

MLL253

Electronic Properties of Materials

$$E(x, t) = E_0 e^{i(\omega t - kx)}$$

phase velocity $v_p = \frac{\omega}{k}$

The avg intensity = $\frac{1}{2} c \epsilon_0 E_0^2$
 \downarrow
 only measurable quantity.

Born interpretation:

~~ψ~~ $\psi(x, y, z, t) \rightarrow$ function

probability to find particle $|\psi(x, y, z, t)|^2$
 at x, y, z at any time t .

\rightarrow Time-independent Schrodinger Egn:

$$\left[\frac{-\hbar^2}{2m} \nabla^2 + V \right] \psi = i\hbar \frac{\partial \psi}{\partial t}$$

$$\frac{-\hbar^2}{2m} \nabla^2 + V = \text{Hamiltonian operator } (\hat{H})$$

\downarrow
 $\text{KE} \quad \text{PE}$

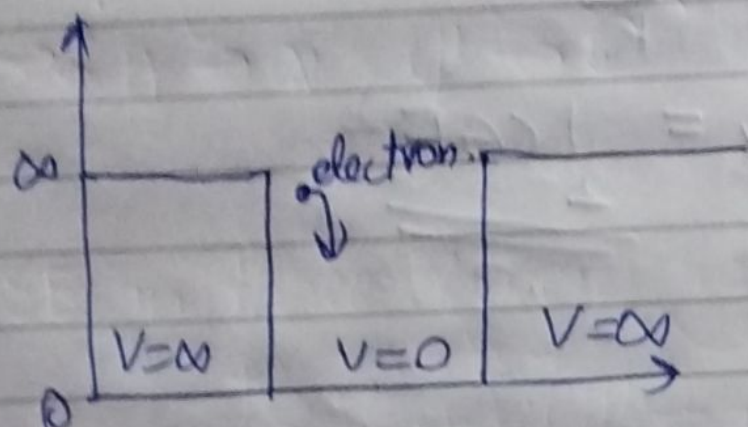
A free particle ($V=0$) everywhere
 $\psi = A \cos\left(\frac{2\pi x}{\lambda} - \omega t\right)$ momentum,
 \rightarrow oscillating everywhere.
 $\frac{2\pi}{\lambda} = \frac{2\pi p}{\hbar} = k$
 $\omega = \frac{h\nu}{\hbar} = \frac{E}{\hbar} \rightarrow$ energy

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$$\psi(x) = Ae^{ikx} \quad \{V=0\}$$

$$\psi^* \psi = A^2$$

Particle In a Box:



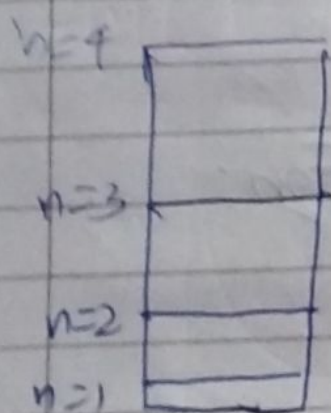
$$\psi(x) = 2A \sin kx$$

$$k = \frac{2\pi n}{a}$$

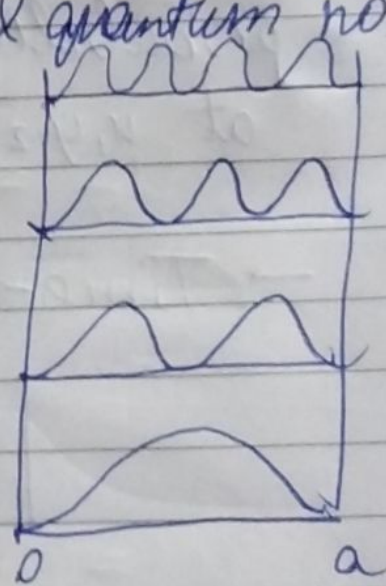
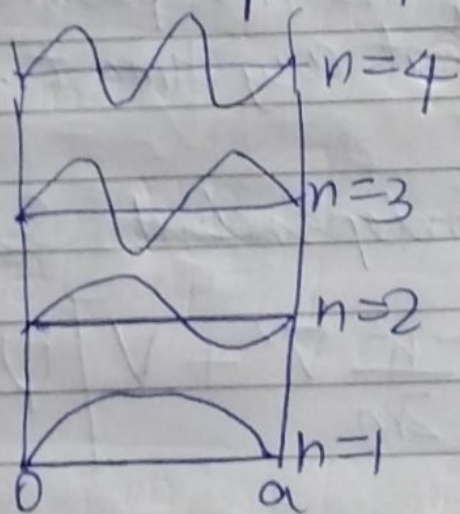
$$E = \frac{\hbar^2 k^2}{2m} = \frac{\hbar^2 n^2}{8ma^2}$$

$$A = \frac{1}{\sqrt{2a}}$$

$n = \text{principal quantum no.}$



Energy levels in well.



→ wavefunctions are no longer continuous.

→ only discrete energy levels exist:

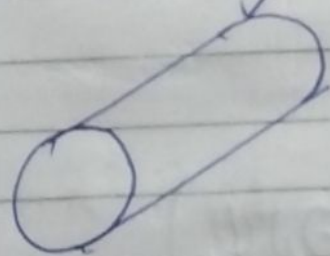
$n=0 \Rightarrow E=0 \Rightarrow$ particle does not exist anywhere
 {Trivial soln}
 \hookrightarrow invalid soln.

Ground state: $n=1$ $E \neq 0$
 non-zero energy
 non-zero velocity

$$\int_{-\infty}^{\infty} \psi^* \psi dr = 1 \quad \text{(Normalization)}$$

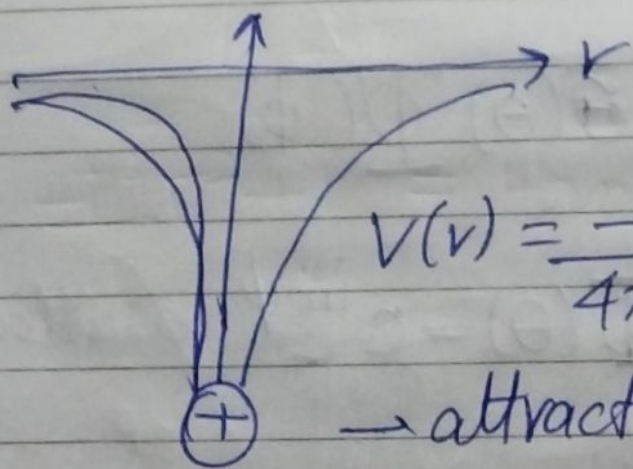
$$\int_0^a 4A^2 \sin^2 kx dx = 1$$

→ looking at the top of cylindrical box.



→ with higher energy, electron starts rotating.

Hydrogen Atom:-



$$V(r) = \frac{-Ze^2}{4\pi\epsilon_0 r} \quad \text{(Coulombic potential)}$$

→ attractive potential from nuclei

→ potential is spherically symmetric

→ $\psi(r, \theta, \phi)$ → univalued.

$$\psi(r, \theta, \phi) = \psi(r, \theta, \phi + 2\pi) = \psi(r, \theta, \phi) \\ = \psi(r, \theta + 2\pi, \phi)$$

→ periodic boundary conditions. ψ univalued

$$r \rightarrow \infty, \psi \rightarrow 0$$

SE in three dimensions:-

$$\left[\frac{\hbar^2}{2m} \nabla^2 + V \right] \psi = i \hbar \frac{\partial \psi}{\partial t} \quad \nabla^2 = \frac{\partial^2}{\partial r^2}$$

In spherical coordinates,

$$\frac{\hbar^2}{2m} \left[\frac{1}{r^2} \frac{\partial}{\partial r} \left(r^2 \frac{\partial \psi}{\partial r} \right) + \frac{1}{r^2 \sin \theta} \frac{\partial}{\partial \theta} \left(\sin \theta \frac{\partial \psi}{\partial \theta} \right) + \frac{1}{r^2 \sin^2 \theta} \frac{\partial^2 \psi}{\partial \phi^2} \right] = i \hbar \frac{\partial \psi}{\partial t}$$

$$\therefore \psi(r, \theta, \phi) = R(r) \Theta(\theta) \Phi(\phi)$$

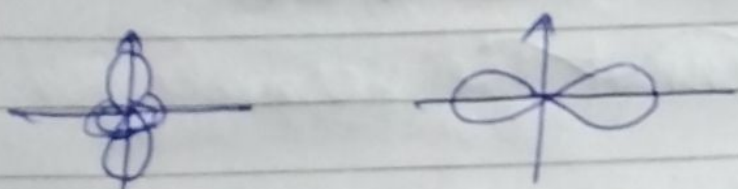
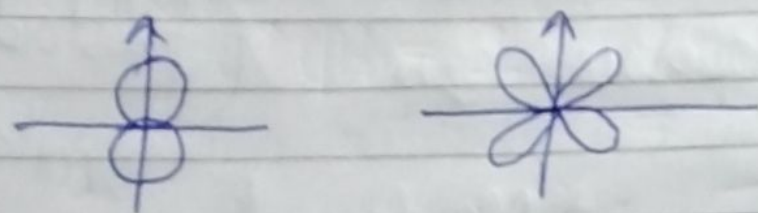
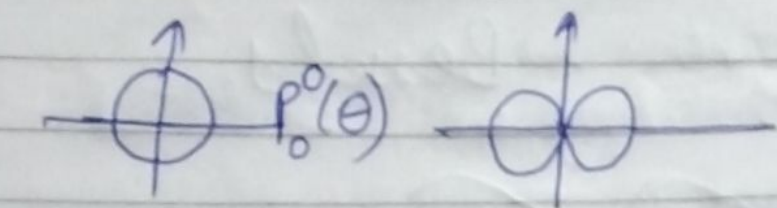
$$\frac{\partial^2 \Phi}{\partial \phi^2} = -m^2 \phi \Rightarrow \Phi(\theta) = e^{im\phi} \text{ oscillating}$$

$$\sin \theta \frac{\partial}{\partial \theta} \left(\sin \theta \frac{\partial \Theta}{\partial \theta} \right) + [l(l+1) \sin^2 \theta - m^2] \Theta = 0$$

sols are associated with Legendre polynomials.

$$P_l^m(u) = (1-u^2)^{\frac{|m|}{2}} \left(\frac{d}{du} \right)^{|m|} P_l^0(u)$$

Associated Legendre's functions P_l^m



Radial:

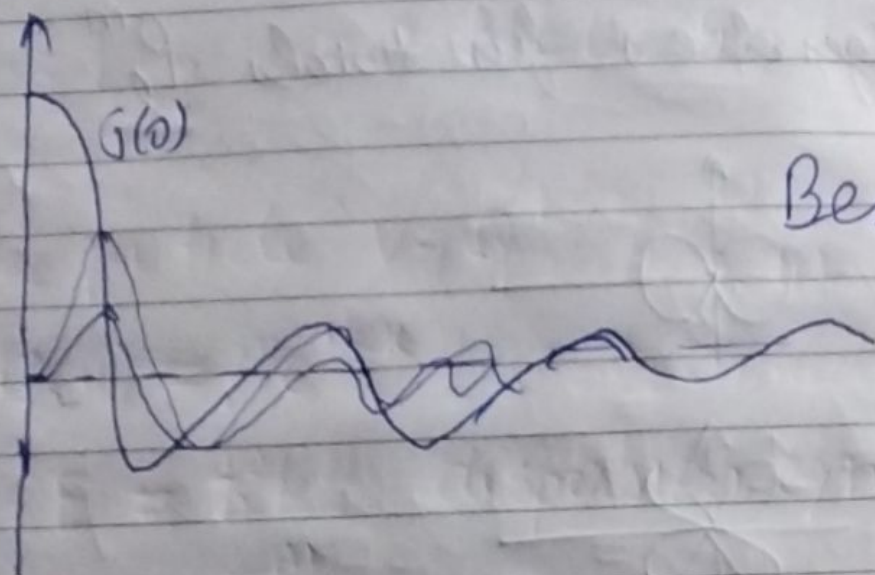
$$\frac{d}{dr} \left(r^2 \frac{dR}{dr} \right) - \frac{2mr^2(V(r) - E)R}{\hbar^2} = l(l+1)R$$

Introduce a new variable: $u = rR$

$$u(r) = Ar j_l(kr) + Br n_l(kr)$$

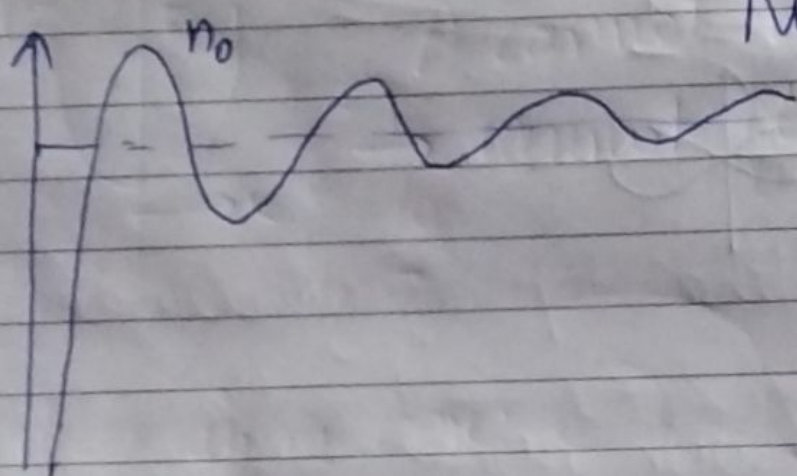
$j_l \Rightarrow$ spherical Bessel's function.

$n_l \Rightarrow$ spherical Neumann's function.



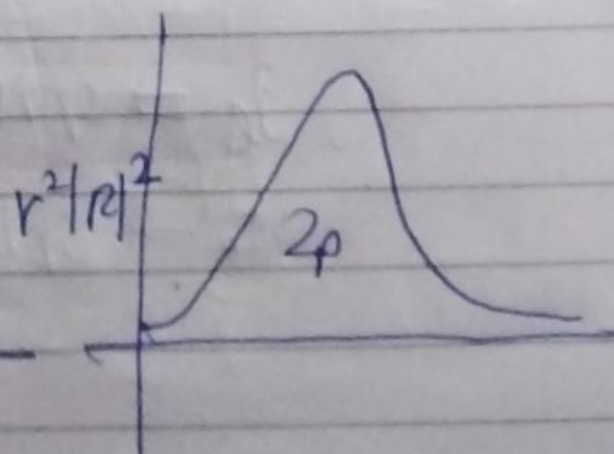
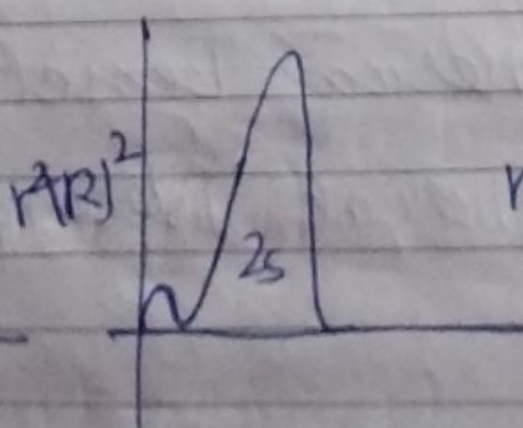
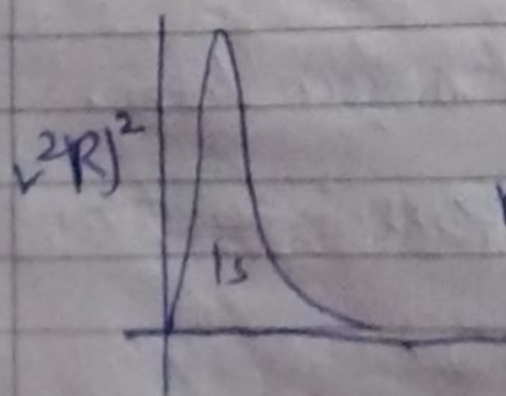
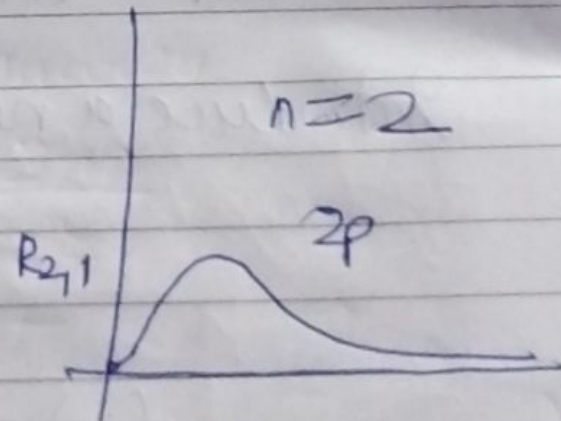
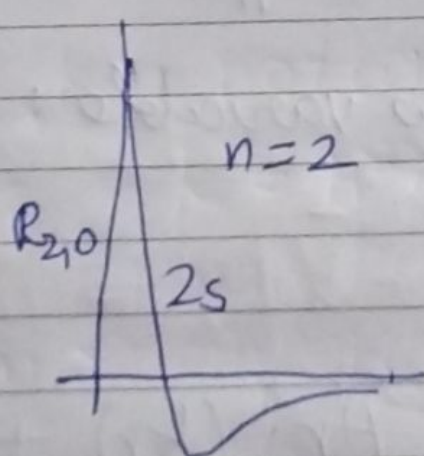
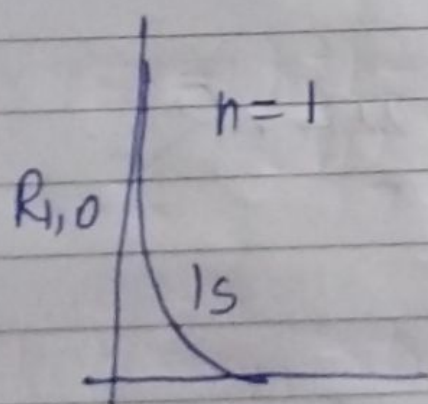
Bessel's

(falstad.com)
orbital viewer



Newman

Radial Wavefunctions:



Energy levels: Only radial dependence!

$$E_n = -\frac{Z^2 E_1}{n^2} \quad E_1 = \frac{me^4}{8\epsilon^2 h^2} = 13.6 \text{ eV}$$

ground state energy of Atom.

→ higher quantum no. orbitals are degenerate.

$$\Psi = R(r) \Theta(\theta) \Phi(\phi)$$

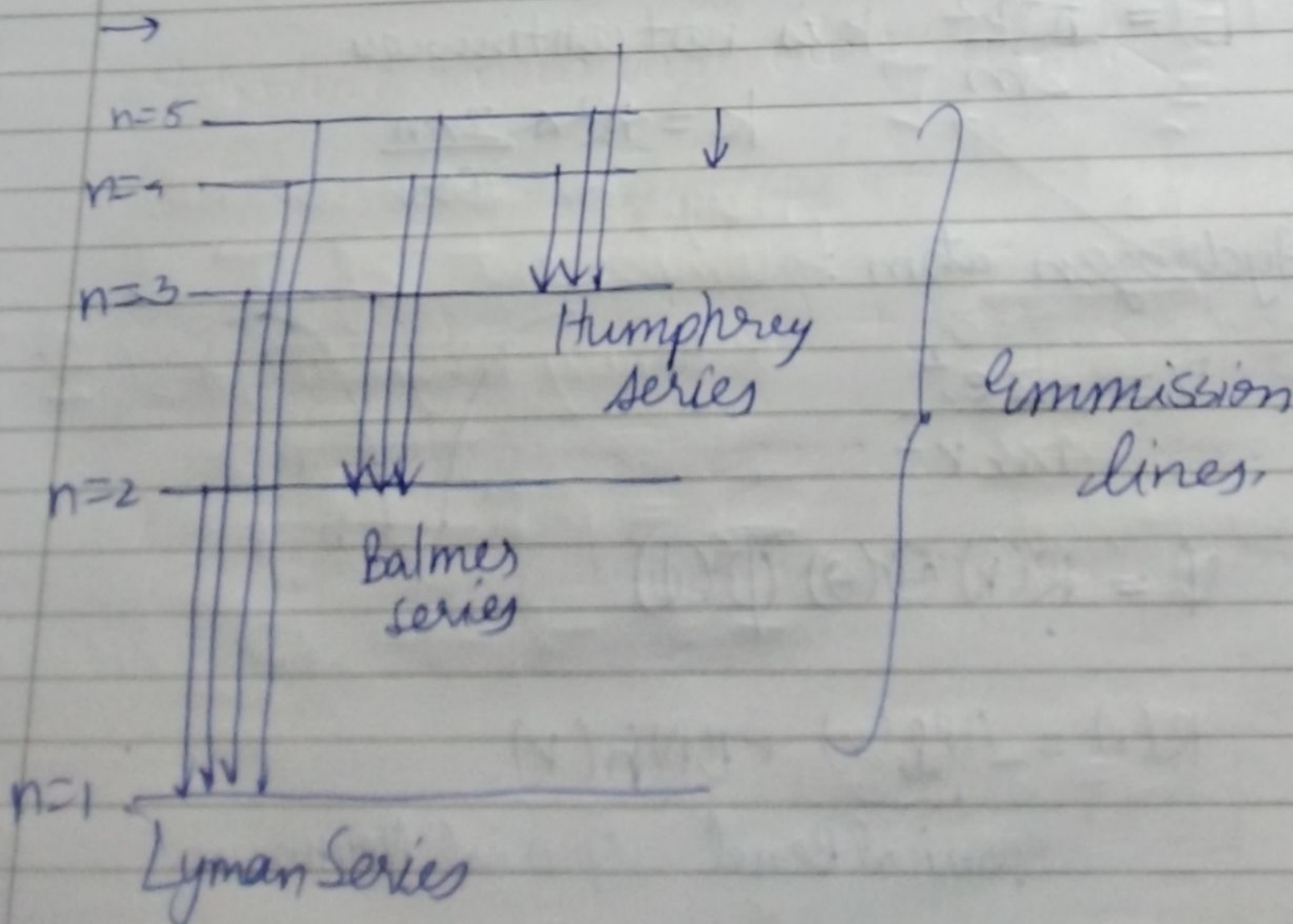
↓
 $AJ_n(r)$

↓
 $P^l(\theta)$

↓
 $e^{-im\phi}$

→ discrete energy levels:

↳ discrete emission ~~lines~~ lines in radiated spectrum.

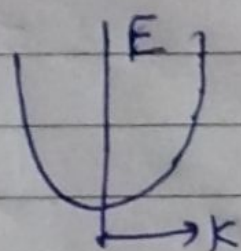


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free particle: $V=0$

$$\psi(x) = Ae^{ikx} \rightarrow k = \text{continuous} = \frac{2\pi}{\lambda}$$

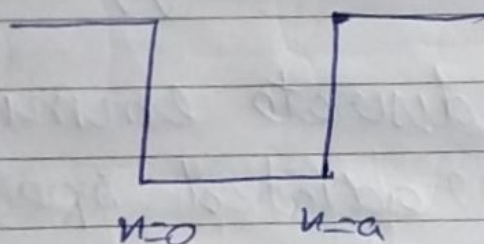
$$E = \frac{\hbar^2 k^2}{2m} \quad \text{dispersion eqn}$$



A cannot be found
→ cannot be normalized.

particle in a box:

$$\psi(x) = 2A \sin kx$$

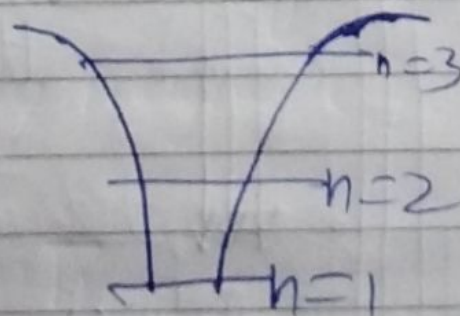


$$E = \frac{\hbar^2 k^2}{2m} \quad k \text{ is not continuous}$$

$$k = \frac{2\pi n}{a}$$

Hydrogen atom:

$$V = \frac{-Ze^2}{4\pi\epsilon r^2}$$



$$\psi = R(r) \Theta(\theta) \Phi(\phi)$$

$$R(r) = A \underline{j_l(r)} + B \underline{N_l(r)}$$

spherical Bessel spherical Neumann.

$$E = -\frac{13.6 \cdot Z^2}{n^2} \quad \text{of Hydrogen like species}$$

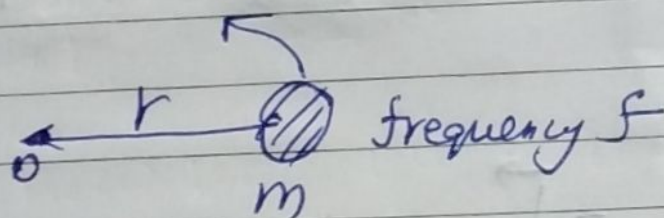
$\Theta_l \rightarrow$ Legendre's function.

$$\phi_m \Rightarrow \underline{m e^{i m \theta}}$$

* State: energy levels that electrons can possibly take.

\rightarrow probability of finding electron at a particular (u, v, z) is univalued.

ANGULAR MOMENTUM:-



$$\vec{L} = I_0 \omega^2 \hat{n} \quad \{\text{classical mechanics}\} = m v r \hat{n}$$

$$\hat{H} \psi = E \psi$$

$$\left[\frac{-\hbar^2 \nabla^2}{2m} + V \right] \psi = E \psi$$

Hamiltonian operator.

energy

$$\hat{L} \psi = L \psi$$

angular momentum operator.

$$L = \hbar \sqrt{l(l+1)}$$

$$* l = 0, 1, \dots, n-1$$

$$n=1, l=0, m=0$$

s orbital
spherically symmetry

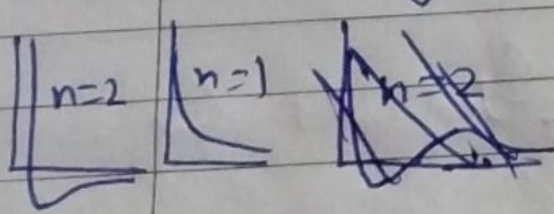
$[L=0] \Rightarrow$ electron is not rotating in s orbital
 \rightarrow it is there in a spherical volume with known probability

$$n=2, l=0$$

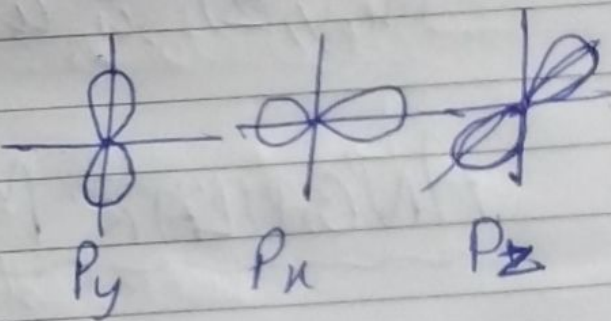
$$l=1$$

1 node radially

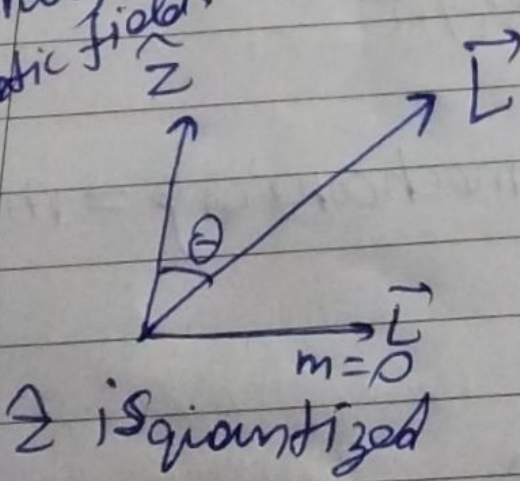
1 node in θ



p orbital
[dumbbell axial symmetry]



arbitrary magnetic field \hat{z}



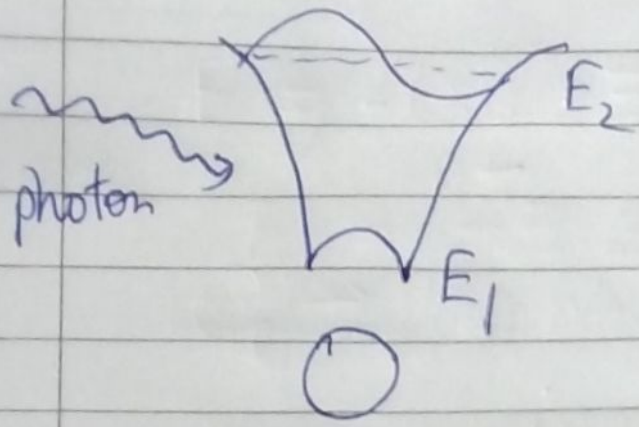
magnetic quantum no.
 \hookrightarrow component of \vec{L} along arbitrary magnetic field \hat{z}

hence θ is also quantised.

$$\vec{L}_z = m\hbar$$

$$m = -l, -l+1, \dots, 0, \dots, l-1, l$$

$$L = \hbar \sqrt{l(l+1)}$$



$$E_2 - E_1 = h\nu$$

$\Delta l = 1$ {in absorption}

- electron cannot be excited from $s \rightarrow s$
- has to excite from $s \rightarrow p$
- explains ~~Fraunhofer~~ atomic spectra lines.