

Glaciology EESCGU4220

Lecture 6: Ice sheet flow

Summary from last lecture

Stress balance equations

Strain rate

Our ice-sheet model:

- 1. Mass conservation
- 2. Stress balance

Summary

- Ice deforms through various mechanisms, dislocation creep in the basal plane dominates.
- Rheology is the relationship between stress and strain.
- Nonlinear rheology of ice is usually approximated by Glen's flow law with n = 3.
- Stress state is described by the "stress tensor".

Stress-balance equations

$$\frac{\partial \sigma_{xx}}{\partial x} + \frac{\partial \tau_{yx}}{\partial y} + \frac{\partial \tau_{zx}}{\partial z} = 0$$
 x-direction

$$\frac{\partial \tau_{xy}}{\partial x} + \frac{\partial \sigma_{yy}}{\partial y} + \frac{\partial \tau_{zy}}{\partial z} = 0$$
 y-direction

$$\frac{\partial \tau_{xz}}{\partial x} + \frac{\partial \tau_{yz}}{\partial y} + \frac{\partial \sigma_{zz}}{\partial z} = \rho g$$
 z-direction

Movie proof of this:

https://github.com/ldeo-glaciology/glaciology_practical/raw/main/derivation_movies/stress%20balance%20z.mov

A hierarchy of ice sheet models based on which terms you neglect

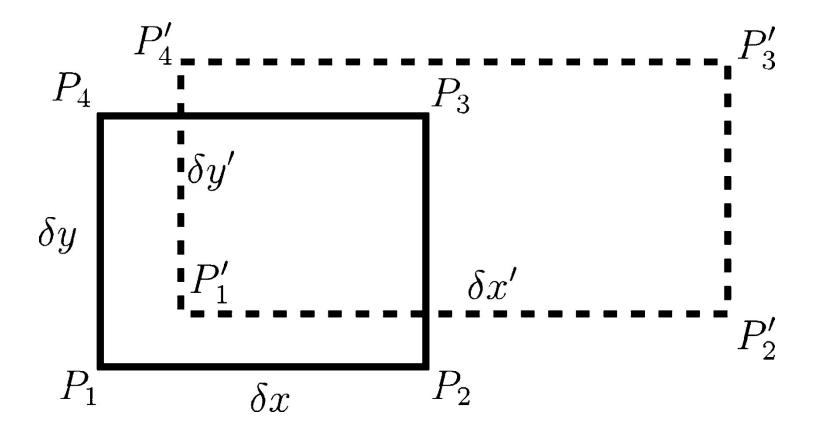
- Shallow Ice Approximation
- Shallow Shelf Approximation
- 'Hybrid'
- 'Higher-order'
- 'Full-stokes'

$$\frac{\partial \sigma_{xx}}{\partial x} + \frac{\partial \tau_{yx}}{\partial y} + \frac{\partial \tau_{zx}}{\partial z} = 0$$

$$\frac{\partial \tau_{xy}}{\partial x} + \frac{\partial \sigma_{yy}}{\partial y} + \frac{\partial \tau_{zy}}{\partial z} = 0$$

$$\frac{\partial \tau_{xz}}{\partial x} + \frac{\partial \tau_{yz}}{\partial y} + \frac{\partial \sigma_{zz}}{\partial z} = \rho g$$

Normal strain and velocity gradients



$$\epsilon_{ij} = \frac{1}{2} \left[\frac{\partial u_j}{\partial x_i} + \frac{\partial u_i}{\partial x_j} \right]$$

$$\boldsymbol{u} = (u, v, w) = (u_x, u_y, u_z)$$

$$\dot{\boldsymbol{\epsilon}} = egin{bmatrix} \dot{\boldsymbol{\epsilon}}_{\chi\chi} & \dot{\boldsymbol{\epsilon}}_{\chi y} & \dot{\boldsymbol{\epsilon}}_{\chi z} \ \dot{\boldsymbol{\epsilon}}_{y\chi} & \dot{\boldsymbol{\epsilon}}_{yz} \ \dot{\boldsymbol{\epsilon}}_{z\chi} & \dot{\boldsymbol{\epsilon}}_{zy} & \dot{\boldsymbol{\epsilon}}_{zz} \end{bmatrix}$$

Strain rate

$$\epsilon_{ij} = \frac{1}{2} \left[\frac{\partial u_j}{\partial x_i} + \frac{\partial u_i}{\partial x_j} \right]$$

For example:

$$\boldsymbol{u} = (u, v, w) = (u_x, u_y, u_z)$$

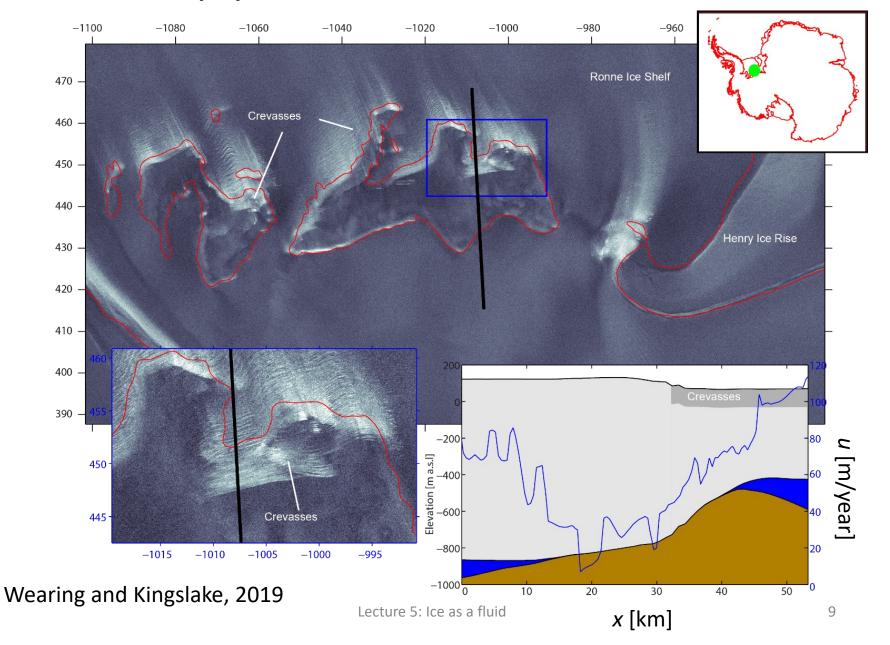
$$\epsilon_{xx}^{\cdot} = \frac{\partial u_x}{\partial x}$$
 or $\epsilon_{xz}^{\cdot} = \frac{1}{2} \left[\frac{\partial u_x}{\partial z} + \frac{\partial u_z}{\partial x} \right]$

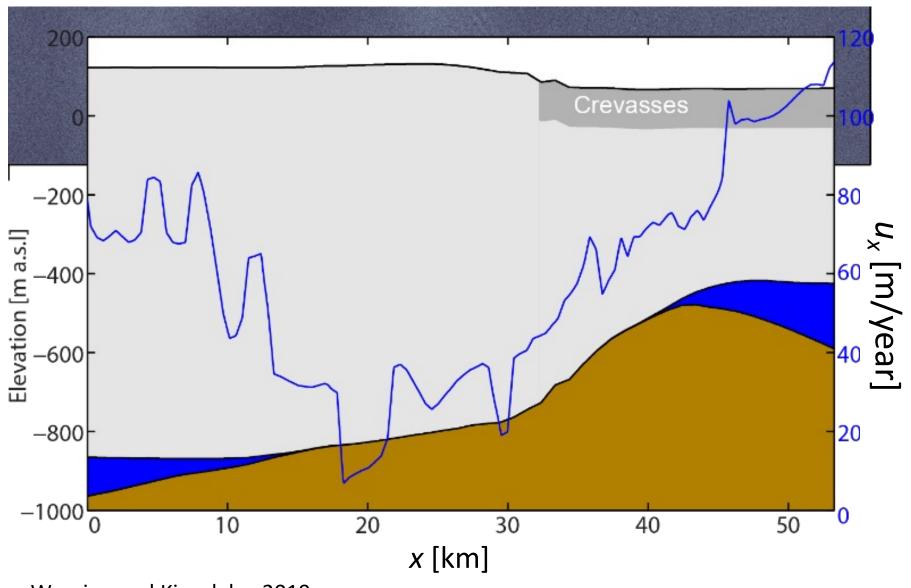
Movie proof of this:

 $https://github.com/ldeo-glaciology/glaciology_practical/raw/main/derivation_movies/strain\%20 and \%20 velocity\%20 gradients. moving the strain of the strai$

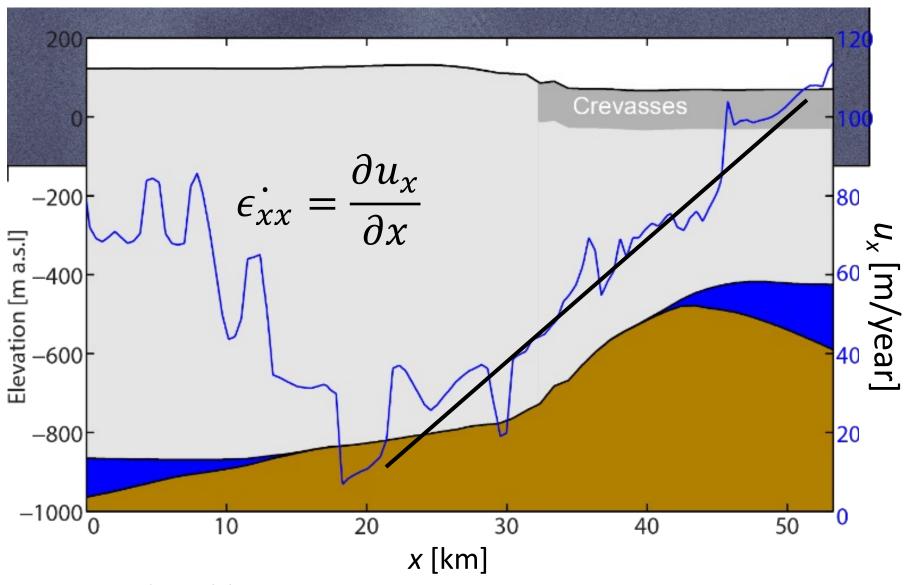
Text proof:

 $https://ldeo-glaciology.github.io/glaciology-intro-book/sections/ice_flow/strain_velocity.html \\$

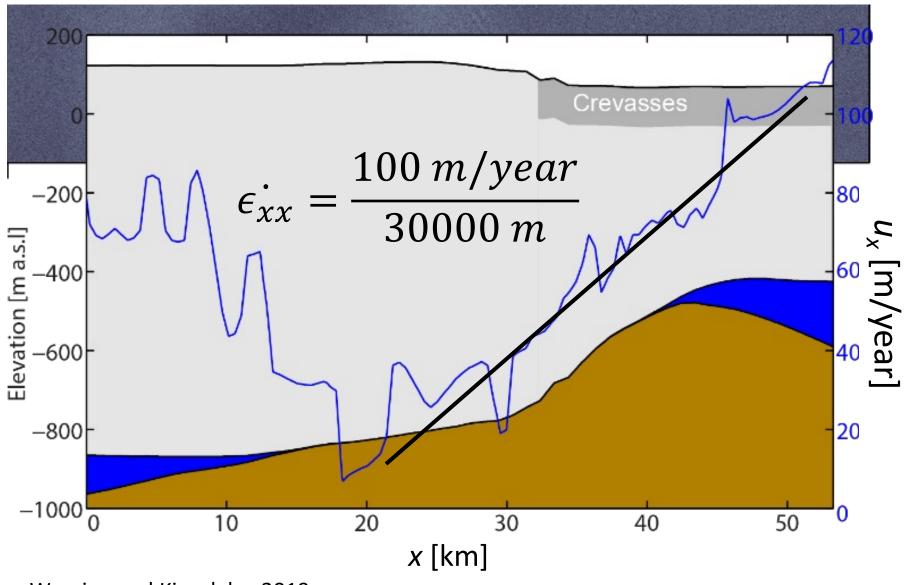




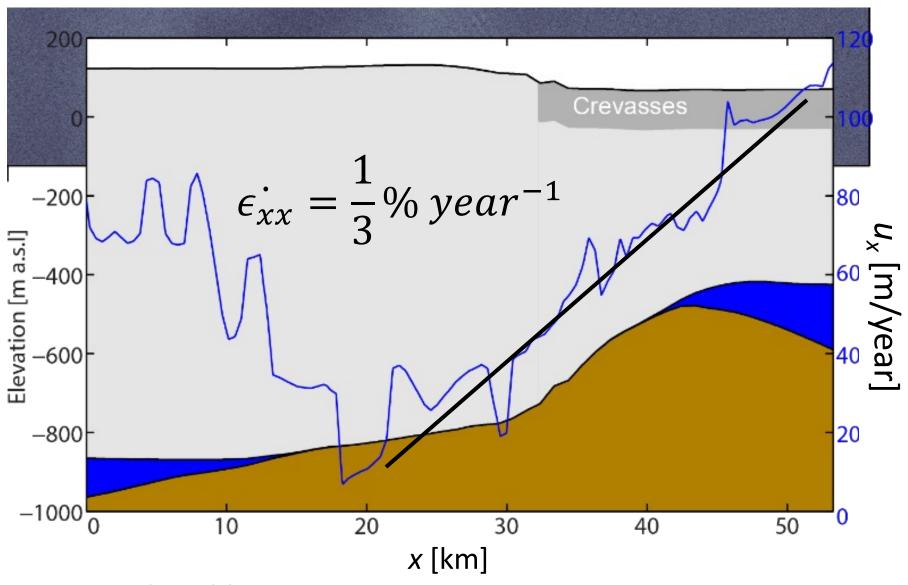
Wearing and Kingslake, 2019



Wearing and Kingslake, 2019



Wearing and Kingslake, 2019



Wearing and Kingslake, 2019

Summary

- Stress balance equations
- Strain rates and velocity gradients
- 'Glen's flow law'
- Ice sheet models make various simplifications
- Our ice-sheet model will only consider τ_{zx}

$$\frac{\partial \sigma_{xx}}{\partial x} + \frac{\partial \tau_{yx}}{\partial y} + \frac{\partial \tau_{zx}}{\partial z} = 0$$

$$\frac{\partial \tau_{xy}}{\partial x} + \frac{\partial \sigma_{yy}}{\partial y} + \frac{\partial \tau_{zy}}{\partial z} = 0$$

$$\frac{\partial \tau_{xz}}{\partial x} + \frac{\partial \tau_{yz}}{\partial y} + \frac{\partial \sigma_{zz}}{\partial z} = \rho g$$

$$\dot{\epsilon_{ij}} = \frac{1}{2} \left[\frac{\partial u_j}{\partial x_i} + \frac{\partial u_i}{\partial x_j} \right]$$

$$\dot{\epsilon}_{zx} = A \tau_{zx}^n$$

Glaciology EESCGU4220

Lecture 7: Ice sheet flow 2

Our ice-sheet model:

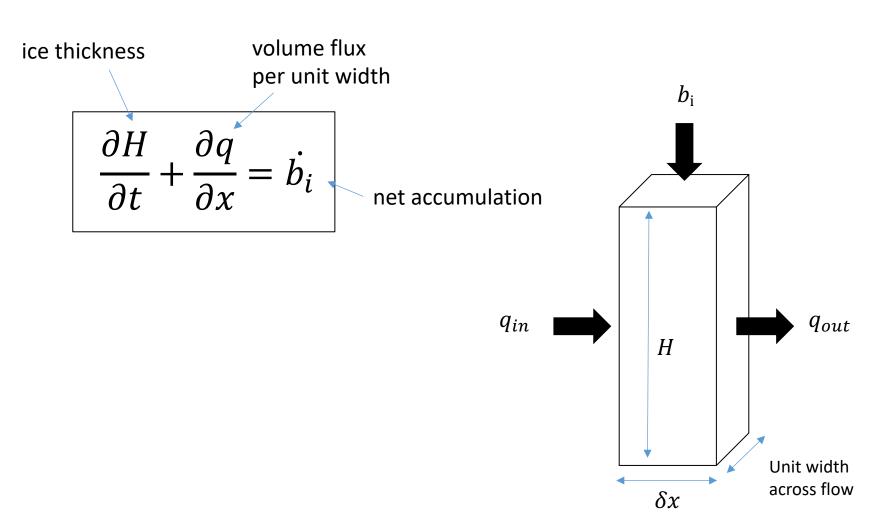
1. Mass conservation

$$\frac{\partial H}{\partial t} + \frac{\partial q}{\partial x} = \dot{b_i}$$

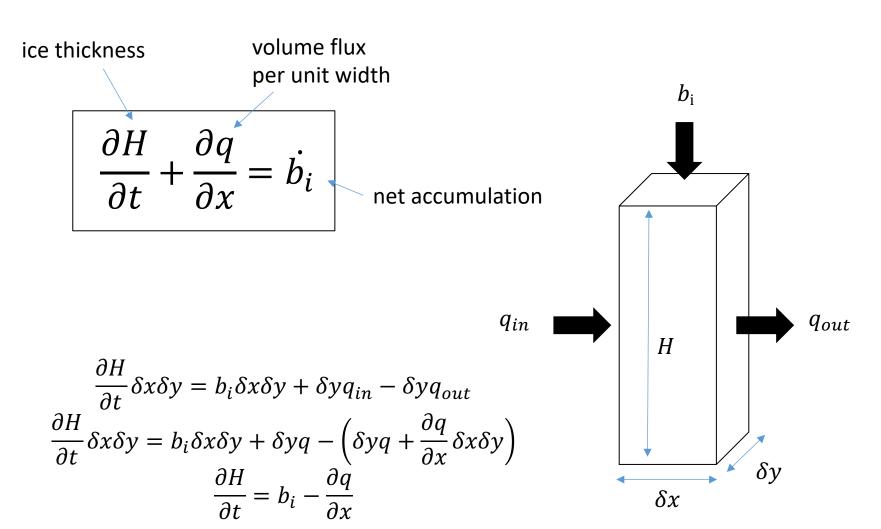
2. Stress balance

$$q = \frac{2A}{n+2} (\rho g \alpha)^n H^{n+2}$$

Mass conservation in a column of ice



Mass conservation in a column of ice



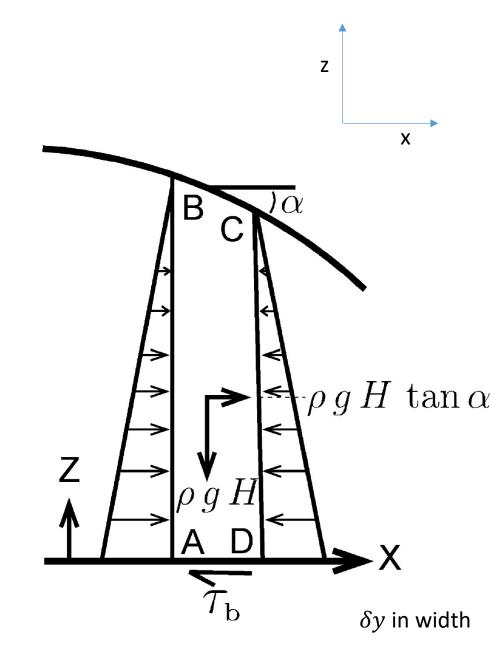
Detailed derivation online here: https://ldeo-glaciology.github.io/glaciology-intro-book/sections/ice-flow/depth-integrated-mass-balance.html

Stress balance

Plan:

- Derive 'driving stress'
- Balance driving stress with viscous stresses in the ice
- Relate these viscous stresses to strain rate with Glen's flow law
- Relate strain rates to gradient in velocity
- Integrate twice to get flux.

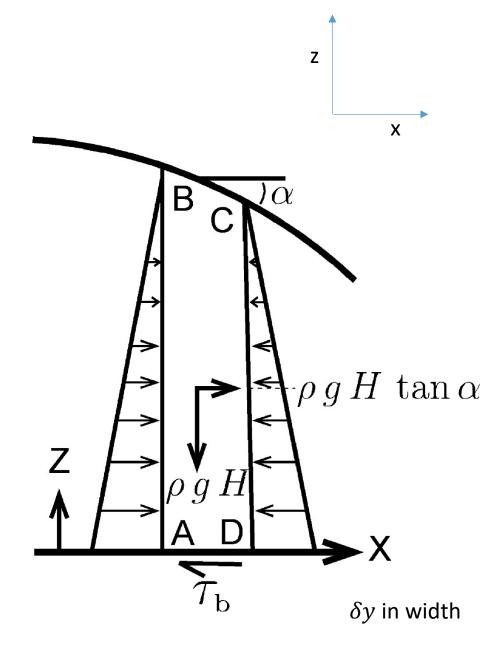
$$\tau_d = -\rho g H \frac{\partial H}{\partial x}$$



$$\tau_d = -\rho g H \frac{\partial H}{\partial x}$$

Total horizontal force acting in one direction integrated over the ice thickness

$$\int_{0}^{\pi} \delta y \rho g(H - z) dz = \frac{1}{2} \delta y \rho g H^{2}$$



$$\tau_d = -\rho g H \frac{\partial H}{\partial x}$$

Total horizontal force acting in one direction integrated over the ice thickness

$$\int_{0}^{H} \delta y \rho g(H - z) dz = \frac{1}{2} \delta y \rho g H^{2}$$

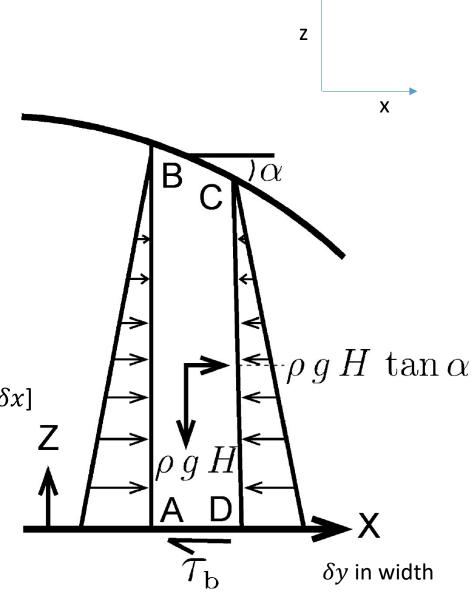
net horizontal force

$$\frac{1}{2}\delta y\rho gH^2 - \left[\frac{1}{2}\delta y\rho gH^2 + \frac{\partial}{\partial x}\left(\frac{1}{2}\delta y\rho gH^2\right)\delta x\right]$$

$$= -\frac{\partial}{\partial x} \left(\frac{1}{2} \delta y \rho g H^2 \right) \delta x$$

Driving force per unit area = driving stress

$$\tau_d = -\rho g H \frac{\partial H}{\partial x}$$



$$\tau_d = -\rho g H \frac{\partial H}{\partial x} \longrightarrow \tau_d = \rho g H \alpha$$

$$\alpha = -\frac{\partial H}{\partial x}$$

• 2-min break

Viscous stresses

We assume the driving stress is balanced by the viscous stresses in the ice.

And we assume that the driving stress is distributed linearly with depth.

Z

driving stress

$$\tau_d = \rho g \alpha H$$

Proof of this: https://ldeo-glaciology.github.io/glaciology- intro-book/sections/ice flow/sia derivation.html#integrate-vertically

viscous stresses in the ice

$$\tau_{zx}(z) = \rho g \alpha H \left(1 - \frac{z}{H} \right)$$

Rheology

Next, we use Glen's flow law to relate these stresses to the shear strain rate ϵ_{zx} .

$$\tau_{zx}(z) = \rho g \alpha H \left(1 - \frac{z}{H} \right)$$

Glen's flow law

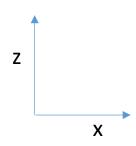
$$\dot{\epsilon_{zx}} = A\tau_{zx}^n$$

$$\epsilon_{zx}^{\cdot} = A(\rho g \alpha H)^n \left(1 - \frac{z}{H}\right)^n$$

Strain rate → velocity gradient

Now we use the definition of strain rate to relate this to the vertical gradient in horizontal velocity.

$$\epsilon_{zx}^{\cdot} = A(\rho g \alpha H)^n \left(1 - \frac{z}{H}\right)^n$$

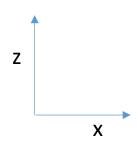


$$\epsilon_{zx}^{\cdot} = \frac{1}{2} \left[\frac{\partial w}{\partial x} + \frac{\partial u}{\partial z} \right]$$

Strain rate → velocity gradient

Now we use the definition of strain rate to relate this to the vertical gradient in horizontal velocity.

$$\epsilon_{zx}^{\cdot} = A(\rho g \alpha H)^n \left(1 - \frac{z}{H}\right)^n$$



$$\dot{\epsilon_{zx}} = \frac{1}{2} \left[\frac{\partial w}{\partial x} + \frac{\partial u}{\partial z} \right]$$

$$\epsilon_{zx}^{\cdot} = \frac{1}{2} \frac{\partial u}{\partial z}$$

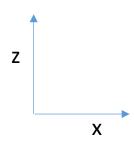
Strain rate \rightarrow velocity gradient

Now we use the definition of strain rate to relate this to the vertical gradient in horizontal velocity.

$$\epsilon_{zx}^{\cdot} = A(\rho g \alpha H)^n \left(1 - \frac{z}{H}\right)^n$$



$$\frac{\partial u}{\partial z} = 2A(\rho g \alpha H)^n \left(1 - \frac{z}{H}\right)^n$$



$$\epsilon_{zx} = \frac{1}{2} \left[\frac{\partial w}{\partial x} + \frac{\partial u}{\partial z} \right]$$

$$\epsilon_{zx}^{\cdot} = \frac{1}{2} \frac{\partial u}{\partial z}$$

Horizontal velocity

We next integrate vertically to get the horizontal velocity as a function of depth.

$$\frac{\partial u}{\partial z} = 2A(\rho g \alpha H)^n \left(1 - \frac{z}{H}\right)^n$$

$$\zeta = 1 - \frac{z}{H}$$

Integrate by substitution

$$u(z) = 2A(\rho g \alpha H)^n \left(\frac{H}{n+1}\right) \left(1 - \left(1 - \frac{z}{H}\right)^{n+1}\right)$$

Ice flux

We next integrate vertically again to get ice flux *q*

$$\frac{\partial u}{\partial z} = 2A(\rho g \alpha H)^n \left(1 - \frac{z}{H}\right)^n$$

$$\zeta = 1 - \frac{z}{H}$$

Integrate by substitution

$$u(z) = 2A(\rho g \alpha H)^n \left(\frac{H}{n+1}\right) \left(1 - \left(1 - \frac{z}{H}\right)^{n+1}\right)$$

Integrate one more time to get depth-integrated volume-flux per unit width:

$$q = \frac{2A}{n+2} (\rho g \alpha)^n H^{n+2}$$

Full derivation here: https://ldeo-

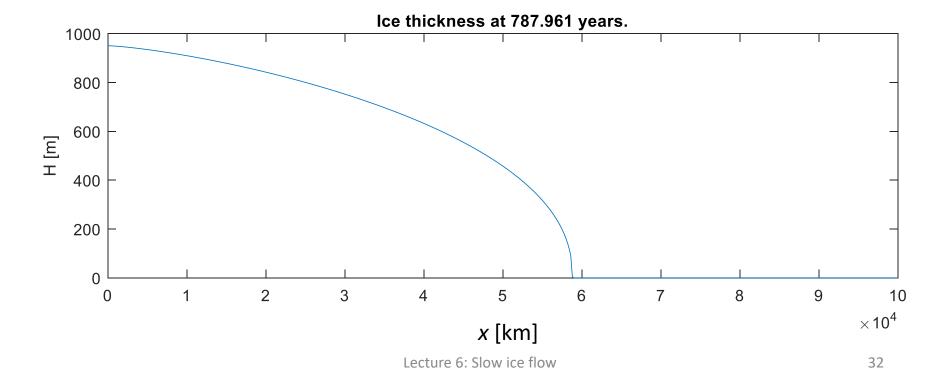
The complete model:

Depth-integrated volume-flux per unit width:

mass conservation:

$$\frac{\partial H}{\partial t} + \frac{\partial q}{\partial x} = \dot{b_i}$$

$$q = \frac{2A}{n+2} (\rho g \alpha)^n H^{n+2}$$



Flux strongly depends on H. Why?

$$q = \frac{2A}{n+2} (\rho g \alpha)^n H^{n+2}$$

$$q \propto H^{n+2}$$

With n = 3 that's H^5 !!!

n comes from the nonlinear rheology and the driving stress being proportional to H

+1 comes from the fact that thicker ice flows faster

+1 comes from the fact that thicker ice moves more ice for a given flow speed

Summary

- Driving stress drives ice flow. It is proportional to ice thickness and surface slope.
- Ice sheet models use depth-integrated mass conservation to evolve the surface up and down.
 - --- a balance between flux coming in, flux going out and net specific accumulation.
- We derived a 'simple' equations for the flux
- Together the mass conservation and flux equations make up our ice sheet model. s

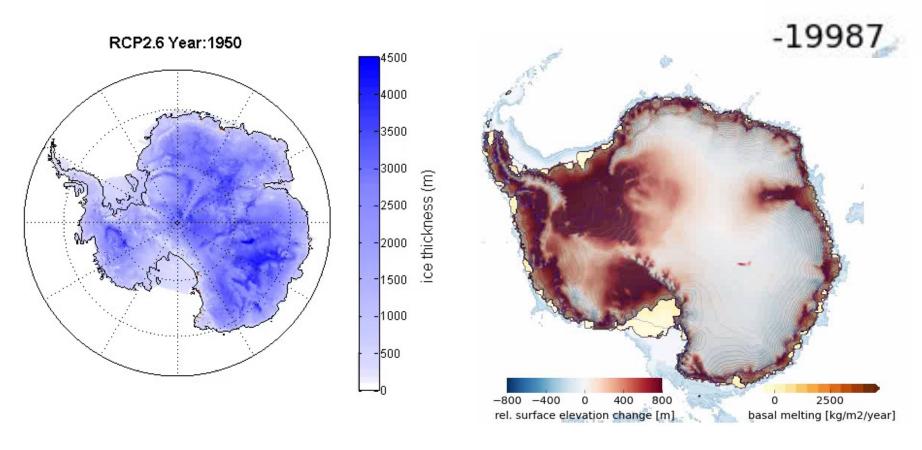
For Tuesday

- (1) go through these derivations and make sure you understand each step (you don't have to memorize anything, just make sure you understand each step).
- (2) revise the "finite-difference method" of solving differential equations. There are many online resources: e.g.

https://www.ljll.math.upmc.fr/frey/cours/UdC/ma691/ma691_ch6.pdf

Two state-of-the-art hydrid models

Take into account GIA, ocean/atmosphere forcing, calving, surface mass balance, spatially-variable bed conditions.



Lecture 6: Slow ice flow

DeConto and Pollard (2013)

Penn State University model

Kingslake et al. (in prep.)

Parallel Ice Sheet model (PISM)₃₆

References

• DeConto, Robert M., and David Pollard. "Contribution of Antarctica to past and future sealevel rise." *Nature* 531, no. 7596 (2016): 591.

• Hindmarsh, R.C.A., 2004. A numerical comparison of approximations to the Stokes equations used in ice sheet and glacier modeling. *Journal of Geophysical Research: Earth Surface*, 109(F1).