

# 1 Tensor-Product Basis

The tensor-product basis is

$$\phi_{(i,j)}(\alpha, \beta) = \tilde{\phi}_{(i)}(\alpha)\tilde{\phi}_{(j)}(\beta), \quad (1)$$

where  $\tilde{\phi}_{(i)}(x)$  denotes the usual 1D GLL basis function at node  $i \in (0, \dots, n_p)$ . For vector fields, the components of the covariant vector field are given by the tensor-product basis (1).

## 2 Hyperviscosity

### 2.1 Vector Hyperviscosity

Fourth-order vector hyperviscosity is implemented using a two stage procedure:

$$\mathbf{f} = -\mathcal{H}(1, 1)\mathbf{u}^n, \quad (2)$$

$$\mathbf{u}^{n+1} = \mathbf{u}^n + \Delta t \mathcal{H}(\nu_d, \nu_v)\mathbf{f}. \quad (3)$$

The hyperdiffusion operator is defined implicitly via

$$\mathbf{f} = \mathcal{H}(\nu_d, \nu_v)\mathbf{u}^n \iff \iint \mathbf{f} \cdot \boldsymbol{\phi} dA = \iint \nu_d (\nabla \cdot \boldsymbol{\phi}) (\nabla \cdot \mathbf{u}^n) + \nu_v (\nabla \times \boldsymbol{\phi})^r (\nabla \times \mathbf{u}^n)_r dA, \quad (4)$$

where  $dA = J d\alpha d\beta$  and

$$(\nabla \cdot \boldsymbol{\phi}) = g^{pq} \nabla_p \phi_q = \frac{1}{J} \frac{\partial}{\partial \alpha} (J g^{\alpha\alpha} \phi_\alpha + J g^{\alpha\beta} \phi_\beta) + \frac{1}{J} \frac{\partial}{\partial \beta} (J g^{\beta\alpha} \phi_\alpha + J g^{\beta\beta} \phi_\beta), \quad (5)$$

$$(\nabla \times \boldsymbol{\phi})^r = \epsilon^{rpq} \nabla_p \phi_q = \epsilon^{rpq} \left[ \frac{\partial \phi_q}{\partial x^p} - \Gamma_{pq}^k \phi_k \right] = \frac{1}{J} \left[ \frac{\partial \phi_\beta}{\partial \alpha} - \frac{\partial \phi_\alpha}{\partial \beta} \right]. \quad (6)$$

Here we assume that  $(\nabla \cdot \mathbf{u})$  and  $(\nabla \times \mathbf{u})_r$  (covariant radial component of curl) have already been computed.

#### 2.1.1 Vector basis with zero $\beta$ component

If  $\phi_{(i,j)\alpha} = \tilde{\phi}_{(i)}(\alpha)\tilde{\phi}_{(j)}(\beta)$  and  $\phi_{(i,j)\beta} = 0$  then

$$\iint \mathbf{f} \cdot \boldsymbol{\phi} dA = \iint f^\alpha \tilde{\phi}_{(i)}(\alpha) \tilde{\phi}_{(j)}(\beta) dA = f_{(i,j)}^\alpha w_i w_j J \Delta\alpha \Delta\beta \quad (7)$$

The divergent term is defined by

$$(\nabla \cdot \boldsymbol{\phi}_{(i,j)}) = \frac{1}{J} \frac{\partial}{\partial \alpha} (J g^{\alpha\alpha} \phi_{(i,j)\alpha}) + \frac{1}{J} \frac{\partial}{\partial \beta} (J g^{\beta\alpha} \phi_{(i,j)\alpha}) \quad (8)$$

$$= \frac{\tilde{\phi}_{(j)}(\beta)}{J} \frac{\partial}{\partial \alpha} (J g^{\alpha\alpha} \tilde{\phi}_{(i)}(\alpha)) + \frac{\tilde{\phi}_{(i)}(\alpha)}{J} \frac{\partial}{\partial \beta} (J g^{\beta\alpha} \tilde{\phi}_{(j)}(\beta)), \quad (9)$$

and

$$\begin{aligned} & \iint (\nabla \cdot \phi_{(i,j)}) (\nabla \cdot \mathbf{u}) dA \\ &= \Delta\alpha\Delta\beta \sum_{m=0}^{n_p-1} \sum_{n=0}^{n_p-1} \left[ \frac{\tilde{\phi}_{(j)}(\beta_n)}{J} \frac{\partial}{\partial\alpha} \left( Jg^{\alpha\alpha} \tilde{\phi}_{(i)}(\alpha) \right) + \frac{\tilde{\phi}_{(i)}(\alpha_m)}{J} \frac{\partial}{\partial\beta} \left( Jg^{\beta\alpha} \tilde{\phi}_{(j)}(\beta) \right) \right] (\nabla \cdot \mathbf{u}) Jw_m w_n \quad (10) \end{aligned}$$

$$\begin{aligned} &= \Delta\alpha\Delta\beta w_j \sum_{m=0}^{n_p-1} \frac{\partial}{\partial\alpha} \left( Jg^{\alpha\alpha} \tilde{\phi}_{(i)}(\alpha) \right) (\nabla \cdot \mathbf{u}) w_m \Big|_{\alpha=\alpha_m, \beta=\beta_j} \\ &\quad + \Delta\alpha\Delta\beta w_i \sum_{n=0}^{n_p-1} \frac{\partial}{\partial\beta} \left( Jg^{\beta\alpha} \tilde{\phi}_{(j)}(\beta) \right) (\nabla \cdot \mathbf{u}) w_n \Big|_{\alpha=\alpha_i, \beta=\beta_n} \quad (11) \end{aligned}$$

$$\begin{aligned} &= \Delta\alpha\Delta\beta w_j \sum_{m=0}^{n_p-1} Jg^{\alpha\alpha} \frac{\partial \tilde{\phi}_{(i)}}{\partial\alpha} (\nabla \cdot \mathbf{u}) w_m \Big|_{\alpha=\alpha_m, \beta=\beta_j} \\ &\quad + \Delta\alpha\Delta\beta w_i \sum_{n=0}^{n_p-1} Jg^{\beta\alpha} \frac{\partial \tilde{\phi}_{(j)}}{\partial\beta} (\nabla \cdot \mathbf{u}) w_n \Big|_{\alpha=\alpha_i, \beta=\beta_n} \quad (12) \end{aligned}$$

Further, the vortical term is defined by

$$(\nabla \times \phi_{(i,j)})^r = -\frac{1}{J} \frac{\partial \phi_{(i,j)\alpha}}{\partial\beta} = -\frac{\tilde{\phi}_{(i)}}{J} \frac{\partial \tilde{\phi}_{(j)}}{\partial\beta} \quad (13)$$

and so

$$\begin{aligned} & \iint (\nabla \times \phi_{(i,j)})^r (\nabla \times \mathbf{u})_r dA \\ &= \Delta\alpha\Delta\beta \sum_{m=0}^{n_p-1} \sum_{n=0}^{n_p-1} \left[ -\frac{\tilde{\phi}_{(i)}(\alpha_m)}{J} \frac{\partial \tilde{\phi}_{(j)}}{\partial\beta} \right] (\nabla \times \mathbf{u})_r Jw_m w_n \Big|_{\alpha=\alpha_m, \beta=\beta_n} \quad (14) \end{aligned}$$

$$= -\Delta\alpha\Delta\beta w_i \sum_{n=0}^{n_p-1} \frac{\partial \tilde{\phi}_{(j)}}{\partial\beta} (\nabla \times \mathbf{u})_r w_n \Big|_{\alpha=\alpha_i, \beta=\beta_n} \quad (15)$$

Combining (7), (12) and (15) then gives

$$\begin{aligned} f_{(i,j)}^\alpha &= \frac{\nu_d}{J(\alpha_i, \beta_j) w_i} \sum_{m=0}^{n_p-1} Jg^{\alpha\alpha} \frac{\partial \tilde{\phi}_{(i)}}{\partial\alpha} (\nabla \cdot \mathbf{u}) w_m \Big|_{\alpha=\alpha_m, \beta=\beta_j} \\ &\quad + \frac{\nu_d}{J(\alpha_i, \beta_j) w_j} \sum_{n=0}^{n_p-1} Jg^{\beta\alpha} \frac{\partial \tilde{\phi}_{(j)}}{\partial\beta} (\nabla \cdot \mathbf{u}) w_n \Big|_{\alpha=\alpha_i, \beta=\beta_n} \\ &\quad - \frac{\nu_v}{J(\alpha_i, \beta_j) w_j} \sum_{n=0}^{n_p-1} \frac{\partial \tilde{\phi}_{(j)}}{\partial\beta} (\nabla \times \mathbf{u})_r w_n \Big|_{\alpha=\alpha_i, \beta=\beta_n} \quad (16) \end{aligned}$$

### 2.1.2 Vector basis with zero $\alpha$ component

If  $\phi_{(i,j)\alpha} = 0$  and  $\phi_{(i,j)\beta} = \tilde{\phi}_{(j)}(\alpha) \tilde{\phi}_{(j)}(\beta)$  then

$$\iint \mathbf{f} \cdot \phi dA = \iint f^\beta \tilde{\phi}_{(i)}(\alpha) \tilde{\phi}_{(j)}(\beta) dA = f_{(i,j)}^\beta w_i w_j J \Delta\alpha \Delta\beta \quad (17)$$

The divergent term is defined by

$$(\nabla \cdot \phi_{(i,j)}) = \frac{1}{J} \frac{\partial}{\partial \alpha} (Jg^{\alpha\beta} \phi_{(i,j)\beta}) + \frac{1}{J} \frac{\partial}{\partial \beta} (Jg^{\beta\beta} \phi_{(i,j)\beta}) \quad (18)$$

$$= \frac{\tilde{\phi}_{(j)}(\beta)}{J} \frac{\partial}{\partial \alpha} (Jg^{\alpha\beta} \tilde{\phi}_{(i)}) + \frac{\tilde{\phi}_{(i)}(\alpha)}{J} \frac{\partial}{\partial \beta} (Jg^{\beta\beta} \tilde{\phi}_{(j)}), \quad (19)$$

and

$$\begin{aligned} & \iint (\nabla \cdot \phi_{(i,j)}) (\nabla \cdot \mathbf{u}) dA \\ &= \Delta\alpha\Delta\beta \sum_{m=0}^{n_p-1} \sum_{n=0}^{n_p-1} \left[ \frac{\tilde{\phi}_{(j)}(\beta_n)}{J} \frac{\partial}{\partial \alpha} (Jg^{\alpha\beta} \tilde{\phi}_{(i)}) + \frac{\tilde{\phi}_{(i)}(\alpha_m)}{J} \frac{\partial}{\partial \beta} (Jg^{\beta\beta} \tilde{\phi}_{(j)}) \right] (\nabla \cdot \mathbf{u}) Jw_m w_n \end{aligned} \quad (20)$$

$$\begin{aligned} &= \Delta\alpha\Delta\beta w_j \sum_{m=0}^{n_p-1} \frac{\partial}{\partial \alpha} (Jg^{\alpha\beta} \tilde{\phi}_{(i)}(\alpha)) (\nabla \cdot \mathbf{u}) w_m \Big|_{\alpha=\alpha_m, \beta=\beta_j} \\ &\quad + \Delta\alpha\Delta\beta w_i \sum_{n=0}^{n_p-1} \frac{\partial}{\partial \beta} (Jg^{\beta\beta} \tilde{\phi}_{(j)}(\beta)) (\nabla \cdot \mathbf{u}) w_n \Big|_{\alpha=\alpha_i, \beta=\beta_n} \end{aligned} \quad (21)$$

$$\begin{aligned} &= \Delta\alpha\Delta\beta w_j \sum_{m=0}^{n_p-1} Jg^{\alpha\beta} \frac{\partial \tilde{\phi}_{(i)}}{\partial \alpha} (\nabla \cdot \mathbf{u}) w_m \Big|_{\alpha=\alpha_m, \beta=\beta_j} \\ &\quad + \Delta\alpha\Delta\beta w_i \sum_{n=0}^{n_p-1} Jg^{\beta\beta} \frac{\partial \tilde{\phi}_{(j)}}{\partial \beta} (\nabla \cdot \mathbf{u}) w_n \Big|_{\alpha=\alpha_i, \beta=\beta_n} \end{aligned} \quad (22)$$

Further, the vortical term is defined by

$$(\nabla \times \phi_{(i,j)})^r = \frac{1}{J} \frac{\partial \phi_{(i,j)\beta}}{\partial \alpha} = \frac{\tilde{\phi}_{(j)}}{J} \frac{\partial \tilde{\phi}_{(i)}}{\partial \alpha} \quad (23)$$

and so

$$\begin{aligned} & \iint (\nabla \times \phi_{(i,j)})^r (\nabla \times \mathbf{u})_r dA \\ &= \Delta\alpha\Delta\beta \sum_{m=0}^{n_p-1} \sum_{n=0}^{n_p-1} \left[ \frac{\tilde{\phi}_{(j)}(\beta_n)}{J} \frac{\partial \tilde{\phi}_{(i)}}{\partial \alpha} \right] (\nabla \times \mathbf{u})_r Jw_m w_n \Big|_{\alpha=\alpha_m, \beta=\beta_n} \end{aligned} \quad (24)$$

$$= \Delta\alpha\Delta\beta w_j \sum_{m=0}^{n_p-1} \frac{\partial \tilde{\phi}_{(i)}}{\partial \alpha} (\nabla \times \mathbf{u})_r w_m \Big|_{\alpha=\alpha_m, \beta=\beta_j} \quad (25)$$

Combining (17), (22) and (25) then gives

$$\begin{aligned} f_{(i,j)}^\beta &= \frac{\nu_d}{J(\alpha_i, \beta_j) w_i} \sum_{m=0}^{n_p-1} Jg^{\alpha\beta} \frac{\partial \tilde{\phi}_{(i)}}{\partial \alpha} (\nabla \cdot \mathbf{u}) w_m \Big|_{\alpha=\alpha_m, \beta=\beta_j} \\ &\quad + \frac{\nu_d}{J(\alpha_i, \beta_j) w_j} \sum_{n=0}^{n_p-1} Jg^{\beta\beta} \frac{\partial \tilde{\phi}_{(j)}}{\partial \beta} (\nabla \cdot \mathbf{u}) w_n \Big|_{\alpha=\alpha_i, \beta=\beta_n} \\ &\quad + \frac{\nu_v}{J(\alpha_i, \beta_j) w_i} \sum_{m=0}^{n_p-1} \frac{\partial \tilde{\phi}_{(i)}}{\partial \alpha} (\nabla \times \mathbf{u})_r w_m \Big|_{\alpha=\alpha_m, \beta=\beta_j} \end{aligned} \quad (26)$$