

MLL253

Electronic Properties of Materials

$$E(v, t) = E_0 e^{i(\omega t - kv)}$$

phase velocity $v_p = \frac{\omega}{k}$

The avg intensity = $\frac{1}{2} c G E_0^2$

only measurable quantity.

Born interpretation:

$\rightarrow \Psi(v, y, z, t) \rightarrow$ function

probability to find particle $\left| \Psi(v, y, z, t) \right|^2$
at v, y, z at any time t .

\rightarrow Time-independent Schrodinger Egn:

$$\left[\frac{-\hbar^2}{2m} \nabla^2 + V \right] \Psi = i\hbar \frac{\partial \Psi}{\partial t}$$

$\frac{-\hbar^2}{2m} \nabla^2 + V$ = Hamiltonian operator (\hat{H})
 \downarrow
 kE PE

A free particle ($V=0$) ^{everywhere} momentum.

$$\Psi = A \cos\left(\frac{2\pi v}{\lambda} - \omega t\right)$$

\rightarrow oscillating everywhere.

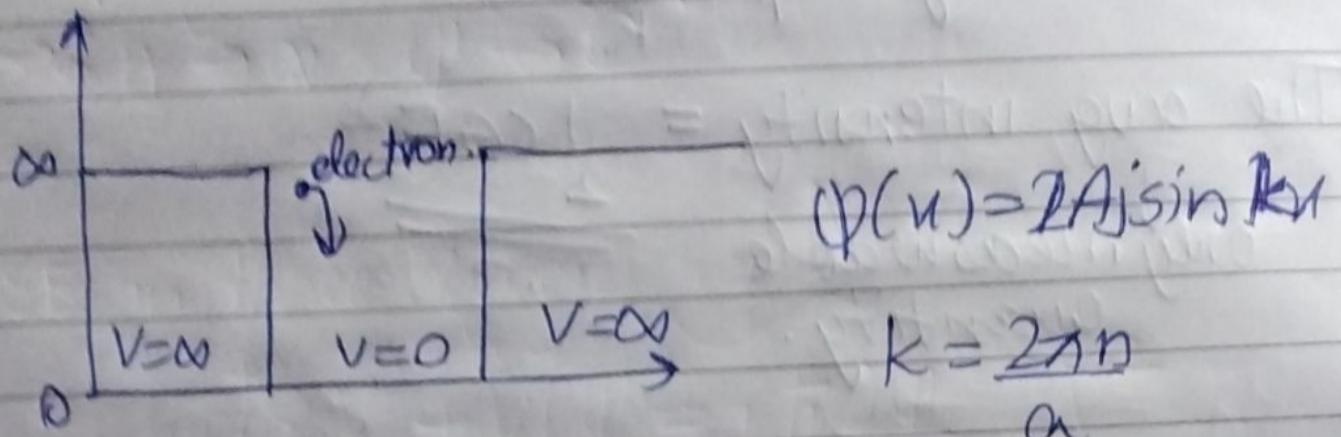
$$\frac{2\pi}{\lambda} = \frac{2\pi p}{h} = k$$

$$\omega = \frac{h\nu}{\lambda} = \frac{E}{\hbar} \rightarrow \text{energy}$$

$$\psi(n) = Ae^{ikx} \quad [V=0]$$

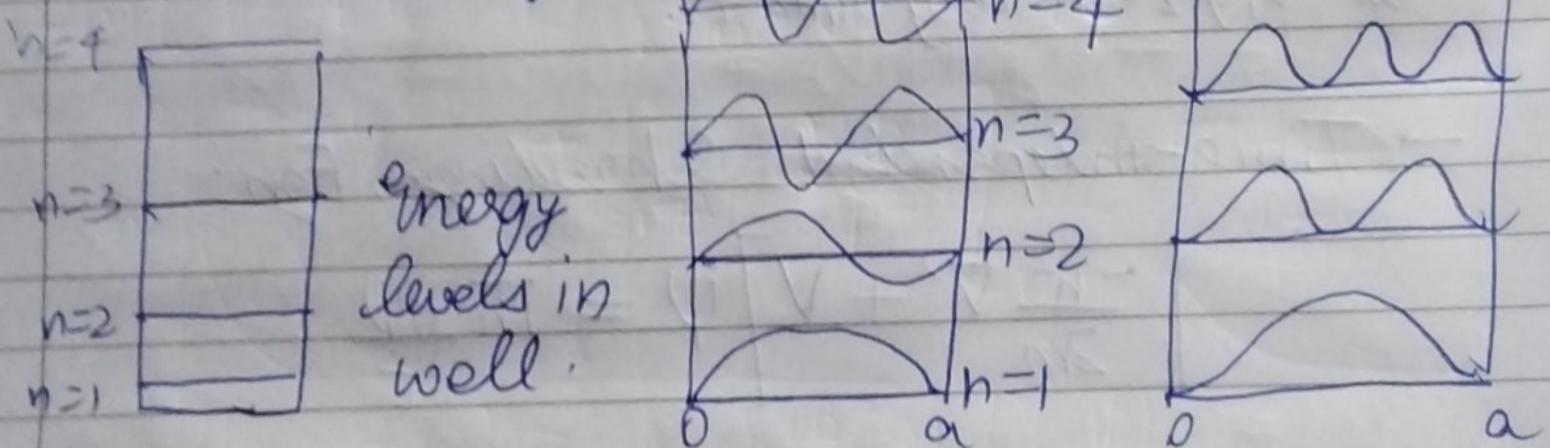
$$|\psi|^2 = A^2$$

Particle In a Box:



$$E = \frac{\hbar^2 k^2}{2m} = \frac{\hbar^2 n^2}{8ma^2}$$

$$A = \frac{1}{\sqrt{2a}} \quad n = \text{principal quantum no.}$$



- wavefunctions are no longer continuous.
- only discrete energy levels exist:

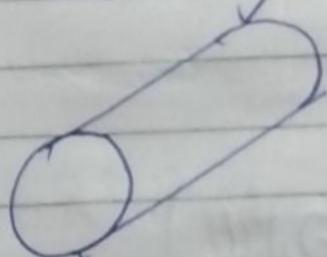
$n=0 \Rightarrow E=0 \rightarrow$ particle does not exist anywhere
 {Trivial soln} \hookrightarrow invalid soln.

Ground state: $n=1$ $E \neq 0$
 non-zero energy
 non-zero velocity

$$\int_{-\infty}^{\infty} \psi^* U \psi dr = 1 \quad \{ \text{Normalization} \}$$

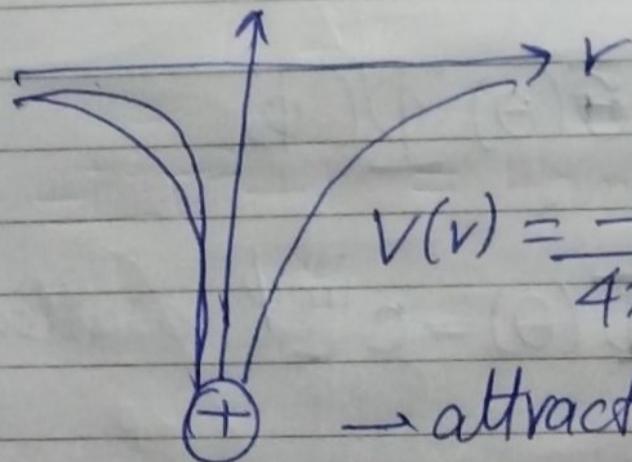
$$\int_0^a 4A^2 \sin^2 kr = 1$$

→ looking at the top of cylindrical box.



→ with higher energy, electron starts rotating.

Hydrogen Atom:-



$$V(r) = \frac{-ze^2}{4\pi\epsilon_0 r} \quad \{ \text{Coulombic potential} \}$$

→ attractive potential from nuclei
 → potential is spherically symmetric

→ $\psi(r, \theta, \phi)$ → univalued

$$\Psi(r, \theta, \phi) = \Psi(r, \theta, \phi + 2\pi) = \Psi(r, \theta, \phi)$$

$$= \Psi(r, \theta + 2\pi, \phi)$$

→ periodic boundary conditions. (univalue)

$$\Rightarrow r \rightarrow \infty, \psi \rightarrow 0$$

SE in three dimensions:-

$$\left[-\frac{\hbar^2}{2m} \nabla^2 + V \right] \Psi = i\hbar \frac{d\Psi}{dt} \quad \nabla^2 = \frac{\partial^2}{\partial r^2}$$

In spherical coordinates.

$$\begin{aligned} -\frac{\hbar^2}{2m} & \left[\frac{1}{r^2} \frac{\partial}{\partial r} \left(r^2 \frac{\partial \Psi}{\partial r} \right) + \frac{1}{r^2 \sin \theta} \frac{\partial}{\partial \theta} \left(\sin \theta \frac{\partial \Psi}{\partial \theta} \right) \right. \\ & \left. + \frac{1}{r^2 \sin^2 \theta} \left(\frac{\partial^2 \Psi}{\partial \phi^2} \right) \right] = i\hbar \frac{d\Psi}{dt} \end{aligned}$$

$$\therefore \Psi(r, \theta, \phi) = R(r) \Theta(\theta) \Phi(\phi)$$

$$\frac{\partial^2 \Phi}{\partial \phi^2} = -m^2 \Phi \Rightarrow \Phi(\phi) = e^{im\phi} \text{ oscillating}$$

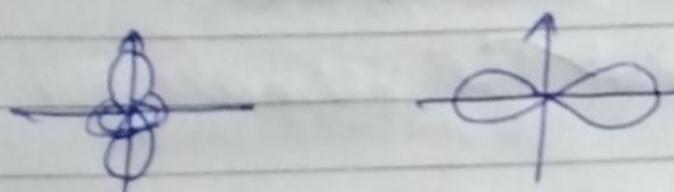
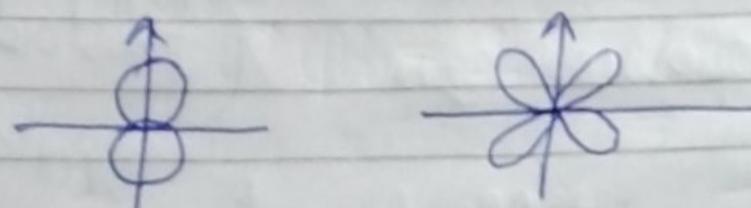
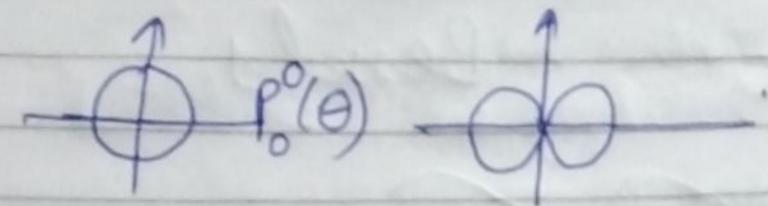
$$\sin \theta \frac{\partial}{\partial \theta} \left(\sin \theta \frac{\partial \Theta}{\partial \theta} \right) + [l(l+1) \sin^2 \theta - m^2] \Theta = 0$$

Sols are associated with Legendre polynomials.

$$P_e^m(\mu) = (1-\mu^2)^{\frac{|m|}{2}} \left(\frac{d}{d\mu} \right)^{|m|}$$

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Associated Legendre's functionals P_l^m



Radial:

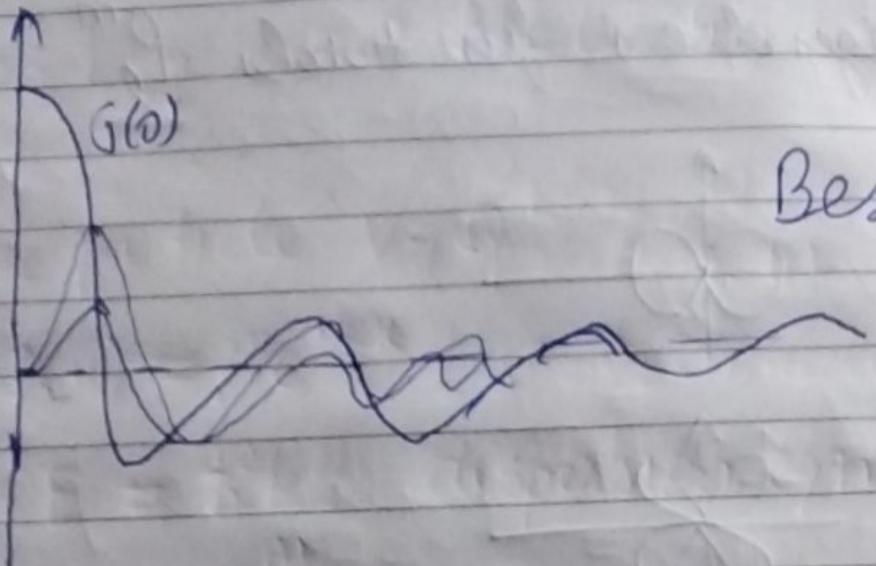
$$\frac{d}{dr} \left(r^2 \frac{dR}{dr} \right) - \frac{2mr^2(V(r)-R)}{\hbar^2} R = l(l+1)R$$

Introduce a new variable: $\mu = rR$

$$u(r) = A r j_\ell(kr) + B r n_\ell(kr)$$

$j_\ell \rightarrow$ spherical Bessel's function.

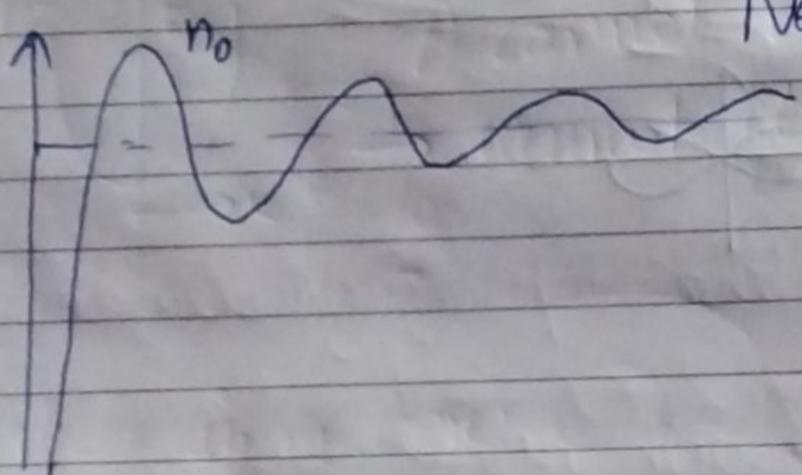
$n_\ell \rightarrow$ spherical Neumann function.



Bessel's

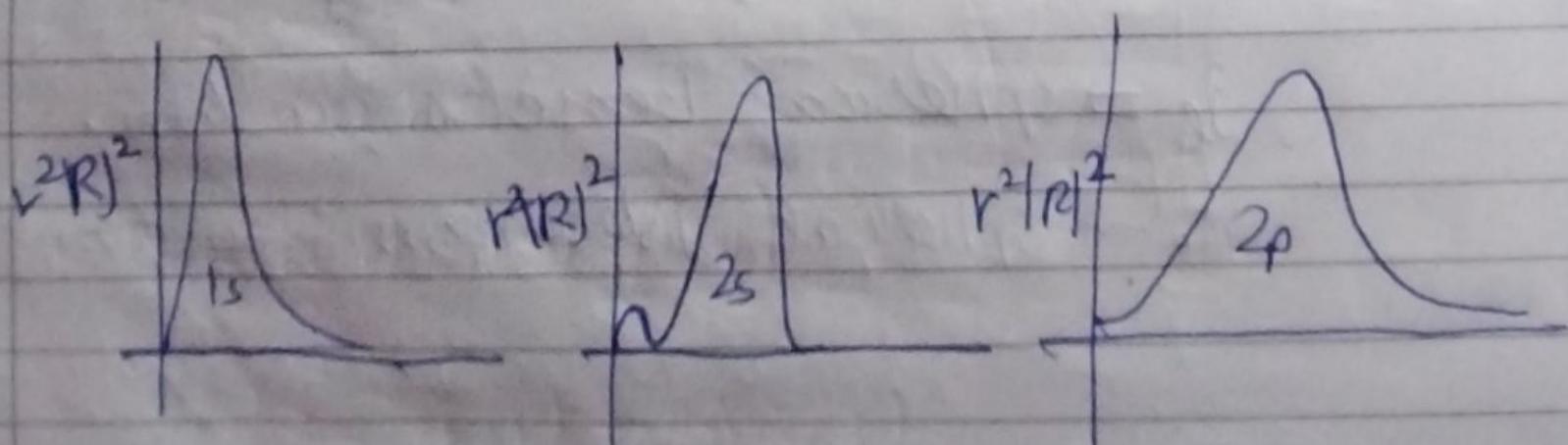
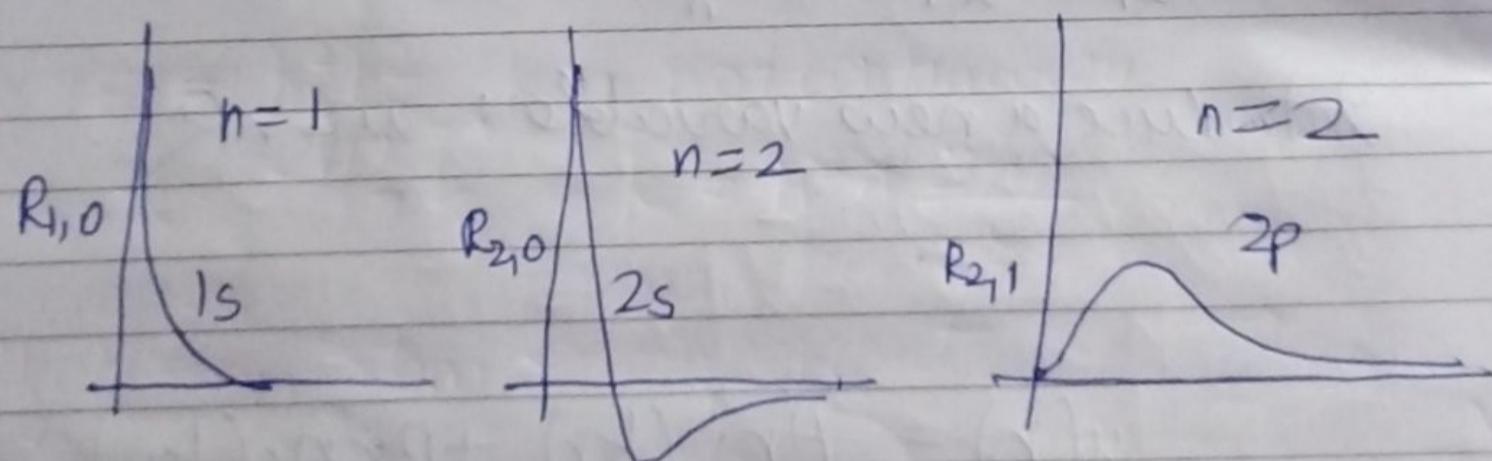
(falstad.com)

orbital viewer



Newman

Radial Wavefunctions:



Energy levels: Only radial dependence!

$$E_n = -\frac{Z^2 E_e}{n^2} \quad E_e = \frac{me^4}{8\epsilon^2 h^2} = 13.6 \text{ eV}$$

ground state energy of H atom.

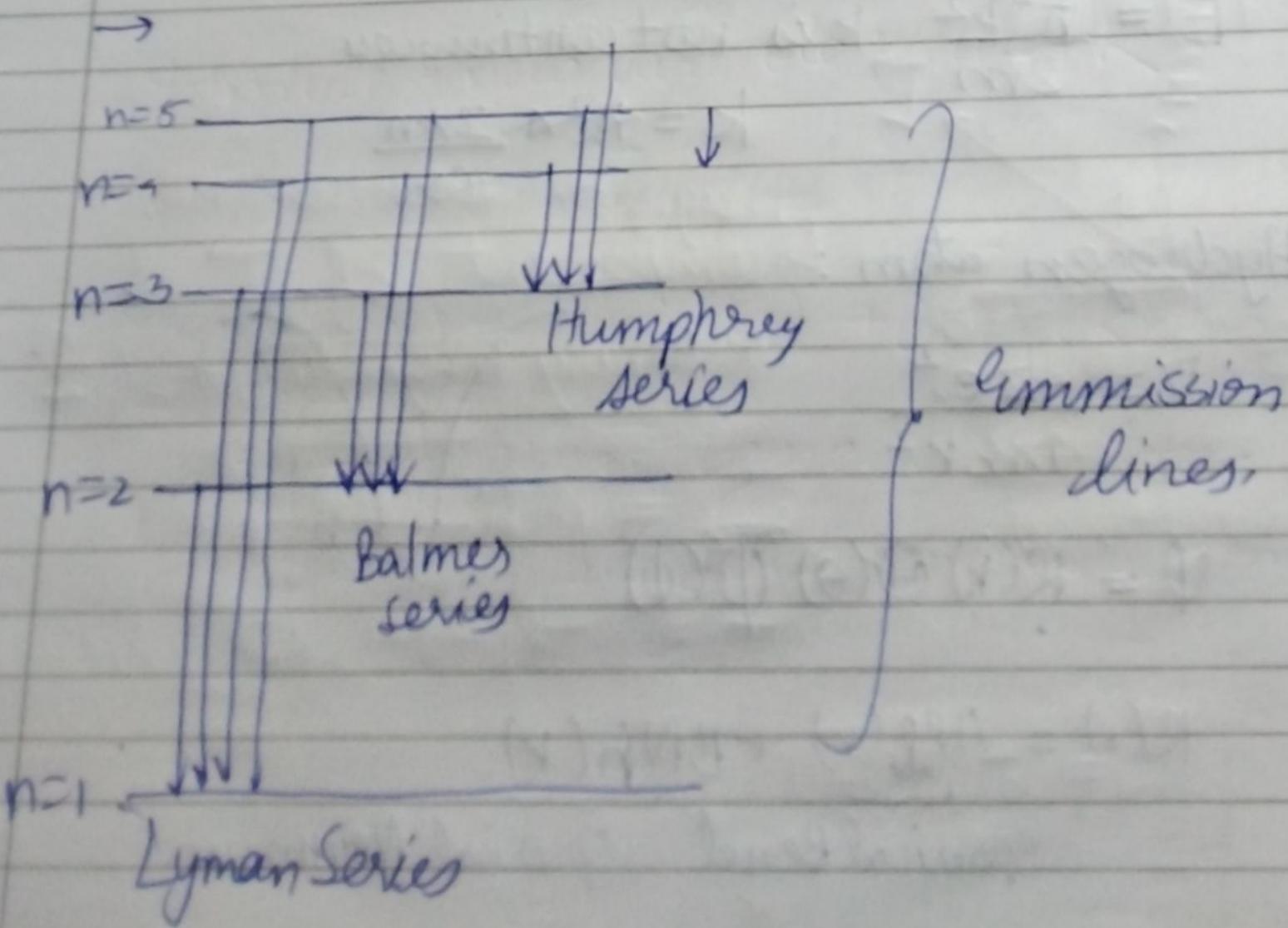
→ higher quantum no. orbitals are degenerate.

$$\Psi = R(r) \Theta(\theta) \Phi(\phi)$$

$\downarrow \quad \downarrow \quad \downarrow e^{-im\phi}$
 $AJ_n(r) \rho^l(\theta)$

→ discrete energy levels:

↳ discrete emmission lines in radiated spectrum.

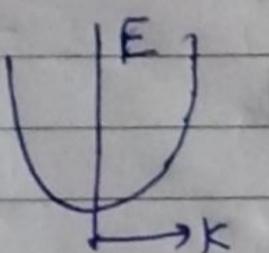


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free particle: $V=0$

$$\Psi(n) = Ae^{ikx} \quad k = \text{continuous} = \frac{2\pi}{\lambda}$$

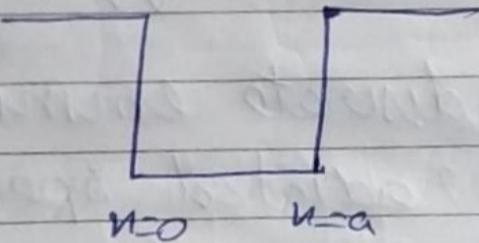
$$E = \frac{\hbar^2 k^2}{2m} \quad \text{dispersion eqn:}$$



A cannot be found
→ cannot be normalized.

Particle in a box:

$$\Psi(n) = 2A_j \sin kx$$



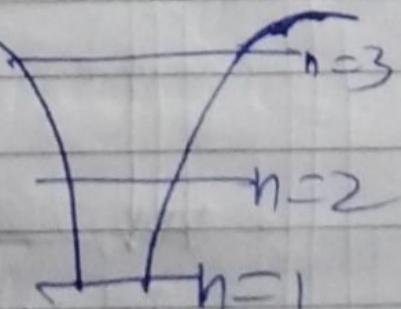
$$E = \frac{\hbar^2 k^2}{2m} \quad k \text{ is not continuous}$$

$$k = \frac{n\pi}{a}$$

Hydrogen atom:

$$V = \frac{-ze^2}{4\pi\epsilon r^2}$$

$$\Psi = R(r) \Theta(\theta) \Phi(\phi)$$



$$R(r) = A \underline{B}_L(r) + B \underline{N}_J(r)$$

spherical Bessel spherical Neumann.

$$E = -\frac{13.6}{n^2} \cdot Z^2 \quad \text{of Hydrogen like species}$$

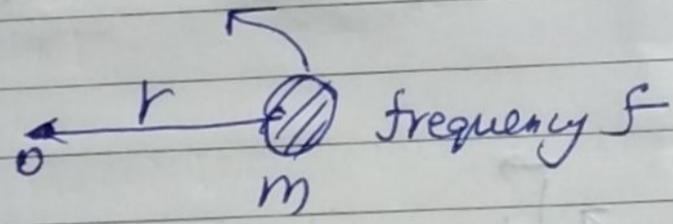
$\Theta_\ell \rightarrow$ Legendre's function.

$$\phi_m \Rightarrow \underline{m e^{im\theta}}$$

* State: energy levels that electrons can possibly take.

→ probability of finding electron at a particular (x, y, z) is univalued.

ANGULAR MOMENTUM:-



$$\vec{L} = I_0 \omega^2 \hat{r} \quad \text{(classical mechanics)} \quad \vec{p} = m v \hat{r}$$

$$\hat{H}\psi = E\psi$$

$$\left[\frac{-\hbar^2}{2m} \nabla^2 + V \right] \psi = E\psi$$

Hamiltonian operator.

$$\hat{L} \psi = L \psi$$

angular momentum operator.

$$L = \hbar \sqrt{\ell(\ell+1)}$$

$$\ell = 0, 1, \dots, n-1$$

$$n=1, l=0, m=0$$

s orbital
spherically symmetric

$[L=0] \Rightarrow$ electron is
not rotating in s orbital

\rightarrow it is there in a spherical volume
with known probability,

$$n=2, l=0$$

$$l=1$$

1 node
radially

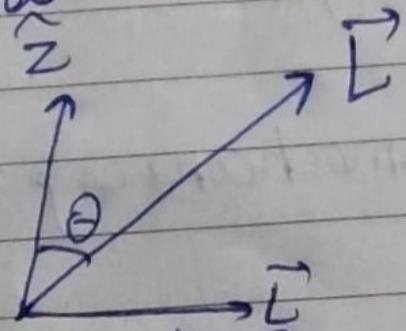
1 node
in θ

p orbital

[dumbbell axial symmetry]



arbitrary
magnetic field.



magnetic quantum no.

Component of L along
arbitrary magnetic field \vec{z}

\vec{z} is quantized

hence θ is also
quantised.

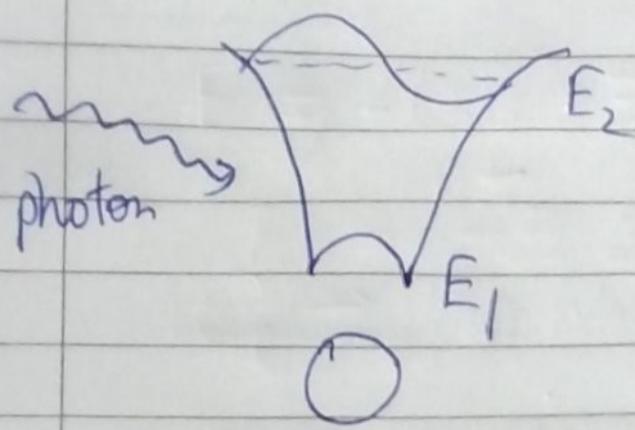
$$L_z = m\hbar$$

$$m = -l, -l+1, \dots, 0, \dots, l$$

$$L = \hbar \sqrt{l(l+1)}$$

wave
photon

\rightarrow ele
 \rightarrow he
 \rightarrow exp



$$E_2 - E_1 = h\nu$$

$\Delta l = 1$ {in absorption}

- electron cannot be excited from $s \rightarrow s$
- has to excite from $s \rightarrow p$
- explains Fraunhofer atomic spectra lines.