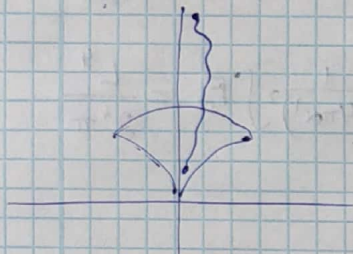


# Asael - Tomasiello: Holographic duals of 3d S-fold CFTs

Dualities in string theory  $\rightarrow$  solutions where transition functions are not coord changes alone.

Ex: F-theory, T-folds

$SL(2, \mathbb{Z})$  in IIB sugra:



Two possibilities for monodromies:

- Around non-contractible path: take path long enough so that fields vary slowly and 2-derivative action suffice
- Around contractible paths: the "long-wavelength" approx will break down near singularity.

Test the non-geometric solutions using holography.

Here:  $AdS_4 \times K_6$  IIB string w/ monodromies in  $K_6$  in  $SL(2, \mathbb{Z})$  (non-contractible) and tested using holography

$\equiv$  S-folds

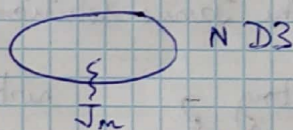
2 classes of S-folds are considered:

①  $AdS_4 \times S^5 \times S^1$  with monodromy  $J$  around  $S^1$ , w/  $\text{Tr} J > 2$ .

Preserves  $\mathcal{Osp}(4|4)$  and has 3d  $cf=4$  dual SCFT.

$\equiv$  Janus S-fold

Ex:  $(N)_m \dashv T[U(N)]$



UV description has only  $cf=3$  but enhanced to  $cf=4$  at low energies (3d)

Related to 4d  $cf=4$   $U(N)$  on a circle with monodromy  $J$  around the circle

Compute the  $S^3$  partition function:

$$F = -\log Z_{S^3} \sim f(J) N^2 \text{ in large } N \text{ limit.}$$

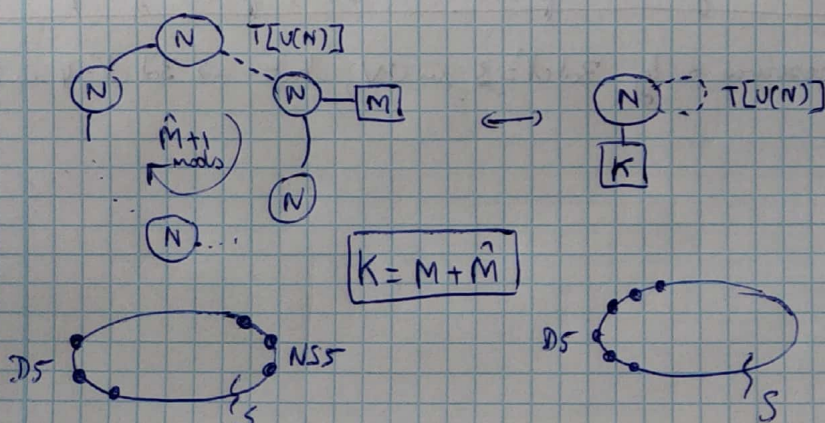
② Quotient of geometrical solution by  $(\text{geom} \times SL(2, \mathbb{Z}))$  involution  $S^5 \times S^1$  with NS5/D5 wrapping  $S^2$ s.  $\equiv$  "S-flip"

Monodromy  $S = \begin{pmatrix} 0 & -1 \\ 1 & 0 \end{pmatrix}$  around non-contractible  $S^1$ .

Preserves  $\mathcal{Osp}(3|4) \rightarrow cf=3$  3d dual SCFT

Impose  $N \geq K$ . Then

$$F \sim \frac{1}{2} N^2 \ln N$$

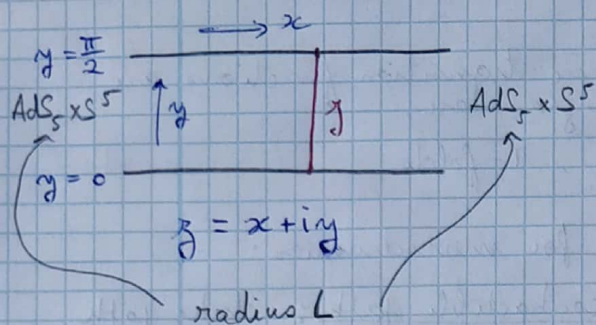




# Janus S-fold

D'Hoker - Estes - Gutperle (2007):  $AdS_4 \times S^2 \times S^2 \times \Sigma$

3d  $\mathcal{N}=4$  SCFT



Riemann Surface  
Infinite strip

Solution parametrized by two functions  $h_1, h_2$ . Then take particular form:

$$\begin{cases} h_1 = \frac{c}{2i}(e^{\tilde{z}} - e^{-\tilde{z}}) \\ h_2 = \frac{c}{2}(e^{-\tilde{z}} + e^{-\bar{\tilde{z}}}) \end{cases} \quad N = \frac{c^2}{8\pi}$$

Geometry has 5-cycle  $S^2 \times S^2 \times J \simeq S^5$  with flux  $N = \frac{1}{(4\pi\alpha')^2} \int F_5 = \frac{L^4}{2^6 \pi}$

Apply transfo  $M \in SL(2, \mathbb{R})$  to that solution given by

$$M = \begin{pmatrix} \frac{1}{1+e^{-T}} & \frac{1}{1-e^{-T}} \\ \frac{-1}{1+e^{-T}} & \frac{1}{1-e^{-T}} \end{pmatrix}$$

$$T = \log \left( \frac{n + \sqrt{n^2 - 4}}{2} \right) \quad n > 2$$

Then  $J_n \tau'(x+T) = \tau'(x)$  with  $J_n = \begin{pmatrix} n & 1 \\ -1 & 0 \end{pmatrix}$  and  $\tau' = M\tau$

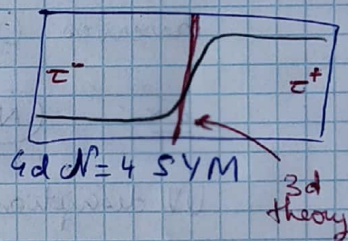
$\Rightarrow$  M-transformed solution invariant under  $\mathcal{G}$  = combination of translation by T along x and  $SL(2, \mathbb{Z})$  transfo  $J_n$ .  $\equiv J_n$ -fold

take  $\mathbb{Z}$  quotient  $\rightarrow AdS_4 \times S^5 \times S^1$

CFT dual . The Janus solution is dual to Janus CFT

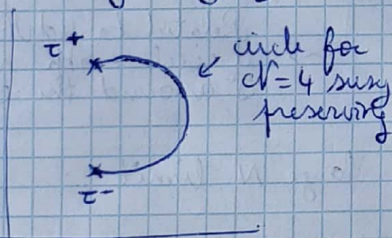
Consider trajectory  $\gamma$

with  $\tau^+, \tau^-$  related by  $J_n$  and take quotient:



4d  $\mathcal{N}=4$  SYM

3d theory



circle for  $\mathcal{N}=4$  sym preserving



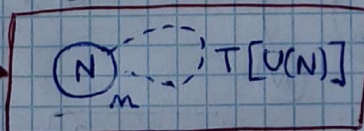
This can be described as 4d  $\mathcal{N}=4$  on circle coupled to 3d theory with quasi-Lagrangian description in UV.

Gaiotto-Witten: this is  $T[U(N)]$  with non-abelian Chern-Simons at level n for one of the two  $U(N)$  flavor groups

because the two  $U(N)$  global sym are not both present in the UV Lagrangian of  $T[U(N)]$ .

Preserves only 3d  $\mathcal{N}=3$  in UV, but  $\rightarrow$  3d  $\mathcal{N}=4$  in IR

In the IR this flows to 3d theory  $\rightarrow$





## Review of $T[U(N)]$ [Gaiotto - Witten, 2008]

$T[SU(N)]$  is the IR limit of  $\textcircled{1} - \textcircled{2} - \dots - \textcircled{N-1} - \boxed{N}$

- \* It is self-mirror (hence Higgs and Coulomb branches are isomorphic)
- \* More generally  $T[G]$  has mirror  $T[G^\vee]$ .
- \*  $T[G]$  has global sym  $G \times G^\vee$  when  $\left. \begin{array}{l} G \text{ acts on Higgs branch} \\ G^\vee \text{ acts on Coulomb branch} \end{array} \right\}$
- \*  $T[G]$  can be constructed as  $\frac{1}{2}$  BPS domain wall with  $G$  gauge theory on one side ( $y < 0$ ) and  $G^\vee$  gauge theory on the other. There are many such walls, but there is a minimal one called the Janus.  
Then you have  $\mathcal{N}=4$  SYM on both sides with gauge coupling varying from very small to very large

## Test of holography

- Regularized on-shell action of  $J_n$ -fold Janus solution [Asrat, Bachas, Estes, Gomis]

$$S_{\text{IB}} = \frac{-1}{(2\pi)^3} \int_0^T dx \int_0^{\pi/2} dy \, h_1 h_2 \partial_{\bar{z}} \partial_{\bar{z}} (h_1 h_2) = \frac{L^8 T}{2^{13} \pi^2} = \frac{1}{2} N^2 \log \left( \frac{n + \sqrt{n^2 - 4}}{2} \right)$$

This sugra result is valid in the regime (large  $N$ , finite  $T$ ).

- Large  $N$  free energy  $F = -\log |\mathbb{Z}_{S^3}|$  of the  $J_n$  theory computed as a matrix model (via localization):

$$F = \frac{N^2}{2} T + \sum_{j=1}^N \log(1 - e^{-jT}) \sim \frac{1}{2} N^2 \log \left( \frac{n + \sqrt{n^2 - 4}}{2} \right)$$

$\Rightarrow$  Results match in the large  $N$  limit

Rg: One can construct other  $J$ -fold theories, using other matrices  $J$ .