$6.02214 \times 10^{23}$  $N_{\rm Av}$ :  $\text{mol}^{-1}$  $1.6605 \times 10^{-27}$ 1 amu: kg  $1.38065 \times 10^{-23}$  $\rm J~K^{-1}$  $8.61734 \times 10^{-5}$  $eV K^{-1}$  $k_{\rm B}$ :  $J K^{-1} mol^{-1}$  $8.2057 \times 10^{-2}$  l atm mol<sup>-1</sup> K<sup>-1</sup> R: 8.314472 ${
m J}~{
m s}^{-1}~{
m m}^{-2}~{
m K}^{-4}$  $5.6704 \times 10^{-8}$  $\sigma_{\mathrm{SB}}$ :  $\rm m\ s^{-1}$  $2.99792458 \times 10^{8}$ c:  $6.62607 \times 10^{-34}$  $4.13566 \times 10^{-15}$ h: J seV s  $1.05457 \times 10^{-34}$  $6.58212 \times 10^{-16}$ eV sJ s $\hbar$ : hc: 1239.8 eV nm $1.60218 \times 10^{-19}$  $\mathbf{C}$ e:  $9.10938215 \times 10^{-31}$  $MeV c^{-2}$ kg 1: 0.5109989  $m_e$ :  $C^2 J^{-1} m^{-1}$  $e^2 \text{ Å}^{-1} \text{ eV}^{-1}$  $8.85419 \times 10^{-12}$  $5.52635 \times 10^{-3}$  $e^2/4\pi\epsilon_0$ :  $2.30708 \times 10^{-28}$ J m 14.39964 eV Å  $0.529177 \times 10^{-10}$ 0.529177Å  $\mathbf{m}$  $a_0$ : 27.212  $E_{\mathrm{H}}$ : Ha eV

Table 1: Key units in Physical Chemistry

## 1 The Classical Foundations

#### 1.1 Lecture 0: Introduction

- 1. Burning lighter
- 2. Foundations of Physical Chemistry
  - (a) Quantum mechanics
  - (b) Statistical mechanics
  - (c) Thermodynamics, kinetics, spectroscopy
  - (d) Physical and chemical properties of matter

#### 1.2 Lecture 1: Basic statistics

- 1. Discrete probability distributions—Coin flip
  - (a) Example of Bernoulli trial,  $2^n$  possible outcomes from n flips
  - (b) Number of ways to get i heads in n flips,  ${}_{n}C_{i} = n!/i!(n-i)!$
  - (c) Probability of *i* heads  $P_i \propto {}_nC_i$
  - (d) Normalized probability,  $\tilde{P}_i = P_i / \sum_i P_i = {}_n C_i / 2^n$
  - (e) Expectation value  $\langle i \rangle = \sum_i i \tilde{P}_i$
- 2. Continuous distributions—temperature
  - (a) Probability density P(x) has units 1/x
  - (b) Normalized  $\tilde{P}(x) = P(x) / \int P(x) dx$
  - (c) (Unitless) probability  $a < x < b = \int_a^b \tilde{P}(x) dx$

- (d) Expectation value  $\langle f(x) \rangle = \int f(x) \tilde{P}(x) dx$
- (e) Mean =  $\langle x \rangle$
- (f) Mean squared =  $\langle x^2 \rangle$
- (g) Variance  $\sigma^2 = \langle x^2 \rangle \langle x \rangle^2$
- (h) Standard deviation  $\Delta x = \sigma$

#### 3. Boltzmann distribution

- (a)  $P(E) \propto e^{-E/k_BT}$ , in some sense the definition of temperature
- (b) Energy and its units
- (c) Absolute temperature and its units
- (d)  $k_BT$  as an energy scale, 0.026 eV at 298 K
- (e) Gravity example
  - i. E(h) = mgh, linear, continuous energy spectrum
  - ii. molecule vs car in a gravitational field (Table 2)
  - iii. Barometric law for gases,  $P = P_0 e^{-mgh/k_BT}$
- (f) Kinetic energy in 1-D example

i. 
$$KE = \frac{1}{2}mv_x^2$$

ii. 
$$P_{1D}(v_x) = \left(\frac{m}{2\pi k_B T}\right)^{1/2} \exp\left(-\frac{m|v_x|^2}{2k_B T}\right)$$

iii. Gaussian distribution, mean  $\mu$ , variance  $\sigma^2$ 

$$G(x) = \frac{1}{\sigma\sqrt{2\pi}} \exp\left(-\frac{(x-\mu)^2}{2\sigma^2}\right)$$

- iv. By inspection,  $\mu = \langle v_x \rangle = 0$ ,  $\sigma^2 = \langle v_x^2 \rangle = k_B T/m$
- v. Molecule vs car again
- (g) Equipartition energy freely exchanged between all degrees of freedom

**Table 2:** Car vs gas molecule at the earth's surface

	car	gas molecule
$\overline{m}$	$1000\mathrm{kg}$	$1 \times 10^{-26}  \mathrm{kg}$
h	$1\mathrm{m}$	1 m
mgh	$9800\mathrm{J}$	$9.8 \times 10^{-26} \mathrm{J}$
	$6.1 \times 10^{22} \mathrm{eV}$	$6.1 \times 10^{-7}  \text{eV}$
T	$298\mathrm{K}$	$298\mathrm{K}$
$k_BT$	$0.026\mathrm{eV}$	$0.026\mathrm{eV}$
$mgh/k_BT$	$2.4 \times 10^{24}$	$2.3\times10^{-5}$
$P(1 {\rm m})/P(0)$	$e^{-2.4 \times 10^{-24}}$	0.99998
$\langle h \rangle$	$0\mathrm{m}$	$42\mathrm{km}$
$\langle v_x \rangle^{1/2}$	$2 \times 10^{-12} \mathrm{m/s}$	$640\mathrm{m/s}$

**Table 3:** Energy conversions and correspondences

	J	eV	Hartree	$kJ \text{ mol}^{-1}$	$\mathrm{cm}^{-1}$
1 J =	1	$6.2415 \times 10^{18}$	$2.2937 \times 10^{17}$	$6.0221 \times 10^{20}$	$5.0340 \times 10^{22}$
1  eV =	$1.6022 \times 10^{-19}$	1	0.036748	96.485	8065.5
1 Ha =	$4.3598 \times 10^{-18}$	27.212	1	2625.6	219474.6
$1 \text{ kJ mol}^{-1} =$	$1.6605 \times 10^{-21}$	0.010364	$3.8087 \times 10^{-4}$	1	83.5935
$1 \text{ cm}^{-1} =$	$1.986410^{-23}$	$1.23984 \times 10^{-4}$	$4.55623 \times 10^{-6}$	0.011963	1

```
import numpy as np
1
2
    import matplotlib.pyplot as plt
3
    R0 = 8.31441
                   J/mol K
5
    mass = 28. /1000 kg/mol N2
6
    def Boltzmann(E,T):
        return np.exp(E/(RO*T))/(RO*T)
8
9
    def MB1D(v,T):
10
        return np.sqrt(mass/(2*np.pi*R0*T))*np.exp((mass*v*v)/(2*R0*T))
11
12
    def MB(c,T):
13
        K = 0.5 * mass * c *c
14
        degeneracy = 4 * np.pi * c * c
15
        normalization = (mass/(2*np.pi*R0*T))**1.5
16
17
        return normalization*degeneracy*Boltzmann(K,T)
18
    energy = np.linspace(0,3000,1500)
19
    velocity = np.linspace(1000,1000,1000)
20
    speed = np.linspace(0,1500,1000)
21
^{22}
    plt.figure()
23
24
    for Temperature in [100,300,1000]:
       Probability = Boltzmann(energy, Temperature)
25
26
       plt.plot(Probability,energy,label=0 K.format(Temperature))
27
    legend = plt.legend()
28
^{29}
    plt.ylabel(Energy (J/mol))
30
    plt.xlabel(Probability (mol/J))
    plt.title(Boltzmann distribution at various temperatures)
32
    plt.savefig(./Images/Boltzmann.png)
33
34
    plt.figure()
35
    for Temperature in [100,200,300,400,500]:
36
        Probability = MB1D(velocity, Temperature)
37
38
        plt.plot(velocity,Probability,label= K.format(Temperature))
39
    legend=plt.legend()
40
41
    plt.xlabel(Velocity (m/s))
    plt.ylabel(Probability (1/(m/s)))
42
43
    plt.title(Boltzmann distribution at various temperatures)
    plt.savefig(./Images/MB1D.png)
44
45
46
    plt.figure()
    for Temperature in [100,200,300,400,500]:
47
        Probability = MB(speed, Temperature)
48
        plt.plot(speed,Probability,label= K.format(Temperature))
49
50
    legend=plt.legend()
51
    plt.xlabel(Speed (m/s))
52
    plt.ylabel(Probability (1/(m/s)))
```

plt.title(Boltzmann distribution at various temperatures)
plt.savefig(./Images/MB.png)

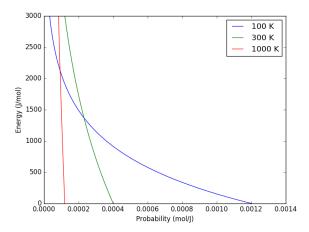


Figure 1: Boltzmann distribution at various temperatures

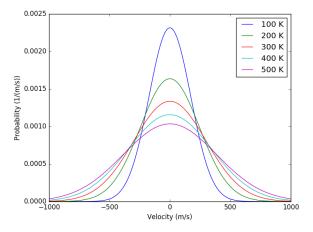


Figure 2: One-dimensional (Gaussian) velocities of  $N_2$  gas

## 1.3 Lecture 2: Kinetic theory of gases

#### 1. Postulates

- (a) Gas is composed of molecules in constant random, thermal motion
- (b) Molecules only interact by perfectly elastic collisions
- (c) Volume of molecules is << total volume

## 2. Maxwell-Boltzmann distribution of molecular speeds

(a) Speed 
$$v = \sqrt{v_x^2 + v_y^2 + v_z^2}$$

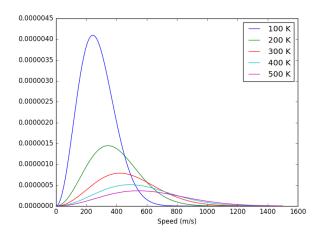


Figure 3: Maxwell-Boltzmann speed distribution of N<sub>2</sub> gas

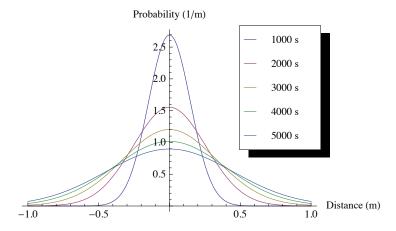
- (b)  $P_{MB}(v)dv = P_{1D}(v_x)P_{1D}(v_y)P_{1D}(v_z) * degeneracy(v)dv$
- (c) mean speeds  $\propto \sqrt{T}$
- (d) mean energy  $U = \frac{3}{2}RT$  and heat capacity  $C_v = \frac{3}{2}R$

#### 3. Flux and pressure

- (a) Velocity flux  $j(v_x)dv_x=v_x\frac{N}{V}P(v_x)dv_x,$  molecules /area /time / $v_x$
- (b) Wall collisions,  $J_w$ , total collisions /area /time
- (c) Momentum exchange, pressure, ideal gas law

#### 4. Collisions and mean free path

- (a) Collision cross section  $\sigma = \pi d^2$ , size of molecule
- (b) Molecular collisions, z per molecule and  $z_{AA}$  per volume
- (c) Mean free path,  $\lambda$ , mean distance between collisions



**Figure 4:** Diffusional spreading,  $\sqrt{\langle x^2 \rangle} = \sqrt{2Dt}$ 

Table 4: Kinetic theory of gases key equations

Boltzmann distribution $(g(E))$ : degeneracy of $E$ )	$P(E) = g(E)e^{-E/k_BT}$
Maxwell-Boltzmann distribution	$P_{\text{MB}}(v) = 4\pi v^2 \left(\frac{m}{2\pi k_B T}\right)^{3/2} \exp\left(-\frac{mv^2}{2k_B T}\right)$
Mean and RMS speeds	$\langle v \rangle = \left(\frac{8k_BT}{\pi m}\right)^{1/2} \qquad \langle v^2 \rangle^{1/2} = \left(\frac{3k_BT}{m}\right)^{1/2}$
Pressure	$\langle P \rangle = \frac{\Delta p}{\Delta t} = m \frac{N}{V} \frac{1}{3} \langle v^2 \rangle = \frac{N k_B T}{V} = \frac{nRT}{V}$
Wall collision frequency	$J_W = \frac{1}{4} \frac{N}{V} \langle v \rangle = \frac{P}{(2\pi m k_B T)^{1/2}}$
Molecular collision frequency	$z = \sqrt{2}\sigma \langle v \rangle \frac{N}{V} = \frac{4\sigma P}{(\pi m k_B T)^{1/2}}$
Total collisions	$z_{AA}=rac{1}{2}rac{N}{V}z$
Mean free path	$\lambda = \frac{\langle v \rangle}{z} = \frac{V}{\sqrt{2}\sigma N}$
Graham's effusion law	$\frac{dN}{dt} = \text{Area} \cdot J_w \propto 1/m^{1/2}$
Effusion from a vessel	$P = P_0 e^{-t/\tau}, \tau = \frac{V}{A} \left(\frac{2\pi m}{k_B T}\right)^{1/2}$
Self-diffusion constant	$D_{11} = \frac{1}{3} \langle v \rangle \lambda$
Diffusion rate	$\langle x^2 \rangle^{1/2} = \sqrt{2Dt}  \langle r^2 \rangle^{1/2} = \sqrt{6Dt}$
Einstein-Smoluchowski equation	$D_{11} = \frac{\delta^2}{2\tau}$
Stokes-Einstein equation for liquids	$D_{11} = \frac{k_B T}{4\pi \eta r}$ "Slip" boundary
	$D_{\text{Brownian}} = \frac{k_B T}{6\pi \eta r}$ "Stick" boundary

## 1.4 Lecture 3: Transport

- 1. Effusion and Graham's law, effusion rate  $\propto MW^{-1/2}$
- 2. Fick's first law: net flux proportional to concentration gradient
  - (a)  $j_x = -D\frac{dc}{dx}$
  - (b) Self-diffusion constant,  $D = \frac{1}{3}\lambda \langle v \rangle$
- 3. Knudsen diffusion,  $D = \frac{1}{3}l\langle v \rangle$
- 4. Fick's second law: time evolution of concentration gradient
  - (a) Continuity with no advection:  $\frac{\partial c}{\partial t} = -\nabla \cdot \vec{j} + \text{gen}$
  - (b) One-dimension:  $\frac{dc}{dt} = D \frac{d^2c}{dx^2}$
  - (c) Diffusion has Gaussian probability distribution:  $c(x,t)/c_0 = [2\sqrt{\pi Dt}]^{-1} \exp(-x^2/4Dt)$
- 5. Seeing is believing—Brownian motion
  - (a) Seemingly random motion of large particles ("dust") due to "kicks" from invisible molecules
  - (b) Einstein receives Nobel Prize for showing:
    - i. Motion follows same Gaussian diffusion behavior
    - ii. From steady-state arguments in a field, diffusion constant is ratio of Boltzmann energy,  $k_BT$ , to mobility
    - iii. Mobility inversely related to viscosity
  - (c) Stokes-Einstein equation
  - (d) Allows measurement of Avogadro's number, final proof of kinetic theory
  - (e) Similar model for diffusion of liquid molecules, slip boundary
- 6. Random walk model of diffusion
  - (a) Binomial distribution
  - (b) Large N and Stirling approximation
  - (c) Einstein-Smoluchowski relation

# 2 Quantum Mechanics: Blurred Lines Between Particles and Waves

#### 2.1 Lecture 4: Duality and demise of classical physics

#### 2.1.1 Properties of waves

- 1. Traveling waves, standing waves
- 2. interference, diffraction
- 3. Expected energy of a classical oscillator,  $\langle \epsilon \rangle_{\nu} = k_B T$  for all  $\nu$

Table 5: Classical waves

The wave equation	$\frac{\partial^2 \Psi(x,t)}{\partial x^2} = \frac{1}{v^2} \frac{\partial^2 \Psi(x,t)}{\partial t^2}$	
General solution	$\Psi(x,t) = A\sin(kx - \omega t)$	
Wavelength (distance)	$\lambda = 2\pi/k$	
Frequency (/time)	$ u = 2\pi/\omega $	
Speed	$v = \lambda \nu$	
Amplitude (distance)	A	
Energy	$E \propto A^2$	
Standing wave	$\Psi(x,t) = A\sin(kx)\cos(\omega t),  k = n\pi/a$	

#### 2.1.2 Blackbody radiation

- 1. Hohlraum spectrum (like the sun) empirically observed to obey:
  - (a) Stefan-Boltzmann law, total irradiance
  - (b) Wien's displacement law
- 2. Rayleigh-Jeans predicts spectrum using classical physics
  - (a) standing waves + classical oscillators  $\rightarrow$  ultraviolet catastrophe
- 3. Planck model
  - (a) Energy spectrum of oscillators are quantized,  $\epsilon_{\nu} = nh\nu$
  - (b) Expected energy of a quantized oscillator,  $\langle \epsilon \rangle_{\nu} = h \nu / \left( e^{h \nu / k_B T} 1 \right)$
  - (c) Correctly reproduces Stefan-Boltzmann and Wien Laws!

### 2.1.3 Heat capacities of solids

- 1. Law of DuLong and Pettite,  $C_v = 3R$ , fails at low T
- 2. Einstein model
  - (a) Atomic vibrations are quantized,  $\epsilon_n = nh\nu$
  - (b) Heat capacity goes to zero at low T

#### 2.1.4 Photoelectric effect

- 1. Stopping potential and work function,  $E_{\text{kinetic}} = h\nu W$
- 2. Kinetic energy varies with light frequency, number of electrons varies with light intensity

#### 2.1.5 Compton effect

- 1. light scattering of electrons changes  $\lambda$
- 2. Photon properties,  $\epsilon = h\nu, p = h/\lambda$

#### 2.1.6 Wave-particle duality

### 2.1.7 Rutherford, planetary model of atom

1. Inconsistent with Maxwell's equations

#### 2.1.8 Bohr model of H atom

- 1. Discrete H energy spectrum and Rydberg formala
- 2. Bohr model (the old quantum mechanics)
  - (a) Stable electron "orbits," quantized angular momentum
  - (b) Light emission corresponds to orbital jumps,  $\nu = \Delta E/h$
  - (c) Bohr equations
  - (d) Comparison with Rydberg formula
  - (e) Failure for larger atoms

#### 2.1.9 de Broglie relation

- 1.  $\lambda = h/p$  universally
- 2. Relation to Bohr orbits
- 3. Davison and Germer experiment,  $e^-$  diffraction off Ni

```
import numpy as np
    import matplotlib.pyplot as plt
    hc = 1239.8
                      eV nm
4
    c = 2.9979e8 * 1.e9
    k = 8.61734e5
                   eV /K
6
    hck = hc/k
    def Irrad(wl,T):
9
          return (8. * np.pi * hc * c * wl**5) / (np.exp(hck/(wl*T))1)
10
    def PlanckEnergy(wl,T):
11
          return (hc/wl) / (np.exp(hck/(wl*T))1)
12
13
   plt.figure()
14
    wl=np.linspace(100,5000,1000)
15
    for T in [1000.,2000.,3000.,4000.,5000.]:
16
17
        Intensity = Irrad(wl,T)
        plt.plot(wl,Intensity,label= K.format(T))
18
19
20
   legend=plt.legend()
21
    plt.xlabel(Wavelength (nm))
    plt.ylabel(Irradiance (eV/nm3/s))
22
     plt.title(Boltzmann distribution at various temperatures)
23
   plt.savefig(./Images/BlackBody.png)
25
    plt.figure()
26
27
    color=[red,orange,green,blue,violet]
    wl=np.linspace(100,20000,1000)
28
    for T in [1000.,2000.,3000.,4000.,5000.]:
        Energy = PlanckEnergy(w1,T)
30
31
        plt.plot(wl,Energy,label= K.format(T),color=color[0])
        kT = k*T
32
        plt.plot([100,max(wl)],[kT,kT],ls=,color=color.pop(0))
33
```

```
legend=plt.legend()
plt.xlabel(Wavelength (nm))
plt.ylabel(Energy (eV))
plt.title(Boltzmann distribution at various temperatures)
plt.savefig(./Images/Planck.png)
```

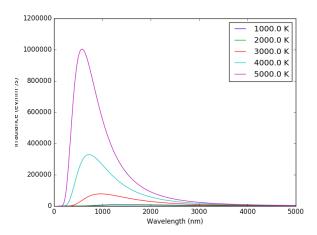


Figure 5: Blackbody irradiance

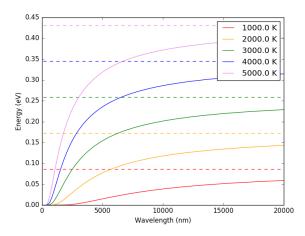


Figure 6: Average energy of a Planck quantized oscillator

## 2.2 Lecture 5: Postulates of quantum mechanics

## 2.2.1 Schrödinger equation describes wave-like properties of matter

### 2.2.2 Born interpretation

- 1. wavefunction is a probability amplitude
- 2. wavefunction squared is probability density

**Table 6:** The new physics

Stefan-Boltzmann Law	$\int I(\lambda, T) d\lambda = \sigma_{\rm SB} T^4$
Wien's Law	$\lambda_{\rm max}T=2897768~{\rm nm~K}$
Rayleigh-Jeans eq	$I(\lambda, T) = \frac{8\pi}{\lambda^4} k_B T c$
Blackbody irradiance	$I(\lambda, T) = \frac{8\pi}{\lambda^5} \frac{hc^2}{e^{hc/\lambda k_B T} - 1}$
Einstein crystal	$C_v = 3R \left(\frac{h\nu}{k_B T}\right)^2 \frac{e^{h\nu/k_B T}}{\left(e^{h\nu/k_B T} - 1\right)^2}$
Photon energy	$\epsilon = h \nu = h c / \lambda$
Rydberg equation	$\nu = R_H c \left( 1/n^2 - 1/k^2 \right)$
Bohr equations $n = 1, 2, \dots$	$l_n = n\hbar$ $r_n = n^2 \left(\frac{4\pi\epsilon_0 \hbar^2}{e^2 m_e}\right) = n^2 a_0$ $E_n = -\frac{m_e e^4}{8\epsilon_0^2 h^2} \frac{1}{n^2} = -\frac{E_H}{2} \frac{1}{n^2}$ $p_n = \frac{e^2}{4\pi\epsilon_0} \frac{m_e}{\hbar} \frac{1}{n} = p_0 \frac{1}{n}$
de Broglie equation	$\lambda = \frac{h}{p}$

#### 2.2.3 Postulates

- 1. Wavefunction contains all information about a system
- 2. Operators used to extract that information
  - (a) QM operators are Hermitian
  - (b) Have eigenvectors and real eigenvalues,  $\hat{O}\psi_i = o\psi_i$
  - (c) Are orthogonal,  $\langle \psi_i | \psi_j \rangle = \delta_{ij}$
  - (d) Always observe an eigenvalue when making an observation
- 3. Expectation values
- 4. Energy-invariant wavefunctions given by Schrödinger equation
- 5. Uncertainty principle

#### 2.2.4 Particle in a box illustrations

- 2.3 Lecture 6: Particle in a box model
- 2.3.1 Particle between infinite walls, electron confined in a wire
- 2.3.2 Classical solution, either stationary or uniform bouncing back and forth

### 2.3.3 One-dimesional QM solutions

- 1. Schrödinder equation and boundary conditions
- 2. discrete, quantized solutions
- 3. standing waves,  $\lambda = 2L/n$ , n-1 nodes, non-uniform probability
- 4. Ho paper, STM of Pd wire
- 5. zero point energy and uncertainty
- 6. correspondence principle
- 7. superpositions

#### 2.3.4 Finite walls and tunneling

- 1. Potential well of finite depth  $V_0$
- 2. Finite number of bound states
- 3. Classical region,  $\psi(x) e^{ikx} + e^{-ikx}, k = \sqrt{2mE}/\hbar$
- 4. "Forbidden" region,  $\psi(x)$   $e^{\kappa x} + e^{-\kappa x}$ ,  $\kappa = \sqrt{2m(V_0 E)}/\hbar$
- 5. Non-zero probability to "tunnel" into forbidden region
- 6. Tunneling between two adjacent wells: chemical bonding, STM, nanoelectronics
- 7. H atom tunneling: NH<sub>3</sub> inversion, H transfer, kinetic isotope effect

#### Table 7: Postulates of Non-relativistic Quantum Mechanics

Postulate 1: The physical state of a system is completely described by its wavefunction  $\Psi$ . In general,  $\Psi$  is a complex function of the spatial coordinates and time.  $\Psi$  is required to be:

- I. Single-valued
- II. continuous and twice differentiable
- III. square-integrable  $(\int \Psi^* \Psi d\tau)$  is defined over all finite domains)
- IV. For bound systems,  $\Psi$  can always be normalized such that  $\int \Psi^* \Psi d\tau = 1$

**Postulate 2:** To every physical observable quantity M there corresponds a Hermitian operator  $\hat{M}$ . The only observable values of M are the eignevalues of  $\hat{M}$ .

Physical quantity	Operator	Expression
Position $x, y, z$	$\hat{x},\hat{y},\hat{z}$	$x\cdot,y\cdot,z\cdot$
		a
Linear momentum $p_x, \dots$	$\hat{p}_x, \dots$	$-i\hbar \frac{\partial}{\partial x}, \dots$
Angular momentum $l_x, \dots$	$\hat{p}_x,\dots$	$-i\hbar \frac{\partial}{\partial x}, \dots$ $-i\hbar \left( y \frac{\partial}{\partial z} - z \frac{\partial}{\partial y} \right), \dots$
Angular momentum $t_x, \dots$	$p_x, \dots$	$-in\left(g\frac{\partial}{\partial z}-z\frac{\partial}{\partial y}\right),\dots$
Kinetic energy $T$	$\hat{T}$	$-\frac{\hbar^2}{2m}\nabla^2$
Potential energy $V$	$\hat{V}$	1/(r +)
Total energy $E$	$\hat{H}$	$-\frac{\hbar^2}{2m}\nabla^2 + V(\mathbf{r},t)$
rotal ellergy E	11	$-\frac{1}{2m}\mathbf{v}^{-}+\mathbf{v}^{-}(\mathbf{I},t)$

**Postulate 3:** If a particular observable M is measured many times on many identical systems is a state  $\Psi$ , the average resuts with be the expectation value of the operator  $\hat{M}$ :

$$\langle M \rangle = \int \Psi^*(\hat{M}\Psi) d\tau$$

Postulate 4: The energy-invariant states of a system are solutions of the equation

$$\hat{H}\Psi(\mathbf{r},t) = i\hbar \frac{\partial}{\partial t}\Psi(\mathbf{r},t)$$

$$\hat{H} = \hat{T} + \hat{V}$$

The time-independent, stationary states of the system are solutions to the equation

$$\hat{H}\Psi(\mathbf{r}) = E\Psi(\mathbf{r})$$

Postulate 5: (The uncertainty principle.) Operators that do not commute  $(\hat{A}(\hat{B}\Psi) \neq \hat{B}(\hat{A}\Psi))$  are called *conjugate*. Conjugate observables cannot be determined simultaneously to arbitrary accuracy. For example, the standard deviation in the measured positions and momenta of particles all described by the same  $\Psi$  must satisfy  $\Delta x \Delta p_x \geq \hbar/2$ .

 $V(x) = \begin{cases} 0 & 0 < x < L \\ \infty & x \le 0 \text{ or } x \ge L \end{cases}$   $\psi_n(x) = \sqrt{\frac{2}{L}} \sin\left(\frac{n\pi x}{L}\right)$   $E_n = \frac{n^2\pi^2\hbar^2}{2mL^2}, n = 1, 2, \dots$  Particle in a box solutions for an electron in a 1 nm box

Table 8: Particle-in-a-box model

#### 2.3.5 Multiple dimensions

1. separation of variables, one quantum number for each dimension

## 2.3.6 Introduce Pauli principle for fermions?

#### 2.4 Lecture 7: Harmonic oscillator

#### 2.4.1 Classical harmonic oscillator

- 1. Hooke's law,  $F = -k(x x_0)$ , k spring constant
- 2. Continuous sinusoidal motion
- 3.  $x(t) = A\sin(\frac{k}{\mu})^{1/2}t, \nu = \frac{1}{2\pi}(\frac{k}{\mu})^{1/2}, E = \frac{1}{2}kA^2$
- 4. Exchanging kinetic and potential energies

#### 2.4.2 Quantum harmonic oscillator

- 1. Solutions like P-I-A-B, waves, nodes, even/odd symmetry
- 2. Zero-point energy

- 3. Expectation values  $\langle x^2 \rangle = \alpha^2(v+1/2), \langle V(x) \rangle = \frac{1}{2}h\nu(v+\frac{1}{2})$
- 4. Classical turning point and tunneling
- 5. Classical limiting behavior

### 2.4.3 HCl example

1. Reduced mass,  $\frac{1}{\mu} = \frac{1}{m_A} + \frac{1}{m_B}$ 

Table 9: Harmonic oscillator model

$$V(x) = \frac{1}{2}kx^2, -\infty < x < \infty$$

$$\psi_v(x) = N_v H_v(x/\alpha) e^{-x^2/2\alpha^2}, v = 0, 1, 2, \dots$$

$$\alpha = (\hbar^2/\mu k)^{1/4}, N_v = (2^v v! \alpha \sqrt{\pi})^{-1/2}$$

$$\frac{\text{Hermite polynomials}}{H_0(y) = 1}$$

$$H_1(y) = 2y$$

$$H_2(y) = 4y^2 - 2$$

$$H_{n+1}(y) = 2yH_n(y) - 2nH_{n-1}(y)$$

$$E_v = (v + \frac{1}{2})h\nu, v = 0, 1, 2, \dots$$
Harmonic oscillator functions

Harmonic oscillator functions

Harmonic oscillator functions

### 2.5 Lecture 8: Rigid Rotor

### 2.5.1 Classical rigid rotor

1. Compare rotation about an axis vs linear motion

- 2. Moment of intertia  $I = \mu r^2$
- 3. Angular momentum,  $l=I\omega=r\times p,\,T=l^2/2I$ 
  - (a) Angular momentum and energy continuous variable

#### 2.5.2 Quantum rotor in a plane

- 1. Angular momentum and kinetic energy operators in polar coordinates,  $\hat{l}_z=-i\hbar\frac{d}{d\phi}$
- 2. Eigenfunctions and energy spectrum
- 3. No zero point energy
- 4. Angular momentum eignefunctions,  $l_z = m_l \hbar$
- 5. Energy superpositions and localization

## 2.5.3 Quantum rotor in 3-D

- 1. Angular momentum and kinetic energy operators in spherical coordinates
- 2. Spherical harmonic solutions,  $Y_{lm_l}$
- 3. Azimuthal QN  $l = 0, 1, \dots$
- 4. Magnetic QN  $m_l = -l, -l+1, ..., l$
- 5. Energy spectrum, 2l + 1 degeneracy
- 6. Vector model can only know total total |L| and  $L_z$
- 7. Wavefunctions look like atomic orbitals, l nodes

#### 2.5.4 Particle angular momentum

- 1. Fermions, mass, half-integer spin
  - (a) Electron,  $s = 1/2, m_s = \pm 1/2$
- 2. Bosons, force-carrying, integer spin

### 2.6 Lecture 9: Spectroscopy

#### 2.6.1 Spectroscopy is quantitative measurement of interaction of light with matter

- 1. Observed  $I(\nu)/I(\nu_0)$
- 2. Bohr condition,  $|E_f E_i|/h = \nu = c\tilde{\nu} = c/\lambda$
- 3. Intensities determined by state populations and transition probabilities

Table 10: 2-D rigid rotor model

$$V(\phi) = 0, 0 \le \phi \le 2\pi$$

$$\hat{H} = -\frac{\hbar^2}{2I} \frac{\partial^2}{\partial \phi^2}, \qquad I = \mu R^2$$

$$\psi_{m_l}(\phi) = \frac{1}{\sqrt{2\pi}} e^{-im_l \phi}, m_l = 0, \pm 1, \pm 2, \dots$$

$$E_{m_l} = \frac{m_l^2 \hbar^2}{2I^2}$$

$$L_z = m_l \hbar$$

$$L_z = m_l \hbar$$

$$U_{20} = \frac{1}{120} \frac{(m=1) + (m=-1)}{120} \frac{(m=1) \cdot (m=-1)}{120} \frac{(m=1) \cdot (m=-1)}{120} \frac{(m=1) \cdot (m=-1)}{120} \frac{(m=1) \cdot (m=-1)}{120} \frac{(m=2) \cdot (m=-2)}{120} \frac{(m=2) \cdot (m=-2)}$$

#### 2.6.2 Einstein coefficients

- 1. Stimulated absorption,  $dn_1/dt = -n_1B\rho(\nu)$
- 2. Stimulated emission,  $dn_2/dt = -n_2B\rho(\nu)$
- 3. Spontaneous emission,  $dn_2/dt = -n_2A$ ,  $A = \left(\frac{8\pi h\nu^3}{c^3}\right)B$
- 4. 1/A = lifetime

### 2.6.3 Transition probability

- 1. Einstein coefficient  $B_{if} = \frac{|\mu_{if}|^2}{6\epsilon_0\hbar^2}$
- 2. Classical electric dipole,  $\overrightarrow{\mu}=q\cdot\overrightarrow{l}$ , quantum dipole operator  $\hat{\mu}=e\cdot\overrightarrow{r}$

- 3. Transition dipole moment,  $\mu_{if} = \left(\frac{d\mu}{dx}\right) \langle \psi_i | \hat{\mu} | \psi_f \rangle$
- 4. Selection rules—conditions that make  $\mu_{if}$  non-zero, "allowed" vs "forbidden" transitions

## 2.7 Lecture 10: Vibrational and rotational spectroscopy

#### 2.7.1 Vibrational transitions

- 1. Gross selection rule: dynamic dipole  $d\mu/dx$  non-zero
- 2. Specific selection rule: dipole integral  $\langle \psi_v | \hat{\mu} | \psi_{v'} \rangle = 0$  unless  $\Delta v = \pm 1$
- 3. Allowed  $\Delta E = h\nu$
- 4. Boltzmann, v = 1 states dominate at normal T

#### 2.7.2 Vibrational spectroscopy

- 1. Diatomics, homo- vs. heteronuclear
- 2. Polyatomics, 3n-6 (3n-5 for linear polyatomic) vibrational modes
- 3.  $CO_2$  example

### 2.7.3 Raman spectroscopy

#### 2.7.4 Diatomic rotational spectroscopy

- 1. Rotational constant  $B = \hbar/4\pi Ic \text{ cm}^{-1}$ ,  $I = \mu R^2$
- 2. Gross selection rule: dipole moment non-zero
- 3. Specific selection rule:  $\Delta l = \pm 1, \ \Delta m_l = 0, \pm 1$
- 4.  $\Delta \tilde{E}_l = 2B(l+1) \text{ cm}^{-1}$
- 5. Rotational state populations

#### 2.7.5 Polyatomic rotational spectroscopy

#### 2.7.6 Vibration-rotation spectroscopy

- 1. Harmonic oscillator + rigid rotor
- 2. Selection rules:  $\Delta v = \pm 1, \Delta l = \pm 1$
- 3. R branch:  $\Delta \tilde{E} = \tilde{\nu} + 2B(l+1), \Delta l = 1$
- 4. P branch:  $\Delta \tilde{E} = \tilde{\nu} 2B(l), \Delta l = -1$

#### 2.7.7 Anharmonicity, Morse potential

## 3 Statistical Mechanics: The Bridge from the Tiny to the Many

Table 11: 3-D rigid rotor model

$$V(\theta,\phi) = 0, 0 \le \phi \le 2\pi, 0 \le \theta < \pi$$

$$\hat{L}^2 = -\hbar^2 \left[ \frac{1}{\sin^2 \theta} \frac{\partial^2}{\partial \phi^2} + \frac{1}{\sin \theta} \frac{\partial}{\partial \theta} \left( \sin \theta \frac{\partial}{\partial \theta} \right) \right]$$

$$\hat{H}_{\text{rot}} = \frac{1}{2I} \hat{L}^2$$

$$Y_{lm_l}(\theta,\phi) = N_l^{|m|} P_l^{|m|} (\cos(\theta)) e^{im_l \phi}$$

$$l = 0, 1, 2, \dots, \qquad m_l = 0, \pm 1, \dots, \pm l$$

$$E_l = \frac{\hbar^2}{2I} l(l+1)$$

$$|L| = \hbar \sqrt{l(l+1)}, L_z = m_l \hbar$$

$$V_{\text{re}}, V_{\text{to}}$$

$$V_{\text{to}}, V_{\text{to}}$$