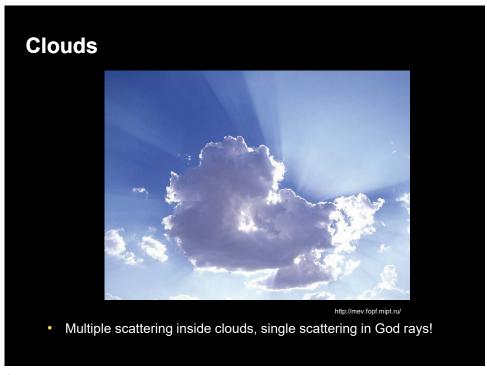
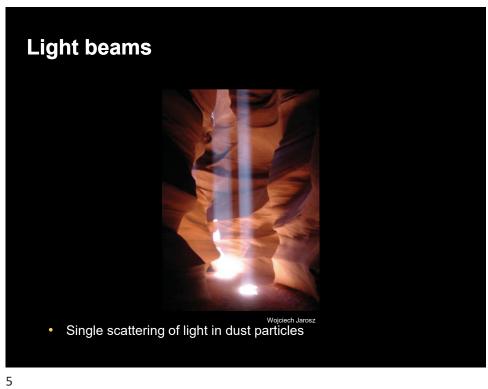


# **Volumetric Scattering**

- Global light transport in participating media
- Single scattering
  - optically thin media
- Multiple scattering
  - optically dense media
- Subsurface scattering









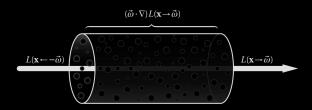
# **Volumetric Scattering**



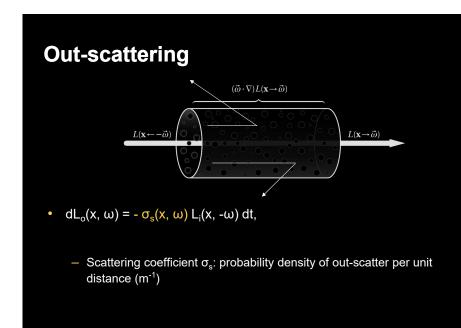
- Participating media
  - particles distributed in a volume of 3D space
- Absorption
- Emission
- Scattering

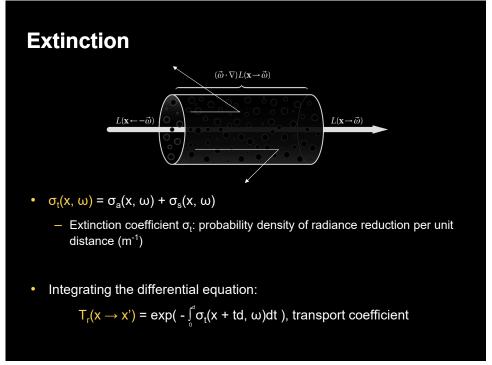
7

# **Absorption**

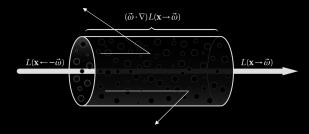


- $L_o(x, \omega) L_i(x, \omega) = dL_o(x, \omega) = -\sigma_a(x, \omega) L_i(x, -\omega) dt$ 
  - Absorption cross-section  $\sigma_a \!\!:\! probability$  density of absorption per unit distance  $(m^{\text{-}1})$





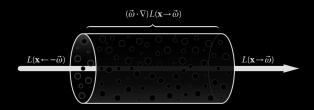
# Homogeneous media



- $\sigma_t(x, \omega)$  constant in homogeneous media
- $T_r(x \rightarrow x') = exp(-\sigma_t d)$ , Beer's Law

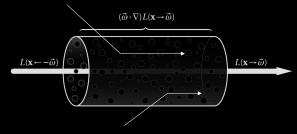
11

# **Emission**



- $dL_o(x, \omega) = Q(x, \omega) dt$ ,
  - $\;\; Q(x,\, \omega)\!\!:$  emitted radiance added per unit distance

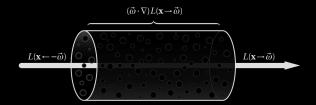
# **In-scattering**



- In-scattering described by phase function  $p(\omega \to \omega')$ 
  - Volumetric analog of BRDF
- $1 \int_{\Omega} p(\omega \to \omega') d\omega' = 1$

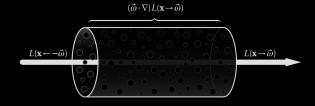
13

# **Added radiance**



- Total added radiance: emission + in-scattering  $dL_o(x,\,\omega) = S(x,\,\omega)\;dt,$
- $\bullet \quad S(x,\,\omega) = Q(x,\,\omega) + \sigma_s(x,\,\omega) \textstyle\int_\Omega p(x,\,-\,\omega' \to \omega) \; L_i(x,\,\omega') \; d\omega'.$

# **Radiative transfer**



• Volume rendering equation :

$$dL_o(x,\,\omega) = -\,\sigma_t(x,\,\omega)\;L_i(x,\,-\omega)\;dt + S(x,\,\omega)\;dt.$$

15

### **Phase Function**

• Isotropic:  $p(\omega \rightarrow \omega') = 1/4\pi$ 

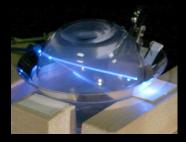


• Anisotropic: Henyey-Greenstein

$$p_{HG}(\cos\theta) = \frac{1 - g^2}{4\pi (1 + g^2 - 2g\cos\theta)^{3/2}}$$

- 0.5 0.5 1 1.5
- $\theta$  is the angle between  $\omega$  &  $\omega'$
- $g \rightarrow [-1, 1],$
- g = + 1 forward scatting, -1 backward scattering, 0 isotropic scattering.

# Measuring scattering in smoke Hawkins et al. 05

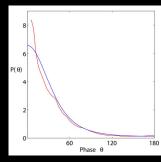


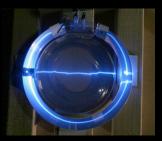


- · Phase function measurement
  - Conical mirror at 45°, laser beam through volume of smoke trapped in glass bowl
  - 1D measurement of the phase function via conical mirror to fit the g parameter of the H-G phase function
  - highly forward scattering g

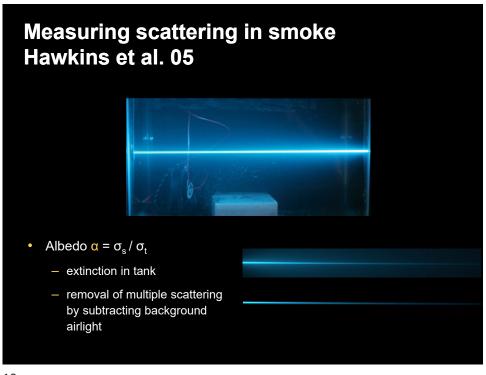
17

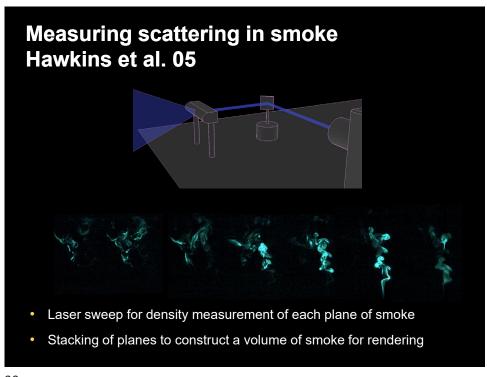
# Measuring scattering in smoke Hawkins et al. 05



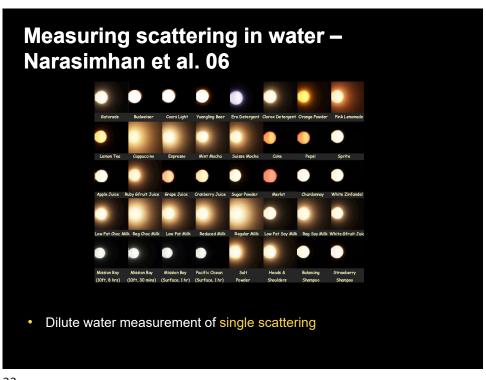


- Phase function measurement
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  - highly forward scattering g



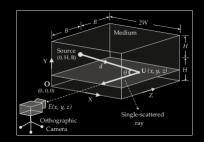






# Measuring scattering in water – Narasimhan et al. 06





• Dilute water measurement of single scattering

23

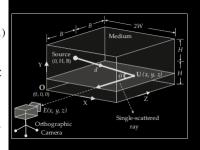
# Measuring scattering in water – Narasimhan et al. 06

$$E(x,y,z) = \frac{I_0}{d^2} \cdot e^{-\sigma d} \cdot \beta P(g,\pi-\theta) \cdot e^{-\sigma z}.$$

$$d = \sqrt{x^2 + (y-H)^2 + (z-B)^2} \cdot \cos \theta = (z-B)/d (1)$$

$$E(x,y) = \int_0^{2B} E(x,y,z) dz$$

$$= \beta \int_0^{2B} \frac{I_0 e^{-\sigma(z+\sqrt{x^2+(y-H)^2+(z-B)^2})}}{x^2 + (y-H)^2 + (z-B)^2} P(g,\pi-\theta) dz.$$



• Dilute water measurement of single scattering

# Measuring scattering in water – Narasimhan et al. 06



measurement concentration

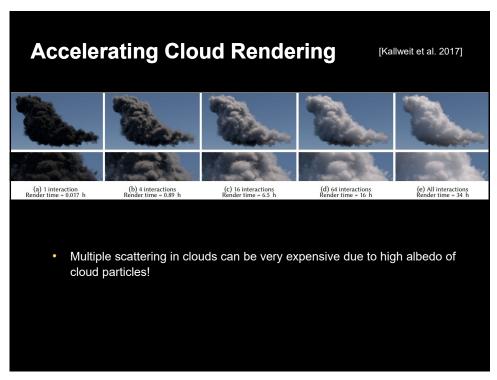


simulated actual concentration

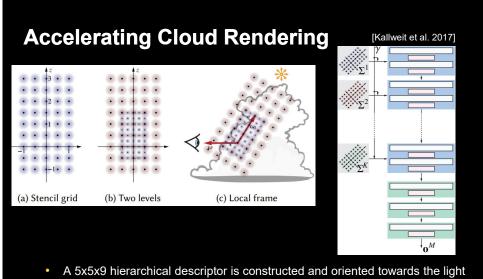
25

# **Radiative Transport**

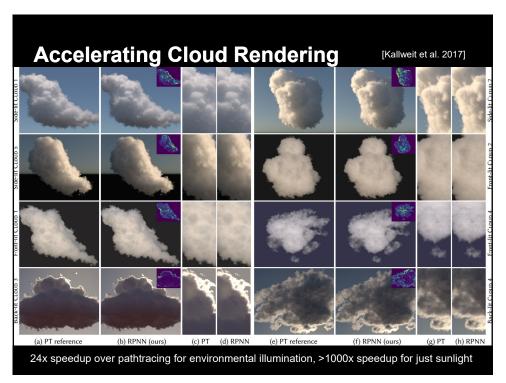
- · Costly computation for rendering of highly scattering materials
- Light also diffuses with multiple scattering!



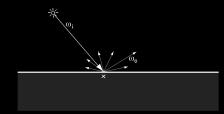
# Accelerating Cloud Rendering Reference RPNN (9 min) Reference Nultiple scattering is clouds predicted using neural networks. Single scattering and direct reflection still rendered using MC rendering. Multiple scattering can very directional in a cloud volume. Hence, a low-resolution 3D descriptor of the cloud volume is employed to correctly predict the spatial and direction distribution of radiance using an MLP.



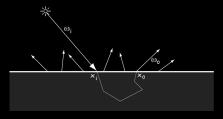
- inside the volume.
- The descriptor is an input to an MLP network for prediction of the directional radiance distribution. The MLP network processes the stencil hierarchically from coarse scale to fine scale for the prediction.



# **Subsurface scattering**



BRDF 
$$f_r(x, \vec{\omega}_i, \vec{\omega}_o) = \frac{dL(x, \vec{\omega}_i)}{dE(x, \vec{\omega}_o)}$$



BSSRDF 
$$S(x_i, \vec{\omega}_i, x_o, \vec{\omega}_o) = \frac{dL(x_o, \vec{\omega}_o)}{d\Phi(x_i, \vec{\omega}_i)}$$

- Bidirectional surface scattering distribution function (BSSRDF), 8D function
  - generalizes the 4D BRDF of surface reflection to transport across material

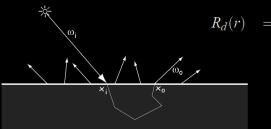
31

Diffuse BSSRDF - Jensen et al. 01

$$R_d(r) = \frac{(\vec{n} \cdot \vec{E}(x_o))}{d\Phi_i(x_i)}$$

- \*\*
- Diffuse BSSRDF Rd given by ratio of radiant exitance and incident flux.
   No dependence on angle of incidence or exitance, only distance of transport!
- R<sub>d</sub>(r) is a 1-D radially symmetric diffusion kernel that depends on distance r of scattering within material, i.e., distance between x<sub>i</sub> and x<sub>o</sub>.

### Diffuse BSSRDF [Jensen et al. 01]



$$R_d(r) = rac{lpha'}{4\pi} \left\{ z_r \left( \sigma_{tr} + rac{1}{d_r} 
ight) rac{e^{-\sigma_{tr}d_r}}{d_r^2} + 
ight.$$

$$\left. z_v \left( \sigma_{tr} + rac{1}{d_v} 
ight) rac{e^{-\sigma_{tr}d_v}}{d_v^2} 
ight\}$$

BSSRDF  $S_d(x_i, \omega_i, x_o, \omega_o) \approx R_d(x_i, x_o)$ 

- $R_d(x_i, x_o)$  models isotropic Gaussian-like diffusion between points  $x_i$  and  $x_o$ 
  - Dipole model for homogeneous semi-infinite medium
  - Sum of two Gaussian like fall-offs

33

# Dipole diffusion - Jensen et al. 01

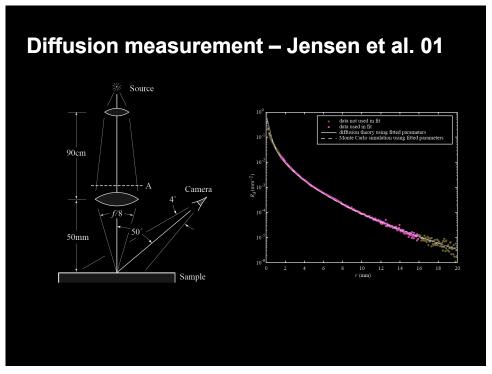
$$R_d(r) = rac{lpha'}{4\pi} \left\{ z_r \left( \sigma_{tr} + rac{1}{d_r} 
ight) rac{e^{-\sigma_{tr}d_r}}{d_r^2} + 
ight.$$
  $z_{\mathcal{V}} \left( \sigma_{tr} + rac{1}{d_{\mathcal{V}}} 
ight) rac{e^{-\sigma_{tr}d_{\mathcal{V}}}}{d_{\mathcal{V}}^2} 
ight\}$ 

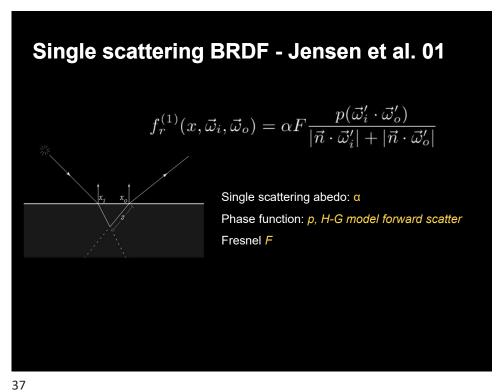


distance of real source from  $x_o$ :  $d_r$  distance of virtual source from  $x_o$ :  $d_v$   $z_r$  &  $z_v$  distance of sources to the surface  $\alpha$ ' albedo,  $\sigma_{tr}$  is extinction coefficient

- Planar homogeneous semi-infinite medium
- Dipoles enforce the net inward flux at the boundary to be zero!
- Think of dipole diffusion as a sum of two Gaussian like fall-offs.

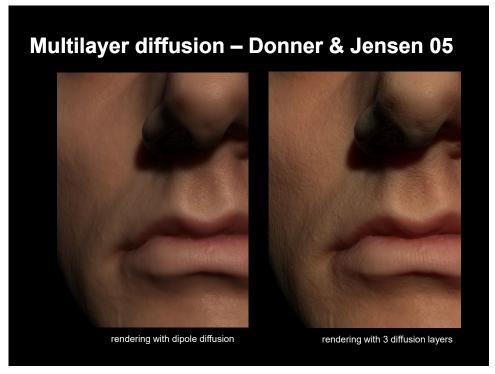


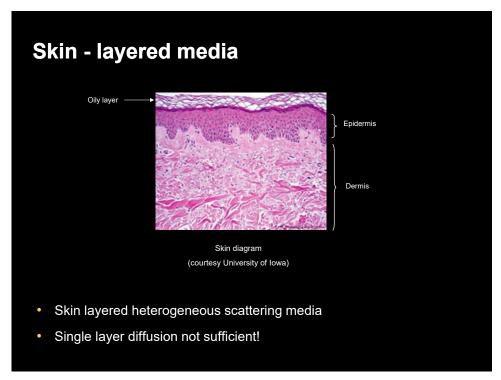




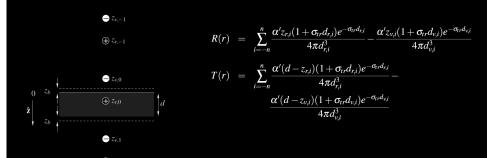








# Multipole model – Donner&Jensen 05



 Multipole models reflectance and transmissions through thin layers more accurately than dipole model

### **Kubelka-Munk theory**

$$T_{12} = T_1 * T_2 + T_1 * R_2 * R_1 * T_2 + T_1 * R_2 * R_1 * R_2 * R_1 * T_2 + \dots$$

$$\oplus z_{\kappa-1}$$

 $\bigoplus z_{r,1}$ 

$$\mathfrak{T}_{12} = \frac{\mathfrak{T}_1 \mathfrak{T}_2}{1 - \mathfrak{R}_2 \mathfrak{R}_1}$$

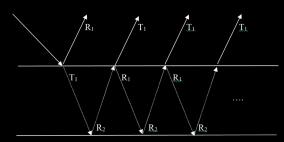
$$\mathcal{R}_{12} = \mathcal{R}_1 + \frac{\mathcal{T}_1 \mathcal{R}_2 \mathcal{T}_1}{1 - \mathcal{R}_2 \mathcal{R}_1}$$

• Derived from geometric series formula!

43

# **Kubelka-Munk theory – geometric series**

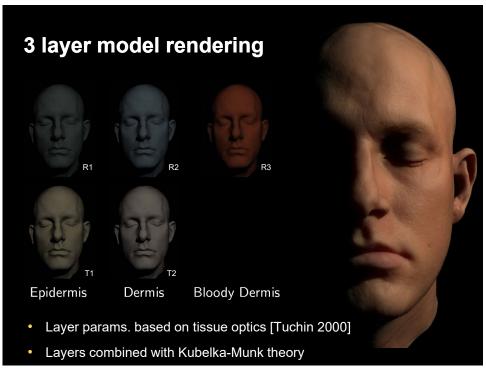
$$\mathcal{R}_{12} = \mathcal{R}_1 + \frac{\mathcal{T}_1 \mathcal{R}_2 \mathcal{T}_1}{1 - \mathcal{R}_2 \mathcal{R}_1}$$

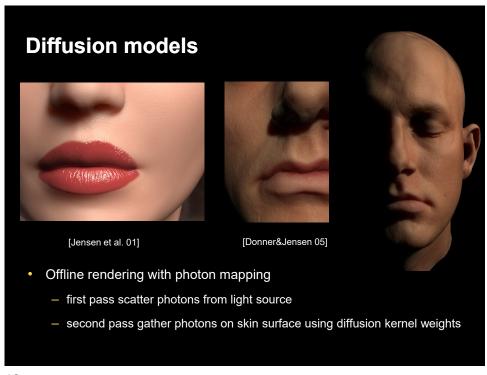


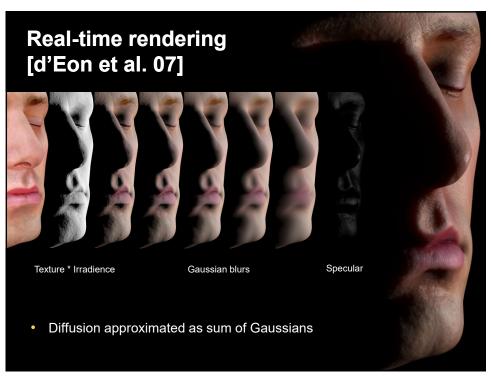
• R<sub>12</sub> = R<sub>1</sub> + T<sub>1</sub> \* R<sub>2</sub> \* T<sub>1</sub> + T<sub>1</sub> \* R<sub>2</sub> \* R<sub>1</sub> \* R<sub>2</sub> \* T<sub>1</sub> + T<sub>1</sub> \* R<sub>2</sub> \* R<sub>1</sub> \* R<sub>2</sub> \* R<sub>1</sub> \* R<sub>2</sub> \* T<sub>1</sub> + ....

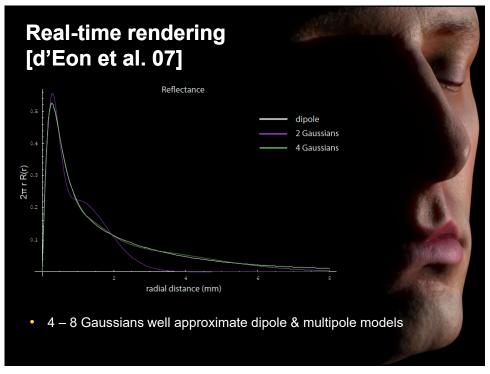
$$= R_1 + T_1 * R_2 * T_1 * (1 + (R_1 * R_2) + (R_1 * R_2)^2 + (R_1 * R_2)^3 + ...)$$

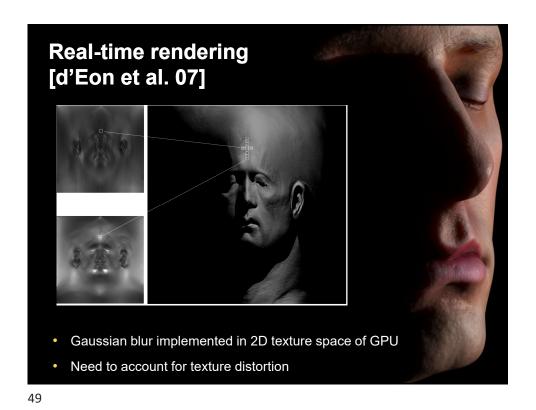
- $= R_1 + T_1 * R_2 * T_1 * (1/(1 (R_1 * R_2)))$
- Derived from geometric series formula!











Real-time rendering
[d'Eon et al. 07]

4 – 8 Gaussians well approximate dipole & multipole models

