

Deriving The 3D QM Solutions From Scratch

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The 3D TISE is

$$-\frac{\hbar^2}{2m}\nabla^2\psi + V\psi = E\psi$$

The spherical Laplacian is

$$\nabla^2 = \frac{1}{r^2} \frac{\partial}{\partial r} \left(r^2 \frac{\partial}{\partial r} \right) + \frac{1}{r^2 \sin \theta} \frac{\partial}{\partial \theta} \left(\sin \theta \frac{\partial}{\partial \theta} \right) + \frac{1}{r^2 \sin^2 \theta} \left(\frac{\partial^2}{\partial \phi^2} \right)$$

Presuming ψ separates in spherical coordinates as $\psi = R(r)Y(\theta, \phi)$, the TISE is

$$\begin{aligned} & -\frac{\hbar^2}{2m} \left[\frac{Y}{r^2} \frac{\partial}{\partial r} \left(r^2 \frac{\partial R}{\partial r} \right) + \frac{R}{r^2 \sin \theta} \frac{\partial}{\partial \theta} \left(\sin \theta \frac{\partial Y}{\partial \theta} \right) + \frac{R}{r^2 \sin^2 \theta} \left(\frac{\partial^2 Y}{\partial \phi^2} \right) \right] + VRY = ERY \\ & \Leftrightarrow \left\{ \frac{1}{R} \frac{\partial}{\partial r} \left(r^2 \frac{\partial R}{\partial r} \right) + \frac{2mr^2}{\hbar^2} [E - V] \right\} + \frac{1}{Y} \left\{ \frac{1}{\sin \theta} \frac{\partial}{\partial \theta} \left(\sin \theta \frac{\partial Y}{\partial \theta} \right) + \frac{1}{\sin^2 \theta} \frac{\partial^2 Y}{\partial \phi^2} \right\} = 0 \end{aligned}$$

This separates the equation; the left side depends solely on r , and the right only on θ, ϕ , implying each term is constant, as if the left term varies in r , the right term cannot vary to compensate and keep the whole equation equal to zero.

We therefore reduce this to the system

$$\frac{1}{R} \frac{\partial}{\partial r} \left(r^2 \frac{\partial R}{\partial r} \right) + \frac{2mr^2}{\hbar^2} [E - V] = \ell(\ell + 1),$$

$$\frac{1}{Y} \left\{ \frac{1}{\sin \theta} \frac{\partial}{\partial \theta} \left(\sin \theta \frac{\partial Y}{\partial \theta} \right) + \frac{1}{\sin^2 \theta} \frac{\partial^2 Y}{\partial \phi^2} \right\} = -\ell(\ell + 1)$$

First, we consider the angular equation: expanding,

$$\sin \theta \left(\cos \theta \frac{\partial Y}{\partial \theta} + \sin \theta \frac{\partial^2 Y}{\partial \theta^2} \right) + \frac{\partial^2 Y}{\partial \phi^2} = -\ell(\ell + 1)Y \sin^2 \theta$$

We separate again. Presuming $Y(\theta, \phi) = \Theta(\theta)\Phi(\phi)$,

$$\sin \theta \Phi(\phi) \left(\cos \theta \frac{\partial \Theta}{\partial \theta} + \sin \theta \frac{\partial^2 \Theta}{\partial \theta^2} \right) + \Theta \frac{\partial^2 \Phi}{\partial \phi^2} = -\ell(\ell + 1)\Theta\Phi \sin^2 \theta$$

Dividing by Y and moving the term on the right over,

$$\left\{ \frac{\sin \theta}{\Theta} \left(\cos \theta \frac{\partial \Theta}{\partial \theta} + \sin \theta \frac{\partial^2 \Theta}{\partial \theta^2} \right) + \ell(\ell + 1) \sin^2 \theta \right\} + \frac{1}{\Phi} \frac{\partial^2 \Phi}{\partial \phi^2} = 0$$

This is clearly separated, so the angular equation reduces to two equations

$$\frac{\sin \theta}{\Theta} \left(\cos \theta \frac{\partial \Theta}{\partial \theta} + \sin \theta \frac{\partial^2 \Theta}{\partial \theta^2} \right) + \ell(\ell + 1) = m^2,$$

$$\frac{\partial^2 \Phi}{\partial \phi^2} = -m^2 \Phi$$

The second is trivial: the ansatz $e^{k\phi}$ yields auxiliary equation $k^2 + m^2 = 0 \Leftrightarrow k = \pm im$. Allowing $m < 0$, this is just $k = im$, so $\Phi(\phi) = e^{im\phi}$.

The first equation is a standard result (I guess),

$$\Theta(\theta) = AP_\ell^m(\cos \theta),$$

where

$$P_\ell^m - (-1)^m (1 - x^2)^{m/2} \left(\frac{d}{dx} \right)^m P_\ell(x),$$

and

$$P_\ell(x) = \frac{1}{2^\ell \ell!} \left(\frac{d}{dx} \right)^\ell (x^2 - 1)^\ell$$

are the Legendre polynomials.