Monte-Carlo and Langevin dynamics

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$$U(x) = \sum_{i,j,i',j'} -Js_{i,j}s_{i',j'}$$

$$X = \{ S_{11}, S_{12} ... \}$$

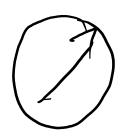
$$S = \{ -1; 1 \}$$

$$S = 5 - 7 - 7 = -27$$

$$S = +3$$

$$U(x) = \sum_{i,j,i',j'}^{neigs.} -Js_{i,j} \cdot s_{i',j'}$$

$$S = \{ A \}$$



















Problem:

$$\langle O \rangle = \sum_{(x)} P(x)O(x)$$

$$P(x) = \frac{1}{Z} e^{-\frac{U(x)}{kT}}$$

$$Z = \sum_{x} e^{-\frac{U(x)}{kT}}$$

$$U(x)$$
 (potential) energy of the system

$$P(x)$$
 probability of state x

$$Z$$
 partition function

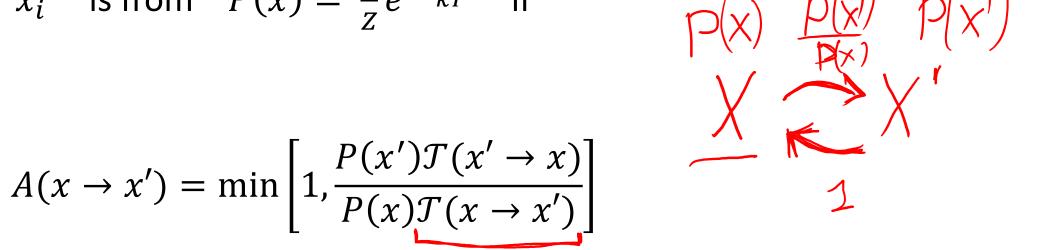
9 spins
$$2$$
 states 2 512

How can we approximate observables?

$$x_i$$
 is from $P(x) = \frac{1}{Z}e^{-\frac{U(x)}{kT}}$

How do we sample? One bit at a time, using Metropolis rejection

$$A(x \to x') = \min \left[1, \frac{P(x')\mathcal{T}(x' \to x)}{P(x)\mathcal{T}(x \to x')} \right]$$



^{*} $A(x \rightarrow x')$ above ensures detailed balance

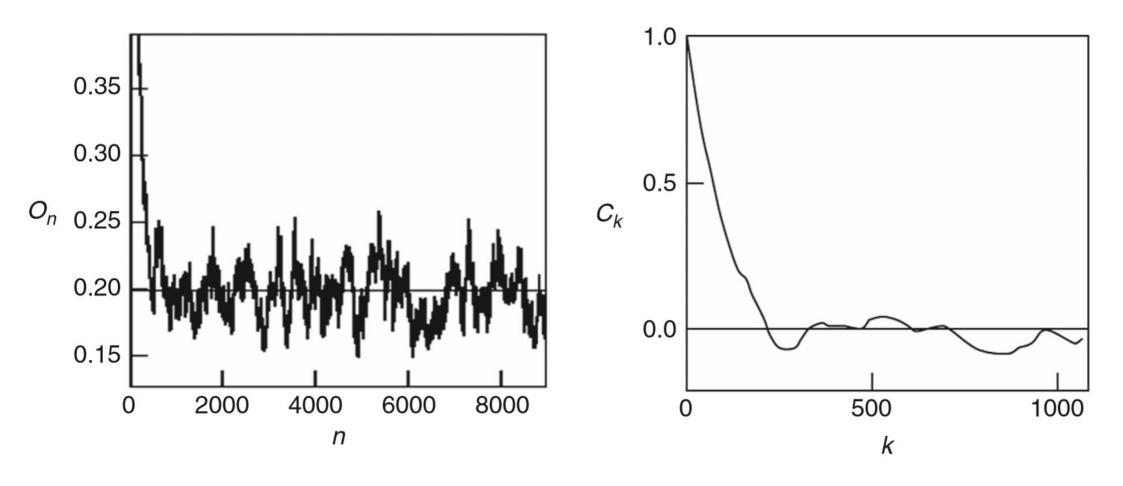
Monte Carlo with Metropolis rejection

- \rightarrow 1. Start from a state x_0
- \rightarrow 2. Iterate the following *N* times
 - \rightarrow Propose a move $x_i \rightarrow x_{i+1}$
 - If $U(x_{i+1}) < U(x_i)$ accept the move else accept the move with probability $p = \exp(-\frac{U(x_{i+1}) U(x_i)}{kT})$
 - Compute and output observables regardless of whether the move is accepted

Demo

```
↓↓↓↓↓↓↓↓↓↓↓↓↓↓
↓↓↓↓↓↓↓↓↓∠∠←<u></u>~~~~~~~~~~~~~↓↓↓↓↓↓
```

Calculation of the observables



From [Martin et al Interacting Electrons ch. 22]

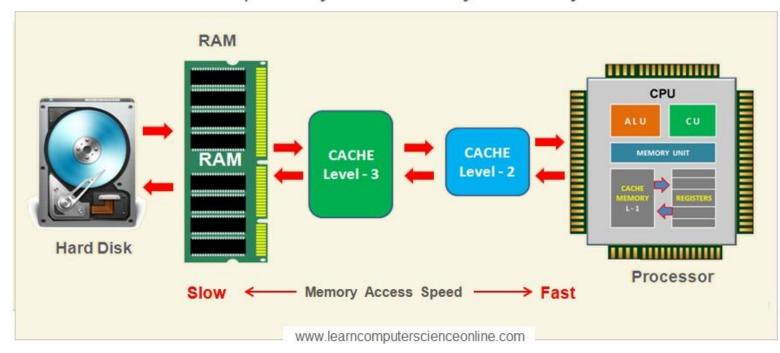
Computer performance I

- Select random spin 2
- Propose change 1
- Calculate Energy 4
- Accept or reject 30

Computer performance I

- Select random spin
- Propose change
- Calculate Energy
- Accept or reject

Computer System Memory Hierarchy



Computer performance II (oversimplification)

```
1 // The worst you can do
2 i = randi(N);
3 j = randi(M);
4 if(s[i,j]==...)
5 ...
1 // The best you can do
2 for(i=0; i<N; i++)
3 for(j=0; j<M; j++)
4 s[i,j] += ...
```

Other sampling methods: Molecular dynamics

$$H(p,x) = U(x) + T(p)$$

$$dx = M^{-1}pdt$$

$$dp = -F(x)dt = -\frac{\partial U(x)}{\partial x}dt$$
 + thermostat

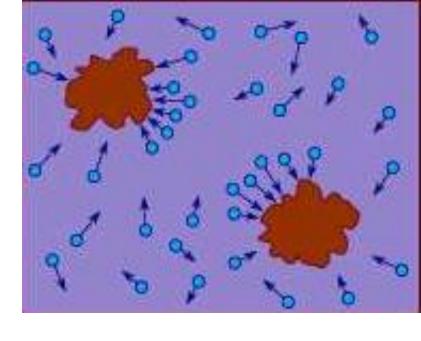
$$P(p,x) = \frac{1}{Z}e^{-\frac{U(x)+T(p)}{kT}}$$

$$P(x) = \sum_{p} P(p, x) = \frac{1}{Z} e^{-\frac{U(x)}{kT}}$$

Other sampling methods

Brownian dynamics

$$dx = M^{-1}F(x)dt + \sqrt{2kT}M^{-1/2}dW$$



Langevin dynamics*

$$dx = M^{-1}pdt dp = [F(x) - \gamma p]dt + \sigma M^{\frac{1}{2}}dW$$

*See [Leimkuhler, Matthews AMRX (2013) v 1 p 34] for implementation details

Brownian dynamics: Implementation

$$dx = M^{-1}F(x)dt + \sqrt{2kT}M^{-1/2}dW$$

Implement as

$$x_{n+1} = x_n + M^{-1}F(x_n)h + \sqrt{2kTh}M^{-1/2}R_n$$

$$R_n = \mathcal{N}(0,1)$$

M- free parameter

h- time step

Brownian dynamics: Implementation

$$dx = M^{-1}F(x)dt + \sqrt{2kT}M^{-1/2}dW$$

Implement as

$$R_n = \mathcal{N}(0,1)$$

$$x_{n+1} = x_n + M^{-1}F(x_n)h + \sqrt{2kTh}M^{-1/2}R_n$$

M- free parameter

or better as*

$$x_{n+1} = x_n + M^{-1}F(x_n)h + \sqrt{2kTh}M^{-1/2}\frac{R_n + R_{n+1}}{2}$$

^{*[}Leimkuhler, Matthews AMRX (2013) v 1 p 34]

Brownian dynamics: Implementation

```
// Calculate forces first, then update is simply
for(i=0; i<N; ++i)
for(j=0; j<M; ++j)
x[i,j] += F[i,j]*h + sqrt(2*T*h)*nrand();</pre>
```

Demonstration

Erratum

Equation

$$dx = M^{-1}F(x)dt + \sqrt{2kT}M^{-1/2}dW$$

does has unit of time missing, correct equation is

$$x_{n+1} = x_n + \frac{\delta t^2}{2} M^{-1} F(x_n) + \frac{\delta t}{2} M^{-1} (p_n + p_{n+1})$$

= $x_n + \frac{\delta t^2}{2} M^{-1} F(x_n) + \frac{\delta t}{2} \sqrt{k_B T} M^{-1/2} (R_n + R_{n+1}),$