Phys 514

Problem Set 3

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Contents

1 Tensors on Manifolds

Consider a D-dimensional manifold. In any coordinate system x^μ, the Kronecker delta symbol δ^μ_ν is defined to be

$$\delta^\mu_\nu = \begin{cases} 1 & \text{if} \quad \mu = \nu \\ 0 & \text{if} \quad \mu \neq \nu \end{cases}$$

1.1 Tensor weight

We need to show that δ^μ_ν is a tensor of weight $(1,1)$. We apply the tensor transformation laws such that $\delta^{\mu'}_{\nu'}$ is a Kronecker delta over a new set of coordinates $x^{\mu'}$ given by

$$\delta^{\mu'}_{\nu'} = \delta^\mu_\nu \frac{\partial x^{\mu'}}{\partial x^\mu} \frac{\partial x^\nu}{\partial x^{\nu'}}$$

We notice that we only consider the terms where $\mu = \nu$, as they will be the only non-zero terms. We can thus define $\mu = \nu = \alpha$ such that:

$$\begin{aligned} \delta^{\mu'}_{\nu'} &= \sum_\alpha \frac{\partial x^{\mu'}}{\partial x^\alpha} \frac{\partial x^\alpha}{\partial x^{\nu'}} \\ &= \sum_\alpha \frac{\partial x^\alpha}{\partial x^{\nu'}} \frac{\partial x^{\mu'}}{\partial x^\alpha} \\ &= \frac{\partial x^{\mu'}}{\partial x^{\nu'}} \\ &= \begin{cases} 1 & \text{if } \mu' = \nu' \\ 0 & \text{if } \mu' \neq \nu' \end{cases} \\ &= \delta^\mu_\nu \end{aligned}$$

Now, considering $\hat{e}_\mu$ as the basis vectors for coordinates x^μ, and $\hat{e}_{\nu'}$ as the basis vectors for coordinates $x^{\nu'}$, it follows that for any transformation of coordinates $\Lambda^{\nu'}_\mu$ we have

$$\hat{e}_\mu = \Lambda^{\nu'}_\mu \hat{e}_{\nu'}$$

And to isolate the basis $\hat{e}_{\nu'}$, we must apply the inverse of $\Lambda^{\nu'}_\mu$, namely $\Lambda^\rho_{\sigma'}$ such that

$$\Lambda^\mu_{\nu'} \Lambda^{\nu'}_\rho = \delta^\mu_\rho, \qquad \Lambda^{\sigma'}_\lambda \Lambda^\lambda_{\tau'} = \delta^{\sigma'}_{\tau'}$$

Which clearly implies that

$$\hat{e}_{\nu'} = \Lambda^\mu_{\nu'} \hat{e}_\mu$$

So δ^μ_ν is a tensor of weight $(1,1)$, as it has a covariant index and a contravariant index.

1.2 Vectors and tensors

We need to show that $\delta^\mu_\nu v^\nu = v^\mu$ and $\delta^\mu_\nu w_\mu = w_\nu$ where v^ν and w_μ are a vector and one form respectively. As seen in the second part of my answer to 1.1, δ^μ_ν represents a transformation multiplied by its own inverse. As such, we can clearly see that

$$\delta^\mu_\nu v^\nu = \Lambda^\mu_{\nu'}\Lambda^{\nu'}_\nu v^\nu = \Lambda^\mu_{\nu'}(\Lambda^{\nu'}_\nu v^\nu) = \Lambda^\mu_{\nu'} v^{\nu'} = v^\mu$$

$$\delta^\mu_\nu w_\mu = \Lambda^\mu_{\nu'}\Lambda^{\nu'}_\nu w_\mu = \Lambda^{\nu'}_\nu(\Lambda^\mu_{\nu'} w_\mu) = \Lambda^{\nu'}_\nu w_{\nu'} = w_\nu$$

As required.

1.3 Computation

First we have δ^μ_μ, which seems very sneaky. A quick guess would assume that

$$\delta^\mu_\mu = \begin{cases} 1 & \text{if} \quad \mu = \mu \\ 0 & \text{if} \quad \mu \neq \mu \end{cases}$$

But aha, this is not what we should see. In this case, we should add the number of dimensions whenever the indices of the Kronecker delta are equal (which they always are), so written in transformations:

$$\delta^\mu_\mu = \Lambda^\mu_{\nu'}\Lambda^{\nu'}_\mu = \sum_i^D \frac{\partial x^i}{\partial x^i} = D$$

Next we have $\delta^\mu_\nu \delta^\mu_\rho$. We again write this down in transformations:

$$\delta^\mu_\nu \delta^\mu_\rho = \Lambda^\mu_{\nu'}\Lambda^{\nu'}_\nu \Lambda^\mu_{\sigma'}\Lambda^{\sigma'}_\rho$$

Which means that

$$\delta^\mu_\nu \delta^\mu_\rho = \begin{cases} D & \text{if} \quad \mu = \nu = \rho \\ 0 & \text{otherwise} \end{cases}$$

So we will only add the elements whenever the indices match, and they will add up to the number of dimensions of the space.

2 General Questions

2.1 Dummy (t h i c c) indices

We must show by manipulating dummy indices that

$$Z_{\mu\nu}v^{\mu}v^{\nu} = \frac{1}{2}\left(Z_{\mu\nu} + Z_{\nu\mu}\right)v^{\mu}v^{\nu}$$

where we assume $Z_{\mu\nu}$, $Z_{\nu\mu}$ are both rank (0,2) tensors that are also symmetric matrices. Hence, we know that $Z_{\mu\nu} = Z_{\nu\mu}^{T}$ (see equation 1.69 in the book for an example, where the tensor is antisymmetric instead of symmetric), so clearly:

$$\frac{1}{2}(Z_{\mu\nu} + Z_{\nu\mu}) = \frac{1}{2}\left(\begin{pmatrix} z_{1,1} & \cdots & z_{1,\nu} \\ \vdots & \ddots & \vdots \\ z_{\mu,1} & \cdots & z_{\mu,\nu} \end{pmatrix} + \begin{pmatrix} z_{1,1} & \cdots & z_{1,\mu} \\ \vdots & \ddots & \vdots \\ z_{\nu,1} & \cdots & z_{\nu,\mu} \end{pmatrix}\right)$$

$$= \frac{1}{2}\begin{pmatrix} 2 \cdot z_{1,1} & \cdots & z_{1,\nu} + z_{1,\mu} \\ \vdots & \ddots & \vdots \\ z_{\mu,1} + z_{\nu,1} & \cdots & \underbrace{2z_{\mu,\nu}}_{=2z_{\nu,\mu}} \end{pmatrix}$$

$$= \begin{pmatrix} z_{1,1} & \cdots & \frac{z_{1,\nu}+z_{1,\mu}}{2} \\ \vdots & \ddots & \vdots \\ \frac{z_{\mu,1}+z_{\nu,1}}{2} & \cdots & z_{\mu,\nu} \end{pmatrix} = \begin{pmatrix} z_{1,1} & \cdots & \frac{z_{1,\mu}+z_{1,\nu}}{2} \\ \vdots & \ddots & \vdots \\ \frac{z_{\nu,1}+z_{\mu,1}}{2} & \cdots & z_{\mu,\nu} \end{pmatrix}$$

Recall that both tensors are symmetric

$$= \begin{pmatrix} z_{1,1} & \cdots & z_{1,\nu} \\ \vdots & \ddots & \vdots \\ z_{\mu,1} & \cdots & z_{\mu,\nu} \end{pmatrix}$$

$$= Z_{\mu\nu}$$

Therefore,

$$Z_{\mu\nu}v^{\mu}v^{\nu} = \frac{1}{2}\left(Z_{\mu\nu} + Z_{\nu\mu}\right)v^{\mu}v^{\nu}$$

and the dummy index is thicc indeed.

2.2 Independent components

For any (0,2) tensor, we have 2 covariant indices and no contravariant indices. Hence, assuming the tensor is non-symmetric to distinguish this case from symmetric and antisymmetric, we can simply represent the number of independent components in any rank $(0,2)$ tensor of dimension D as the total number of components in the tensor, so this is given by D^2

Now if we consider symmetric tensors, then we know that $T_{ij} = T_{ji}$, so we may only consider $\sum_{k=1}^{D} k$ independent elements. In 3D, this yields 1+2+3 = 6, in 4D we have 1+2+3+4 = 10, etc.

For an antisymmetric tensor where $T_{ij} = -T_{ji}$ for $i \neq j$ and $T_{ij} = 0$ for $i = j$, we consider the same amount of independent components as the symmetric case but without the diagonal, so we have $\sum_{k=1}^{D-1} k$ independent components. In 3D we have 1+2 =3 independent component, in 4D we have 1+2+3 = 6, and so on.

2.3 Sum of tensors

As we saw in part 2.2, an antisymmetric tensor will contain independent elements $T_{ij} = -T_{ji}$ except $T_{ij} = 0$ for $i = j$, while a symmetric tensor need only have the property that $T_{ij} = T_{ji}$ and $T_{ij} \neq 0$ for $i = j$. Thus, we can build a (0,2) tensor with D^2 independent components by adding a symmetric matrix and an antisymmetric matrix together. This will yield a new tensor where we will have the component property $T_{ij} \neq 0$ for $i = j$ and also $T_{ij} \neq T_{ji} \neq -T_{ji}$.

2.4 Independent components pt. 2

This case, we are talking about an outer product of two (0,1) tensors, which yields a (0,2) tensor as desired. Thus, we can evidently see that we would have D^2 independent components again, as in the general case. So for any $D \geq 2$ dimensions, we will have a (0,2) tensor resulting from 2 (0,1) one forms.

2.5 Independent components (bonus)

For $(0,p)$ tensors, we would have D^p independent components in the general case, $\sum_{k=1}^{D} k^{p-1}$ for the symmetric case, and $\sum_{k=1}^{D-1} k^{p-1}$ for the antisymmetric case. We extend the logic from 2.2 such that in a symmetric tensor we have $T_{ijk} = T_{jik} = T_{jki} = T_{kji} = T_{ijk}$ and also equals all permutations of ijk in T_{ijk}. In an antisymmetric tensor, we'd have $T_{ijk} = -T_{jik} = T_{jki} = -T_{kji} = T_{kij} = -T_{ikj}$, so clearly, when adding a symmetric and antisymmetric tensor we would yield a (0,3) nonsymmetric tensor. And this can be extended to any $(0,p)$ tensor as well, as every general nonsymmetric tensor can be decomposed into an antisymmetric and a symmetric part.

3 Higher-dimensional Space

3.1 Transformations

3.2 Transformations 2

3.3 Metric on Cartesian Coordinates

3.4 Equivalent metrics

4 Minkowski metric problem