Chapter 2 Problems

William Arnold

May 19, 2020

1 Problem 2.1

Prove the following theorems:

(a) For normalizable solutions, the separation constant E must be *real*

Assume E is complex, meaning $E = a + ib$ for some $a, b \in \mathbb{R}$. For a separable solution $\Psi(x,t) = \psi(x)e^{-iEt/\hbar}$, we have that

$$\begin{aligned} |\Psi(x,t)|^2 &= |\psi|^2 \left| e^{-i(a+ib)t/\hbar} \right|^2 \\ &= |\psi|^2 \left| e^{-iat/\hbar} e^{bt/\hbar} \right|^2 \\ &= |\psi|^2 e^{2bt/\hbar} \end{aligned}$$

For $t = 0$, we know that $\int_{-\infty}^{\infty} |\Psi(x,0)|^2 dx = 1$ so

$$\int_{-\infty}^{\infty} |\Psi(x,0)|^2 dx = \int_{-\infty}^{\infty} |\psi|^2 e^0 dx = \int_{-\infty}^{\infty} |\psi|^2 dx = 1$$

Now for any $t > 0$, we have

$$\int_{-\infty}^{\infty} |\Psi(x,t)|^2 dx = \int_{-\infty}^{\infty} |\psi|^2 e^{2bt/\hbar} dx = e^{2bt/\hbar} \int_{-\infty}^{\infty} |\psi|^2 dx = e^{2bt/\hbar} = 1$$

Taking the natural log of both sides we see

$$2bt/\hbar = 0$$

and hence $b = 0$ and E is composed of only the real component.

(b) The time independenent wave function ψ can always be taken to be *real* (unlike Ψ which is necessarily complex). That doesn't mean that every solution to the time-independent Schroedinger equation is real; what it says is that if you've got one that is not, it can always be expressed as a linear combination of solutions (with the same energy) that are real. So you might as well stick to ψs that are real.

Assume $\psi = a + ib$ is a solution to the time-independent Schroedinger equation

$$-\frac{\hbar^2}{2m}\frac{\partial^2\psi}{\partial x^2} + \psi V = \psi E$$

Since both ψ and ψ^* solve this, we can plug both in to this equation:

$$\begin{aligned}
-\frac{\hbar^2}{2m}\left(\frac{\partial^2 a}{\partial x^2} + i\frac{\partial^2 b}{\partial x^2}\right) + (a+ib)V = (a+ib)E \\
-\frac{\hbar^2}{2m}\left(\frac{\partial^2 a}{\partial x^2} - i\frac{\partial^2 b}{\partial x^2}\right) + (a-ib)V = (a-ib)E
\end{aligned}$$

Adding these equations together we get

$$\begin{aligned}
&-\frac{\hbar^2}{2m}\left(2\frac{\partial^2 a}{\partial x^2} + i\frac{\partial^2 b}{\partial x^2} - i\frac{\partial^2 b}{\partial x^2}\right) + (2a+ib-ib)V \quad = (a+ib-ib)E \\
&-\frac{\hbar^2}{2m}\left(2\frac{\partial^2 a}{\partial x^2}\right) + 2aV = 2aE \\
&-\frac{\hbar^2}{2m}\left(\frac{\partial^2 a}{\partial x^2}\right) + aV = aE
\end{aligned}$$

And hence $a = \frac{1}{2}(\psi + \psi^*)$ is also a solution, and has no imaginary component. It is worth noting that we should probably use $\frac{1}{\sqrt{2}}(\psi + \psi^*)$ instead, so that we still have a normalized solution.

(c) If $V(x)$ is an even function, then ψ can always be taken to be either even or odd.

Assume we have some solution $\psi(x)$. Let $\gamma(x) = \psi(-x)$. We then have

$$\begin{aligned}
\frac{\partial\gamma}{\partial x}(x) = -\frac{\partial\psi}{\partial x}(-x) \\
\frac{\partial^2\gamma}{\partial x^2}(x) = \frac{\partial^2\psi}{\partial x^2}(-x)
\end{aligned}$$

If we look at γ in the time-independent equation we get

$$\begin{aligned}
-\frac{\hbar^2}{2m}\frac{\partial^2\gamma}{\partial x^2}(x) + \gamma(x)V(x) = \gamma(x)E \\
-\frac{\hbar^2}{2m}\frac{\partial^2\psi}{\partial x^2}(-x) + \psi(-x)V(x) = \psi(-x)E \\
-\frac{\hbar^2}{2m}\frac{\partial^2\psi}{\partial x^2}(-x) + \psi(-x)V(-x) = \psi(-x)E
\end{aligned}$$

Where the last equation uses the evenness of V. This equation also holds since the time independent equation holds for all $x \in \mathbb{R}$. Thus, $\psi(-x)$ is also a solution if V is even.

Now, since any linear combination of solutions is also a solution, we can use the even combination $\psi(x) + \psi(-x)$ or the odd combination $\psi(x) - \psi(-x)$ in order to get an even or odd solution to the equation.

2 Problem 2.2

Prove that E must exceed the minimum value of $V(x)$ at some point.
Assume $E < V(x)$ for all x. Then we can solve for $\frac{\partial^2 \psi}{\partial x^2}$ to see that

$$\frac{\partial^2 \psi}{\partial x^2} = \frac{2m}{\hbar^2} \psi \left(V(x) - E \right)$$

This says that ψ and $\frac{\partial^2 \psi}{\partial x^2}$ always have the same sign, taking ψ as a real solution. Assuming ψ is C^2, if it ever takes a positive value, then it must continue to curve upwards in the positive direction and hence would grow unbounded. Symmetrically if ψ takes on some negative value, it will grow unboundedly negative in the negative direction. Hence, since $\psi \neq 0$ we have a contradiction and $E > V(x)$ for some x.

3 Problem 2.3

Calculate $\langle x \rangle, \langle x^2 \rangle, \langle p \rangle, \langle p^2 \rangle, \sigma_x$, and σ_p for the n-th stationary state of the infinite square well. Check that the uncertainty principle is satisfied. Which state comes closest to the uncertainty limit?

Recall that

$$\psi_n = \sqrt{\frac{2}{a}} \sin(\frac{n\pi}{a} x)$$

Thus we compute

$$\begin{aligned}
\langle x^2 \rangle &= \int_0^a x^2 |\psi_n|^2 dx \\
&= \left[\frac{1}{12} \left(\frac{3\left(a^2 - 2\pi^2 n^2 x^2\right) \sin\left(\frac{2\pi n x}{a}\right) - 6\pi a n x \cos\left(\frac{2\pi n x}{a}\right)}{\pi^3 n^3} + \frac{4x^3}{a} \right) \right]_0^a \\
&= \frac{a^2 \left(4\pi^3 n^3 + 3\sin(2\pi n) - 6\pi n(\pi n \sin(2\pi n) + \cos(2\pi n))\right)}{12\pi^3 n^3} \\
&= \frac{a^2 (4\pi^3 n^3 - 6\pi n)}{12\pi^3 n^3} \\
&= \frac{a^2}{3} - \frac{a^2}{2\pi^2 n^2} \\
\langle x \rangle &= \int_0^a x |\psi_n|^2 dx \\
&= \left[\frac{a\left(-2\pi^2 n^2 + 2\pi n \sin(2\pi n) + \cos(2\pi n) - 1\right)}{4\pi^2 n^2} \right] \\
&= \frac{a}{2} \text{ (I used mathematica)}
\end{aligned}$$

Thus we have

$$\begin{aligned}
\sigma_x^2 &= \left\langle x^2 \right\rangle - \left\langle x \right\rangle^2 \\
&= \frac{a^2}{3} - \frac{a^2}{2\pi^2 n^2} - \frac{a^2}{4} \\
&= a^2 \left(\frac{1}{12} - \frac{1}{2n^2\pi^2} \right)
\end{aligned}$$

It makes sense that the standard deviation would scale linearly with the area.

$$\begin{aligned}
\langle p \rangle &= \int_0^a \psi^* \left[-i\hbar \frac{\partial \psi}{\partial x} \right] dx \\
&= \int_0^a \frac{-2\hbar i n\pi \cos(\frac{n\pi x}{a}) \sin(\frac{n\pi x}{a})}{a^2} dx \\
&= \left[\frac{\hbar i \cos(\frac{2n\pi x}{a})}{2a} \right]_0^a \\
&= 0 \\
\left\langle p^2 \right\rangle &= \int_0^a \psi^* \left[-\hbar^2 \frac{\partial^2 \psi}{\partial x^2} \right] dx \\
&= \int_0^a \frac{1}{a^3} 2\hbar^2 n^2 \pi^2 \sin^2(\frac{n\pi x}{a}) \\
&= \frac{\hbar^2 n^2 \pi^2}{a^2}
\end{aligned}$$

Thus $\sigma_p^2 = \left\langle p^2 \right\rangle = \frac{\hbar^2 n^2 \pi^2}{a^2}$

We can use this to check the uncertainty principle:

$$\begin{aligned}
\sigma_x \sigma_p &= \left(\frac{\hbar n \pi}{a} \right) \left(a \sqrt{\frac{1}{12} - \frac{1}{2n^2\pi^2}} \right) \\
&> \hbar n \pi \frac{1}{2\sqrt{3}} \\
&> \hbar/2
\end{aligned}$$

This quantity is minimized for $n = 1$