

# 1 Continuity equation in a periodic electrodynamic medium

We consider a periodic cell  $Z$  of a random heterogeneous material. The conductivity at  $\mathbf{x} \in Z$  is  $\boldsymbol{\sigma}(\mathbf{x})$  (symmetric, positive definite, second-order tensor);  $\mathbf{E}(\mathbf{x})$ ,  $\phi(\mathbf{x})$  and  $\mathbf{j}(\mathbf{x})$  denote the electric field, the electric potential and the volumic current, respectively, at point  $\mathbf{x}$ . The permittivity is noted  $\boldsymbol{\epsilon}(\mathbf{x})$ .

$$Z : \quad \operatorname{div} \left( \mathbf{j} + \frac{d}{dt} \mathbf{D} \right) = 0, \quad (1)$$

$$Z : \quad \mathbf{j} = \boldsymbol{\sigma} \cdot \mathbf{E}, \quad (2)$$

$$Z : \quad \mathbf{D} = \boldsymbol{\epsilon} \cdot \mathbf{E}, \quad (3)$$

$$Z : \quad \mathbf{E} = \bar{\mathbf{E}} + \mathbf{grad} \phi^{\text{per}}, \quad (4)$$

where  $\phi^{\text{per}}$  is a  $Z$ -periodic field,  $\mathbf{j}$  is a  $Z$ -antiperiodic field, and  $\bar{\mathbf{E}}$  is a prescribed constant vector. Eq. (4) ensures that the electric field is curl-free, with  $\bar{\mathbf{E}} = \langle \mathbf{E} \rangle$ , where angle brackets denote volume averages over the cell  $Z$ .

## 2 The time dependence

We consider the following electric field as a data:

$$\bar{\mathbf{E}} = t\bar{\mathbf{E}}_1 + \frac{t^2}{2}\bar{\mathbf{E}}_2 + \frac{t^3}{6}\bar{\mathbf{E}}_3 \quad (5)$$

and we assume the following development for  $\phi^{\text{per}}$

$$\phi^{\text{per}} = \sum_{i=1}^{\infty} \phi_i \frac{t^i}{i!} \quad (6)$$

and we note  $\mathbf{grad} \phi^{\text{per}} = \mathbf{e}$ , and also  $\mathbf{grad} \phi_i = \mathbf{e}_i$  (for all  $i$ ). Hence, we have the following problems:

$$\begin{aligned} \operatorname{div} (\boldsymbol{\epsilon} \cdot (\bar{\mathbf{E}}_1 + \mathbf{e}_1)) &= 0 \\ \operatorname{div} (\boldsymbol{\sigma} \cdot (\bar{\mathbf{E}}_1 + \mathbf{e}_1) + \boldsymbol{\epsilon} \cdot (\bar{\mathbf{E}}_2 + \mathbf{e}_2)) &= 0 \\ \operatorname{div} (\boldsymbol{\sigma} \cdot (\bar{\mathbf{E}}_2 + \mathbf{e}_2) + \boldsymbol{\epsilon} \cdot (\bar{\mathbf{E}}_3 + \mathbf{e}_3)) &= 0 \\ \operatorname{div} (\boldsymbol{\sigma} \cdot (\bar{\mathbf{E}}_3 + \mathbf{e}_3) + \boldsymbol{\epsilon} \cdot \mathbf{e}_4) &= 0 \\ \operatorname{div} (\boldsymbol{\sigma} \cdot \mathbf{e}_4 + \boldsymbol{\epsilon} \cdot \mathbf{e}_5) &= 0 \\ \operatorname{div} (\boldsymbol{\sigma} \cdot \mathbf{e}_5 + \boldsymbol{\epsilon} \cdot \mathbf{e}_6) &= 0 \\ \dots \end{aligned} \quad (7)$$

These problems can be resolved in an iterative way.

### 3 Resolution of the problems

#### 3.1 first problem

We consider a biphasic cell with a matrix whose conductivity is isotropic  $\sigma_0$  and another phase whose conductivity is isotropic  $\sigma_1$ . The permittivity is also isotropic,  $\epsilon_0 = \tau_0 \sigma_0$  for medium 0 and  $\epsilon_1 = \tau_1 \sigma_1$  for medium 1. Hence, the first problem to solve is

$$\operatorname{div} (\epsilon(\bar{\mathbf{E}}_1 + \mathbf{e}_1)) = 0 \quad (8)$$

and its solution can be expressed as

$$\bar{\mathbf{E}}_1 + \mathbf{e}_1 = \bar{\mathbf{E}}_1 - \mathbf{A}(\mathbf{x}) \cdot \mathbf{G}^{-1} \cdot \mathbf{U}^1 \quad (9)$$

where the subblocks of  $\mathbf{G}$  are:

$$\mathbf{G}_{ik} = \int_Z \epsilon {}^t \mathbf{A}^i \cdot \mathbf{A}^k \quad (10)$$

and

$$\mathbf{U}_i^1 = \int_Z \epsilon {}^t \mathbf{A}^i \cdot \bar{\mathbf{E}}_1 \quad \text{and} \quad \mathbf{A}_{ij}^k(\mathbf{x}) = \Gamma_i^{0,\text{per}} \left[ (-\chi \epsilon_0 \Gamma_0^{\text{per}})^{k-1} (\chi \mathbf{e}_j) \right] (\mathbf{x}) \quad (11)$$

#### 3.2 second problem

The second problem to solve is

$$\operatorname{div} (\epsilon(\bar{\mathbf{E}}_2 + \mathbf{e}_2) + \mathbf{P}_1) = 0 \quad (12)$$

with  $\mathbf{P}_1 = \sigma(\bar{\mathbf{E}}_1 + \mathbf{e}_1) = \sigma(\bar{\mathbf{E}}_1 - \mathbf{A}(\mathbf{x}) \cdot \mathbf{G}^{-1} \cdot \mathbf{U}^1)$ . Its solution can be expressed as

$$\bar{\mathbf{E}}_2 + \mathbf{e}_2 = \bar{\mathbf{E}}_2 - \mathbf{A}(\mathbf{x}) \cdot \mathbf{G}^{-1} \cdot \mathbf{U}^2 - \Gamma_0^{\text{per}}(\mathbf{P}_1) \quad (13)$$

where

$$\mathbf{U}_i^2 = \int_Z {}^t \mathbf{A}^i \cdot [\epsilon(\bar{\mathbf{E}}_2 - \Gamma_0^{\text{per}}(\mathbf{P}_1)) + \mathbf{P}_1] \quad (14)$$

#### 3.3 third problem

The third problem to solve is

$$\operatorname{div} (\epsilon(\bar{\mathbf{E}}_3 + \mathbf{e}_3) + \mathbf{P}_2) = 0 \quad (15)$$

with  $\mathbf{P}_2 = \sigma(\bar{\mathbf{E}}_2 + \mathbf{e}_2) = \sigma(\bar{\mathbf{E}}_2 - \mathbf{A}(\mathbf{x}) \cdot \mathbf{G}^{-1} \cdot \mathbf{U}^2 - \Gamma_0^{\text{per}}(\mathbf{P}_1))$ . Its solution can be expressed as

$$\bar{\mathbf{E}}_3 + \mathbf{e}_3 = \bar{\mathbf{E}}_3 - \mathbf{A}(\mathbf{x}) \cdot \mathbf{G}^{-1} \cdot \mathbf{U}^3 - \Gamma_0^{\text{per}}(\mathbf{P}_2) \quad (16)$$

where

$$\mathbf{U}_i^3 = \int_Z {}^t \mathbf{A}^i \cdot [\epsilon(\bar{\mathbf{E}}_3 - \Gamma_0^{\text{per}}(\mathbf{P}_2)) + \mathbf{P}_2] \quad (17)$$

### 3.4 fourth problem

The fourth problem to solve is

$$\operatorname{div}(\epsilon \mathbf{e}_4 + \mathbf{P}_3) = 0 \quad (18)$$

with  $\mathbf{P}_3 = \sigma(\bar{\mathbf{E}}_3 + \mathbf{e}_3) = \sigma(\bar{\mathbf{E}}_3 - \mathbf{A}(\mathbf{x}) \cdot \mathbf{G}^{-1} \cdot \mathbf{U}^3 - \mathbf{\Gamma}_0^{\text{per}}(\mathbf{P}_2))$ . Its solution can be expressed as

$$\mathbf{e}_4 = -\mathbf{A}(\mathbf{x}) \cdot \mathbf{G}^{-1} \cdot \mathbf{U}^4 - \mathbf{\Gamma}_0^{\text{per}}(\mathbf{P}_3) \quad (19)$$

where

$$\mathbf{U}_i^4 = \int_Z {}^t \mathbf{A}^i \cdot [-\epsilon \mathbf{\Gamma}_0^{\text{per}}(\mathbf{P}_3) + \mathbf{P}_3] \quad (20)$$

### 3.5 $k^{\text{th}}$ problem with $k \geq 5$

The  $k^{\text{th}}$  problem to solve is

$$\operatorname{div}(\epsilon \mathbf{e}_k + \mathbf{P}_{k-1}) = 0 \quad (21)$$

with  $\mathbf{P}_{k-1} = \sigma \mathbf{e}_{k-1} = \sigma(-\mathbf{A}(\mathbf{x}) \cdot \mathbf{G}^{-1} \cdot \mathbf{U}^{k-1} - \mathbf{\Gamma}_0^{\text{per}}(\mathbf{P}_{k-2}))$ . Its solution can be expressed as

$$\mathbf{e}_k = -\mathbf{A}(\mathbf{x}) \cdot \mathbf{G}^{-1} \cdot \mathbf{U}^k - \mathbf{\Gamma}_0^{\text{per}}(\mathbf{P}_{k-1}) \quad (22)$$

where

$$\mathbf{U}_i^k = \int_Z {}^t \mathbf{A}^i \cdot [-\epsilon \mathbf{\Gamma}_0^{\text{per}}(\mathbf{P}_{k-1}) + \mathbf{P}_{k-1}] \quad (23)$$

## 4 Numerical application

We choose  $\sigma_0 = 1$ ,  $\tau_0 = 1e - 6$ , so that  $\epsilon_0 = 1e - 6$ ; and  $\sigma_1 = 1e6$ ,  $\tau_1 = 1e - 6$ , so that  $\epsilon_1 = 1$ . For the remote field, we choose  $\bar{\mathbf{E}}_1 = 4e5$ ,  $\bar{\mathbf{E}}_2 = -4e5$ ,  $\bar{\mathbf{E}}_3 = 0$  and hence the time of interest is between  $1e - 6$  and  $1e - 5$ .