Why Symplectic Singularities?

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Based on long time collaboration with P. Argyres, S. Cabrera, A. Collinucci, S. Giacomelli, J. Grimminger, A. Hanany, R. Kalveks, M. Martone, S. Schäfer-Nameki, M. Sperling, G. Zafrir, Z. Zhong.

[AB, Grimminger, 2209.15016]

Why Symplectic Singularities?

- Symplectic singularities are a basic kind of geometric spaces.
- They feature in moduli space of vacua of supersymmetric theories.
- They provide a geometric perspective to understand
 - the phases of a given theory
 - the landscape of all theories
 - and the relations between them.

Why Symplectic Singularities?

- Introduction and Examples
- 2 Definition of Conical Symplectic Singularities
- Moduli Space of Vacua
 - SCFTs and Moduli Space of Vacua
 - Classification of SCFTs
 - Phases of SCFTs
- CSS stratification and fibrations
 - General Strategy
 - Phase diagrams from Magnetic Quivers
 - The quest for the elementary building blocks
- Conclusion

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Example 1 : Gravitational Instantons

Start from Asymptotically Locally Euclidean HyperKähler 4-manifolds. For instance Eguchi-Hanson space: $\mathcal{T}^{\star}\mathbb{P}^{1}$ with metric

$$\mathrm{d}s^2 = \frac{\mathrm{d}r^2}{1 - \frac{a}{r^4}} + \frac{r^2}{4} \left(1 - \frac{a}{r^4} \right) \sigma_3^2 + \frac{r^2}{4} (\sigma_1^2 + \sigma_2^2)$$

with

$$\sigma_1 = \sin \psi d\theta - \cos \psi \sin \theta d\phi
\sigma_2 = \cos \psi d\theta + \sin \psi \sin \theta d\phi
\sigma_3 = d\psi + \cos \theta d\phi$$

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When $a \rightarrow 0$:

$$T^{\star}\mathbb{P}^1 o \mathbb{R}^4/\mathbb{Z}_2$$
 .

This is the A_1 singularity.

Example 2 : Moduli Space of Gauge Instantons

Let $\mathcal{M}_{k,N}$ be the moduli space of k SU(N) framed instantons on \mathbb{R}^4 . It is a singular HyperKähler manifold,

$$\dim_{\mathbb{R}} \mathcal{M}_{k,N} = 4kN$$
.

Example 2 : Moduli Space of Gauge Instantons

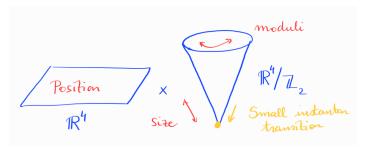
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For k = 1, N = 2:

$$\mathcal{M}_{1,2} = \mathbb{R}^4 \times \mathbb{R}^4/\mathbb{Z}_2 = \mathbb{H} \times A_1 \,.$$

Singularity: small instanton transition.



Supersymmetric QFT in spacetime dimension $3 \le d \le 6$ with 8 supercharges have *Higgs branches*.

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Exemple: SQED with gauge group U(1) and $N_f = 2$:

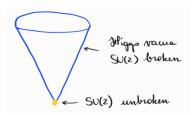
$$\mathbb{C}[\mathcal{H}] = \left(rac{\mathbb{C}[Q_{2 imes 1}, ilde{Q}_{1 imes 2}]}{(ilde{Q}Q = 0\,,\, |Q|^2 - | ilde{Q}|^2 = 0)}
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Supersymmetric QFT in spacetime dimension $3 \le d \le 6$ with 8 supercharges have *Higgs branches*.

Exemple: SQED with gauge group U(1) and $N_f = 2$:

$$\mathbb{C}[\mathcal{H}] = \left(\frac{\mathbb{C}[Q_{2\times 1}, \tilde{Q}_{1\times 2}]}{(\tilde{Q}Q = 0, |Q|^2 - |\tilde{Q}|^2 = 0)}\right)^{U(1)} \qquad \mathcal{H} = \mathbb{C}^4 / / U(1) = \mathbb{C}^2 / \mathbb{Z}_2.$$

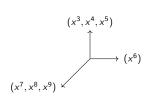
$$\mathcal{H} = \{M_{2\times 2}|M^2 = 0 \,,\, \mathrm{tr}(M) = 0\} \,,\,\, \mathrm{with}\,\, M = Q\tilde{Q} \,.$$

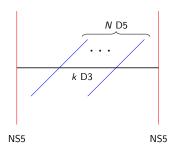


Resolution / Deformation by Fayet-Iliopoulos terms.

Brane system for 3d $\mathcal{N}=4$ system (Hanany-Witten):

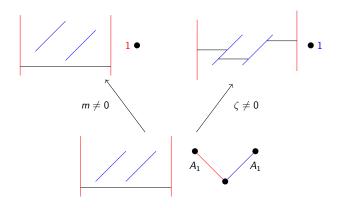
	x ⁰	x^1	x ²	<i>x</i> ³	<i>x</i> ⁴	<i>x</i> ⁵	x ⁶	x ⁷	x ⁸	x ⁹
NS5	х	х	х	х	х	х				
D3							Х			
D5	х	х	х					х	х	х







Full moduli space of vacua for 3d $\mathcal{N}=4$ SQED with $N_f=2$:



Nilpotent cone:

$$\mathcal{N}(\mathfrak{sl}_N)=\{M\in \mathrm{Mat}(N,\mathbb{C})\,|\,M^N=0\}$$

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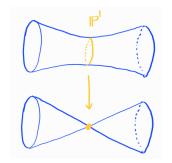
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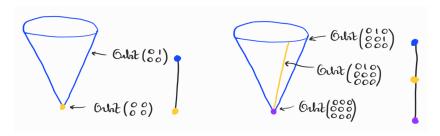
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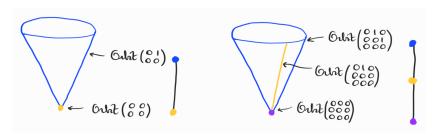
Springer resolution: $T^*\operatorname{Flag}_N \to \mathcal{N}(\mathfrak{sl}_N)$



Nilpotent orbits and stratification of the nilpotent cone:



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Minimal nilpotent orbit closure $\overline{\mathcal{O}_{\min}(\mathfrak{g})}$:

$$\dim_{\mathbb{R}} \overline{\mathcal{O}_{\min}(\mathfrak{g})} = 4(\mathit{h}^{\vee}(\mathfrak{g}) - 1)$$

$$\mathcal{M}_{1,G} = \mathbb{R}^4 imes \overline{\mathcal{O}_{\min}(\mathfrak{g})}$$

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Common features / Definition

- Presence of singularities;
- Complex symplectic structure on smooth locus;
- Singular limit of a HyperKähler manifolds;
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A Poisson variety $(X, \{,\})$ is a Conical Symplectic Singularity if:

- **1** The Poisson structure is non-degenerate on X^{reg} : there is a symplectic form ω ;
- **③** For some / any resolution of singularities $\pi: Y \to X$, $\pi^\star \omega$ extends to a regular 2-form on Y;
- **3** There is a \mathbb{C}^* action that contracts X to a point under which ω has positive integer weight.

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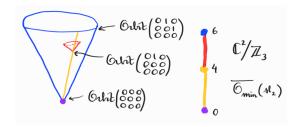
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A Conical Symplectic Singularity is isolated (ICSS) if the singular locus is a point.

Stratification and Transverse Slices

ICSS are the building blocks of all CSS:

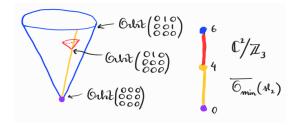
- ullet A CSS has a stratification into symplectic leaves o Hasse diagram
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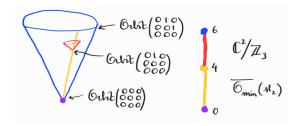


Open question: is there a classification of ICSS?¹

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Open question: is there a classification of ICSS?¹

Only known example is $\overline{\mathcal{O}_{\min}(\mathfrak{g})}$ [Beauville, 2000]

But see later...



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Dimension	Susy	Bosonic subalgebra		SCA
d = 6	$\mathcal{N}=(1,0)$	$\mathfrak{so}(6,2)\oplus\mathfrak{su}(2)_H$	\subset	$\mathfrak{osp}(6,2 1)$
d = 5	$\mathcal{N}=1$	$\mathfrak{so}(5,2)\oplus\mathfrak{su}(2)_H$	\subset	f(4)
d = 4	$\mathcal{N}=2$	$\mathfrak{so}(4,2) \oplus \mathfrak{su}(2)_H \oplus \mathfrak{u}(1)_C$	\subset	$\mathfrak{su}(2,2 2)$
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SCFTs are

- "Rare" in 6d / 5d isolated, rely on exceptional isomorphisms, non Lagrangian.
- More common in 4d (some Lagrangian; existence of conformal manifolds).
 Classification?
- Very large number in 3d.

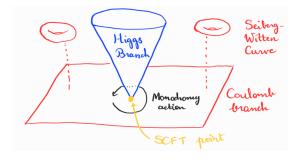
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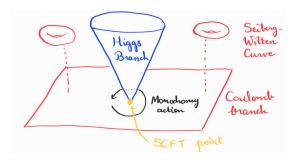
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Existence of Moduli space of vacua, always contains Higgs branch.

Example: SU(2) gauge theory with $N_f = 4$:



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Lesson 1

The Moduli Space of Vacua can be used to "solve" certain sectors of the theory, and go beyond perturbation theory.

Idea: classify possible MSV (bottom-up approach).

Example: "simplest" 4d $\mathcal{N}=2$ SCFTs moduli spaces

• Coulomb branch: dimension 1.

$$II^{\star}$$
, III^{\star} , IV^{\star} , I_0^{\star} , IV , III , II .

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Coulomb branch : dimension 1.

$$II^{\star}\,,\quad III^{\star}\,,\quad IV^{\star}\,,\quad I_0^{\star}\,,\quad IV\,,\quad III\,,\quad II\,.$$

 Higgs branch: Isolated Conical Symplectic Singularities that appear as associated varieties to VOAs:

$$\mathfrak{e}_8$$
, \mathfrak{e}_7 , \mathfrak{e}_6 , \mathfrak{f}_4 , \mathfrak{d}_4 , \mathfrak{g}_2 , \mathfrak{a}_2 , \mathfrak{a}_1 .

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Coulomb branch : dimension 1.

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One needs to take into account deformations! Example: SU(2) $N_f=4$ vs $\mathcal{N}=4$.

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To match with CB classification, needs to also consider non-isolated singularities!

Classification of 4d rank 1 $\mathcal{N}=2$ SCFT Coulomb branch geometries:

Flavor	CB geometry and deformation	$\Delta(u)$
E ₈	$II^* ightarrow \{I_1^{10}\}$	6
E_7	$III^* ightarrow \{\hat{m{I}}_1^9\}$	4
E_6	$IV^* o \{I_1^8\}$	3
D_4	$I_0^* \rightarrow \{I_1^6\}$	2
A_2	$IV \rightarrow \{I_1^4\}$	3/2
A_1	$III \rightarrow \{I_3^{\clip{3}}\}$	4/3
Ø	$H o \{I_1^{\frac{1}{3}}\}$	6/5
C ₅	$II^* \to \{I_1^6, I_4\}$	6
C_3A_1	$III^* \to \{I_1^5, I_4\}$	4
C_2U_1	$IV^* ightarrow \{I_1^4, I_4\}$	3
C_1	$I_0^* \to \{I_1^2, I_4\}$	2
$A_3 \rtimes \mathbb{Z}_2$	$II^* \to \{I_1^3, I_1^*\}$	6
$A_1U_1 \rtimes \mathbb{Z}_2$	$III^* o \{ \vec{I_1^2}, \vec{I_1^*} \}$	4
U_1	$IV^* o \{\hat{I_1^1}, \hat{I_1^*}\}$	3
$A_2 \rtimes \mathbb{Z}_2$	$II^* \to \{I_1^2, IV_{Q=1}^*\}$	6
$U_1 \rtimes \mathbb{Z}_2$	$III^* \rightarrow \{I_1, IV_{Q=1}^*\}$	4
Ø	$IV_{O=1}^*$	3
C_1	$I_0^* o \overline{\{I_2^3\}}$	2

[Argyres, Lotito, Lü, Martone 18]

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Lesson 2

"Conjecture": SCFTs with 8 supercharges in dimensions $3 \le d \le 6$ are entirely characterized by their full moduli space of vacua.

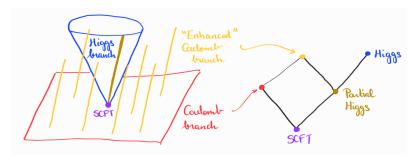
Example: C_5 theory.

Classification of SCFTs

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Phases of SCFTs

Lesson 3

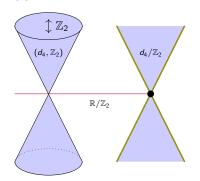
Phases of the theory \equiv Smooth loci of the MSV Phase transitions \equiv Singularities

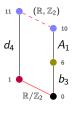
Phases of SCFTs

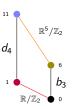
Lesson 3

Phases of the theory ≡ Smooth loci of the MSV
Phase transitions ≡ Singularities

Example: d=6, $\mathcal{N}=(1,0)$ SCFT from SU(2) with $N_f=4$. At the SCFT point, Spin(7) symmetry! [AB, Grimminger, 2022]







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General Strategy

Fibration of the full moduli space:

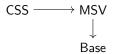
$$\begin{array}{c} \mathsf{CSS} & \longrightarrow \mathsf{MSV} \\ & \downarrow \\ \mathsf{Base} \end{array}$$

- For each fiber, compute singular stratification.
- The phase diagram is obtained by deleting non-singular transitions.

[AB, Grimminger, 22]

General Strategy

Fibration of the full moduli space:

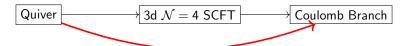


- For each fiber, compute singular stratification. How to do this?
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[AB, Grimminger, 22]

Magnetic Quivers

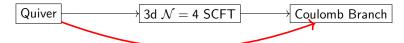
The Coulomb branch of a 3d $\mathcal{N}=4$ SCFT is a CSS due to $\mathfrak{su}(2)_{\mathcal{C}}$.



[Cremonesi, Hanany, Zaffaroni 14] , [Bullimore, Dimofte, Gaiotto 15] , [Braverman, Finkelberg, Nakajima 15]

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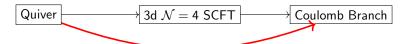
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Let X be a CSS. We say that the (generalized) quiver Q is a **magnetic quiver** for X if

$$\mathcal{C}^{\mathrm{3d}\ \mathcal{N}=4}(Q)=X\,.$$

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Lesson 4

A magnetic quiver is a simple combinatorial object that encodes CSS and which gives access to its singular stratification.

Rank-1 4d $\mathcal{N}=2$ magnetic quivers

Flavor	$dim_{\mathbb{H}}(\mathit{HB})$	Magnetic Quive
E ₈	29	
E_7	17	Affine
E_6	11	
D_4	5	Dynkin
A_2	2	
A_1	1	Diagrams
Ø	0	

[AB, Grimminger, Hanany, Sperling, Zafrir, Zhong 20]

Flavor	$dim_{\mathbb{H}}(\mathit{HB})$	Magnetic Quiver
C ₅	16	0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0
C_3A_1	8	
C_2U_1	4	
C_1	1	
$A_3 \rtimes \mathbb{Z}_2$	9	0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0
$A_1U_1 \rtimes \mathbb{Z}_2$	3	
U_1	1	
$A_2 \rtimes \mathbb{Z}_2$	5	O ● 1 2 3
$U_1 \rtimes \mathbb{Z}_2$	1	
Ø	0	
<i>C</i> ₁	1	

Quiver Subtraction

From a magnetic quiver, the stratification is obtained via a quiver algorithm.

Physically corresponds to Higgsing graph between SCFTs.

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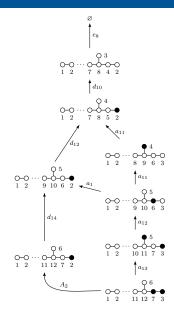
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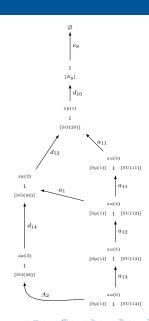
Different algorithms:

- Quiver subtraction [AB, Cabrera, Grimminger, Hanany, Sperling, Zajac, Zhong,
 2019]: needs a list of ICSS as an input / or alternatively a brane realization.
- Quiver fission and decay [AB, Sperling, Zhong, 2023]: no input needed. ICSS correspond to stable quivers (which can't fission or decay).

Example : 6d $\mathcal{N}=(1,0)$ theory from SU(6) gauge theory with $N_f=14$ and one antisymmetric hyper.

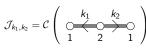
Quiver Subtraction





New elementary slices

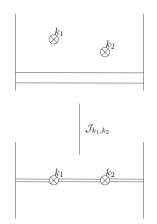
Using brane systems / quiver fission and decay, one can construct **new Isolated Conical Symplectic Singularities**!



	k_2	≥ 4	3	2	1
	≥ 4	$\mathcal{Y}(k_1) \qquad \mathcal{Y}(k_2)$ $A_1 \qquad A_1$ $\operatorname{su}(2)^2$	$\mathcal{Y}(k_1)$ A_1 $\operatorname{su}(2)^2$	$\mathcal{Y}(k_1)$ A_1 $\operatorname{su}(2)^2$	$\mathcal{Y}(k_1)$ A_1 $\mathfrak{su}(2)^2$
	3		\mathcal{J}_{33} su(2) ²	\mathcal{J}_{32} su(2) ²	g ₂ G ₂
	2			d ₃ so(6)	⊞ ³ ● sp(3)
	1				bad

See also [Bellamy, Bonafé, Fu, Levy, Juteau, Sommers 2022]

New elementary slices



Example:

$$\operatorname{HS}(\mathcal{J}_{33}) = \frac{1 + 3t^2 + 18t^4 + 14t^6 + 18t^8 + 3t^{10} + t^{12}}{(1 - t^2)^3(1 - t^4)^3}.$$

Outline

- Introduction and Examples
- 2 Definition of Conical Symplectic Singularities
- Moduli Space of Vacua
 - SCFTs and Moduli Space of Vacua
 - Classification of SCFTs
 - Phases of SCFTs
- CSS stratification and fibrations
 - General Strategy
 - Phase diagrams from Magnetic Quivers
 - The quest for the elementary building blocks
- 5 Conclusion



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Not covered today:

- Magnetic quivers can be seen as *intersection graphs* in brane systems.
- Links with (generalized) toric polygons [AB, Collinucci, Schäfer-Nameki, 22]
- RG flows between SCFTs (e.g. [AB, Giacomelli, Grimminger, 22])
- Monopole formula and Hilbert series [Cremonesi, Hanany, Zaffaroni, 14]



Future Directions

- How to construct magnetic quivers in general (e.g. from pure geometry) ?
- List all ICSS from stable quivers.
- Monopole Formula for arbitrary transverse slices.
- Boostrap the geometry to non-isolated CSS.
- Understand the corresponding dynamics of QFTs.
- Beyond classical geometry: scheme and stack aspects.

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Thank you for your attention!