearranging the terms gives

$$\begin{aligned} \frac{d^2\phi}{dx^2} &= -\frac{2m}{\hbar^2} E \left(\hbar \frac{d\psi}{dx} + m\omega x \psi \right) + \frac{2m}{\hbar^2} \hbar \omega \left(\hbar \frac{d\psi}{dx} + m\omega x \psi \right) \\ &\quad + \frac{2m}{\hbar^2} \frac{1}{2} m\omega^2 x^2 \left(\hbar \frac{d\psi}{dx} + m\omega x \psi \right) \\ &= -\frac{2m}{\hbar^2} E \phi + \frac{2m}{\hbar^2} \hbar \omega \phi + \frac{2m}{\hbar^2} \frac{1}{2} m\omega^2 x^2 \phi \\ -\frac{\hbar^2}{2m} \frac{d^2\phi}{dx^2} + \frac{1}{2} m\omega^2 x^2 \phi &= \underbrace{(E - \hbar\omega)}_{\text{energy of } \phi} \phi \end{aligned} \quad (7)$$

Therefore ϕ satisfies Schrödinger's equation with an energy of $(E - \hbar\omega)$, which is lower than the energy of ψ by a value of $\hbar\omega$.

(b) Ground state ψ_0 has the lowest possible energy E_0 . Therefore

$$-\frac{\hbar^2}{2m} \frac{d^2\psi_0}{dx^2} + \frac{1}{2} m\omega^2 x^2 \psi_0 = E_0 \psi_0, \quad (8)$$

Since E_0 is the lowest, a state $\phi = \left(\hbar \frac{d}{dx} + m\omega x\right) \phi$ where

$$-\frac{\hbar^2}{2m} \frac{d^2\phi}{dx^2} + \frac{1}{2}m\omega^2 x^2 \phi = \underbrace{(E_0 - \hbar\omega)}_{\text{lower than } E_0} \phi \quad (9)$$

with a lower energy cannot exist. Therefore we conclude that $\phi = 0$:

$$\begin{aligned} \phi &= \hbar \frac{d\psi_0}{dx} + m\omega x \psi_0 = 0 \\ \frac{d\psi_0}{\psi_0} &= -\frac{m\omega}{\hbar} x dx \\ \int \frac{d\psi_0}{\psi_0} &= -\frac{m\omega}{\hbar} \int x dx \\ \psi_0 &= C e^{-m\omega x^2/2\hbar}, \end{aligned} \quad (10)$$

where C is an integration constant.

Question 2: Step potential

There are two distinct regions, as shown in the figure below:

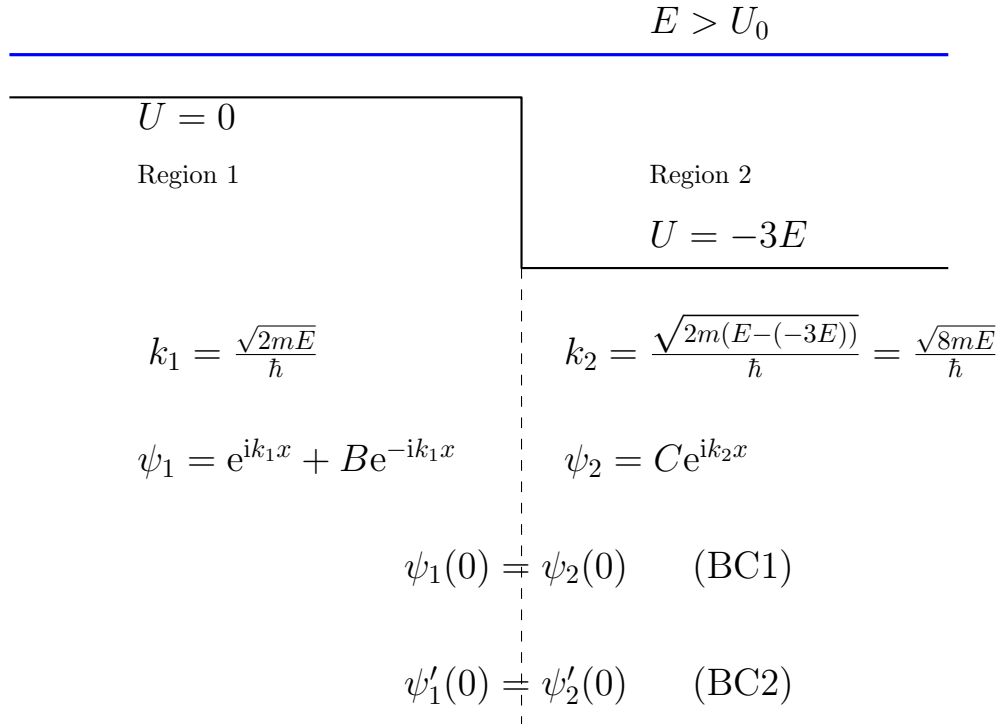


Figure 1: Step potential

- (a) Classically, the particle would drop to a lower potential and have a gain in kinetic energy by a magnitude of $3E$.

(b) Based on Fig. 1, our wavefunction is

$$\psi(x) = \begin{cases} e^{ik_1x} + Be^{-ik_1x}, & x < 0, \\ Ce^{ik_2x}, & x > 0. \end{cases} \quad (11)$$

The values of k_1 and k_2 that satisfies Schrödinger's equation at their respective regions are

$$k_1 = \frac{\sqrt{2mE}}{\hbar}, \quad k_2 = \frac{\sqrt{2m[E - (-3E)]}}{\hbar} = \frac{\sqrt{8mE}}{\hbar}. \quad (12)$$

Applying boundary conditions: At BC1, $\psi_1(0) = \psi_2(0)$ leads to

$$1 + B = C. \quad (13)$$

BC2, $\psi'_1(0) = \psi'_2(0)$ gives

$$\begin{aligned} ik_1(1 - B) &= ik_2C \\ 1 - B &= \frac{k_1}{k_2}C = \frac{\sqrt{8mE}/\hbar}{\sqrt{2mE}/\hbar}C = 2C. \end{aligned} \quad (14)$$

Taking Eq. (13)+(14) gives

$$C = \frac{2}{3}, \quad (15)$$

and Eq. (13)-(14) gives

$$B = -\frac{1}{3}. \quad (16)$$

Therefore, the full solution is

$$\psi(x) = \begin{cases} e^{i\sqrt{2mE}x/\hbar} - \frac{1}{3}e^{-i\sqrt{2mE}x/\hbar}, & x < 0, \\ \frac{2}{3}e^{i\sqrt{8mE}x/\hbar}, & x > 0. \end{cases} \quad (17)$$

(c) The reflection probability is

$$R = \frac{|B|^2}{|A|^2} = \frac{(1/3)^2}{1} = \frac{1}{9}. \quad (18)$$

Question 3: Transmission resonance

There are three regions, as shown in Fig. 2

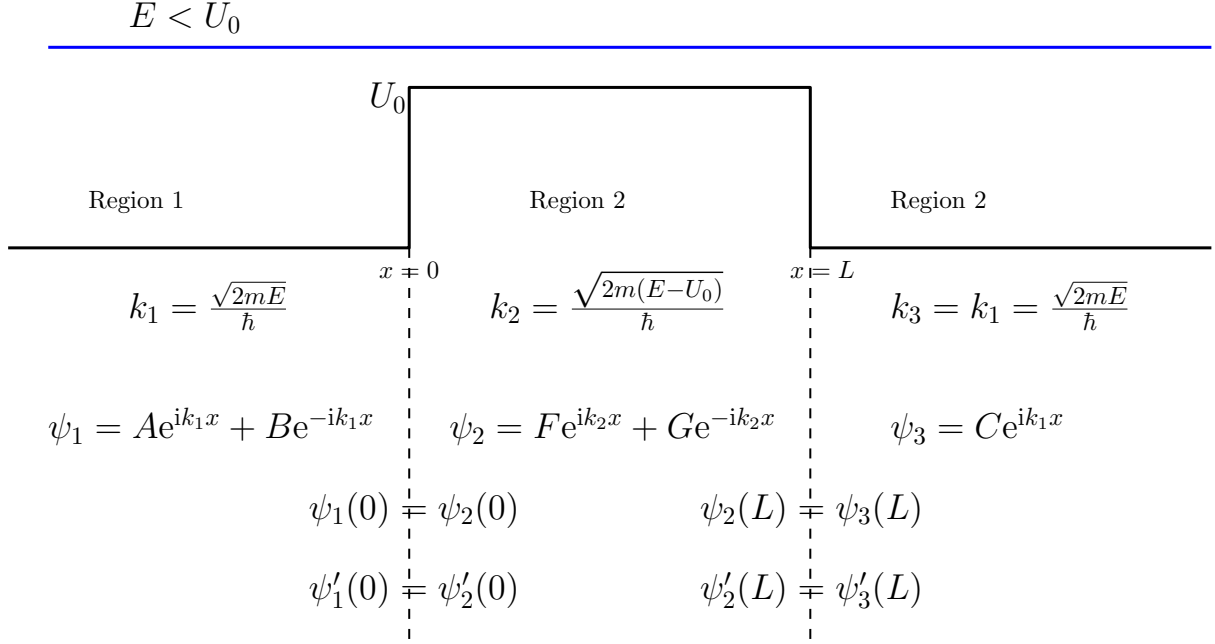


Figure 2: Tunneling through a barrier.

Since we are given $E - U_0 = \frac{\pi^2 \hbar^2}{2mL^2}$, we get

$$k_2 = \frac{\sqrt{2m(E - U_0)}}{\hbar} = \frac{\pi}{L}. \quad (19)$$

The boundary conditions give

$$\psi_1(0) = \psi_2(0) \rightarrow A + B = F + G, \quad (20)$$

$$\psi_2(L) = \psi_3(L) \rightarrow -(F + G) = Ce^{ik_1L}, \quad (21)$$

$$\psi_1'(0) = \psi_2'(0) \rightarrow ik_1(A - B) = ik_2(F - G), \quad (22)$$

$$\psi_2'(L) = \psi_3'(L) \rightarrow ik_2(Fe^{i\pi} - Ge^{-i\pi}) = ik_1Ce^{ik_1L}. \quad (23)$$

Using Eqs. (20) and (21) to eliminate $(F + G)$ gives

$$A + B = -Ce^{ik_1L}. \quad (24)$$

Equation (23) gives

$$\begin{aligned} k_2(Fe^{i\pi} - Ge^{-i\pi}) &= ik_1Ce^{ik_1L} \\ -\frac{\pi}{L}(F - G) &= k_1Ce^{ik_1L} \end{aligned} \quad (25)$$

Using Eqs. (22) and (25) to eliminate $(C - D)$ gives

$$A - B = -Ce^{ik_1L}. \quad (26)$$

Finally, taking (24)–(26) will lead to

$$B = 0. \quad (27)$$

Question 4: Two-dimensional well

Let

$$\psi(x, y) = F(x)G(y). \quad (28)$$

Substitute this into the 2D Schrödinger equation (in the region where $U = 0$) to get

$$\begin{aligned} G \frac{d^2 F}{dx^2} + F \frac{d^2 G}{dy^2} &= -\frac{2mE}{\hbar^2} FG \\ \frac{1}{F} \frac{d^2 F}{dx^2} + \frac{1}{G} \frac{d^2 G}{dy^2} &= -\frac{2mE}{\hbar^2}. \end{aligned} \quad (29)$$

For this equation to be true, each term on the left hand side must individually be constants:

$$\frac{1}{F} \frac{d^2 F}{dx^2} = C_x, \quad \frac{1}{G} \frac{d^2 G}{dy^2} = C_y. \quad (30)$$

Look at the general solution for the x equation first:

$$\frac{d^2 F}{dx^2} = C_x F, \quad \rightarrow \quad F = A_x \sin k_x x + B_x \cos k_x x. \quad (31)$$

We can see that

$$\frac{d^2 F}{dx^2} = -k^2 (A_x \sin k_x x + B_x \cos k_x x) = -k^2 F, \quad (32)$$

so that comparing Eq. (32) with (31), we conclude that

$$-k^2 = C_x. \quad (33)$$

For boundary conditions,

$$\begin{aligned} F(0) = 0 &= 0 + B_x, \quad \rightarrow \quad B_x = 0 \\ F(L) = 0 &= A_x \sin k_x L + \underbrace{B_x}_{=0} \cos k_x L, \quad \rightarrow k_x L = n\pi, \quad n = \text{integer}. \end{aligned} \quad (34)$$

Therefore,

$$k_x^2 = \frac{n^2 \pi^2}{L^2}. \quad (35)$$

Do the similar calculation for y , and substitute the results into Eq. (30):

$$\begin{aligned} \frac{n_x^2 \pi^2}{L^2} + \frac{n_y^2 \pi^2}{L^2} &= \frac{2mE}{\hbar^2} \\ E &= \frac{\pi^2 \hbar^2}{2mL^2} (n_x^2 + n_y^2). \end{aligned} \quad (36)$$

Normalization requires

$$1 = \int dx dy |\psi|^2 = \int dx dy A_x^2 A_y^2 \sin^2 \frac{n_x \pi x}{L} \sin^2 \frac{n_y \pi y}{L}. \quad (37)$$

This should give $A_x A_y = 2/L$. Therefore the final wavefunction should be

$$\psi(x, y) = F(x)G(y) = \frac{2}{L} \sin \frac{n_x \pi x}{L} \sin \frac{n_y \pi y}{L}. \quad (38)$$

The five lowest energy levels are

| n_x | n_y | Energy, E | Degeneracy, d_n |
|-------|-------|-------------|-------------------|
| 1 | 1 | $2K$ | 1 |
| 1 | 2 | $5K$ | 2 |
| 2 | 1 | | |
| 2 | 2 | $8K$ | 1 |
| 3 | 1 | $10K$ | 2 |
| 1 | 3 | | |

where we have introduced the abbreviation

$$K = \frac{\pi^2 \hbar^2}{2mL^2}. \quad (39)$$

Question 5: Tunneling

Given $E = 5.5 \text{ eV} = 8.8 \times 10^{-19} \text{ J}$, $U_0 = 10.0 \text{ eV} = 1.6 \times 10^{-18} \text{ J}$. Accordingly, for an electron,

$$\alpha = \frac{\sqrt{2m(U_0 - E)}}{\hbar} = 1.086 \times 10^{10} \text{ m}^{-1}. \quad (40)$$

The tunneling formula:

$$T \simeq \frac{16E}{U_0} \left(1 - \frac{E}{U_0}\right) e^{-2\alpha L}, \quad (41)$$

Tunneling probability is $T = 0.10\% = 0.001$, therefore

$$\begin{aligned} 0.001 &= 16 \frac{5.5}{10} \left(1 - \frac{5.5}{10}\right) e^{-2(1.086 \times 10^{10})L} \\ 2.525 \times 10^{-4} &= e^{-(2.72 \times 10^{10})L} \\ L &= \mathbf{3.82 \times 10^{-10} \text{ m}}. \end{aligned} \quad (42)$$