

4125 HW 4

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3.10a

The heat transfer in this process is, using 334 J/g as the latent heat of fusion of water,

$$Q = (334 \text{ J/g})(30 \text{ g}) = 10 \text{ kJ}$$

The change in entropy is therefore

$$\Delta S = \frac{Q}{T} = \frac{10 \text{ kJ}}{273 \text{ K}} = 36.7 \text{ J/K}$$

3.10b

The heat capacity of the melted ice cube is $C_V = (4.2 \text{ J/g} \cdot \text{K})(30 \text{ g}) = 126 \text{ J/K}$ The change in entropy is therefore, presuming (justifiably) a constant heat capacity,

$$\Delta S = C_V \ln \left(\frac{T_f}{T_i} \right) = (126 \text{ J/K}) \ln \left(\frac{298 \text{ K}}{273 \text{ K}} \right) = 137.5 \text{ J/K}$$

3.10c

The heat required to melt the ice cube is $C_V(25 \text{ K}) = 3150 \text{ J}$ The temoerature of the kitchen doesn't change much, so the change in its entropy is

$$\Delta S = -\frac{Q}{T} = -\frac{3150 \text{ J}}{298 \text{ K}} = -10.6 \text{ J/K}$$

where the negative sign is added because Q is flowing out of the kitchen.

3.10d

Adding the changes in entropy,

$$\Delta S_{total} = 137.5 \text{ J/K} - 10.6 \text{ J/K} = 126.9 \text{ J/K}$$

3.16a

When the memory was initialized, it is in one of $2^{2^{33}}$ microstates. For each final state of the system, the previous state could have been any of those initial microstates, so the multiplicity is the above number. The entropy is then

$$S = k \ln \Omega = (1.38 \times 10^{-23} \text{ J/K}) \ln \left(2^{2^{33}} \right) = 8.22 \times 10^{-14} \text{ J/K}$$

3.16b

The corresponding heat would be

$$Q = ST = (8.22 \times 10^{-14} \text{ J/K})(300 \text{ K}) = 2.56 \times 10^{-11} \text{ J}$$

which is a likely immesurable amount of additional energy.

3.23

The expression for the temperature of a two-state paramagnet is

$$\frac{1}{T} = \frac{k}{2\mu B} \ln \left(\frac{N - U/\mu B}{N + U/\mu B} \right) \Leftrightarrow x = \frac{\mu B}{kT} = \frac{1}{2} \ln \left(\frac{N - U/\mu B}{N + U/\mu B} \right)$$

Writing $U/\mu B = N - 2N_{\uparrow}$ from the definition of U ,

$$x = \frac{1}{2} \ln \left(\frac{N_{\uparrow}}{N - N_{\uparrow}} \right) = \frac{1}{2} \ln \left(\frac{N_{\uparrow}}{N_{\downarrow}} \right)$$

Plugging this in to the given expression,

$$\begin{aligned} S &= Nk [\ln (2 \cosh x) - x \tanh x] \\ Nk \left\{ \ln \left(2 \cosh \left[\frac{1}{2} \ln \left(\frac{N_{\uparrow}}{N_{\downarrow}} \right) \right] \right) - \frac{1}{2} \ln \left(\frac{N_{\uparrow}}{N_{\downarrow}} \right) \tanh \left[\frac{1}{2} \ln \left(\frac{N_{\uparrow}}{N_{\downarrow}} \right) \right] \right\} \\ &= Nk \left\{ \ln \left(2 \cdot \frac{1}{2} \left[\frac{\exp \left(\ln \left[\frac{N_{\uparrow}}{N_{\downarrow}} \right] \right) + 1}{\exp \left(\frac{1}{2} \ln \left[\frac{N_{\uparrow}}{N_{\downarrow}} \right] \right)} \right] \right) \right. \\ &\quad \left. - \frac{1}{2} \ln \left(\frac{N_{\uparrow}}{N_{\downarrow}} \right) \left[\frac{\exp \left(\ln \left[\frac{N_{\uparrow}}{N_{\downarrow}} \right] \right) - 1}{\exp \left(\ln \left[\frac{N_{\uparrow}}{N_{\downarrow}} \right] \right) + 1} \right] \right\} \\ &= Nk \left\{ \ln \left(\frac{\frac{N_{\uparrow}}{N_{\downarrow}} + 1}{\frac{N_{\uparrow}}{N_{\downarrow}} \sqrt{e}} \right) - \frac{1}{2} \ln \left(\frac{N_{\uparrow}}{N_{\downarrow}} \right) \left[\frac{\frac{N_{\uparrow}}{N_{\downarrow}} - 1}{\frac{N_{\uparrow}}{N_{\downarrow}} + 1} \right] \right\} \\ &= Nk \left\{ \ln(N_{\uparrow} + N_{\downarrow}) - \ln(N_{\downarrow}) - \ln(N_{\uparrow}) + \ln(N_{\downarrow}) - \frac{1}{2} - \frac{1}{2} \frac{N_{\uparrow} - N_{\downarrow}}{N_{\uparrow} + N_{\downarrow}} (\ln(N_{\uparrow}) - \ln(N_{\downarrow})) \right\} \end{aligned}$$

$$\begin{aligned}
&= Nk \left\{ \ln N - \ln N_{\uparrow} - \frac{1}{2} - \frac{1}{2} \frac{N_{\uparrow}}{N} \ln N_{\uparrow} + \frac{1}{2} \frac{N - N_{\uparrow}}{N} \ln(N - N_{\uparrow}) \right\} \\
&= Nk \ln N - Nk \ln N_{\uparrow} - \frac{N_{\uparrow}}{2} k \ln N_{\uparrow} + \frac{Nk}{2} \ln(N - N_{\uparrow}) - \frac{N_{\uparrow}k}{2} \ln(N - N_{\uparrow}) - \frac{Nk}{2}
\end{aligned}$$

Neglecting the merely large terms,

3.25a

$$\begin{aligned}
S = k \ln \Omega &= k \ln \left[\left(\frac{q+N}{q} \right)^q \left(\frac{q+N}{N} \right)^N \right] = k [q \ln(q+N) - q \ln q + N \ln(q+N) - N \ln N] \\
&= k [N \ln N - q \ln q + (q+N) \ln(q+N)]
\end{aligned}$$

The neglected term is a division of the multiplicity expression by $\sqrt{2\pi q(q+N)/N}$ which, after a logarithm, is asymptotic to $-\log N + \log q + \log(q+N)$ as $q, N \rightarrow \infty$ which is $o(N \log N + q \log q + (q+N) \log(q+N))$ in the same limit. It therefore is acceptable to neglect this for large q and N .

3.25b

Since $U = \epsilon q \Leftrightarrow q = \frac{U}{\epsilon}$, the entropy may be written

$$S = k \left[N \ln N - \frac{U}{\epsilon} (\ln U - \ln \epsilon) + \left(\frac{U}{\epsilon} + N \right) \ln \left(\frac{U}{\epsilon} + N \right) \right]$$

Differentiating with respect to internal energy at constant N and V ,

$$\begin{aligned}
\frac{1}{kT} &= \frac{1}{k} \left(\frac{\partial S}{\partial U} \right)_{N,V} = \frac{\ln \epsilon}{\epsilon} - \left(\frac{1}{\epsilon} + \frac{\ln U}{\epsilon} \right) + \left(\frac{1}{\epsilon} + \frac{1}{\epsilon} \ln \left[\frac{U}{\epsilon} + N \right] \right) \\
\Leftrightarrow T &= \frac{\epsilon}{k \ln \left(\frac{U+\epsilon N}{U} \right)} = \frac{\epsilon}{k \ln \left(1 + \frac{\epsilon N}{U} \right)}
\end{aligned}$$

3.25c

Solving for U ,

$$1 + \frac{\epsilon N}{U} = e^{\epsilon/kT} \Leftrightarrow U = \frac{\epsilon N}{e^{\epsilon/kT} - 1}$$

Differentiating with respect to temperature,

$$C_V = \left(\frac{\partial U}{\partial T} \right)_{N,V} = \frac{\epsilon^2 N e^{\epsilon/kT}}{kT^2 (e^{\epsilon/kT} - 1)^2}$$

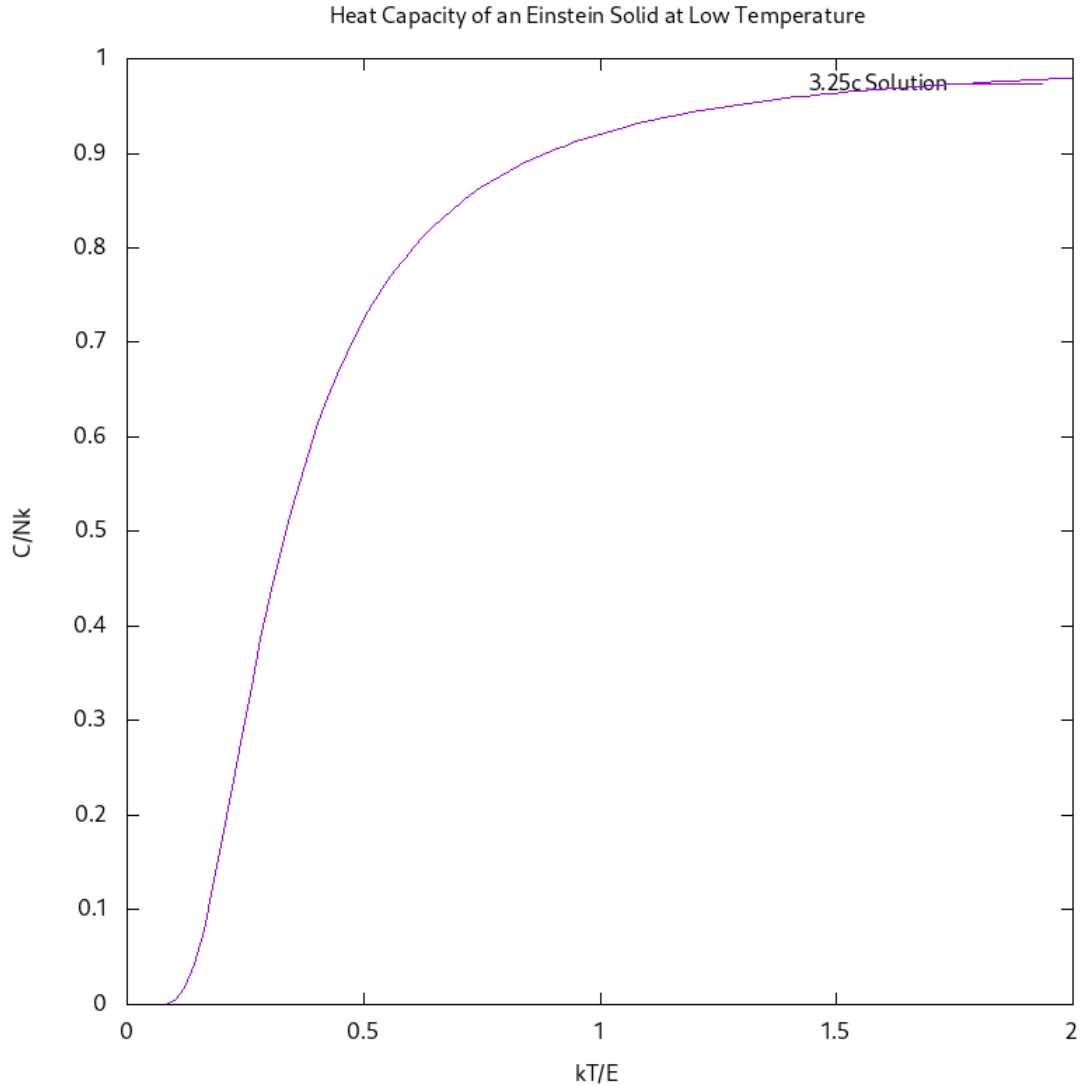
3.25d

Using $e^x = 1 + x + O(x^2)$ as $x \rightarrow 0$ and $\frac{1}{T} \rightarrow 0$ as $T \rightarrow \infty$,

$$C_V \approx \frac{\epsilon^2 N}{k} \frac{1 + \frac{\epsilon}{kT}}{T^2 \left(1 + \frac{\epsilon}{kT} - 1 \right)^2} = k \epsilon^2 N \left(\frac{1}{\epsilon^2} + \frac{1}{\epsilon kT} \right) \approx Nk$$

3.25e

Using gnuplot,



This predicts an exponential falloff of the heat capacity at low temperature, which is corroborated qualitatively by the experimental data presented in figure 1.14. From this plot, $\frac{kT}{\epsilon} = 0.5$ is a little past the inflection point of the curve. Looking at the experimental data, the analogous point is around 40 K for lead, 110 K for aluminum, and 400 K for diamond (the x range of the graph is too short to really tell, but it appears the slope is decreasing at 400 K, so it's a decent estimate). Solving for ϵ in the initial equation,

$$\epsilon_{Pb} = 2kT_{Pb} = 2(1.38 \times 10^{-23} \text{ J/K})(40 \text{ K}) = 1.1 \times 10^{-21} \text{ J} = 0.0069 \text{ eV}$$

$$\epsilon_{Al} = 2kT_{Al} = 2(1.38 \times 10^{-23} \text{ J/K})(110 \text{ K}) = 3.04 \times 10^{-21} \text{ J} = 0.019 \text{ eV}$$

$$\epsilon_{Dia} = 2kT_{Dia} = 2(1.38 \times 10^{-23} \text{ J/K})(400 \text{ K}) = 1.10 \times 10^{-20} \text{ J} = 0.069 \text{ eV}$$

Couldn't progress further due to the odd deadline of 6pm instead of the standard midnight time.