

Rodrigues' Formula, and Detailed Examples

A ChatGPT compilation

0. Functional-Analytic Background (brief)

Let $w(x) > 0$ on an interval (a, b) . Define the weighted Hilbert space

$$L_w^2(a, b) = \{f : (a, b) \rightarrow \mathbb{R} \mid \|f\|^2 = \langle f, f \rangle = \int_a^b |f(x)|^2 w(x) dx < \infty\}.$$

Inner product:

$$\langle f, g \rangle = \int_a^b f(x)g(x)w(x) dx.$$

A sequence of polynomials $\{p_n\}$ with $\deg p_n = n$ is orthogonal w.r.t. w if

$$\langle p_n, p_m \rangle = 0 \quad (n \neq m).$$

Orthogonal polynomials often arise as eigenfunctions of a self-adjoint Sturm–Liouville operator and therefore serve as natural bases for expansions of physical wavefunctions.

1. Rodrigues' Formula — motivation and statement

Rodrigues' formula gives an explicit differential operator representation of classical orthogonal polynomials. A common pattern: the weight w satisfies a Pearson-type identity

$$\frac{d}{dx}[\sigma(x)w(x)] = \tau(x)w(x),$$

with $\sigma(x)$ (usually polynomial degree ≤ 2) and $\tau(x)$ (degree ≤ 1). For the classical families there exists a constant A_n such that

$$p_n(x) = \frac{1}{A_n w(x)} \frac{d^n}{dx^n} [\sigma(x)^n w(x)]$$

produces a polynomial of degree n which satisfies a second-order differential equation of Sturm–Liouville type and is orthogonal w.r.t. w . Below we demonstrate this concretely for Legendre, Hermite and (generalized) Laguerre polynomials.

3. Worked Example: Legendre polynomials

Data

Weight $w(x) = 1$ on $[-1, 1]$. Take $\sigma(x) = 1 - x^2$. Rodrigues formula (classical form):

$$P_n(x) = \frac{1}{2^n n!} \frac{d^n}{dx^n} [(x^2 - 1)^n].$$

(i) Polynomial degree and leading term

$(x^2 - 1)^n$ has degree $2n$. Differentiating n times gives degree n . The leading term: from $(x^2)^n = x^{2n}$,

$$\frac{d^n}{dx^n} x^{2n} = \frac{(2n)!}{n!} x^n,$$

so prefactor $1/(2^n n!)$ yields leading coefficient $\frac{(2n)!}{2^n (n!)^2}$ (conventional normalization).

(ii) Verify the Legendre ODE

Legendre ODE:

$$\frac{d}{dx} [(1 - x^2) P'_n(x)] + n(n + 1) P_n(x) = 0.$$

To verify, substitute $P_n = \frac{1}{2^n n!} Q^{(n)}$ with $Q(x) = (x^2 - 1)^n$. We compute $(1 - x^2) P'_n$ as

$$(1 - x^2) P'_n(x) = \frac{1}{2^n n!} (1 - x^2) Q^{(n+1)}(x).$$

Differentiate once more:

$$\frac{d}{dx} [(1 - x^2) P'_n] = \frac{1}{2^n n!} [(1 - x^2) Q^{(n+2)} - 2x Q^{(n+1)}].$$

One then uses repeated Leibniz expansions of derivatives of $Q = (x^2 - 1)^n$. Because Q contains factor $(x^2 - 1)^n$, many lower-order terms vanish or cancel, and the leftover simplifies to $-n(n + 1) P_n(x)$. (This algebra is mechanical but lengthy; for teaching you may expand for small n to see the cancellations explicitly.)

(iii) Orthogonality (integration by parts)

Let $m < n$. Consider

$$\int_{-1}^1 P_n(x) P_m(x) dx = \frac{1}{2^n n!} \int_{-1}^1 P_m(x) Q^{(n)}(x) dx.$$

Integrate by parts n times, moving derivatives onto P_m . Boundary terms vanish because $Q = (x^2 - 1)^n$ and its first $n - 1$ derivatives are zero at $x = \pm 1$ (factor of $(x \pm 1)^n$ kills them). After n integrations we obtain an integral involving $P_m^{(n)}(x)$, but $\deg P_m = n < m$ so $P_m^{(n)} \equiv 0$. Therefore the integral is zero: orthogonality proven.

(iv) Explicit small n check

Compute directly:

$$\begin{aligned} P_0 &= \frac{1}{1} \frac{d^0}{dx^0} (x^2 - 1)^0 = 1, \\ P_1 &= \frac{1}{2 \cdot 1!} \frac{d}{dx} (x^2 - 1) = x, \\ P_2 &= \frac{1}{4 \cdot 2!} \frac{d^2}{dx^2} (x^2 - 1)^2 = \frac{1}{2} (3x^2 - 1). \end{aligned}$$

These match the standard Legendre polynomials.

4. Worked Example: Hermite polynomials

Data

Weight $w(x) = e^{-x^2}$ on $(-\infty, \infty)$. The Rodrigues formula (physicist's Hermite) is

$$H_n(x) = (-1)^n e^{x^2} \frac{d^n}{dx^n} (e^{-x^2}).$$

This family satisfies

$$H_n''(x) - 2xH_n'(x) + 2nH_n(x) = 0,$$

and orthogonality

$$\int_{-\infty}^{\infty} e^{-x^2} H_n(x) H_m(x) dx = \sqrt{\pi} 2^n n! \delta_{nm}.$$

(i) Polynomial degree and small n

Let $R(x) = e^{-x^2}$. Then $H_n(x) = (-1)^n e^{x^2} R^{(n)}(x)$. Differentiating R n times produces e^{-x^2} times a polynomial of degree n ; multiplication by e^{x^2} cancels the exponential leaving a polynomial of degree n .

Check:

$$H_0 = 1, \quad H_1 = -e^{x^2} \frac{d}{dx} e^{-x^2} = 2x, \quad H_2 = e^{x^2} \frac{d^2}{dx^2} e^{-x^2} = 4x^2 - 2.$$

(ii) Verify the Hermite ODE (sketch)

Using $R' = -2xR$ and induction on derivatives, one shows

$$e^{x^2} \frac{d}{dx} \left(e^{-x^2} \frac{d}{dx} H_n \right) = -2nH_n,$$

which expands to the ODE above. For a direct check, you can substitute the small- n polynomials above into the ODE.

(iii) Orthogonality via integration by parts

For $m < n$,

$$\int_{-\infty}^{\infty} e^{-x^2} H_n(x) H_m(x) dx = (-1)^n \int_{-\infty}^{\infty} H_m(x) (R^{(n)}(x)) dx,$$

with $R = e^{-x^2}$. Integrate by parts n times moving derivatives onto H_m . The boundary terms at $\pm\infty$ vanish because R decays faster than any polynomial; after n steps one gets an integral with $H_m^{(n)}$, which is zero if $m < n$. Thus orthogonality holds.

(iv) Physics connection

Hermite polynomials appear in the quantum harmonic oscillator wavefunctions:

$$\psi_n(x) = \mathcal{N}_n e^{-x^2/2} H_n(x),$$

where the Gaussian factor is part of the weight and H_n supplies the polynomial part of the n -th excited eigenstate.

5. Worked Example: Generalised Laguerre polynomials

Data

Weight $w(x) = x^\alpha e^{-x}$ on $[0, \infty)$ with $\alpha > -1$. Rodrigues formula:

$$L_n^{(\alpha)}(x) = \frac{1}{n!} x^{-\alpha} e^x \frac{d^n}{dx^n} (e^{-x} x^{n+\alpha}).$$

They satisfy

$$xy'' + (\alpha + 1 - x)y' + ny = 0,$$

and orthogonality

$$\int_0^\infty x^\alpha e^{-x} L_n^{(\alpha)}(x) L_m^{(\alpha)}(x) dx = \frac{\Gamma(n + \alpha + 1)}{n!} \delta_{nm}.$$

(i) Polynomial degree and small n

Inside the n -th derivative there is a factor $x^{n+\alpha}$; differentiating n times produces a leading term proportional to x^α . Multiplying by $x^{-\alpha} e^x$ cancels the power and exponential leaving a polynomial of degree n .

Check for α generic:

$$L_0^{(\alpha)}(x) = 1, \quad L_1^{(\alpha)}(x) = -x + (\alpha + 1).$$

(ii) Verify the Laguerre ODE (sketch)

Set $S(x) = e^{-x} x^{n+\alpha}$, so

$$L_n^{(\alpha)}(x) = \frac{1}{n!} x^{-\alpha} e^x S^{(n)}(x).$$

Using product-rule expansions and the Pearson identity for the weight, algebraic cancellation yields the Laguerre ODE with eigenvalue n . As with Legendre, the full expansion is mechanical; checking small n explicitly is instructive.

(iii) Orthogonality

For $m < n$ consider

$$\int_0^\infty x^\alpha e^{-x} L_n^{(\alpha)}(x) L_m^{(\alpha)}(x) dx.$$

Substitute the Rodrigues form of $L_n^{(\alpha)}$ and integrate by parts n times, moving derivatives onto the other polynomial times the weight. Boundary terms at $x \rightarrow \infty$ vanish due to e^{-x} ; at $x = 0$ the behavior is controlled by x^α with $\alpha > -1$, ensuring boundary terms vanish as well. After n integrations the integrand contains derivatives of order n of $L_m^{(\alpha)}$ which vanish if $m < n$. Hence orthogonality holds.

(iv) Quantum application

Laguerre polynomials (with $\alpha = 2\ell + 1$ or related values) appear in radial parts of hydrogen-like wavefunctions, where the radial equation reduces to a Laguerre-type ODE after separation of variables.

6. Practical remarks and exercises

Practical remarks

- The Rodrigues representation is extremely useful for deriving explicit expressions, recurrence relations, and normalization constants.
- The three-term recurrence (Favard) can also be derived from orthogonality and gives a tridiagonal (Jacobi) representation of the multiplication operator M_x in the polynomial basis.
- Generating functions are complementary: once you have Rodrigues (or recurrence) you can often derive the generating function and vice versa.

Exercises (suggested)

1. Using Gram–Schmidt on $\{1, x, x^2, x^3\}$ with weight $w(x) = 1$ on $[-1, 1]$, compute P_0, \dots, P_3 explicitly (show integrals).
2. Verify that Hermite polynomials obtained from Rodrigues satisfy the orthogonality constant $\sqrt{\pi}2^n n!$ by evaluating the Gaussian integrals for small n .
3. Using the Rodrigues form for Laguerre with $\alpha = 0$, compute L_0, L_1, L_2 and verify the orthogonality numerically (a short Python/Mathematica check).
4. Derive the three-term recurrence for Legendre polynomials from the differential equation and show it matches the standard $(n+1)P_{n+1} = (2n+1)xP_n - nP_{n-1}$.