

2231 HW 6

Duncan Wilkie

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The multipole expansion is given in general by

$$V(r) = \frac{1}{4\pi\epsilon_0} \sum_{n=0}^{\infty} \frac{1}{r^{n+1}} \int (r')^n P_n(\cos \alpha) \rho(r') dV$$

where r is the distance from the origin to the point of consideration, r' is the distance from the origin to the differential volume element of consideration, α is the angle between r and r' , and $\rho(r')$ is the charge density at r' . The first three terms of this sum are, since $\rho(r') = \delta(\rho - R)\delta(z)\lambda$ where ρ and z are the radial and z components of r' respectively in cylindrical coordinates.

$$V_0(r) = \frac{1}{4\pi\epsilon_0 r} \int \rho(r') dV = \frac{\lambda R}{2\epsilon_0 r}$$

$$\begin{aligned} V_1(r) &= \frac{1}{4\pi\epsilon_0 r^2} \int r' \cos \alpha \rho(r') dV = \frac{1}{4\pi\epsilon_0 r^2} \int_{-\infty}^{\infty} \int_0^{2\pi} \int_0^{\infty} r' \cos \alpha \delta(r' - R) \delta(z) \lambda r' dr' d\theta' dz' \\ &= \frac{\lambda R^2}{8\pi\epsilon_0 r^2} \int_0^{2\pi} \cos \alpha d\theta \end{aligned}$$

Noting

$$\begin{aligned} \vec{r} &= r \cos \theta \hat{i} + r \sin \theta \hat{j} + z \hat{k}, \vec{r}' = R \cos \theta' \hat{i} + R \sin \theta' \hat{j} \\ \Rightarrow \vec{r} \cdot \vec{r}' &= rR \cos \alpha = rR \cos \theta \cos \theta' + rR \sin \theta \sin \theta' \\ \Rightarrow V_1(r) &= \frac{\lambda R^2}{8\pi\epsilon_0 r^2} \int_0^{2\pi} \cos \theta \cos \theta' + \sin \theta \sin \theta' d\theta' = 0 \end{aligned}$$

$$\begin{aligned} V_2(r) &= \frac{1}{4\pi\epsilon_0 r^3} \int (r')^2 P_2(\cos \alpha) \rho(r') dV = \frac{1}{4\pi\epsilon_0 r^3} \int_{-\infty}^{\infty} \int_0^{2\pi} \int_0^{\infty} (r')^2 \left(\frac{3}{2} \cos^2 \alpha - \frac{1}{2} \right) \delta(r' - R) \delta(z) \lambda r' dr' d\theta' dz \\ &= \frac{\lambda R^3}{8\pi\epsilon_0 r^3} \int_0^{2\pi} \left(\frac{3}{2} \cos^2 \alpha - \frac{1}{2} \right) d\theta' = \frac{\lambda R^3}{8\pi\epsilon_0 r^3} \int_0^{2\pi} \left(\frac{3}{2} (\cos \theta \cos \theta' + \sin \theta \sin \theta')^2 - \frac{1}{2} \right) d\theta' \\ &= \frac{\lambda R^3}{8\pi\epsilon_0 r^3} \int_0^{2\pi} \left(\frac{3}{2} [\cos^2 \theta \cos^2 \theta' + 2 \cos \theta \cos \theta' \sin \theta \sin \theta' + \sin^2 \theta \sin^2 \theta'] - \frac{1}{2} \right) d\theta' \\ &= \frac{\lambda R^3}{8\pi\epsilon_0 r^3} \left[\frac{3}{2} \left(\cos^2 \theta \int_0^{2\pi} \cos^2 \theta' d\theta' + \frac{\sin(2\theta)}{2} \int_0^{2\pi} \sin(2\theta') d\theta' + \sin^2 \theta \int_0^{2\pi} \sin^2 \theta' d\theta' \right) - \int_0^{2\pi} \frac{1}{2} \right] \\ &= \frac{\lambda R^3}{8\pi\epsilon_0 r^3} \left[\frac{3}{2} (\pi) - \pi \right] = \frac{\lambda R^3}{16\epsilon_0 r^3} \end{aligned}$$

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By symmetry, \vec{p} is in the \hat{z} direction, so

$$\begin{aligned} |\vec{p}| &= \int \vec{r} \rho(\vec{r}) dV = \int z k \cos \theta dA = k \int (R \cos \theta) \cos \theta R^2 \sin \theta d\theta d\phi = 2\pi k R^3 \int \cos^2 \theta \sin \theta d\theta \\ &= -2\pi k R^3 \frac{\cos^3(\theta)}{3} \Big|_0^\pi = -\frac{4k\pi R^3}{3} \end{aligned}$$

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The dipole moment of this arrangement should be in the \hat{z} direction once again. This allows computation of

$$\begin{aligned} |\vec{p}| &= \int \vec{r} \rho(\vec{r}) dV = \int_N z \rho_0 dV - \int_S z \rho_0 dV \\ &= \int_0^{\frac{\pi}{2}} \int_0^{2\pi} \int_0^R \rho_0 R \cos \theta R^2 \sin \theta dr d\phi d\theta - \int_{\frac{\pi}{2}}^\pi \int_0^{2\pi} \int_0^R \rho_0 R \cos \theta R^2 \sin \theta dr d\phi d\theta \\ &= 2\pi \rho_0 R^3 \left(-\frac{\cos(2\theta)}{4} \Big|_0^{\frac{\pi}{2}} + \frac{\cos(2\theta)}{4} \Big|_{\frac{\pi}{2}}^\pi \right) = 2\pi \rho_0 R^3 \end{aligned}$$

The corresponding electric field is

$$E_{dip}(r, \theta) = \frac{p}{4\pi\epsilon_0 r^3} (2 \cos \theta \hat{r} + \sin \theta \hat{\theta}) = \frac{\rho_0 R^3}{2\epsilon_0 r^3} (2 \cos \theta \hat{r} + \sin \theta \hat{\theta})$$

The force is attractive.

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The dipole moment of the atom due to the field of the charge is $\vec{p} = \alpha \frac{kq}{r^2} \hat{r}$. The force of attraction is

$$F = (\vec{p} \cdot \nabla) \vec{E} = \frac{\alpha k q}{r^2} \frac{\partial}{\partial r} \frac{k q}{r^2} = -\frac{2\alpha k^2 q^2}{r^5}$$

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We apply the method of images. A dipole $-\vec{p}$ at $-z$ instead of the plane is equivalent to this problem. The field due to this image dipole is

$$\vec{E} = \frac{p}{4\pi\epsilon_0 (2z)^3} (2 \cos \theta \hat{r} + \sin \theta \hat{\theta})$$

and the associated torque is, with the origin at the image dipole,

$$\begin{aligned} \tau &= \vec{p} \times \vec{E} = (p \cos \theta \hat{r} + p \sin \theta \hat{\theta}) \times \left(\frac{p}{4\pi\epsilon_0 r^3} (2 \cos \theta \hat{r} + \sin \theta \hat{\theta}) \right) \\ &= \frac{p^2}{4\pi\epsilon_0 r^3} \left(-\frac{\sin(2\theta)}{2} \hat{\phi} \right) = -\frac{p^2 \sin(2\theta)}{8\pi\epsilon_0 r^3} \hat{\phi} \end{aligned}$$