

# 7380 HW 1

Duncan Wilkie

12 September 2021

## 1

Presuming  $\epsilon$  is everywhere nonzero, we may rewrite (2) as  $\frac{i}{\omega\epsilon}(\nabla \times H) = E$ . Taking the curl of both sides and substituting the right side of (1) in for  $\nabla \times E$  yields

$$\nabla \times \left( \frac{i}{\omega\epsilon}(\nabla \times H) \right) = i\omega\mu H$$

Since  $H$  is always perpendicular to the  $(x_1, x_2)$  plane, it only has a component in the  $x_3$  direction, and so its curl is computed as  $\left( \frac{\partial H}{\partial x_2}, -\frac{\partial H}{\partial x_1}, 0 \right)$  using the abusive notation  $H = |H|$ . Moving the constants to the other side, the subsequent curl is

$$\left( -\frac{1}{\epsilon} \frac{\partial^2 H}{\partial x_3 \partial x_1} - \frac{1}{\epsilon^2} \frac{\partial \epsilon}{\partial x_3}, \frac{1}{\epsilon^2} \frac{\partial \epsilon}{\partial x_3} \frac{\partial H}{\partial x_2} - \frac{1}{\epsilon} \frac{\partial^2 H}{\partial x_3 \partial x_2}, \frac{1}{\epsilon^2} \left( \frac{\partial \epsilon}{\partial x_1} \frac{\partial H}{\partial x_1} + \frac{\partial \epsilon}{\partial x_2} \frac{\partial H}{\partial x_2} \right) + \frac{1}{\epsilon} \left( \frac{\partial^2 H}{\partial x_1^2} + \frac{\partial^2 H}{\partial x_2^2} \right) \right)$$

Since neither  $H$  nor  $\epsilon$  depend on  $x_3$ , all the terms containing those partials are zero. This expression then reduces to

$$\left( 0, 0, \frac{1}{\epsilon^2} \left( \frac{\partial \epsilon}{\partial x_1} \frac{\partial H}{\partial x_1} + \frac{\partial \epsilon}{\partial x_2} \frac{\partial H}{\partial x_2} \right) + \frac{1}{\epsilon} \left( \frac{\partial^2 H}{\partial x_1^2} + \frac{\partial^2 H}{\partial x_2^2} \right) \right)$$

On the right hand side, we similarly have a nonzero component only in the  $x_3$  direction since  $H$  is perpendicular to the  $(x_1, x_2)$  plane. Thus, we obtain a single scalar equation. Applying the same abusive notation to the right side, this equation is

$$\frac{1}{\epsilon^2} \left( \frac{\partial \epsilon}{\partial x_1} \frac{\partial H}{\partial x_1} + \frac{\partial \epsilon}{\partial x_2} \frac{\partial H}{\partial x_2} \right) + \frac{1}{\epsilon} \left( \frac{\partial^2 H}{\partial x_1^2} + \frac{\partial^2 H}{\partial x_2^2} \right) = \omega^2 \mu H$$

This is clearly a scalar second-order PDE for  $H$ , as desired.

## 5

Let  $\tau = t - t_0$ . Then  $\tilde{E}(x, t) = E(x, \tau)$ ,  $\tilde{H}(x, t) = H(x, \tau)$  and subsequently

$$\begin{aligned} \nabla \times \tilde{E}(x, t) &= -\frac{\partial}{\partial t}(\mu * \tilde{H}(x, t)) \Leftrightarrow \nabla \times E(x, \tau) = -\left( \frac{\partial}{\partial \tau}(\mu * H(x, \tau)) \right) \frac{\partial \tau}{\partial t} \\ &\Leftrightarrow \nabla \times E(x, \tau) = -\frac{\partial}{\partial \tau}(\mu * H(x, \tau)) \end{aligned}$$

This is identical to the assumption that  $E$  and  $H$  satisfy this equation in  $t$ . An identical argument holds for the second equation of the system.