

# Report For Discontinuous Galerkin Method

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## 1 Problem Statement

Using the third-order DG scheme to solve the Burgers' equation

$$\begin{cases} u_t + uu_x = 0, \\ u(x, 0) = g(x) = 1/2 + \sin(x), \quad 0 \leq x \leq 2\pi, \end{cases} \quad (1)$$

with periodic boundary condition. This report will contain the following contents:

- Derive the exact solution of Burgers equation (1).
- Algorithm design and some related properties.
- Numerical results.
- Discussions.

## 2 Burgers' equation

In this section, we will introduce some theoretical results for Burgers' equation.

Choose the **flux**  $f(u) := \frac{u^2}{2}$ , Burgers' equation is a special type of the 1D conservation law:

$$\begin{cases} u_t + (f(u))_x = 0, \\ u(x, 0) = g(x). \end{cases} \quad (2)$$

Consider the curve  $y = \eta(t; x_0)$  satisfying:

$$\begin{aligned} \eta'(t; x_0) &= u(\eta(t; x_0), t), \\ \eta(0) &= x_0, \end{aligned} \quad (3)$$

along the curve  $\eta(t; x_0)$ , equation (1) is equivalent to

$$\begin{cases} \frac{d}{dt} u(\eta(t; x_0), t) = 0, \\ u(x_0, 0) = g(x_0). \end{cases} \quad (4)$$

It means that  $u(x, t)$  remains constant along the curve  $\eta(t; x_0)$ , and the curve  $y = \eta(t; x_0)$  is the **characteristic line** for equation (1).

So, we find the solution  $u(x, t)$  by the following two steps:

- Derive a characteristic line  $y = \eta(t; \xi)$  from equation (3), which passes through the point  $(x, t)$ .
- Set  $u(x, t) := g(\xi)$ .

By (3),  $\eta(t; \xi) = \xi + g(\xi)t$ , i.e.

$$x = \xi + ut.$$

Then we can derive the solution as an implicit form:

$$u(x, t) = g(x - u(x, t)t). \quad (5)$$

If the characteristic lines don't intersect for  $t \in [0, T]$ , the equation (1) is well-posed. Unfortunately, it isn't still true for each  $T > 0$ . Here is the characteristic lines if we choose  $T = 1.2$ .

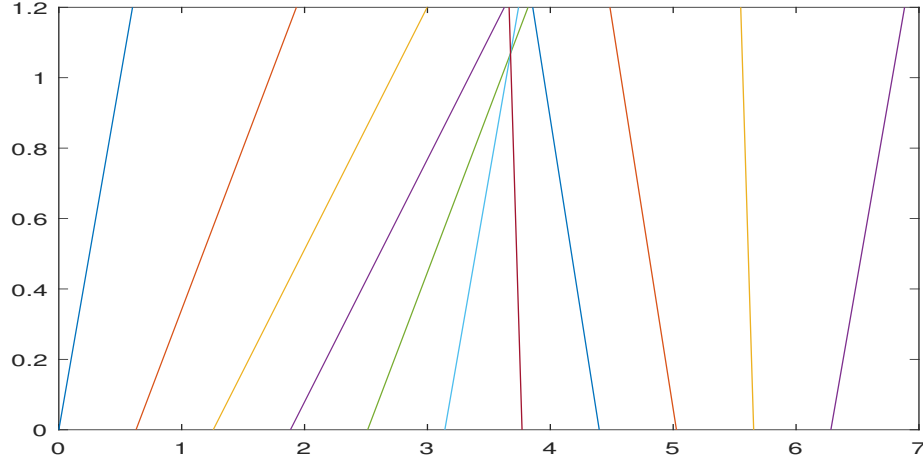


Figure 1: Characteristic lines.

If the characteristic lines intersect with each other, it means that equation (1) has no strong solutions, and its weak solutions are **incontinuous**. By [1, Exercise 3.3], the solution will break at time

$$T_b = \frac{-1}{\min g'(x)} = 1. \quad (6)$$

And the incontinuous point will move forward as well, which forms a **shock wave**. The entropy weak solution is a strong solution on both sides of the shock wave, and incontinuous on the shock wave points. In this problem, the shock wave originates from  $x_0 = \pi$  with the speed

$$s = \frac{g(x_0^-) + g(x_0^+)}{2} = 0.5, \quad (7)$$

see [1, (3.26)].

### 3 Algorithm

In this section, we introduce the DG scheme briefly with some tricks for implementation.

### 3.1 Mesh and finite element space

In this assignment, we choose equidistant grids. Assume there are  $N$  elements on the interval  $[0, 2\pi]$ , mark  $h := \frac{2\pi}{N}$ , the  $j$ -th element is

$$I_j := [x_{j-\frac{1}{2}}, x_{j+\frac{1}{2}}] = \left[ \frac{2\pi(j-1)}{N}, \frac{2\pi j}{N} \right],$$

and the finite element space

$$V_h^2 := \{u \in L^2([0, 2\pi]) : u|_{I_j} \in P^2(I_j)\}.$$

$P_2(I_j)$  is the vector space of all polynomials on  $I_j$  with degree less than or equal to 2. We choose the Legendre-form basis for  $P_2(I_j)$ , i.e.

$$\begin{aligned} P_2(I_j) &= \text{span}\{p_j^0, p_j^1, p_j^2\}, \\ p_j^0 &= 1, \quad p_j^1 = \frac{1}{h} \left( x - \frac{x_j + x_{j+1}}{2} \right), \\ p_j^2 &= \frac{3}{2} \left( \frac{2x - x_j - x_{j+1}}{h} \right)^2 - \frac{1}{2}. \end{aligned} \tag{8}$$

It's an orthogonal basis for  $P_2(I_j)$ .

### 3.2 Weak form and finite element approximation

Now derive the weak form for (2). Choose  $v \in V_h^2$ , integration by parts on  $I_j$ , we get:

$$\int_{I_j} u_t v dx - \int_{I_j} f(u) v_x dx + f(u)v|_{x_{j+\frac{1}{2}}^-} - f(u)v|_{x_{j-\frac{1}{2}}^+}, \forall u \in V_h^2. \tag{9}$$

Since  $u$  maybe incontinuous on  $x_{j\pm\frac{1}{2}}$ , we use **numerical flux** to substitute  $f(u)|_{x_{j\pm\frac{1}{2}}}$ . In this assignment, we choose **Lax-Friedrichs flux**:

$$\hat{f}^{LF}(u^+, u^-) := \frac{1}{2} (f(u^-) + f(u^+) - \alpha(u^+ - u^-)), \quad \alpha = \max_u |f'(u)|. \tag{10}$$

Then, choose  $u_j := c_j^0(t)p_j^0 + c_j^1(t)p_j^1 + c_j^2(t)p_j^2$ ,  $v_{j,i} := p_j^i$ , (9) and (10) derive a semi-discretize numerical scheme.

The Lax-Friedrichs flux (10) is **consistent**, **Lipschitz continuous** and **monotone**, so this scheme satisfies discretize entropy inequality, see [2].

In fact, since  $(p_i^2, p_j^2)_{L^2(I_j)} = k\delta_{ij}$ , the mass matrix of (9) is just a diagonal matrix. So this scheme is efficient.

### 3.3 Time Integration

In this assignment, we use **TVD Runge-Kutta scheme** to solve the ODE system  $U_t = \mathcal{L}(U)$ :

$$\begin{aligned} U^{(1)} &= U^n + \Delta t \mathcal{L}(U^n), \\ U^{(2)} &= \frac{3}{4}U^n + \frac{1}{4}(U^{(1)} + \Delta t \mathcal{L}(U^{(1)})), \\ U^{n+1} &= \frac{1}{3}U^n + \frac{2}{3}(U^{(2)} + \Delta t \mathcal{L}(U^{(2)})). \end{aligned} \tag{11}$$

## 4 Numerical Result

In this section, we will give the numerical results.

### 4.1 Before break time

First, we test the  $L^1$ ,  $L^\infty$  error and convergence rates for this numerical scheme. We choose  $T = 0.2$ ,  $N = 20, 40, 80, 160, 320$  and  $k = 0.1h$ . Here are the results.

	20	rate	40	rate	80	rate	160	rate	320
$L^1$ norm	8.18e-4	2.82	1.16e-4	2.76	1.71e-5	2.75	2.53e-6	2.75	3.77e-7
$L^\infty$ norm	1.10e-3	2.60	1.81e-4	2.62	2.95e-5	2.66	4.69e-6	2.67	7.38e-7

Table 1: The error and convergence rate

In this test, we choose  $\alpha = 1.5$  for the Lax-Friedrichs flux, and use Newton's method to derive the solution of Burgers' equation. We write the modified equation

$$w = \sin(x - wt - \frac{t}{2}), \quad u = w + \frac{1}{2},$$

and set the initial value  $w = 1$  beyond the shock wave, i.e.  $x < \pi + 0.5t$ , the initial value  $w = -1$  after the shock wave, i.e.  $x > \pi + 0.5t$ .

### 4.2 After break time

In this section, we give some figures for  $T = 0, 0.5, 1, 1.5$  to show the continuous and incontinuous solutions.

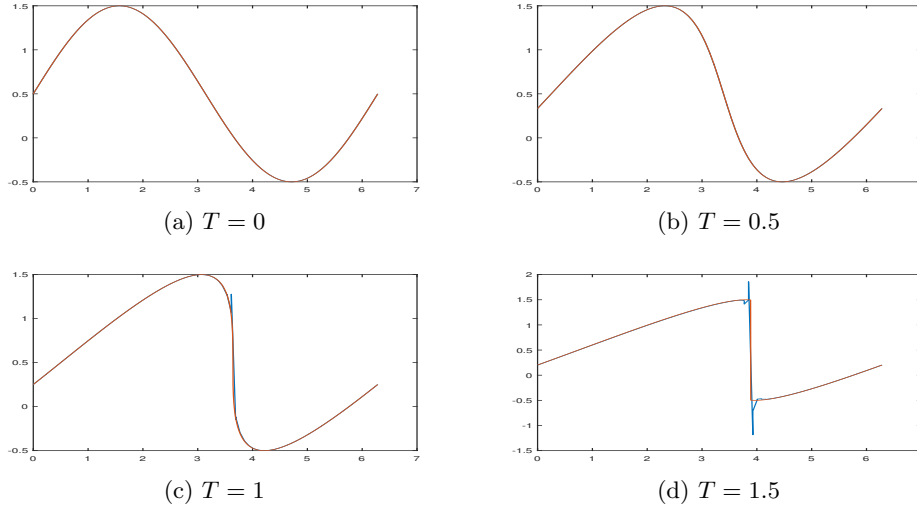


Figure 2: The numerical solution of Burgers equation

The red lines are the exact solutions, and the blue lines are the numerical solutions. Severe oscillations occur for  $T > 1$ , which means the  $P_2$  element DG method isn't TV (Total Variation) stable.

Now, we try to operate the limiter introduced in [2, Section 3.2.2]. After we operate the limiter, we eliminate the oscillations. It shows:

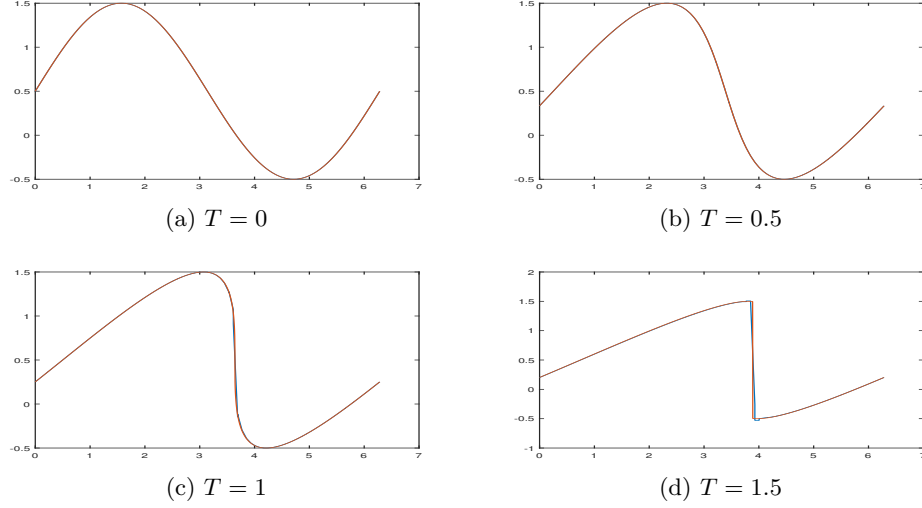


Figure 3: The numerical solution of Burgers equation after limiter

These figures show that the numerical solutions fit the actual solution of Burgers' equation well.

## 5 Discussions

Now, we make some discussions for the numerical results.

- By (10), we should choose  $\alpha = 1.5$ . But for  $\alpha = 0$  or  $\alpha = 1$ , the scheme also works, even gives better numerical result, why?
- How to make a prior estimation for the parameter  $\alpha$ ?
- In this assignment, I generate the initial value by interpolation method on  $x_{j-\frac{1}{2}}, x_{j+\frac{1}{2}}, \frac{x_{j-\frac{1}{2}} + x_{j+\frac{1}{2}}}{2}$ . If I try  $L^2$ -projection method to generate the initial value, things will different or not?

## References

- [1] Randall J. Leveque, Numerical Methods for Conservation Laws.
- [2] Chi-Wang Shu, Discontinuous Galerkin Methods: General Approach and Stability.