Application of a Discontinuous Petrov-Galerkin method to compressible flow problems

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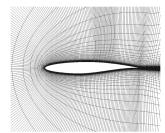
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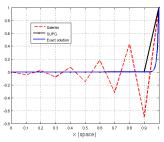
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Compressible Navier-Stokes equations

Numerical difficulties:

- Resolving solution features (sharp, localized viscous-scale phenomena)
 - Shocks
 - Boundary layers resolution needed for drag/load
- Stability of numerical schemes
 - Coarse/adaptive grids
 - Higher order
- Stabilized methods
 - Stabilization parameters
 - Higher-order





DPG: a minimum residual method via optimal testing

Given a trial space U and Hilbert test space V,

$$b(u,v) = \ell(v) \iff Bu = \ell, \qquad \begin{cases} \langle Bu,v \rangle_V & \coloneqq b(u,v) \\ \langle \ell,v \rangle_V & \coloneqq \ell(v). \end{cases}$$

We seek to minimize the dual residual over $U_h \subset U$

$$J(u_h) = \frac{1}{2} \|Bu_h - \ell\|_{V'}^2 = \frac{1}{2} \|R_V^{-1}(Bu_h - \ell)\|_V^2.$$

where $R_V:V\to V'$ is the isometric Riesz map

$$\langle R_V v, \delta v \rangle_V := (v, \delta v)_V, \quad \forall \delta v \in V.$$

Conditions for minimization of the convex functional give

$$b(u_h, v_{\delta u}) = \ell(v_{\delta u}), \quad \forall \delta u \in U_h, \quad v_{\delta u} := R_V^{-1} B \delta u.$$

Practical details of DPG

Computation of $v_{\delta u} := R_V^{-1} B \delta u$ is **global** and **infinite-dimensional**.

- By choosing a broken test space V and localizable norm $\|v\|_V^2 = \sum_K \|v\|_{V(K)}^2$, test functions can be determined locally.
- In practice, we use an enriched space $V_h \subset V$, where $\dim(V_h) > \dim(U_h)$ elementwise, and optimal test functions are approximated by computing $v_{\delta u} := R_{V_h}^{-1} B \delta u \in V_h$ through¹

$$(v_{\delta u}, \delta v)_V = b(\delta u, \delta v), \quad \delta u \in U_h, \quad \forall \delta v \in V_h$$

Typically, if $U_h = \mathcal{P}^p(\mathbb{R}^n)$, $V_h = \mathcal{P}^{p+\triangle p}(\mathbb{R}^n)$, where $\triangle p \geq n$.²

¹L. Demkowicz and J. Gopalakrishnan. A class of discontinuous Petrov-Galerkin methods. II. Optimal test functions. *Num. Meth. for Partial Diff. Eq.* 27:70–105, 2011

²J. Gopalakrishnan and W. Qiu. An analysis of the practical DPG method. Technical report. IMA, 2011. Submitted

Properties of DPG

■ Stiffness matrices are symmetric positive-definite. For trial/test bases $\{\phi_j\}_{j=1}^m$ and $\{v_i\}_{i=1}^n$, with $B_{ji} = b(\phi_j, v_i)$ and $I_i = \ell(v_i)$. DPG solves

$$\left(B^T R_V^{-1} B\right) u = \left(B^T R_V^{-1}\right) I,$$

For localizable norms and discontinuous testing, R_V is block diagonal.

■ DPG provides the best approximation in the energy norm

$$||u||_{E} = ||Bu||_{V'} = \sup_{||v||_{V}=1} |b(u,v)|.$$

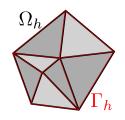
■ The energy error is computable through the error representation function e defined through $(e, \delta v)_V = \ell(v) - b(u_h, \delta v)$ for all $\delta v \in V$.

$$||u - u_h||_E = ||B(u - u_h)||_{V'} = ||R_V^{-1}(I - Bu_h)||_V = ||e||_V$$

Ultra-weak formulation for convection-diffusion

The first order convection-diffusion system:

$$A(u,\sigma) := \left[\begin{array}{c} \nabla \cdot (\beta u - \sigma) \\ \frac{1}{\epsilon} \sigma - \nabla u \end{array} \right] = \left[\begin{array}{c} f \\ 0 \end{array} \right].$$



The variational formulation is

$$b\left(\left(u,\sigma,\widehat{u},\widehat{f}_{n}\right),\left(v,\tau\right)\right) = \left(u,\nabla_{h}\cdot\tau - \beta\cdot\nabla_{h}v\right)_{\Omega_{h}} + \left(\sigma,\epsilon^{-1}\tau + \nabla_{h}v\right)_{\Omega_{h}} - \left\langle \left[\tau\cdot n\right],\widehat{u}\right\rangle_{\Gamma_{h}} + \left\langle\widehat{f}_{n},\left[v\right]\right\rangle_{\Gamma_{h}},$$

where $\widehat{f}_n := \beta_n u - \sigma_n$ and $\left\langle \widehat{f}_n, \llbracket v \rrbracket \right\rangle_{\Gamma}$ is defined

$$\left\langle \widehat{f}_n, \llbracket v \rrbracket \right\rangle_{\Gamma_h} := \sum_K \int_{\partial K} \operatorname{sgn}(\vec{n}) \, \widehat{f}_n v.$$

Graph norm under convection-diffusion

The graph norm³ for convection-diffusion gives exceptional stability.

$$\|(v,\tau)\|_{V(K)}^2 = \|\nabla \cdot \tau - \beta \cdot \nabla v\|_{L^2(K)}^2 + \|\epsilon^{-1}\tau + \nabla v\|_{L^2(K)}^2 + \|v\|_{L^2(K)}^2.$$

Problem with this test norm: approximability of test functions.

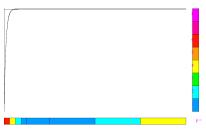


Figure: Components of optimal test functions for flux \hat{f}_n on the *right-hand* side of a unit element for $\epsilon = 0.01$.

³T. Bui-Thanh, L. Demkowicz, and O. Ghattas, A unified discontinuous Petrov-Galerkin method and its analysis for Friedrichs' systems. Submitted to SIAM J. Numer. Anal., 2011. Also ICES report 11-34, November 2011

Determining an alternative test norm

$$b(\mathbf{U}, \mathbf{V}) = (u, \nabla \cdot \tau - \beta \cdot \nabla v)_{\Omega_h} + (\sigma, \epsilon^{-1}\tau + \nabla v)_{\Omega_h} + \text{boundary terms}$$

Recover $\|u\|_{L^2(\Omega)}^2$ with conforming (v,τ) satisfying the *adjoint equations*

$$\begin{array}{rcl}
\nabla \cdot \tau - \beta \cdot \nabla v &= u \\
\frac{1}{\epsilon} \tau + \nabla v &= 0
\end{array}, \quad \langle \llbracket \tau \cdot n \rrbracket, \widehat{u} \rangle_{\Gamma}, \left\langle \widehat{f}_{n}, \llbracket v \rrbracket \right\rangle_{\Gamma} = 0$$

"Necessary" conditions for robustness —

$$\|u\|_{L^{2}(\Omega)}^{2} = b(\mathbf{U}, (v, \tau)) = \frac{b(\mathbf{U}, (v, \tau))}{\|(v, \tau)\|_{V}} \|(v, \tau)\|_{V} \le \|\mathbf{U}\|_{E} \|(v, \tau)\|_{V}$$

Let \lesssim denote a robust bound - if $\|(v,\tau)\|_V \lesssim \|u\|_{L^2(\Omega)}$, then we have that

$$||u||_{L^2(\Omega)} \lesssim ||\mathbf{U}||_E$$

Main idea: the test norm should measure adjoint solutions robustly.

Dirichlet inflow condition: issues as $\epsilon \to 0$

Standard choice of boundary condition: $u = u_0$ on inflow boundary Γ_{in} , induces boundary layers in adjoint problems, $\|\beta \cdot \nabla v\|_{L^2} = O(\epsilon^{-1})$.

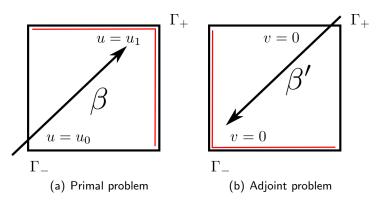


Figure: For the standard Dirichlet inflow condition, the solution to the adjoint problem can develop strong boundary layers at the outflow of the adjoint problem.

Solution: New inflow boundary condition on \widehat{f}_n

Non-standard choice of boundary condition: $\hat{f}_n = \beta_n u_0$ on $\Gamma_{\rm in}$, induces smoother adjoint problems, $\|\beta \cdot \nabla v\|_{L^2} = O(1)$.

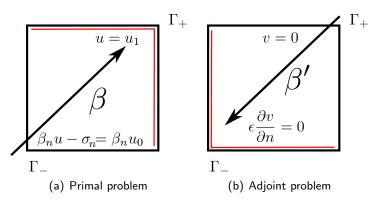


Figure: Under the new inflow condition, the wall-stop boundary condition is relaxed to a zero-stress condition at the outflow boundary of the adjoint problem.

Test norms and adjoint solutions

Intuition: the effectiveness of DPG under a test norm is governed by how a **specific test norm** measures the **solutions of the adjoint problem**.

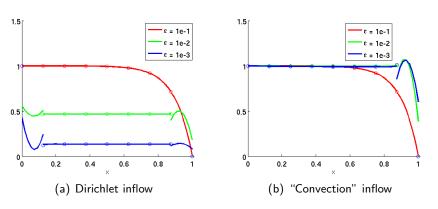


Figure: DPG solutions to convection-diffusion for both inflow conditions using an H^1 test norm.

Mesh-scaled test norms

For solutions (v, τ) of the adjoint equations, we derive quantities that are robustly bounded from above by $||u||_{L^2(\Omega)}$. Our test norm, as defined over a single element K, is now

$$\| (v,\tau) \|_{V,K}^2 = \min \left\{ \frac{\epsilon}{|K|}, 1 \right\} \|v\|^2 + \epsilon \|\nabla v\|^2 + \|\beta \cdot \nabla v\|^2 + \|\nabla \cdot \tau\|^2 + \min \left\{ \frac{1}{\epsilon}, \frac{1}{|K|} \right\} \|\tau\|^2.$$

which induces the proven robust bound⁴

$$\|u\|_{L^{2}(\Omega)} + \|\sigma\|_{L^{2}(\Omega)} + \epsilon \|\widehat{u}\| + \sqrt{\epsilon} \|\widehat{f}_{n}\| \lesssim \|(u, \sigma, \widehat{u}, \widehat{f}_{n})\|_{E}.$$

⁴J. Chan, N. Heuer, T. Bui Thanh, and L. Demkowicz. Robust DPG method for convection-diffusion problems II: natural inflow conditions. Technical Report 12-21, ICES, June 2012. Submitted

Numerical verification: Erikkson-Johnson problem

Experiments done using Camellia⁵ and a boundary layer solution.

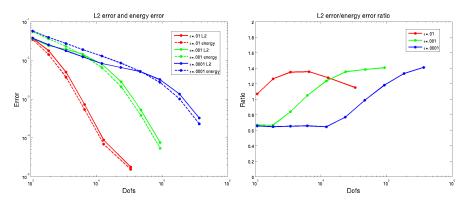


Figure: L^2 /energy errors and their ratio for $\epsilon = .01, .001, .0001$.

⁵N. Roberts, D. Ridzal, P. Bochev, and L. Demkowicz. A Toolbox for a Class of Discontinuous Petrov-Galerkin Methods Using Trilinos. Technical Report SAND2011-6678, Sandia National Laboratories, 2011

2D test case: Burgers equation

$$\frac{\partial \left(u^2/2\right)}{\partial x} + \frac{\partial u}{\partial y} + \epsilon \Delta u = f$$

Burgers equation can be written with $\beta(u)=(u/2,1)$

$$\nabla \cdot (\beta(u)u - \sigma) = f$$
$$\frac{1}{\epsilon}\sigma - \nabla u = 0.$$

i.e. nonlinear convection-diffusion on domain $[0,1]^2$.

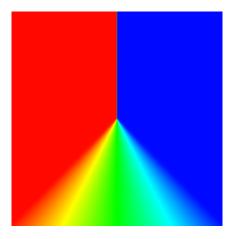


Figure: Shock solution for Burgers' equation, $\epsilon=1e-4$, using Newton-Raphson.

Adaptivity begins with a cubic 4×4 mesh.

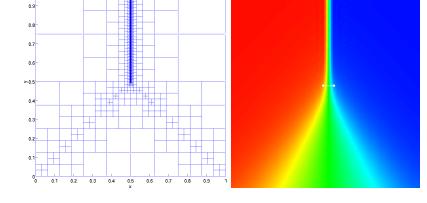


Figure: Adaptive mesh after 9 refinements, and zoom view at point (.5,.5) with shock formation and 1e-3 width line for reference.

2D Compressible Navier-Stokes equations (ideal gas)

Given density ρ , velocities $\mathbf{u} = (u_1, u_2)$ and temperature T,

$$\nabla \cdot \begin{bmatrix} \rho u_1 \\ \rho u_2 \end{bmatrix} = 0,$$

$$\nabla \cdot \left(\begin{bmatrix} \rho u_1^2 + p \\ \rho u_1 u_2 \end{bmatrix} - \sigma_1 \right) = 0,$$

$$\nabla \cdot \left(\begin{bmatrix} \rho u_1 u_2 \\ \rho u_2^2 + p \end{bmatrix} - \sigma_2 \right) = 0,$$

$$\nabla \cdot \left(\begin{bmatrix} ((\rho e) + p) u_1 \\ ((\rho e) + p) u_2 \end{bmatrix} - \sigma \mathbf{u} + \vec{q} \right) = 0,$$

$$\frac{1}{2\mu} \boldsymbol{\sigma} - \frac{\lambda}{4\mu(\mu + \lambda)} \text{tr}(\boldsymbol{\sigma}) \mathbf{I} = \nabla \mathbf{u} - \text{Re} \, \mathbf{w},$$

$$\kappa^{-1} \vec{q} = \nabla T,$$

where **w** represents the antisymmetric part of ∇u

Extrapolation of test norms

Convection-diffusion:

$$\nabla \cdot (\frac{\beta u}{\epsilon} - \sigma) = f$$
$$\frac{1}{\epsilon} \sigma - \nabla u = 0.$$

Navier-Stokes: defining vector variables $U = \{\rho, u_1, u_2, T\}$ and $\Sigma = \{\sigma, \mathbf{q}, w\}$,

$$\nabla \cdot (A_{\text{invisc}} \frac{U}{U} - A_{\text{visc}} \Sigma) = R_{\text{conserv}}(U, \Sigma)$$
$$E_{\text{visc}} \Sigma - \nabla U = R_{\text{constit}}(U, \Sigma)$$

where $R_{\rm conserv}(U,\Sigma)$ and $R_{\rm constit}(U,\Sigma)$ are the conservation/constitutive residuals.

Extrapolation of test norms

Convection-diffusion:

$$\nabla \cdot (\beta u - \sigma) = f$$
$$\frac{1}{\epsilon} \sigma - \nabla u = 0.$$

Navier-Stokes: defining vector variables $U = \{\rho, u_1, u_2, T\}$ and $\Sigma = \{\sigma, \mathbf{q}, w\}$,

$$\nabla \cdot (A_{\text{invisc}} U - A_{\text{visc}} \Sigma) = R_{\text{conserv}}(U, \Sigma)$$
$$\frac{E_{\text{visc}} \Sigma - \nabla U}{E_{\text{constit}}(U, \Sigma)}$$

where $R_{\rm conserv}(U,\Sigma)$ and $R_{\rm constit}(U,\Sigma)$ are the conservation/constitutive residuals.

Carter's flat plate problem

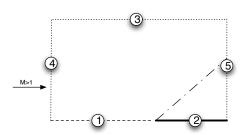
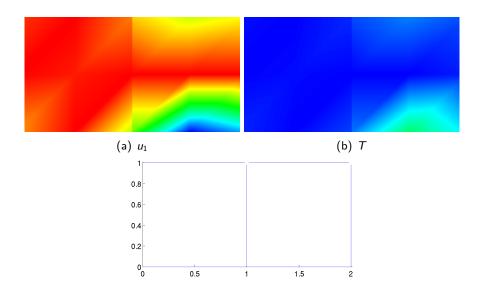
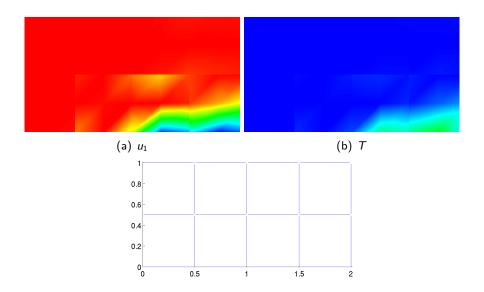


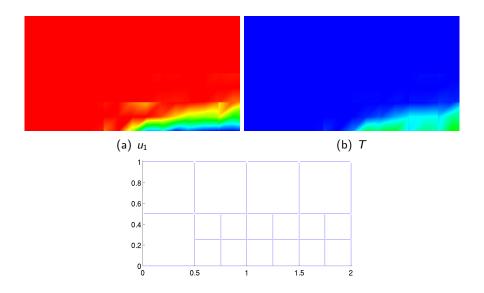
Figure: Carter flat plate problem on domain $[0,2] \times [0,1]$. Plate begins at x = 1, Re = 1000.

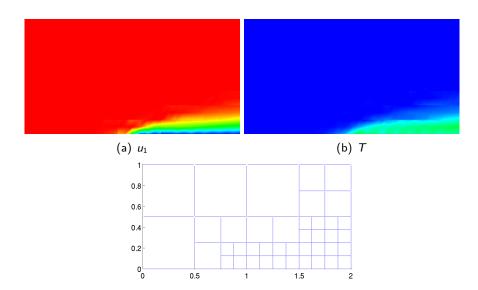
- Symmetry boundary conditions.
 Prescribed temperature and
- wall stagnation conditions.
- 3 Symmetry boundary conditions.
- 4 Inflow: conserved quantities specified using far-field values.
- 5 No outflow condition set.

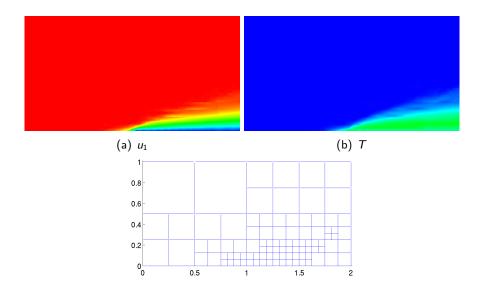
Stress/heat flux boundary conditions are set in terms of the momentum and energy fluxes.

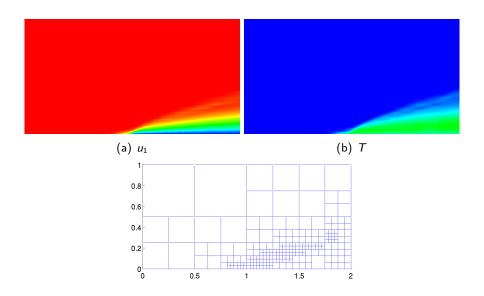


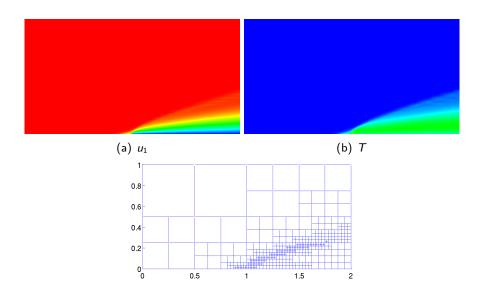


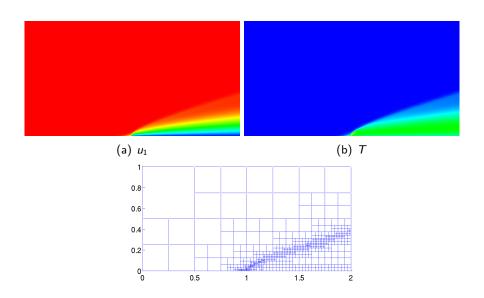


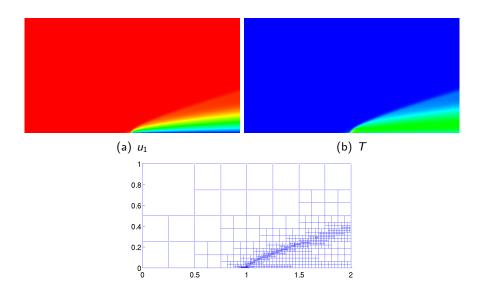


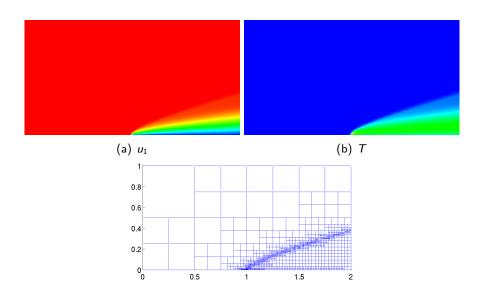


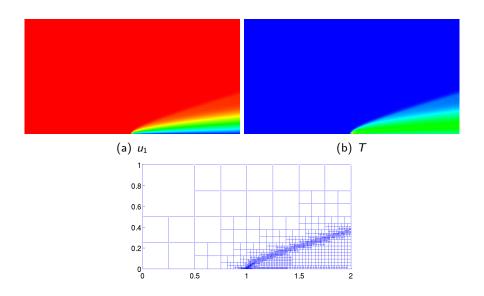




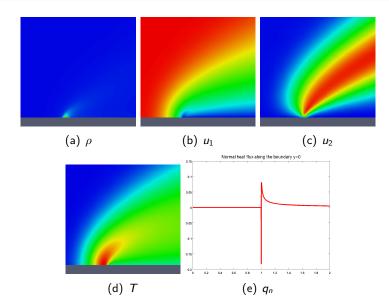








Zoomed solutions at plate/stagnation point



Thank you!

Questions?



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