



Technische Universität München

School of Computation, Information and Technology



Master's Thesis

Residual-based Methods for Robust Spectral Analysis in Dynamical Systems

April Herwig

april.herwig@tum.de

Supervisor: Prof Oliver Junge

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I assure the single handed composition of this master's thesis only supported by declared resources.

Garching,

Abstract

Koopmania

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1 Introduction

Dynamical systems theory has seen many revolutions. One recent such revolution is to use operator theoretic tools to analyze the global behavior of chaotic systems. Concieved in the 1930's by Koopman and Von Neumann [7], the study of global dynamics under a mapping S is conducted by studying the action of the composition $\psi \circ S$ of S with an observable ψ . Observables which act as eigenvalues for the composition operator and its dual hold information on slowly decaying structures in phase space.

Currently the most popular method by far for analyzing composition operators (named Koopman operators) is dynamic mode decomposition (DMD). The past 2 decades have seen massive growth in this topic, and DMD has recieved many evolutions [3]. One primary concern of nearly all DMD methods is *spectral pollution*. The Koopman operator often has unfavorable spectral qualities such as continuous spectrum which are unstable, and are destroyed by finite approximation. The task presented in the present paper is to identify which candidate eigenvalues (computed e.g. by DMD) are approximations of eigenvalues of the true Koopman operator (or its adjoint), and which arise due to discretization.

2 Background

We consider a discrete dynamical system generated by a map $S : \Omega \rightarrow \Omega$ defined on a measure space $(\Omega, \mathcal{A}, dx)$ which is *nonsingular*, that is $\int_{S^{-1}(A)} dx = 0$ for any $A \in \mathcal{A}$ with $\int_A dx = 0$.

2.1 Koopman and Perron-Frobenius Operators

Koopman operator theory shifts the focus from dynamics of points in state space \mathcal{X} to dynamics of observables $g : \Omega \rightarrow \mathbb{C}$. A state $x \in \Omega$ evolves by iteratively applying the map S , and similarly observables evolve under the action of the *Koopman operator*

$$\mathcal{K}g = g \circ S. \tag{2.1}$$

The primary benefit of this alternative viewpoint of dynamics is that \mathcal{K} is *linear* and can therefore be analyzed using algebraic and functional analytic tools. However, the Koopman operator acts on function spaces which are typically infinite-dimensional. Even the Krylov space $\{g, \mathcal{K}g, \mathcal{K}^2g, \dots\}$ is generically infinite-dimensional: consider e.g. an indicator function $g = \mathbb{1}_{[0,1]}$ and a translation $S(x) = x - 1$.

It should be noted that until now we have not declared a function space to act as a domain for \mathcal{K} . This is because there are many potential domains for which \mathcal{K} is not only well-defined, but has interesting properties worth studying. For the current section we consider $\mathcal{K} : L^\infty \rightarrow L^\infty$. On this space \mathcal{K} is obviously bounded, in fact it is a contraction.

The Koopman operator tracks how observables evolve with S . There is however a dual viewpoint: that of *densities* and *pushforwards*. From the duality pairing

$$\langle f, \mathcal{K}g \rangle = \langle \mathcal{L}f, g \rangle \quad (2.2)$$

for $f \in L^1$, $g \in L^\infty$ we may deduce the form of an adjoint operator, known as the *transfer operator* or *Perron-Frobenius operator*. Consider an indicator $g = \mathbb{1}_A$ for an $A \in \mathcal{A}$ and let $f \geq 0$. The left hand side of 2.2 is

$$\int f(x) \mathbb{1}_A(S(x)) \, dx = \int f(x) \mathbb{1}_{S^{-1}(A)}(x) \, dx = \int_{S^{-1}(A)} f(x) \, dx, \quad (2.3)$$

and the right hand side

$$\int \mathcal{L}f(x) \mathbb{1}_A(x) \, dx = \int_A \mathcal{L}f(x) \, dx. \quad (2.4)$$

Hence $\mathcal{L}f \in L^1$ should satisfy the equation

$$\int_A \mathcal{L}f(x) \, dx = \int_{S^{-1}(A)} f(x) \, dx. \quad (2.5)$$

A short exercise in measure theory shows that

$$A \mapsto \int_{S^{-1}(A)} f(x) \, dx \quad (2.6)$$

defines a finite absolutely continuous measure, and hence by the Radon-Nikodym theorem the measure has a unique density which by 2.5 is precisely $\mathcal{L}f$. For general $f \in L^1$ decompose $f = f^+ - f^-$ with $f^+, f^- \geq 0$ and set $\mathcal{L}f = \mathcal{L}f^+ - \mathcal{L}f^-$.

2.6 shows that \mathcal{L} is precisely the action of the pushforward $S_\# \mu = \mu \circ S^{-1}$ for densities. This provides a useful intuition for the transfer operator: we move away from the viewpoint of single points $x \in \Omega$, and instead consider *distributions* of point density. If x is chosen from a random distribution with density f , then the distribution of $S(x)$ has density $\mathcal{L}f$.

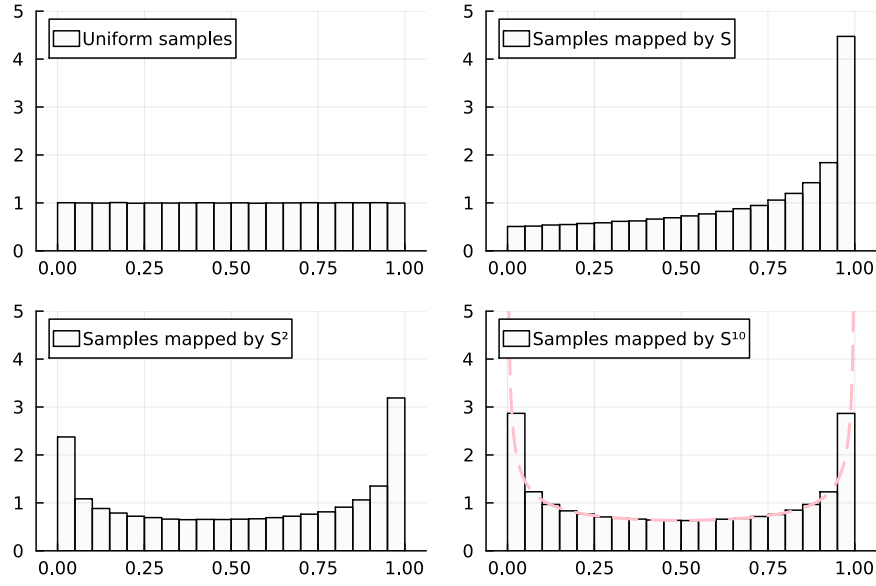


Figure 2.1: Illustration of the pushforward intuition of the Perron-Frobenius operator. Action of the \mathcal{L} associated with the quadratic map $S : [0, 1] \rightarrow [0, 1]$, $x \mapsto 4x(1 - x)$. One million points are sampled uniformly in $[0, 1]$ and mapped forward under S , S^2 , and S^{10} . The respective distributions are shown in histograms.

2.2 Spectral Properties

2.2.1 The Spectrum

We begin with abstract definitions for parts of the spectrum.

Definition 2.1. Let $M : D(M) \rightarrow \mathcal{X}$ be a linear operator on a Banach space $(X, \|\cdot\|)$. The domain $D(M) \subset \mathcal{X}$ of M may be a dense subset of \mathcal{X} . M is called closed if: whenever $x_k \rightarrow x$ is a convergent sequence with Mx_k convergent, then $Mx_k \rightarrow Mx$.

Closedness is weaker than continuity due to the extra assumption that Mx_k is convergent.

Definition 2.2. The spectrum $\sigma(M)$ of a closed operator M is the set of numbers $\lambda \in \mathbb{C}$ for which $M - \lambda I$ does not have a bounded inverse. The complement $\rho(M) := \mathbb{C} \setminus \sigma(M)$ is called resolvent set and $(M - \lambda I)^{-1}$ is called the resolvent.

Proposition 2.3. The spectrum admits a decomposition

$$\sigma(M) = \sigma_p(M) \oplus \sigma_c(M) \oplus \sigma_r(M) \quad (2.7)$$

into point-, continuous-, and residual-spectrum with

$$\sigma_p(M) = \{\lambda \in \mathbb{C} \mid M - \lambda I \text{ is not injective}\}, \quad (2.8)$$

$$\sigma_c(M) = \{\lambda \in \mathbb{C} \mid M - \lambda I \text{ is injective, but its range is a dense subset of } \mathcal{X}\}, \quad (2.9)$$

$$\sigma_r(M) = \{\lambda \in \mathbb{C} \mid M - \lambda I \text{ is injective, but does not have dense range}\}. \quad (2.10)$$

Remark. In some literature, 0 is never part of the spectrum. We do not take this approach.

An element $\lambda \in \sigma_p(M)$ is a *true* eigenvalue in the sense that there exists a vector $x \in \mathcal{X}$ such that $(M - \lambda I)x = 0$. The continuous spectrum has a similar characterization as the set of $\lambda \in \mathbb{C}$ for which $M_\lambda I$ is injective but *not* bounded from below, i.e. there exists a sequence $(x_k)_k$, $\|x_k\| = 1$, for which $(M - \lambda I)x_k \rightarrow 0$. Equivalently, $(M - \lambda I)^{-1}$ cannot be extended to a bounded linear operator, but is still a (densely defined) closed linear operator. (See the proof of theorem 2.16 for a proof of this characterization.)

One might ask why such detail is required in defining the spectrum. The answer lies in the subtle consequences of infinite-dimensionality. The following example from [6] shows the subtle interplay between residual spectrum and spectrum of the adjoint operator, all of which requires infinite-dimensionality.

Example 2.4 (Shift operators on $\ell^2(\mathbb{N})$). The canonical right-shift operator

$$R : (a_1, a_2, \dots) \mapsto (0, a_1, a_2, \dots) \quad (2.11)$$

is adjoint to the left shift

$$L : (a_1, a_2, \dots) \mapsto (a_2, a_3, \dots). \quad (2.12)$$

We claim:

$$\sigma_p(L) = \{\lambda \in \mathbb{C} \mid |\lambda| < 1\} \quad \sigma_p(R) = \emptyset \quad (2.13)$$

$$\sigma_c(L) = \{\lambda \in \mathbb{C} \mid |\lambda| = 1\} \quad \sigma_c(R) = \{\lambda \in \mathbb{C} \mid |\lambda| = 1\} \quad (2.14)$$

$$\sigma_r(L) = \emptyset \quad \sigma_r(R) = \{\lambda \in \mathbb{C} \mid |\lambda| < 1\}. \quad (2.15)$$

Indeed, to see 2.13 consider $|\lambda| < 1$, $a^\lambda = (\lambda, \lambda^2, \dots)$. Then clearly $La^\lambda = \lambda a^\lambda$. On the other hand, suppose $Ra = \lambda a$ for some $a \neq 0$, $\lambda \in \mathbb{C}$. Then letting n be the first index such that $a_n \neq 0$, we must have $0 = a_{n-1} = (Ra)_n = \lambda a_n$ so $\lambda = 0$. But the kernel of R is clearly $\{0\}$.

To see 2.14, take $|\lambda| = 1$ and construct *approximate* (L, λ) -eigenvectors $v^n = (\lambda, \lambda^2, \dots, \lambda^n, 0, \dots)$. Then $(L - \lambda I)v^n = (0, \dots, 0, \lambda^{n+1}, 0, \dots)$ so $\|(L - \lambda I)v^n\| = 1$ but $\|v^n\| = \sqrt{n}$. It remains to show that λ is of continuous spectrum, as opposed to residual spectrum. Indeed, for any w such that for all $v \in \ell^2$, $0 = \langle (L - \lambda I)v, w \rangle = \langle v, (R - \bar{\lambda} I)w \rangle$ so that w is an eigenvector of R which by 2.13 is a contradiction. An entirely analogous argument shows that the circle is also continuous spectrum for R .

We now reverse the discussion to see that any $|\lambda| < 1$ is in the residual spectrum of R . For $v \in \ell^2$ and w a $\bar{\lambda}$ -eigenvector of L we have $0 = \langle v, 0 \rangle = \langle v, (L - \bar{\lambda} I)w \rangle = \langle (R - \lambda I)v, w \rangle$ so that the range of $R - \lambda I$ is orthogonal to w .

Finally, $\sigma(M)$ is bounded by $\|M\|$ for any operator M , and clearly $\|L\| = 1$. But $\sigma_p(L) \cup \sigma_c(L) = \{\lambda \in \mathbb{C} \mid |\lambda| \leq 1\}$ so $\sigma_r(L) = \emptyset$. Combined with the previous paragraph, this shows 2.15.

The example already suggests some important relationships regarding the spectrum of adjoint operators.

Theorem 2.5. *Let M be a densely defined closed linear operator on a Hilbert space.*

1. $\lambda \in \sigma(M)$ iff $\bar{\lambda} \in \sigma(M^*)$.
2. If $\lambda \in \sigma_r(M)$, then $\bar{\lambda} \in \sigma_p(M^*)$.
3. Conversely, if $\lambda \in \sigma_p(M)$, then $\bar{\lambda} \in \sigma_p(M^*) \cup \sigma_r(M^*)$.
4. $\sigma_c(M) = \sigma_c(M^*)$.

Proof. 1. Suppose $M - \lambda I$ has a bounded inverse B . Then $(M - \lambda I)B = B(M - \lambda I) = I$. Equivalently $B^*(M^* - \bar{\lambda} I) = (M^* - \bar{\lambda} I)B^* = I^* = I$.

2. The range of $M - \lambda I$ is not dense in \mathcal{X} . Hence there exists a $v \perp \text{ran}(M - \lambda I)$. But this implies $v \in \ker(M^* - \bar{\lambda} I)$.

3. There exists a $v \in \ker(M - \lambda I)$ which implies $v \perp \text{ran}(M^* - \bar{\lambda} I)$.

4. This follows from 1, 2, and 3.

□

2.2.2 Spectra of Koopman and Perron-Frobenius Operators

The Koopman and Perron-Frobenius operator spectrum holds information about the long-term mixing rates of structures in phase space. We write (with some abuse of notation) $\mathcal{L}|_{\mathcal{X}}$ and $\mathcal{K}|_{\mathcal{X}}$ to denote the Perron-Frobenius / Koopman operators acting on the domain \mathcal{X} .

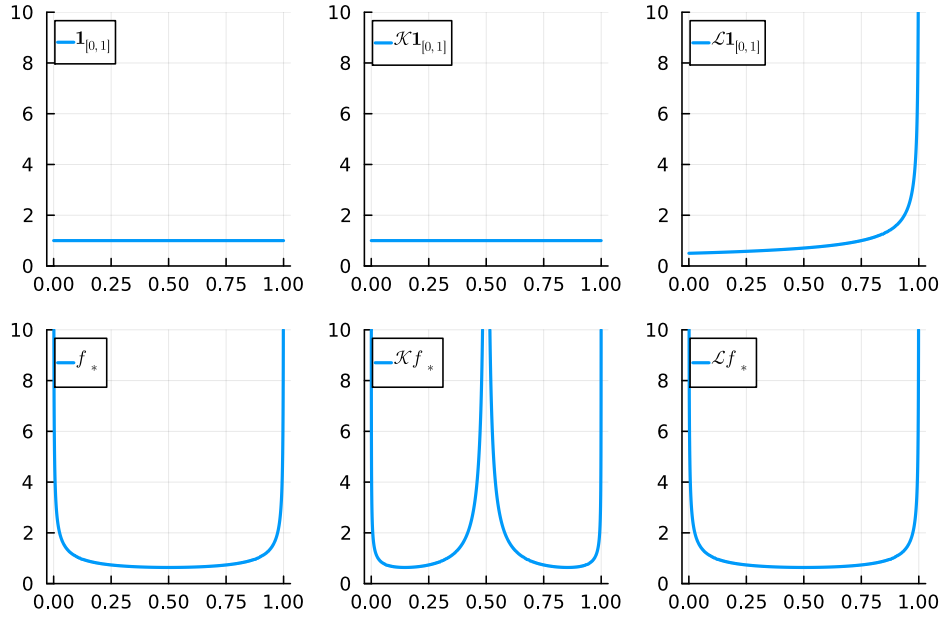


Figure 2.2: Actions of the Koopman and Perron-Frobenius operators associated with the quadratic map $S : [0, 1] \curvearrowright, x \mapsto 4x(1-x)$. The constant density $\mathbb{1}_{[0,1]}$ is fixed under \mathcal{K} and the density $f_*(x) = \frac{1}{\pi\sqrt{x(1-x)}}$ is fixed under \mathcal{L} (c.f. figure 2.1).

Definition 2.6. An eigenfunction f for $\mathcal{L}|_{L^1}$ with eigenvalue 1 is the density of a (signed) invariant measure. We say S preserves the measure $f dx$.

The following are well-known facts about the Koopman and Perron-Frobenius operator. They can be found in many sources e.g. [9].

- Theorem 2.7.**
1. If S is ergodic¹, then there is at most one invariant density. Conversely, if there is a unique invariant density which is $(dx-)$ almost everywhere positive, then S is ergodic.
 2. If S is ergodic, then every eigenvalue of $\mathcal{K}|_{L^1}$ is simple.
 3. Suppose S is invertible. Then S is weak-mixing iff 1 is the only eigenvalue of $\mathcal{K}|_{L^1}$.

The following theorem is from [5].

Definition 2.8. A set $A \subset \Omega$ in phase space is called δ -almost-invariant if

$$\int_{S^{-1}(A) \cap A} dx = (1 - \delta) \cdot \int_A dx. \quad (2.16)$$

¹Ergodicity and mixing describe how observations become decorrelated over time. S is ergodic if for all $A, B \in \mathcal{A}$ we have $\frac{1}{n} \sum_{j=0}^{n-1} \int_{S^{-j}(A) \cap B} dx \rightarrow \int_A dx \int_B dx$ as $n \rightarrow \infty$. S is weak-mixing if we have $\frac{1}{n} \sum_{j=0}^{n-1} \left| \int_{S^{-j}(A) \cap B} dx - \int_A dx \int_B dx \right| \rightarrow 0$ as $n \rightarrow \infty$.

Theorem 2.9. *Let $\mathbb{R} \ni \lambda < 1$ be an eigenvalue corresponding to a real-valued normalized eigenfunction f of $\mathcal{L}|_{L^1}$. Let further $A \subset \Omega$ be such that $\int_A f dx = \frac{1}{2}$. Then*

$$\delta + \eta = \lambda + 1 \quad (2.17)$$

if A is δ -almost-invariant and $\Omega \setminus A$ is η -almost-invariant.

These are only a few interesting properties of the spectrum of these operators, but there are many more which can be studied.

2.3 Pseudospectra

The reader will have likely noticed the sensitivity required in understanding the spectrum for infinite-dimensional operators. In particular the spectral types can be unstable w.r.t. perturbations of the operator. For example, for any self-adjoint operator M (e.g. the generator for the Koopman semigroup in a continuous-time dynamical system) there exists a compact operator E with arbitrarily small norm such that the perturbation $M + E$ has purely point spectrum.

The situation is even worse when one considers *perturbations of the dynamics* instead of perturbations of the Koopman / Perron-Frobenius operators. Consider a circle rotation $S : \mathbb{T} \rightarrow \mathbb{T}$, $e^{2\pi i \theta} \mapsto e^{2\pi i (\theta + \alpha)}$. We have $\sigma(\mathcal{L}|_{L^2}) = \sigma_p(\mathcal{L}|_{L^2}) = \alpha^{\mathbb{N}_0}$. If the rotation is rational then the spectrum is discrete, but if the rotation is irrational then the spectrum is dense in the unit circle.

In the present paper we tackle the instability problem for operator perturbations. One might wonder whether operators which are "close" (in some sense) also have spectra which are "close". Unfortunately, this is in general not true. Even in finite dimensions this breaks down, as the following example shows.

Example 2.10 ([12]). Let $\mathcal{X} = \mathbb{R}^n$ and let

$$M = \begin{bmatrix} 0 & 1 & & \\ & \ddots & \ddots & \\ & & & 1 \\ & & & 0 \end{bmatrix}. \quad (2.18)$$

M is nilpotent so $\sigma(M) = \{0\}$. However, for an arbitrarily small $\epsilon > 0$, the perturbation

$$M + E = \begin{bmatrix} 0 & 1 & & \\ & \ddots & \ddots & \\ & & & 1 \\ \epsilon & & & 0 \end{bmatrix}. \quad (2.19)$$

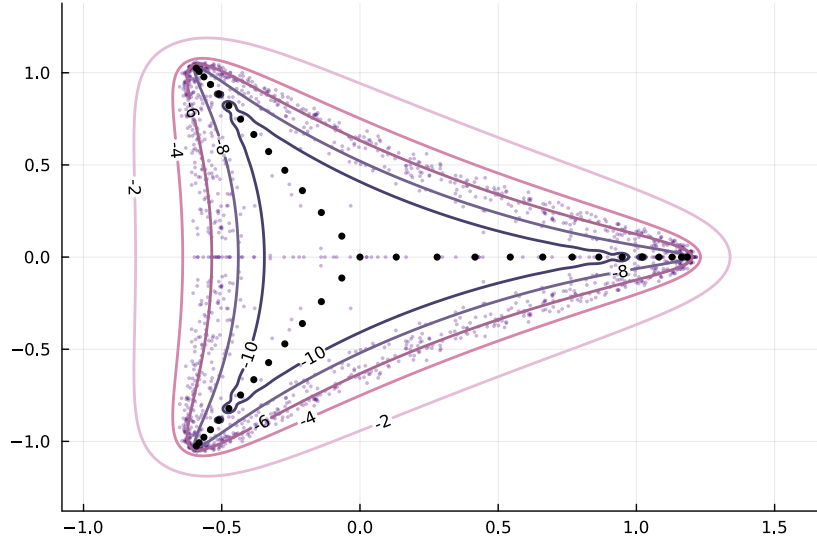


Figure 2.3: Pseudospectra for the 40×40 Toeplitz matrix T with 0 on the main diagonal, 1 on the first upper off-diagonal, and $1/4$ on the second lower off-diagonal. Contour lines show the boundary of the ϵ -pseudospectrum for $\epsilon = 10^{-2}, 10^{-4}, 10^{-6}, 10^{-8}, 10^{-10}$. Black: Spectrum of T . Purple: Spectra of $T + E$ for 40 randomly sampled perturbations E with $\|E\| < 10^{-4}$.

has characteristic polynomial $(-\lambda)^n - (-1)^n \epsilon$ so that $\sigma(M+E) = \{\lambda \in \mathbb{C} \mid \lambda^n = \epsilon\}$. For growing n , $\sigma(M + E)$ comes asymptotically close to filling the unit circle.

2.3.1 Definitions of the Pseudospectrum

Definition 2.11. Let $M : D(M) \rightarrow \mathcal{X}$ be a closed linear operator. The ϵ -pseudospectrum of M is the smallest set in \mathbb{C} which contains the spectrum of all perturbations of M with norm less than ϵ .²

$$\sigma_\epsilon(M) = \bigcup_{\|E\| < \epsilon} \sigma(M + E). \quad (2.20)$$

Theorem 2.12. We have the following equivalent formulation of the pseudospectrum:

$$\sigma_\epsilon(M) = \left\{ \lambda \in \mathbb{C} \mid \|(M - \lambda I)^{-1}\| > \frac{1}{\epsilon} \right\} \quad (2.21)$$

where we use the convention that $\|(M - \lambda I)^{-1}\| = \infty$ if $M - \lambda I$ is not invertible.

Remark. While this result is practically pseudospectrum folklore, a correct proof in the general case of closed linear operators on Banach spaces is difficult to find. Virtually all sources which cite this result, cite an unpublished

²Some authors will define the pseudospectrum with " \leq " instead of " $>$ ". This makes $\sigma_\epsilon(M)$ a closed set, but breaks theorem 2.12.

technical report by Chaitin-Chatelin and Harrabi [2]. The following proof is based on [1] and extended slightly from bounded operators to closed (potentially) unbounded operators.

Proof. We prove " \subset " by contraposition: assume $\|(M - \lambda I)^{-1}\| > \frac{1}{\epsilon}$ and let $\|E\| < \epsilon$. Then $\|(M - \lambda I)^{-1}E\| < 1$ and hence $I + (M - \lambda I)^{-1}E$ is invertible. This implies

$$M - \lambda I + E = (M - \lambda I) (I + (M - \lambda I)^{-1}E) \quad (2.22)$$

is invertible.

Conversely, we prove " \supset " by showing there exists an operator E with $\|E\| < \epsilon$ such that $M - \lambda I + E$ is not invertible. Since $\|(M - \lambda I)^{-1}\| > \frac{1}{\epsilon}$ there exists a $u \in \mathcal{X}$ with $\|u\| = 1$ and $(M - \lambda I)^{-1}u = v \in D(M)$ with $\|v\| = \frac{1}{\delta} > \frac{1}{\epsilon}$.³ The Hahn-Banach theorem provides a $v^* \in \mathcal{X}^*$ with $\|v^*\| = 1$ and $v^*v = \|v\| = \frac{1}{\delta}$. Set $E = -\delta uv^*$. Then $\|E\| = \delta < \epsilon$ and

$$Ev = -\delta uv^*v = -u = -(M - \lambda I)(M - \lambda I)^{-1}u = -(M - \lambda I)v. \quad (2.23)$$

Rearranging yields $(M + E - \lambda I)v = 0$. \square

The proof of the above theorem also shows that one can consider only rank-one perturbations in 2.20:

Corollary 2.13. *We can restrict 2.20 to*

$$\sigma_\epsilon(M) = \bigcup_{\substack{\|E\| < \epsilon \\ \text{rank } E = 1}} \sigma(M + E). \quad (2.24)$$

The following formulation of the pseudospectrum in the Hilbert space case will be the main tool we use in the section on numerical methods.

Definition 2.14. *The function*

$$\sigma_{\text{inf}}(M) = \inf_{\|x\|=1} \|Mx\| \quad (2.25)$$

is known as the injection modulus.

Lemma 2.15. *Let $M : D(M) \rightarrow \mathcal{X}$ be a closed linear operator on a Hilbert space. If $\lambda \in \rho(M)$ then*

$$\sigma_{\text{inf}}(M - \lambda I) = \sigma_{\text{inf}}(M^* - \bar{\lambda}I), \quad (2.26)$$

but this is not necessarily true if $\lambda \in \sigma(M)$.

³At this point we require the strict inequality. Otherwise, the existence of such a pair u, v is not guaranteed.

Proof. 2.26 for $\lambda \in \rho(M)$ follows from the fact that for bounded operators A we have $\|A\| = \|A^*\|$, applied to $A = (M - \lambda I)^{-1}$. To see that 2.26 does not hold for $\lambda \in \sigma(M)$, consider again the right shift from example 2.4. Clearly $\|Rx\| = \|x\|$ so R is bounded from below but 0 is part of the spectrum. \square

Theorem 2.16. *For a Hilbert space operator M ,*

$$\frac{1}{\|(M - \lambda I)^{-1}\|} = \min \{ \sigma_{\inf}(M - \lambda I), \sigma_{\inf}(M^* - \bar{\lambda} I) \}. \quad (2.27)$$

where we use the convention that $1/\|(M - \lambda I)^{-1}\| = 0$ when $\lambda \in \sigma(M)$.

Proof. When $\lambda \in \sigma_p(M)$, then obviously $\sigma_{\inf}(M - \lambda I) = 0$. Due to theorem 2.5 part 2, $\lambda \in \sigma_r(M)$ has $\sigma_{\inf}(M^* - \bar{\lambda} I) = 0$. Finally, let $\lambda \in \sigma_c(M)$ so $M - \lambda I$ is injective. Assume $\sigma_{\inf}(M - \lambda I) > 0$. Then $(M - \lambda I)^{-1} : \text{ran}(M - \lambda I) \rightarrow D(M)$ is bounded. But since $\text{ran}(M - \lambda I) \subset \mathcal{X}$ is dense, $(M - \lambda I)^{-1}$ has a unique bounded extension to \mathcal{X} , a contradiction. \square

Corollary 2.17. *In the Hilbert space setting, the pseudospectrum can be formulated as:*

$$\sigma_{\epsilon}(M) = \{ \lambda \in \mathbb{C} \mid \exists u \in \mathcal{X} : \|u\| = 1 \text{ and either } \|(M - \lambda I)u\| < \epsilon \text{ or } \|(M^* - \bar{\lambda} I)u\| < \epsilon \}. \quad (2.28)$$

In the general Banach space setting,

$$\sigma_{\epsilon}(M) = \sigma(M) \cup \{ \lambda \in \mathbb{C} \mid \exists u \in \mathcal{X} : \|u\| = 1, \|(M - \lambda I)u\| < \epsilon \}. \quad (2.29)$$

One sees from the proofs of lemma 2.15 and theorem 2.16 that $\sigma(M)$ in equation 2.29 can be replaced with $\sigma_r(M)$. The residual spectrum causes many common and unintuitive issues in spectral theory. In the study of pseudospectra, this issue is alleviated in one of three ways: (1): consider only classes of operators which have no residual spectrum (e.g. finite-dimensional or compact operators), (2): compute a slightly different object known as the *approximate-point spectrum* (see below), (3): ignore the problem entirely and write proofs which are all slightly incorrect.

Definition 2.18. *The approximate-point pseudospectrum is the set of points for which there exists an ϵ -approximate pseudoeigenvector*

$$\sigma_{ap,\epsilon}(M) = \{ \lambda \in \mathbb{C} \mid \exists u \in \mathcal{X} : \|u\| = 1, \|(M - \lambda I)u\| < \epsilon \}. \quad (2.30)$$

The $\epsilon \rightarrow 0$ limit of such sets is the approximate-point spectrum

$$\sigma_{ap}(M) = \bigcap_{\epsilon > 0} \sigma_{ap,\epsilon}(M) \quad (2.31)$$

$$= \left\{ \lambda \in \mathbb{C} \mid \exists (u_k)_k : \|u_k\| = 1 \ \forall k, \lim_{k \rightarrow \infty} \|(M - \lambda I)u_k\| = 0 \right\}. \quad (2.32)$$

2.3.2 Properties

To gain some intuition for the pseudospectrum, we derive some properties for specific types of operators.

Definition 2.19. *An operator M on a Hilbert space is called normal if $M^*M = MM^*$. Special classes of normal operator include unitary operators ($M^* = M^{-1}$) and self-adjoint operators ($M^* = M$).*

Proposition 2.20. *For a normal operator M , $\sigma_r(M) = \emptyset$.*

Proof. Let $\lambda \in \sigma(M)$, $A = M - \lambda I$. If λ is not an eigenvalue then $\ker A = \{0\}$ and

$$\{0\} = \ker A = \ker A^* = (\operatorname{ran} A)^\perp \quad (2.33)$$

where for the second equality we used the normality. Hence $\operatorname{ran} A$ is dense. \square

Lemma 2.21. *For a closed linear operator M we have*

$$\operatorname{dist}(\lambda, \sigma(M)) \geq \frac{1}{\|(M - \lambda I)^{-1}\|}. \quad (2.34)$$

Moreover, if M is a normal operator on a Hilbert space, then we have equality.

Proof. Let $\eta \in \sigma(M)$. Denote $A = M - \lambda I$, $B = M - \eta I$. Then A is invertible and

$$B = A - (A - B) = A(I - A^{-1}(A - B)). \quad (2.35)$$

Now $\|A^{-1}(A - B)\| \leq \|A^{-1}\| |\lambda - \eta|$ so if $|\lambda - \eta| < \frac{1}{\|A^{-1}\|}$ then the right hand side of 2.35 is invertible, a contradiction to $\eta \in \sigma(M)$.

To prove equality in the case of normal operators requires some heavy machinery from operator theory, namely:

1. Every normal operator is unitarily similar to a multiplication operator on some semi-finite measure space (i.e. there exists a unitary U such that $M = U^{-1}M_a U$ where $(M_a f)(x) = a(x)f(x)$ is a multiplication operator).
2. For multiplication operators, equation 2.34 holds with equality.

These results can be found in many sources e.g. [11]. \square

Theorem 2.22. *We have*

1. $\sigma_\delta(M) \subset \sigma_\epsilon(M)$ whenever $\delta \leq \epsilon$,
2. $\bigcap_{\epsilon > 0} \sigma_\epsilon(M) = \sigma(M)$,
3. $\sigma(M) + B_\epsilon \subset \sigma_\epsilon(M)$, where B_ϵ is the ball with radius ϵ centered at 0 and $+$ refers to pointwise summation $C + D = \{x + y \mid x \in C, y \in D\}$.
4. When M is normal, $\sigma(M) + B_\epsilon = \sigma_\epsilon(M)$.

Remark. Theorem 2.22 effectively says that the Pseudospectra are nested sets which grow *at least* as quickly as ϵ -balls around the spectrum.

Proof. 1 and 2 are consequences of the representation 2.20 of the pseudospectrum. 3 and 4 are consequences of lemma 2.21. \square

Theorem 2.23. *The function which maps $(\epsilon, M) \mapsto \sigma_\epsilon(M)$ which sends an $\epsilon > 0$ and a bounded M acting on a Hilbert space \mathcal{X} to the set $\sigma_\epsilon(M)$ is continuous using the metric*

$$d((\epsilon_1, M_1), (\epsilon_2, M_2)) = |\epsilon_1 - \epsilon_2| + \|M_1 - M_2\| \quad (2.36)$$

in the domain and the Hausdorff metric in the codomain.

Remark. 1. The idea for the proof is taken from [4], though it originally included two errors: it does not take into account residual spectrum, and makes a bound which is unfortunately not correct. These issues can be lifted in the Hilbert space setting, so that the theorem is still correct as it is stated in [4].

2. If one replaces $\sigma_\epsilon(M)$ with $\sigma_{ap,\epsilon}(M)$ in theorem 2.23, then the proof in [4] is correct and even works for general Banach spaces. To the best of this author's knowledge, there is no published correct proof for the general Banach space setting with $\sigma_\epsilon(M)$, despite it being considered common knowledge in the study of pseudospectra.

Proof. Let (M, ϵ) and (M', ϵ') be such that $d((M, \epsilon), (M', \epsilon')) < \delta$ for some $0 < \delta < \epsilon/2$. Without loss of generality let $\epsilon \leq \epsilon'$.

Let $\lambda \in \sigma_{\epsilon'}(M')$. By 2.28 there exists a normalized $v \in \mathcal{X}$ such that either $\|(M' - \lambda I)v\| < \epsilon'$ or $\|(M'^* - \bar{\lambda}I)v\| < \epsilon'$. Assume the former is true. It follows

$$\|(M - \lambda I)v\| \leq \|(M - M')v\| + \|(M' - \lambda I)v\| < \delta - |\epsilon - \epsilon'| + \epsilon' \leq \delta + \epsilon \quad (2.37)$$

where for the last inequality we used that $\epsilon \geq \epsilon' - |\epsilon' - \epsilon|$. By 2.28, $\lambda \in \sigma_{\epsilon+\delta}(M)$ and hence $\sigma_{\epsilon'}(M') \subset \sigma_{\epsilon+\delta}(M)$.

For the case that $\|(M'^* - \bar{\lambda}I)v\| < \epsilon'$, equation 2.37 can be done exactly the same with $\|(M^* - \bar{\lambda}I)v\|$ on the left hand side.

Now let $\lambda \in \sigma_{\epsilon-\delta}(M)$. Again this implies there exists a normalized $v \in \mathcal{X}$ such that either $\|(M - \lambda I)v\| < \epsilon - \delta$ or $\|(M^* - \bar{\lambda}I)v\| < \epsilon - \delta$. Assume again the former is true. Then

$$\|(M' - \lambda I)v\| \leq \|(M' - M)v\| + \|(M - \lambda I)v\| < \delta + \epsilon - \delta \leq \epsilon' \quad (2.38)$$

and hence $\lambda \in \sigma_{\epsilon'}(M')$ so that $\sigma_{\epsilon-\delta}(M) \subset \sigma_{\epsilon'}(M')$ (again the adjoint case can be done exactly the same with adjoints on the left hand side of the equation).

We now have $\sigma_{\epsilon-\delta}(M) \subset \sigma_{\epsilon'}(M') \subset \sigma_{\epsilon+\delta}(M)$ which implies for the Hausdorff distance between $\sigma_{\epsilon}(M)$ and $\sigma_{\epsilon'}(M')$,

$$\mathcal{H}(\sigma_{\epsilon}(M), \sigma_{\epsilon'}(M')) \quad (2.39)$$

$$\leq \max \{ \mathcal{H}(\sigma_{\epsilon}(M), \sigma_{\epsilon-\delta}(M)), \mathcal{H}(\sigma_{\epsilon}(M), \sigma_{\epsilon+\delta}(M)) \} \quad (2.40)$$

$$\leq \mathcal{H}(\sigma_{\epsilon-\delta}(M), \sigma_{\epsilon+\delta}(M)) \quad (2.41)$$

so we must investigate the difference between $\sigma_{\epsilon-\delta}(M)$ and $\sigma_{\epsilon+\delta}(M)$. But

$$\sigma_{\epsilon+\delta}(M) \setminus \sigma_{\epsilon-\delta}(M) = \left\{ \lambda \in \mathbb{C} \mid \frac{1}{\epsilon+\delta} < \|(M - \lambda I)^{-1}\| < \frac{1}{\epsilon-\delta} \right\} \quad (2.42)$$

which is just a difference of sublevel sets for the continuous function $\lambda \mapsto \|(M - \lambda I)^{-1}\|$. Hence 2.42 collapses continuously to \emptyset as $\delta \rightarrow 0$, which implies 2.41 converges to 0 as $\delta \rightarrow 0$. \square

With an understanding of the spectral properties of Koopman and Perron-Frobenius operators as well as the pseudospectrum, we are prepared to apply pseudospectral methods to some of the most common Koopman approximation methods, known under the umbrella term *dynamic mode decomposition*.

3 Numerical Methods

3.1 Petrov-Galerkin Methods

The original Ritz-Galerkin method is described as follows: we are given a PDE problem in its *weak formulation*:

$$\text{find } u \in \mathcal{X} \text{ such that } q(v, u) = \langle f, v \rangle \quad \forall v \in \mathcal{W} \quad (3.1)$$

where q is some elliptic sesquilinear form, $f \in \mathcal{X}$ given, $\langle \cdot, \cdot \rangle$ some inner product, and \mathcal{W} some *test set*. The form q is typically derived from some minimization problem for a functional representing energy.

The idea behind the Ritz-Galerkin method is to *solve 3.1 on a finite-dimensional subspace*: let $\mathcal{W} = \text{span} \{ \psi_1, \dots, \psi_N \}$ be q -linearly independent.

Then writing

$$\Psi(x) = [\psi_1(x) \mid \dots \mid \psi_N(x)], \quad (3.2)$$

we make the approximation $u \approx \Psi c_u$ and have $v = \Psi c_v$ for $c_u, c_v \in \mathbb{C}^N$. Now 3.1 reduces to

$$\text{find } c_u \in \mathbb{C}^N \text{ such that } q(\Psi c_v, \Psi c_u) = \langle f, \Psi c_v \rangle \quad \forall c_v \in \mathbb{C}^N. \quad (3.3)$$

One quickly verifies that c_u is the unique solution to the matrix equation $Ax = b$ with

$$A_{ij} = q(\Psi e_i, \Psi e_j) = q(\psi_i, \psi_j), \quad b_i = \langle f, \Psi e_i \rangle = \langle f, \psi_i \rangle. \quad (3.4)$$

where $e_i \in \mathbb{C}^N$ is the i -th standard unit vector. A is known as the *stiffness matrix*.

One can extend the Ritz-Galerkin formulation by allowing the test space to differ from the basis space: keep $u \approx \Psi c_u$ but let $\mathcal{W} = \text{span}\{\psi_1^*, \dots, \psi_M^*\}$, then A and b become

$$A_{ij} = q(\psi_i^*, \psi_j), \quad b_i = \langle f, \psi_i^* \rangle. \quad (3.5)$$

In the regime $M > N$ this equation is *overdetermined* so it is solved in a least squares sense

$$\|Ax - b\|_2^2 = \min_{x \in \mathbb{C}^N}! \quad (3.6)$$

where $\|\cdot\|_2$ is the vector ℓ^2 norm. This is known as the Petrov-Galerkin method.

One can extend the method further by asking that multiple solutions for multiple right hand sides $(f_j)_{j=1}^N$ are computed simultaneously, that is

$$\|AX - B\|_F^2 = \min_{X \in \mathbb{C}^{N \times N}}!, \quad B_{ij} = \langle f_j, \psi_i^* \rangle \quad (3.7)$$

where $\|\cdot\|_F$ denotes the Frobenius norm.

Numerically, 3.7 can be solved using the Moore-Penrose inverse

$$X = A^\dagger B. \quad (3.8)$$

An exercise in matrix calculus shows that the solution can also be written in the form

$$X = (A^* A)^{-1} A^* B. \quad (3.9)$$

Both forms will prove to be useful later.

3.2 Extended Dynamic Mode Decomposition (EDMD)

3.2.1 The Galerkin Ansatz

We apply the Petrov-Galerkin Ansatz to obtain a matrix approximation K for $\mathcal{K}|_{\mathcal{X}}$. Let $q(\cdot, \cdot) = \langle \cdot, \cdot \rangle$ and consider a linearly independent family of functions $\{\psi_1, \dots, \psi_N\} \subset \mathcal{X}$. Take delta distributions, that is $\langle \delta_x, \psi \rangle = \psi(x)$, for the test function(s): let $(w_i, x_i)_{i=1}^M$ represent a quadrature scheme and set $\psi_i^* = \sqrt{w_i} \delta_{x_i}$.

We then solve the Galerkin equation 3.7 for $f_j = \mathcal{K}\psi_j$:

$$\left\| \sqrt{W} \Psi_X K - \sqrt{W} \Psi_Y \right\|_F^2 = \min_{K \in \mathbb{C}^{N \times N}}! \quad (3.10)$$

with $W = \text{diag}(w_1, \dots, w_M)$ and

$$(\Psi_X)_{ij} = \langle \delta_{x_i}, \psi_j \rangle = \psi_j(x_i), \quad (\Psi_Y)_{ij} = \langle \delta_{x_i}, \mathcal{K}\psi_j \rangle = \psi_j(S(x_i)). \quad (3.11)$$

This results in the *EDMD matrix* $K = \Psi_X^\dagger \Psi_Y = (\Psi_X^* W \Psi_X)^{-1} (\Psi_X^* W \Psi_Y)$. Inverting Ψ_X involves computing the Moore-Penrose inverse of an $N \times M$ matrix, whereas inverting $\Psi_X^* \Psi_X$ involves computing the inverse of a symmetric $N \times N$ matrix. Depending on the relationship between M and N in the particular usecase, either formulation might be cheaper.

Another look at the second formulation of the EDMD matrix shows that

$$G_{ij} := (\Psi_X^* W \Psi_X)_{ij} = \sum_{k=1}^M w_k \overline{\psi_i(x_k)} \psi_j(x_k) \xrightarrow{M \rightarrow \infty} \langle \psi_i, \psi_j \rangle =: \mathbb{G}_{ij}, \quad (3.12)$$

$$A_{ij} := (\Psi_X^* W \Psi_Y)_{ij} = \sum_{k=1}^M w_k \overline{\psi_i(x_k)} \psi_j(S(x_k)) \xrightarrow{M \rightarrow \infty} \langle \psi_i, \mathcal{K} \psi_j \rangle =: \mathbb{A}_{ij}, \quad (3.13)$$

which can be computed with constant memory requirement.

3.2.2 Functional Minimization

We could have arrived at equation 3.10 completely differently: equation 3.10 is precisely a quadrature approximation of the functional least squares minimization:

$$\|\Psi K - \mathcal{K} \Psi\|_{\mathcal{X}^{1 \times N}}^2 = \min_{K \in \mathbb{C}^{N \times N}} \quad (3.14)$$

where $\mathcal{K} \Psi$ is understood elementwise and $\mathcal{X}^{1 \times N}$ is the space of (row) vector-valued functions with each component in \mathcal{X}

$$\|[f_1 \mid \dots \mid f_N]\|_{\mathcal{X}^{1 \times N}}^2 = \|f_1\|_{\mathcal{X}}^2 + \dots + \|f_N\|_{\mathcal{X}}^2. \quad (3.15)$$

Let us investigate $\Psi : \mathbb{C}^N \rightarrow \mathcal{X}$ a bit closer. Decompose the Hilbert space \mathcal{X} into

$$\mathcal{X} = \text{span} \{\psi_j\}_{j=1}^N \oplus \mathcal{V}, \quad \mathcal{V} = \left(\{\psi_j\}_{j=1}^N \right)^\perp \quad (3.16)$$

and let $\{\psi_j\}_{j=N+1}^\infty$ be a basis of \mathcal{V} .

A short calculation using the orthogonality of $\{\psi_j\}_{j=1}^N$ and $\{\psi_j\}_{j=N+1}^\infty$ shows that for $\phi \in \mathcal{X}$,

$$\langle \phi, \Psi c \rangle = \sum_{j=1}^N \langle \phi, \psi_j \rangle c_j. \quad (3.17)$$

Hence the adjoint quasi-matrix $\Psi^* : \mathcal{X} \rightarrow \mathbb{C}^N$ acts as

$$\Psi^* \phi = \begin{pmatrix} \langle \phi, \psi_1 \rangle \\ \vdots \\ \langle \phi, \psi_1 \rangle \end{pmatrix}. \quad (3.18)$$

Lemma 3.1. *For an arbitrary basis $\{\psi_j\}_{j=1}^N$, the orthogonal projector P onto the space spanned by the basis is*

$$P = \Psi(\Psi^*\Psi)^{-1}\Psi^*. \quad (3.19)$$

Proof. The solution of a linear least squares problem $\|Ax - b\| = \min!$ is the result of orthogonally projecting b onto the range of A , that is, $Ax = Pb$. From equation 3.9 we know $x = (A^*A)^{-1}A^*b$. Therefore $Pb = Ax = A(A^*A)^{-1}A^*b$. Since this holds for arbitrary b , the result is proven. \square

Using equation 3.9 we see that in the infinite-data limit $M \rightarrow \infty$, K has the form

$$K = (\Psi^*\Psi)^{-1}\Psi^*(\mathcal{K}\Psi) \quad (3.20)$$

Inserting the definitions of Ψ and Ψ^* yields $K = \mathbb{G}^{-1}\mathbb{A}$ exactly as in equations 3.12 and 3.13.

If we view the result of applying K as in equation 3.20 to a vector c as an object in \mathcal{X} , that is, ΨKc , we see that

$$\Psi Kc = \Psi(\Psi^*\Psi)^{-1}\Psi^*(\mathcal{K}\Psi)c = PK\Psi c = PKP\Psi c. \quad (3.21)$$

Since this holds for arbitrary c we have proven that (when viewed as an operator on \mathcal{X}) K is precisely the action of PKP , the finite section of \mathcal{K} over span $\{\psi_j\}_{j=1}^N$. In this way, EDMD can just as well be viewed as a finite section method.

Algorithm 1 Extended Dynamic Mode Decomposition (EDMD)

Require: Dictionary $\{\psi_j\}_{j=1}^N$, data points and weights $\{(w_i, x_i)\}_{i=1}^M$

- 1: Construct G , A as in 3.12, 3.13
 - 2: Set $K = G^{-1}A$ (or $L = G^{-1}A^*$)
 - 3: Compute an eigendecomposition $KV = V\Lambda$ (or $LV = V\Lambda$)
 - 4: **return** Eigenvalues and eigenmodes Λ , V
-

3.2.3 EDMD for the Perron-Frobenius Operator

From equation 3.20, and noting the form of \mathbb{A} ,

$$\langle \psi_i, \mathcal{L}\psi_j \rangle = \langle \mathcal{K}\psi_i, \psi_j \rangle = \overline{\langle \psi_j, \mathcal{K}\psi_i \rangle} = \overline{\mathbb{A}_{ji}} \quad (3.22)$$

yields an equivalent Galerkin method for \mathcal{L} :

$$L = (\Psi^*\Psi)^{-1}\Psi^*(\mathcal{L}\Psi) = \mathbb{G}^{-1}\mathbb{A}^* \quad (3.23)$$

or the finite analogue:

$$L = G^{-1}A^*. \quad (3.24)$$

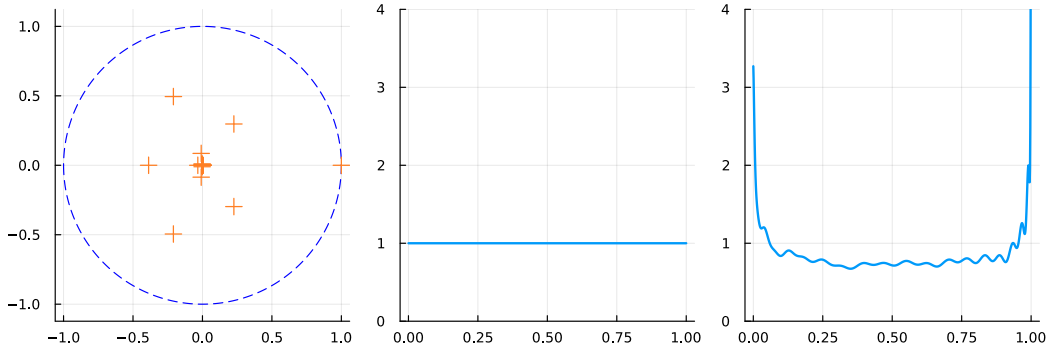


Figure 3.1: Algorithm 1 applied to the quadratic map (c.f. figure 2.1) performed with $M = 100$ Gauß-Legendre quadrature nodes and weights, and $N = 40$ Legendre polynomials transplanted to the interval $[0, 1]$. Left: spectrum of K . Middle: (normalized) eigenfunction of K for the eigenvalue $\lambda = 1$. Right: (normalized) eigenfunction of L for the eigenvalue $\lambda = 1$. Compare with figure 2.2.

3.3 Residual EDMD (ResDMD)

The formulation 3.21 shows that K (when viewed as a operator on \mathcal{X}) converges weakly to \mathcal{K} . However, from example 2.10 we know that the spectrum is already unstable for operators which are close in the *strong* operator topology, let alone in the *weak* topology. It is therefore entirely unclear *a priori* that the eigenvalues and eigenvectors of K actually represent eigenvalues and eigenfunctions of \mathcal{K} .

In the Perron-Frobenius operator community, a common solution to the mentioned issue is *stochastic smoothing*. Instead of considering a deterministic dynamical system generated by S , one consider a stochastic system: $x \in \Omega$ is assigned a *distribution* of possible image points instead of being assigned to the point $S(x)$. The resulting Markov process has an associated (stochastic) Perron-Frobenius operator which (under some conditions on the type of stochastic smoothing) is Hilbert-Schmidt on $L^2(\Omega)$. This way, the finite sections converge strongly to the (stochastic) Perron-Frobenius and (since it is Hilbert-Schmidt) so do the eigenvalues [5].

We take a different approach using the theory of pseudospectra. We wish to know which of our eigenvalues are *spurious*, that is, caused by the reduction to a finite section, and which eigenvalues are accurate. To determine this, we consider equation 2.28: if a sequence $(\lambda_N, c_N)_N$ of (normalized) eigenpairs for $K = K(\{\psi_j\}_{j=1}^N)$ converges to a true $\lambda \in \sigma(\mathcal{K})$ then we must have

$$\lim_{N \rightarrow \infty} \|(\mathcal{K} - \lambda_N I)\Psi c_N\|_{\mathcal{X}}^2 = 0 \quad (3.25)$$

(where $\Psi = \Psi_N$ is as in 3.2) or

$$\lim_{N \rightarrow \infty} \|(\mathcal{L} - \lambda_N I) \Psi c_N\|_{\mathcal{X}}^2 = 0. \quad (3.26)$$

Conversely, if neither of these tend to 0 as N grows, then we can rule out λ_N as a candidate eigenvalue.

From section 3.2.2 we know that the Galerkin equation 3.10 is a quadrature approximation of equation 3.14. Similarly, the *regression error*

$$\text{res}(\lambda_N, c_N) := \left\| \sqrt{W}(\Psi_Y - \lambda_N \Psi_X) c_N \right\|_2^2 \quad (3.27)$$

is precisely a quadrature approximation of equation 3.25.

Proposition 3.2. *Let λ and $g = \Psi c$ be a candidate eigenpair for \mathcal{K} . Then*

$$\lim_{M \rightarrow \infty} \text{res}(\lambda, c) = \|(\mathcal{K} - \lambda I)g\|_{\mathcal{X}}^2 \quad (3.28)$$

Proof. Denote $J = \Psi_Y^* W \Psi_Y$ and observe that

$$\lim_{M \rightarrow \infty} J_{ij} = \langle \mathcal{K} \psi_i, \mathcal{K} \psi_j \rangle =: \mathbb{J}_{ij}. \quad (3.29)$$

Consider the action of Ψ on standard unit vectors:

$$\langle \Psi e_i, \Psi e_j \rangle = \langle \psi_i, \psi_j \rangle = \mathbb{G}_{ij} = e_i^* \mathbb{G} e_j. \quad (3.30)$$

Analogously $e_i^* \mathbb{A} e_j = \langle \psi_i, \mathcal{K} \psi_j \rangle$, $e_i^* \mathbb{J} e_j = \langle \mathcal{K} \psi_i, \mathcal{K} \psi_j \rangle$. Sesquilinearity of $\langle \cdot, \cdot \rangle$ yields

$$\langle g, g \rangle = c^* \mathbb{G} c, \quad \langle g, \mathcal{K} g \rangle = c^* \mathbb{A} c, \quad \langle \mathcal{K} g, \mathcal{K} g \rangle = c^* \mathbb{J} c. \quad (3.31)$$

The proof is now simply a calculation. Indeed,

$$\begin{aligned} & \left\| \sqrt{W}(\Psi_Y - \lambda \Psi_X) c \right\|_2^2 \\ &= ((\Psi_Y - \lambda \Psi_X) c)^* W ((\Psi_Y - \lambda \Psi_X) c) \\ &= c^* (\Psi_Y^* W \Psi_Y - \bar{\lambda} \Psi_X^* W \Psi_Y - \lambda \Psi_Y^* W \Psi_X + |\lambda|^2 \Psi_X^* W \Psi_X) c \\ &= c^* J c - \bar{\lambda} c^* A c - \lambda c^* A^* c + |\lambda|^2 c^* G c. \end{aligned} \quad (3.32)$$

Taking the infinite-data limit,

$$\begin{aligned} & \lim_{M \rightarrow \infty} \left\| \sqrt{W}(\Psi_Y - \lambda \Psi_X) c \right\|_2^2 \\ &= c^* \mathbb{J} c - \bar{\lambda} c^* \mathbb{A} c - \lambda c^* \mathbb{A}^* c + |\lambda|^2 c^* \mathbb{G} c \\ &= \langle \mathcal{K} g, \mathcal{K} g \rangle - \bar{\lambda} \langle g, \mathcal{K} g \rangle - \lambda \langle \mathcal{K} g, g \rangle + |\lambda|^2 \langle g, g \rangle \\ &= \langle (\mathcal{K} - \lambda I) g, (\mathcal{K} - \lambda I) g \rangle \\ &= \|(\mathcal{K} - \lambda I) g\|_{\mathcal{X}}^2. \end{aligned} \quad (3.33)$$

□

From the proof we also directly see the following corollaries.

Corollary 3.3. *Let $\lambda \in \mathbb{C}$, define*

$$\text{res}(\lambda) = \text{res}_N(\lambda) := \min_{c^* G c = 1} \left\| \sqrt{W}(\Psi_Y - \lambda \Psi_X) c \right\|_2. \quad (3.34)$$

Then

$$\lim_{M \rightarrow \infty} \text{res}(\lambda) = \min_{\substack{g \in \text{span}\{\psi_1, \dots, \psi_N\} \\ \|g\|=1}} \|(\mathcal{K} - \lambda I)g\|_{\mathcal{X}}. \quad (3.35)$$

Corollary 3.4. *Assume $\overline{\text{span}\{\psi_j\}_{j=1}^\infty} = \mathcal{X}$. Then*

$$\lim_{N \rightarrow \infty} \lim_{M \rightarrow \infty} \text{res}(\lambda) = \sigma_{\inf}(\mathcal{K} - \lambda I). \quad (3.36)$$

Moreover the outer limit $N \rightarrow \infty$ is monotonically decreasing. Hence

$$\sigma_{ap, \epsilon}(\mathcal{K}) = \bigcap_{N > 0} \left\{ \lambda \in \mathbb{C} \mid \lim_{M \rightarrow \infty} \text{res}_M(\lambda) < \epsilon \right\}. \quad (3.37)$$

Corollary 3.4 suggests a way to compute the ϵ -approximate point pseudospectrum of \mathcal{K} . This is summarized in algorithm 2.

The computation of $\text{res}(\lambda)$ reduces to a generalized eigenvalue problem. Let

$$U (= U(\lambda)) = J - \bar{\lambda} A - \lambda A^* + |\lambda|^2 G. \quad (3.38)$$

Then computing $\text{res}(\lambda)^2$ is equivalent to solving the minimization problem

$$\min_{c \in \mathbb{C}^N} c^* U c \quad \text{such that} \quad c^* G c = 1. \quad (3.39)$$

Let ξ be a Lagrange multiplier, that is, a necessary condition for a solution of 3.39 is

$$U c - \xi G c = 0 \quad (3.40)$$

since U and G are symmetric. Inserting such a c into the objective yields

$$c^* U c = c^* \xi G c = \xi \quad (3.41)$$

since $c^* G c = 1$. Hence 3.39 is solved by computing the smallest generalized eigenvalue from 3.40 (symmetry of U and G guarantees that all such eigenvalues are real).

3.4 kernel EDMD (kEDMD)

3.5 kernel ResDMD (kResDMD)

Algorithm 2 Residual DMD

Require: Dictionary $\{\psi_j\}_{j=1}^N$, data points and weights $\{(w_i, x_i)\}_{i=1}^M$, grid $\{z_\nu\}_{\nu=1}^T \subset \mathbb{C}$, tolerance ϵ

- 1: Perform algorithm 1 to obtain G , A , K
- 2: Construct J as in 3.29
- 3: **for** z_ν **do**
- 4: Compute $\text{res}(z_\nu)$ as in 3.39
- 5: **return** $\{z_\nu \mid \text{res}(z_\nu) < \epsilon\}$ as an approximation for $\sigma_{ap,\epsilon}(\mathcal{K})$

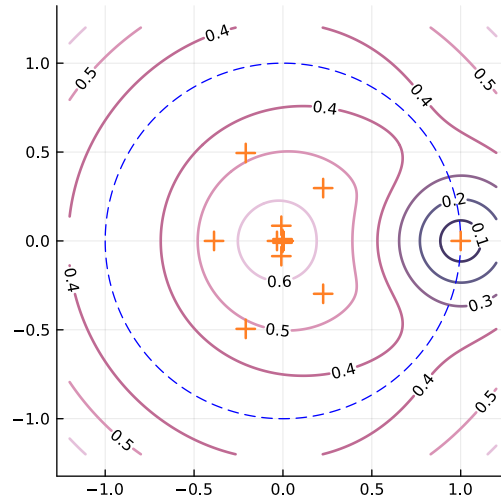


Figure 3.2: Algorithm 2 applied to the quadratic map (c.f. figure 2.1) performed with the same parameters as figure 3.1. Contours of $\lambda \mapsto \text{res}(\lambda)$ are shown with the spectrum of K . The quadratic map is ergodic [10] which is verified by the residuals. The eigenvalues of K other than 1 are spurious.

4 Benchmark Examples

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5 Appendix

5.1 Transfer Operator for Blaschke Products

We consider the analytic mapping on the unit circle in \mathbb{C} ,

$$\tau : \mathbb{T} \ni z \mapsto z \frac{z - \mu}{1 - \bar{\mu}z}$$

dependent on a parameter $\mu \in \mathbb{D}$ in the open unit disk.

τ is a Lebesgue-measure preserving 2-to-1 map. The spectrum $\sigma(\mathcal{K}|_{L^2(d\theta)}) = \overline{\mathbb{D}}$ is the whole unit disk, point spectrum $\sigma_p(\mathcal{K}|_{L^2(d\theta)}) = \{1\}$.

When considering instead the adjoint $\mathcal{L} = \mathcal{K}^*$, the point spectrum has the interesting property that $\{\mu^k\}_{k=0}^\infty \cup \{\bar{\mu}^k\}_{k=0}^\infty \subset \sigma_p(\mathcal{L}|_{L^2(d\theta)})$.

More specifically, let $\mathbb{T} \subset A$ be a (suitably chosen) open annulus containing the unit circle, and $H^2(A)$ be the Hardy Hilbert space of holomorphic functions on A which are square integrable on the edge ∂A . Then on this space, $\sigma(\mathcal{L}|_{H^2(A)}) = \{\mu^k\}_{k=0}^\infty \cup \{\bar{\mu}^k\}_{k=0}^\infty$.

EDMD with monomials $\{z \mapsto z^k\}_{k=-N}^N$ (Fourier basis) results in eigendecompositions which converge exponentially to $\sigma(\mathcal{L}|_{H^2(A)})$ as $N \rightarrow \infty$.

5.2 Some scribbles about kernels and residuals

EDMD: EDMD is a Petrov-Galerkin method as written in Péter's 2016 paper. Considering a dictionary $\{\psi_i\}_{i=1}^N \subset \mathcal{H}$ ($\mathcal{H} = L^2$ mostly) and snapshots $\{x_j\}_{j=1}^M \subset X$, EDMD makes the following construction:

$$\Psi(x) = [\psi_1 \dots \psi_N] \tag{5.1}$$

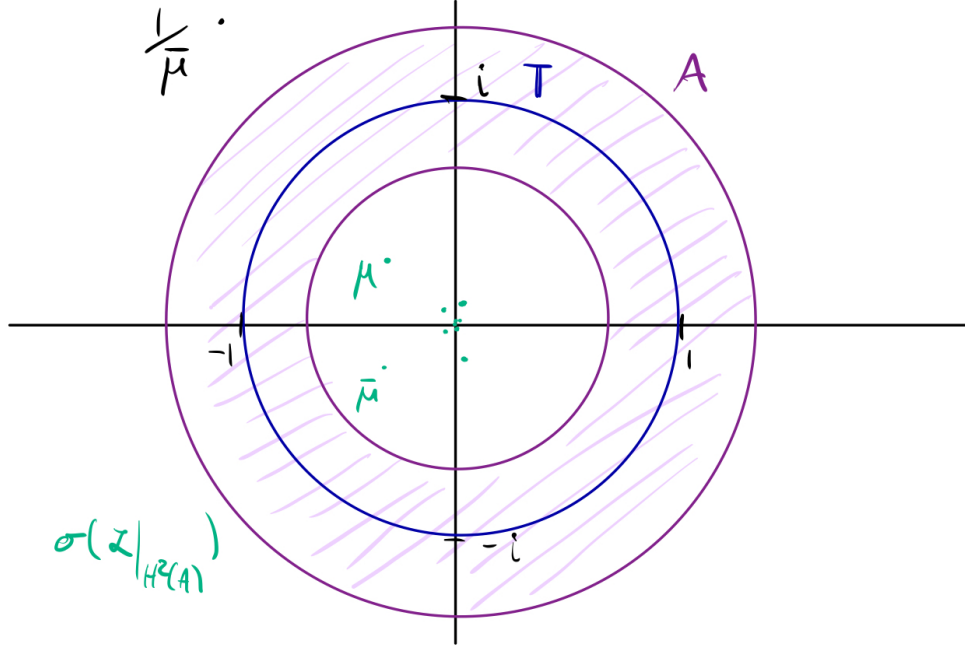
We find a matrix $K \in \mathbb{C}^{N \times N}$ that minimizes

$$\|\mathcal{K}\Psi - \Psi K\|_{\mathcal{H}(X, \mathbb{C}^{1 \times N})} \tag{5.2}$$

The \mathcal{H} norm is approximated by a quadrature:

$$\Psi_X = \Psi \begin{pmatrix} x_1 \\ \vdots \\ x_M \end{pmatrix}, \quad \Psi_Y = \Psi \circ S \begin{pmatrix} x_1 \\ \vdots \\ x_M \end{pmatrix} \tag{5.3}$$

$$\|\mathcal{K}\Psi - \Psi K\|_{\mathcal{H}(X, \mathbb{C}^{1 \times N})} \approx \|\Psi_Y - \Psi_X K\|_F. \tag{5.4}$$

Figure 5.1: Unit circle \mathbb{T} and Annulus A .

5.2 is minimized by

$$K = \Psi_X^\dagger \Psi_Y = G^\dagger A \quad (5.5)$$

where $G = \Psi_X^* \Psi_X$, $A = \Psi_X^* \Psi_Y$.

EDMD for Perron-Frobenius: Since

$$A_{ij} \stackrel{M \rightarrow \infty}{\approx} \langle \psi_i, \mathcal{K} \psi_j \rangle = \langle \mathcal{L} \psi_i, \psi_j \rangle, \quad (5.6)$$

an equivalent Galerkin approximation for $\mathcal{L} = \mathcal{K}^*$ is given by

$$P = G^\dagger A^* = \Psi_X^\dagger \Psi_X^* \Psi_Y^\dagger \Psi_Y^* \Psi_X. \quad (5.7)$$

kEDMD: Let $k : X \times X \rightarrow \mathbb{C}$ be a kernel so that

$$k(x_i, x_j) = \langle \Psi(x_i), \Psi(x_j) \rangle_{\ell^2} \quad (= \Psi(x_i) \Psi(x_j)^*). \quad (5.8)$$

(Note that $\Psi(x)$ might be infinite-dimensional, i.e. $N = \infty$, but we think of $N < \infty$ for clarity of the matrix manipulations.)

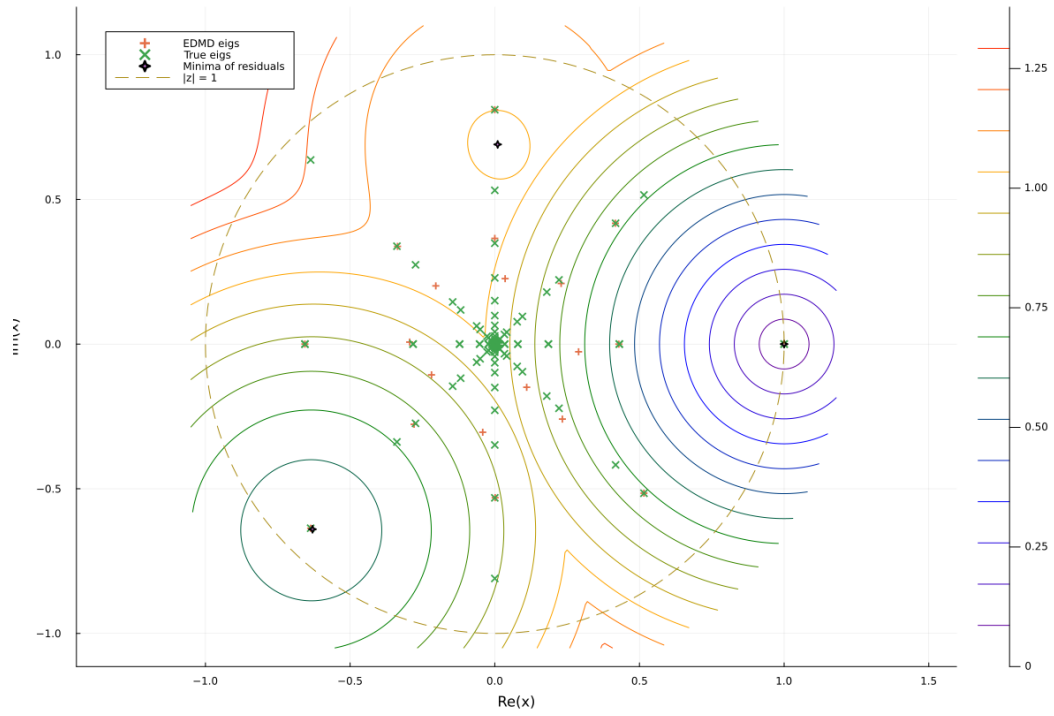


Figure 5.2: ResDMD over the space $L^2(\mathbb{T})$ for a Blaschke map $\tau(z) = z \frac{z-\mu}{1-\bar{\mu}z}$ with $\mu = 0.9e^{\pi i/4}$. True eigenvalues are known in the Hardy space $H^2(A)$ of functions holomorphic on the annulus $A = \{z \in \mathbb{C} \mid r < |z| < R\}$, $R = 11/10 - 1/32$, $r = 1/R$. ResDMD matrices G, A, L are computed w.r.t. this Hilbert space, using a dictionary of $N = 20$ monomials and $M = 200,000$ equally spaced quadrature nodes.

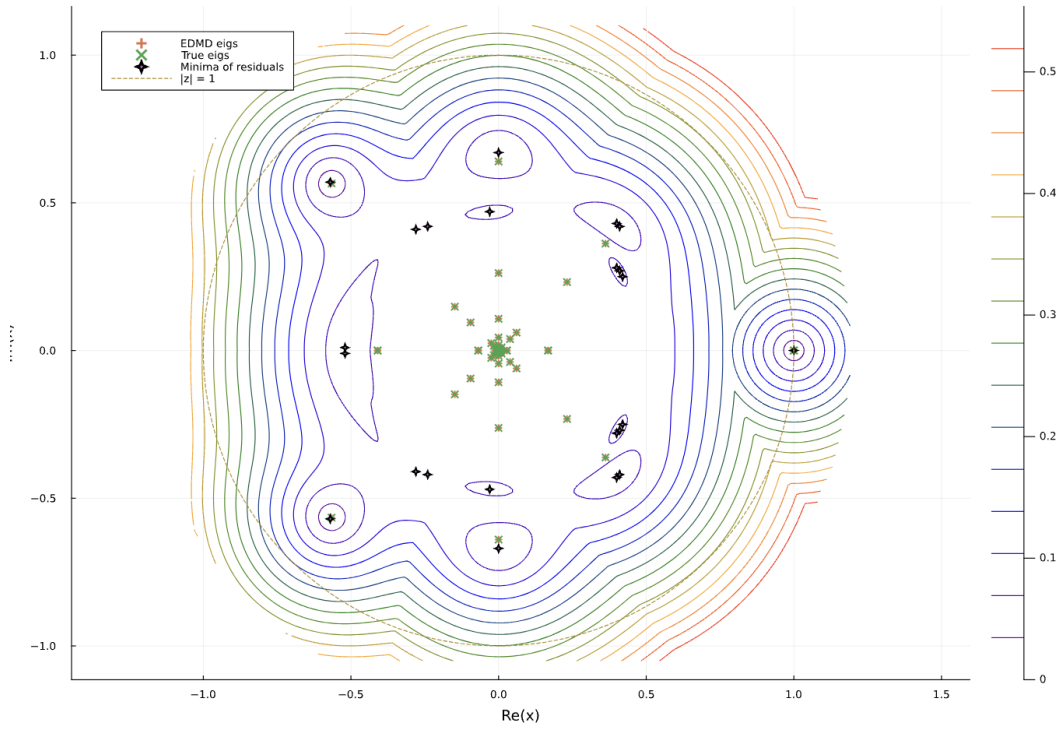


Figure 5.3: ResDMD for $\mu = 0.8e^{\pi i/4}$ over the space $H^2(A)^*$, where $H^2(A) \subset L^2$ is seen as embedded in L^2 and $(\cdot)^*$ the Banach space adjoint. This dual space is isomorphic to the space $H^2(\mathbb{D}_r) \oplus H^2(\mathbb{D}_{1/r}^\infty)$ and has a norm $\|f\|_{H^2(A)^*}^2 = \sum_{n \in \mathbb{Z}} |\hat{f}|^2 r^{|n|}$.

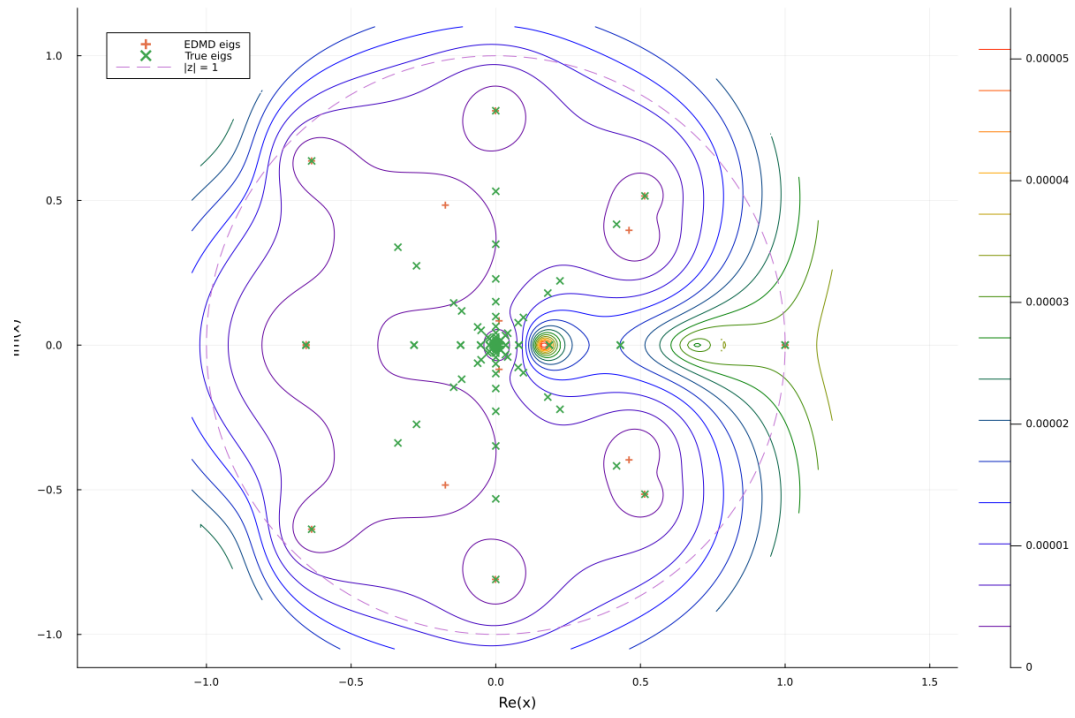


Figure 5.4: Kernelized ResDMD for the Blaschke map using a degree 20 polynomial kernel $k(x, y) = (1 + x'y)^{20}$ and $M = 2,000$ equally spaced quadrature nodes.

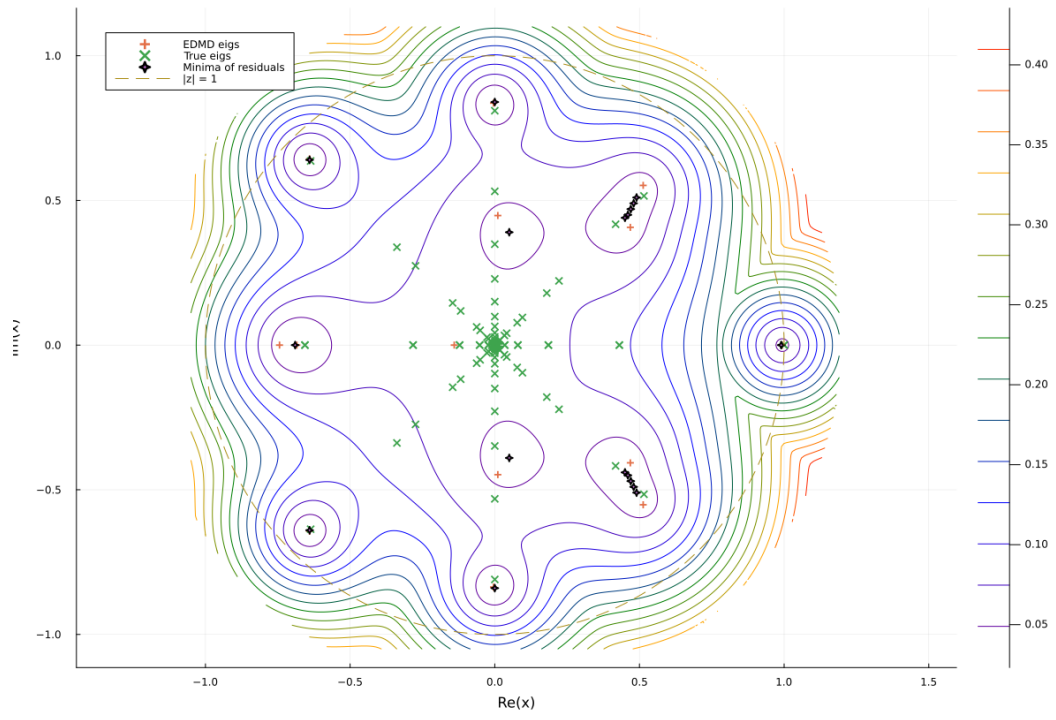


Figure 5.5: Kernelized ResDMD using an RBF kernel $k(x, y) = e^{-|x-y|^2/c}$ with $c = 0.01$ and $M = 2,000$ equally spaced quadrature nodes.

Let further $\Psi_X \stackrel{\text{SVD}}{\approx} Q\Sigma Z^*$ and

$$\hat{G} = (k(x_i, x_j))_{ij} = \Psi_X \Psi_X^* \stackrel{\text{diagonalize}}{=} Q(\Sigma^* \Sigma) Q^*, \quad (5.9)$$

$$\hat{A} = (k(y_i, x_j))_{ij} = \Psi_Y \Psi_X^* \quad (5.10)$$

Then

$$\hat{K} = (\Sigma^\dagger Q^*) \hat{A} (Q \Sigma^\dagger) \in \mathbb{C}^{M \times M} \quad (5.11)$$

has the same eigenvalues as K .

ResDMD: The key idea behind ResDMD is, for an element $z \in \mathbb{C}^N$, $g = \Psi z$:

$$\text{res}(\lambda, z) := \|(\mathcal{K} - \lambda)g\|_{\mathcal{H}(X, \mathbb{C})}^2 \quad (5.12)$$

$$\stackrel{M \rightarrow \infty}{\approx} \|(\Psi_Y - \lambda \Psi_X)z\|_{\ell^2}^2 \quad (5.13)$$

$$= z^* (\Psi_Y^* \Psi_Y - \lambda \Psi_Y^* \Psi_X - \bar{\lambda} \Psi_X^* \Psi_Y + |\lambda|^2 \Psi_X^* \Psi_X) z \quad (5.14)$$

$$= z^* (L - \lambda A^* - \bar{\lambda} A + |\lambda|^2 G) z. \quad (5.15)$$

We interpret the first equation as

$$\text{regression error} \stackrel{M \rightarrow \infty}{\approx} \text{operator residual}. \quad (5.16)$$

The function $\lambda \mapsto \text{res}(\lambda) = \min_z \text{res}(\lambda, z)$ is shown in 5.2

kResDMD: We apply idea 5.16 to \hat{K} :

$$\hat{K}^* = (\Psi_X^* Q \Sigma^\dagger)^\dagger (\Psi_Y^* Q \Sigma^\dagger) =: \hat{\Psi}_X^\dagger \hat{\Psi}_Y \quad (5.17)$$

so that \hat{K}^* solves the regression problem

$$\min_M \|\hat{\Psi}_Y - \hat{\Psi}_X M\|_F. \quad (5.18)$$

The regression error reduces to

$$\text{r\hat{e}s}(\lambda, z) := \|(\hat{\Psi}_Y - \lambda \hat{\Psi}_X)z\|_{\ell^2}^2 \quad (5.19)$$

$$= (Q\Sigma^\dagger z)^* (\Psi_Y \Psi_Y^* - \lambda \Psi_Y \Psi_X^* - \bar{\lambda} \Psi_X \Psi_Y^* + |\lambda|^2 \Psi_X \Psi_X^*) (Q\Sigma^\dagger z) \quad (5.20)$$

$$= (Q\Sigma^\dagger z)^* \left(\hat{L} - \lambda \hat{A}^* - \bar{\lambda} \hat{A} + |\lambda|^2 \hat{G} \right) (Q\Sigma^\dagger z) \quad (5.21)$$

Note at this point it is unclear whether this regression error in the case of \hat{K} has a physical meaning.

ResDMD for Perron-Frobenius: Consider $\tilde{P} = Z^* P Z$, which has the same eigenvalues as P . Then

$$\tilde{P} = (\Psi_X^* Q \Sigma)^\dagger (\Psi_Y^* Q \Sigma) =: \tilde{\Psi}_X^\dagger \tilde{\Psi}_Y \quad (5.22)$$

so that \hat{P} solves the regrestion problem

$$\min_M \|\tilde{\Psi}_Y - \tilde{\Psi}_X M\|_F. \quad (5.23)$$

The interesting thing to note is that

$$\tilde{\text{r\hat{e}s}}(\lambda, z) := \|(\tilde{\Psi}_Y - \lambda \tilde{\Psi}_X)z\|_{\ell^2}^2 \quad (5.24)$$

$$= (Q\Sigma z)^* (\Psi_Y \Psi_Y^* - \lambda \Psi_Y \Psi_X^* - \bar{\lambda} \Psi_X \Psi_Y^* + |\lambda|^2 \Psi_X \Psi_X^*) (Q\Sigma z) \quad (5.25)$$

$$= (Q\Sigma z)^* \left(\hat{L} - \lambda \hat{A}^* - \bar{\lambda} \hat{A} + |\lambda|^2 \hat{G} \right) (Q\Sigma z) \quad (5.26)$$

and hence whenever Ψ_X has the same number of zero singular values (we assume there are none i.e. $s_{\min}(\Psi_X) > 0$),

$$\text{r\hat{e}s}(\lambda) := \min_z \text{r\hat{e}s}(\lambda, z) = \min_\xi \tilde{\text{r\hat{e}s}}(\lambda, \xi) =: \tilde{\text{r\hat{e}s}}(\lambda). \quad (5.27)$$

The function $\text{r\hat{e}s}$ is shown in 5.4, 5.5.

Interpreting the regression error: Let $y = Q\Sigma z$. For convenience we conjugate λ , i.e. let $\tilde{P}z = \bar{\lambda}z$ be a candidate eigenpair for \mathcal{L} .

What follows is a rough, purely formal sketch of the idea that we want to show, from what I calculated this morning (no promises for correctness):

$$\|(\Psi_Y^* - \bar{\lambda}\Psi_X^*)y\|_{\ell^2}^2 \quad (5.28)$$

$$= \sum_{i=1}^N \left| \sum_{j=1}^M \overline{(\mathcal{K}\psi_i - \lambda\psi_i)(x_j)} y_j \right|^2 \quad (5.29)$$

$$\stackrel{M \rightarrow \infty}{\approx} \sum_{i=1}^N \left| \int \overline{(\mathcal{K}\psi_i - \lambda\psi_i)(\xi)} y(\xi) d\xi \right|^2 \quad (5.30)$$

$$= \sum_{i=1}^N |\langle (\mathcal{K} - \lambda)\psi_i, y \rangle|^2 \quad (5.31)$$

$$= \sum_{i=1}^N |\langle \psi_i, (\mathcal{L} - \bar{\lambda})y \rangle|^2 \quad (5.32)$$

$$\stackrel{N \rightarrow \infty}{\approx} \|(\mathcal{L} - \bar{\lambda})y\|_{\mathcal{H}(X, \mathbb{C})}^2. \quad (5.33)$$

where for 5.30 we say that the vector $y = (y_j)_j$ converges to a function $y : \xi \mapsto y(\xi)$, and for 5.33 we use Plancherel's theorem (hence we need to assume Ψ is an orthonormal family that in the limit has span dense in $\mathcal{H}(X, \mathbb{C})$).

It seems likely that these convergences will work, under the crucial assumption that the singular values stay bounded from above and below as $N \rightarrow \infty$, which is given if Ψ is an orthonormal family. There is some literature on kernel *QR* algorithms, which are essentially Gram-Schmidt for kernel weight matrices. However, I would much prefer if this would work without the orthonormal requirement, as most kernel functions and snapshot datasets will not make orthonormal Gram matrices.

For e.g. Gauß kernels $k(x, y) = \exp(-\|x - y\|^2/c)$, the sharpness parameter c has to be changed as M changes in order for the singular value condition to hold since if c is kept constant and M grows, the Gram matrix slowly becomes the all-ones matrix.

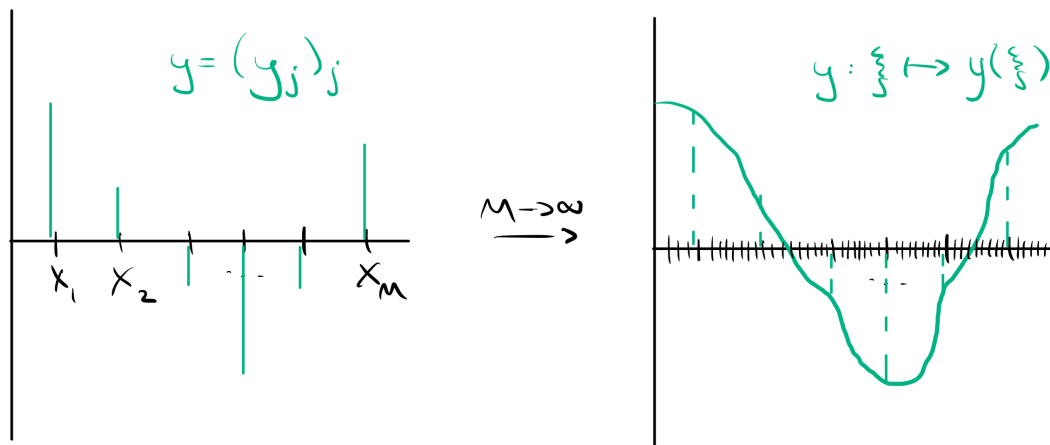


Figure 5.6: Idea behind equation 5.30