

# **CPEN 400Q / EECE 571Q Lecture 02**

## **Quantum circuits and PennyLane**

Thursday 13 January 2022

# Announcements

- Classes are online until 7 Feb
- Piazza has been setup for the class
- Assignment 0 due on Tuesday before class
  - Instructions have been updated
  - Submit GitHub username/student ID as text response
  - Please update forked repo permissions
  - Make PR to master branch on *your* copy of the repo
- Assignment 1 will be available tomorrow (due in 2 weeks; lots of time)

We learned that qubits are physical systems whose states are represented by complex-valued vectors that are linear combinations of two **basis states**  $|0\rangle$  and  $|1\rangle$ :

$$|\psi\rangle = \alpha|0\rangle + \beta|1\rangle, \quad |\alpha|^2 + |\beta|^2 = 1.$$

## Last time

A qubit lives in a 2-dimensional complex vector space with an **inner product** called a **Hilbert space**. The inner product tells us about the *overlap* between two states.

## Last time

The coefficients in the linear combination (**amplitudes**) tell us the probability of observing a particular basis state  $|\psi_i\rangle$  when we **measure** a qubit.

We can compute these probabilities by projecting onto basis states using the inner product.

$$\Pr(\text{outcome } i) = |\langle \psi_i | \varphi \rangle|^2$$

## Last time

In between state preparation and measurement, we apply  $2 \times 2$  **unitary matrices** (gates/operations) to modify the qubit's state.

A matrix  $U$  is unitary if

$$UU^\dagger = U^\dagger U = \mathbb{1}.$$

Unitaries preserve the length of qubit state vectors, and the angles between them.

## Last time

We wrote some NumPy code to do all this:

```
def ket_0():  
    return np.array([1, 0])  
  
def ket_1():  
    return np.array([0, 1])  
  
def superposition(alpha, beta):  
    return alpha * ket_0() + beta * ket_1()  
  
def apply_op(U, state):  
    return np.dot(U, state)  
  
def apply_ops(list_U, state):  
    for U in list_U:  
        state = np.dot(U, state)  
    return state
```

## Last time

```
def measure(state, num_samples):  
    # Compute using the inner product method  
    prob_0 = np.abs(np.vdot(ket_0(), state)) ** 2  
    prob_1 = np.abs(np.vdot(ket_1(), state)) ** 2  
  
    samples = np.random.choice(  
        [0, 1], size=num_samples, p=[prob_0, prob_1]  
    )  
  
    return samples
```



## Last time

Quantum computing involves preparing a qubit in a particular state, applying one or more unitary operations, and performing a measurement.

```
def quantum_algorithm(alpha, beta, list_U):  
    initial_state = superposition(alpha, beta)  
    state = apply_ops(initial_state, list_U)  
    return measure(state)
```

But doing all of this both by hand or using pure NumPy can be tedious, so today we will shift from NumPy to the quantum software framework PennyLane.

- Implement single-qubit quantum algorithms in PennyLane
- Describe the behaviour of common single-qubit gates
- Calculate the expectation value of an observable
- Perform measurements in other bases

## Quantum functions

Recall three of our quantum gates from last time:

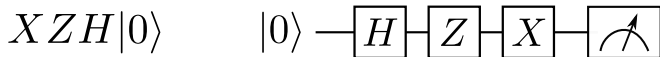
$$H = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 & 1 \\ 1 & -1 \end{pmatrix}, \quad X = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}, \quad Z = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}.$$

We can apply these gates to a qubit and express the computation in matrix form, or as a quantum circuit.

$$XZH|0\rangle \qquad |0\rangle \text{ --- } \boxed{H} \text{ --- } \boxed{Z} \text{ --- } \boxed{X} \text{ --- } \boxed{\text{Measurement}}$$

# Quantum functions

We can also express this circuit as a **quantum function** in PennyLane.



```
import pennylane as qml

def my_quantum_function():
    qml.Hadamard(wires=0)
    qml.PauliZ(wires=0)
    qml.PauliX(wires=0)
    return qml.sample()
```

# Quantum functions

Quantum functions are like normal Python functions, with two special properties:

1. Apply one or more quantum operations

```
import pennylane as qml

def my_quantum_function():
    qml.Hadamard(wires=0) # Apply Hadamard gate to qubit 0
    qml.PauliZ(wires=0)   # Apply Pauli Z gate to qubit 0
    qml.PauliX(wires=0)   # Apply Pauli X gate to qubit 0
    return qml.sample()
```

Q: Why wires? A: PennyLane can be used for continuous-variable quantum computing, which does not use qubits.

# Quantum functions

Quantum functions are like normal Python functions, with two special properties:

1. Apply one or more quantum operations
2. Return a measurement on one or more qubits

```
import pennylane as qml

def my_quantum_function():
    qml.Hadamard(wires=0)
    qml.PauliZ(wires=0)
    qml.PauliX(wires=0)
    return qml.sample() # Return measurement samples
```

Quantum functions are executed on **devices**. These can be either *simulators*, or *actual quantum hardware*.

```
import pennylane as qml  
  
dev = qml.device('default.qubit', wires=1, shots=100)
```

This creates a device of type **'default.qubit'** with 1 qubit that returns 100 measurement samples for anything that is executed.

# Quantum functions

A **QNode** (quantum node) is an object that binds a quantum function to a device, and executes it.

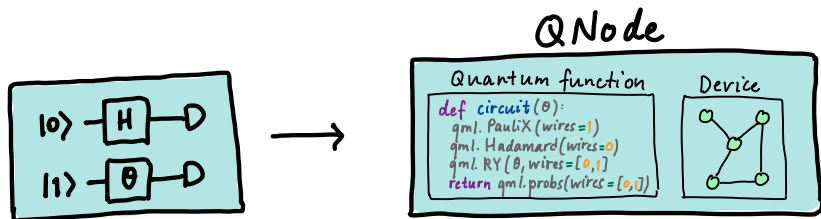


Image credit: [https://pennylane.ai/qml/glossary/quantum\\_node.html](https://pennylane.ai/qml/glossary/quantum_node.html)



# Quantum nodes

```
import pennylane as qml

dev = qml.device('default.qubit', wires=1, shots=100)

def my_quantum_function():
    qml.Hadamard(wires=0)
    qml.PauliZ(wires=0)
    qml.PauliX(wires=0)
    return qml.sample()
```

With these two components, we can create and execute a QNode.

```
# Create a QNode
my_qnode = qml.QNode(my_quantum_function, dev)

# Execute the QNode
result = my_qnode()
```

# Hands-on with QNodes

You probably have some questions...

1. Where's the state?

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1. Where's the state?
  - Inside the device!
2. What happens to the gates?
  - Operations are recorded onto a “tape”
  - The QNode constructs the tape when it is called
  - The tape is then executed on the device.



## Single-qubit unitary operations

## More quantum gates

So far, we know 3 gates that do the following:

$$\begin{aligned} X|0\rangle &= |1\rangle, & X|1\rangle &= |0\rangle, \\ Z|0\rangle &= |0\rangle, & Z|1\rangle &= -|1\rangle, \\ H|0\rangle &= \frac{1}{\sqrt{2}}(|0\rangle + |1\rangle), & H|1\rangle &= \frac{1}{\sqrt{2}}(|0\rangle - |1\rangle). \end{aligned}$$

But a general qubit state looks like

$$|\psi\rangle = \alpha|0\rangle + \beta|1\rangle,$$

where  $\alpha$  and  $\beta$  are *complex numbers* (such that  $|\alpha|^2 + |\beta|^2 = 1$ ).

How do we make the rest?

## Z rotations

Consider the operation  $Z$ :

$$Z|0\rangle = |0\rangle, \quad Z|1\rangle = -|1\rangle.$$

Apply this to a superposition:

The *sign* of the amplitude on the  $|1\rangle$  state has changed.

We know that  $-1 = e^{i\pi}$ :

What if instead of  $\pi$ , we used an arbitrary angular parameter?

The extra  $e^{i\theta}$  is called a **relative phase**.

## $Z$ rotations

The “proper” form of this rotation is

$$RZ(\theta) = e^{-i\frac{\theta}{2}Z} = \begin{pmatrix} e^{-i\frac{\theta}{2}} & 0 \\ 0 & e^{i\frac{\theta}{2}} \end{pmatrix}$$

Exercise: expand out the exponential of  $Z$  to obtain the matrix representation.

## $S$ and $T$

Two other special cases:  $\theta = \pi/2$ , and  $\theta = \pi/4$ .

$$S = RZ(\pi/2) = \begin{pmatrix} e^{-i\frac{\pi}{4}} & 0 \\ 0 & e^{i\frac{\pi}{4}} \end{pmatrix} \sim \begin{pmatrix} 1 & 0 \\ 0 & i \end{pmatrix}$$

$$T = RZ(\pi/4) = \begin{pmatrix} e^{-i\frac{\pi}{8}} & 0 \\ 0 & e^{i\frac{\pi}{8}} \end{pmatrix} \sim \begin{pmatrix} 1 & 0 \\ 0 & e^{i\frac{\pi}{4}} \end{pmatrix}$$

$S$  is part of a special group called the **Clifford group**.

$T$  is used in universal gate sets for fault-tolerant QC.

## $X$ and $Y$ rotations

$RZ$  changes the phase, but not the magnitudes of the amplitudes.  
How do we change those?

$RX$ , and  $RY$  rotations...

## “Rotations”?

There is a reason we are calling these rotations.

$$|\psi\rangle = \alpha|0\rangle + \beta|1\rangle$$

We can rewrite  $\alpha = ae^{i\phi}$  and  $\beta = be^{i\omega}$  where  $a, b$  are real-valued numbers:

Factor out the  $e^{i\phi}$  (a **global phase**):



## “Rotations”?

The global phase doesn't matter though!

It does not affect the measurement outcome probabilities.

## “Rotations”?

If the global phase doesn't matter...

$$|\psi\rangle = e^{i\phi} \left( a|0\rangle + be^{i(\omega-\phi)}|1\rangle \right) \sim a|0\rangle + be^{i(\omega-\phi)}|1\rangle$$

Relabel  $\varphi = \omega - \phi$ :

$$|\psi\rangle = a|0\rangle + be^{i\varphi}|1\rangle$$

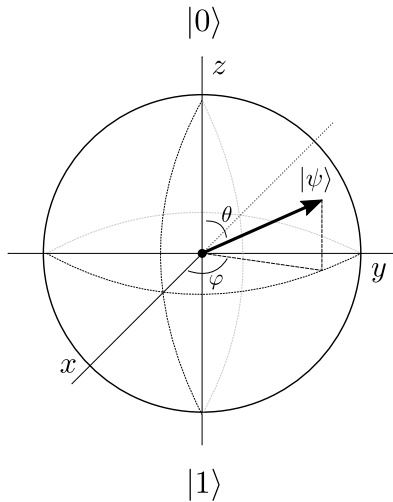
## “Rotations”?

Normalization tells us that  $a^2 + b^2 = 1$ . What else has this relationship?

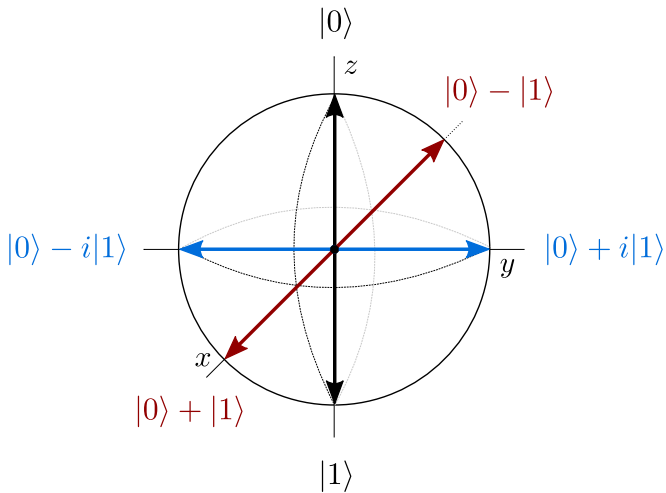
We can rewrite as:

So any single-qubit state can be specified by two angular parameters... just like points on a sphere!

## Rotations: the Bloch sphere



## Rotations: the Bloch sphere



# Rotations: the Bloch sphere

$RX$ ,  $RY$ , and  $RZ$  correspond visually to rotations about their respective axes.

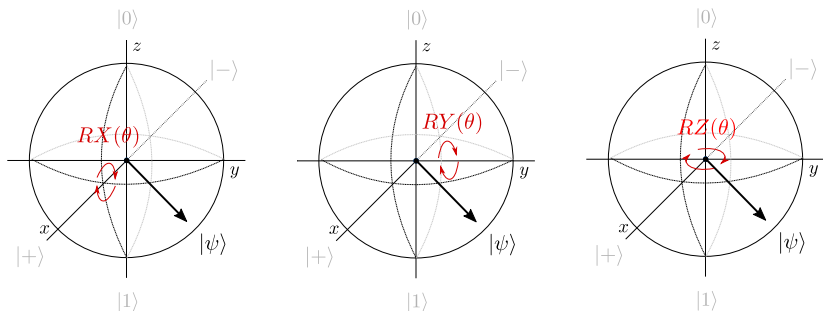
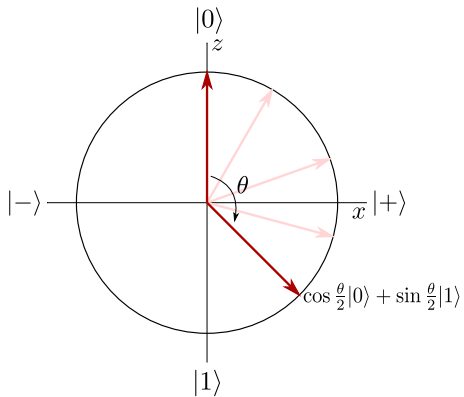
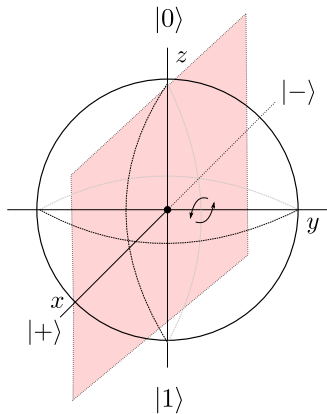


Image credit: Codebook node I.6

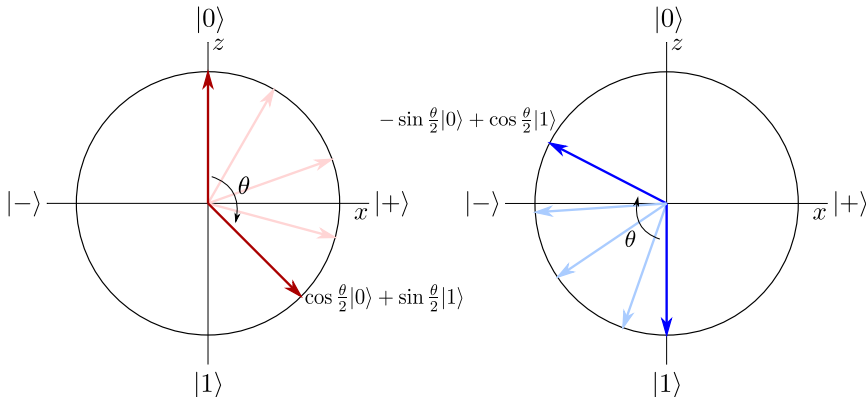
# Rotations: $RY$



## Rotations: $RY$

The matrix representation of  $RY$  is

$$RY(\theta) = \begin{pmatrix} \cos \frac{\theta}{2} & -\sin \frac{\theta}{2} \\ \sin \frac{\theta}{2} & \cos \frac{\theta}{2} \end{pmatrix}$$

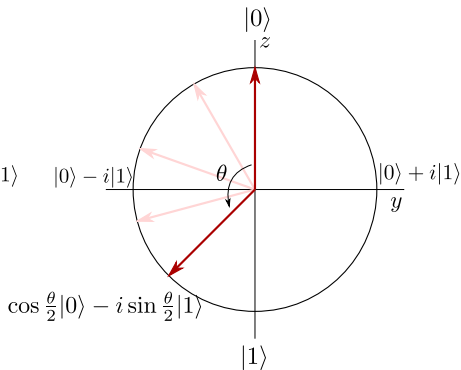
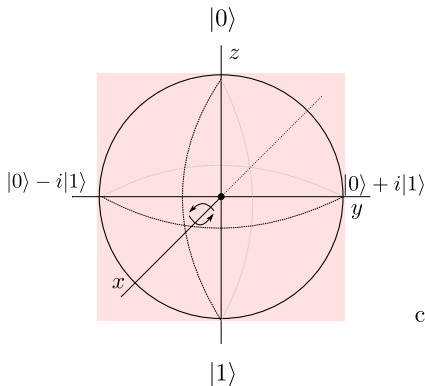




## Rotations: $RX$

$RX$  is similar but has complex components:

$$RX(\theta) = \begin{pmatrix} \cos \frac{\theta}{2} & -i \sin \frac{\theta}{2} \\ -i \sin \frac{\theta}{2} & \cos \frac{\theta}{2} \end{pmatrix}$$



## Pauli rotations

These unitary operations are called **Pauli rotations**.

	Math	Matrix	Code	Special cases
$RZ$	$e^{-i\frac{\theta}{2}Z}$	$\begin{pmatrix} e^{-i\frac{\theta}{2}} & 0 \\ 0 & e^{i\frac{\theta}{2}} \end{pmatrix}$	<code>qml.RZ</code>	$Z(\pi), S(\pi/2), T(\pi/4)$
$RY$	$e^{-i\frac{\theta}{2}Y}$	$\begin{pmatrix} \cos \frac{\theta}{2} & -\sin \frac{\theta}{2} \\ \sin \frac{\theta}{2} & \cos \frac{\theta}{2} \end{pmatrix}$	<code>qml.RY</code>	$Y(\pi)$
$RX$	$e^{-i\frac{\theta}{2}X}$	$\begin{pmatrix} \cos \frac{\theta}{2} & -i \sin \frac{\theta}{2} \\ i \sin \frac{\theta}{2} & \cos \frac{\theta}{2} \end{pmatrix}$	<code>qml.RX</code>	$X(\pi), SX(\pi/2)$

# Adjoints

We can rotate forwards, or backwards by negating the angle. But there is a more general way of rotating backwards. In PennyLane, we can compute adjoints of operations *and* entire quantum functions using `qml.adjoint`:

```
def some_function(x):  
    qml.RZ(Z, wires=0)  
  
def apply_adjoint(x):  
    qml.adjoint(qml.S)(wires=0)  
    qml.adjoint(some_function)(x)
```

`qml.adjoint` is a special type of function called a **transform**. We will cover transforms in more detail around beginning of week 4.

Hands-on time...

What about  $H$ ?

$$H = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 & 1 \\ 1 & -1 \end{pmatrix}$$

This does not have the form of  $RX$ ,  $RY$ , or  $RZ$ .

But, we can use a combination of these to make an  $H$  (actually, just need two of the three).

## Deep dive: unitary operations

The  $n \times n$  unitary matrices are a mathematical group under matrix multiplication,  $U(n)$ :

1. Closure: for  $U, V$  unitary,  $UV$  is also unitary
2. Associativity:  $(UV)W = U(VW)$
3. Identity:  $\mathbb{1}$
4. Inverses:  $U^{-1} = U^\dagger$

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Any unitary matrix can be written in terms of a finite set of real-valued parameters:

$$U(\phi, \theta, \omega) = e^{i\alpha} \begin{pmatrix} e^{-i(\phi+\omega)/2} \cos(\theta/2) & -e^{i(\phi-\omega)/2} \sin(\theta/2) \\ e^{-i(\phi-\omega)/2} \sin(\theta/2) & e^{i(\phi+\omega)/2} \cos(\theta/2) \end{pmatrix}$$

## Universal gate sets: Pauli rotations

With just  $RZ$  and  $RY$  (or  $RZ/RX$ ,  $RY/RX$ ), we can implement *any single-qubit unitary operation*<sup>1</sup>:

$$U = e^{i\alpha} RZ(\omega) RY(\theta) RZ(\phi)$$

$\{RZ, RY\}$  is **universal** for single-qubit quantum computing.

Hands-on...

For more fun: do text exercises in Codebook node I.3 and I.7.

---

<sup>1</sup>Note that the  $\alpha$  technically doesn't matter.

## Universal gate sets: $H$ and $T$

With just  $H$  and  $T$ , we can approximate any single-qubit rotation up to arbitrary accuracy. For example, we can implement  $RZ(0.1)$  up to accuracy  $10^{-10}$ :

```
→ gridsynth 0.1 -d 10  
HTHTHTHTHTSHTHTHTHTSHTSHTHTHTSHTSHTHTSHTHTHTSHTSHTHTSHTSHTSHTS  
HTHTHTHTHTHTHTHTHTHTSHTSHTSHTSHTSHTSHTSHTHTSHTSHTSHTSHTHTHTSHTSHTHT  
SHTSHTSHTHTHTHTSHTHTHTSHTSHTHTHTHTSHTHTHTSHTSHTSHTSHTSHTHTSHTHTHT  
HTHTHTHTHTSHTHTHTSHTHTSHTHTHTSHTSHTHTSHTSHTHTXWWW
```

This was generated using the newsynth Haskell package:  
<https://www.mathstat.dal.ca/~selinger/newsynth/>



## Universal gate sets: $H$ and $T$

Or to accuracy  $10^{-100}$ :

[illegible]

...we'll talk more about this in a few weeks when we discuss *quantum compilation*.

Measurement: observables and expectation values

# Sampling

So far, we've learned how take measurement samples in the computational basis.

```
dev = qml.device('default.qubit', wires=1, shots=100)

@qml.qnode(dev)
def rotate_with_rz(theta):
    qml.Hadamard(wires=0)
    qml.RZ(theta, wires=0)
    return qml.sample()
```

What else can we do?

# Measurement outcome probabilities

Compute the measurement outcome probabilities from the results:

```
dev = qml.device('default.qubit', wires=1, shots=100)

@qml.qnode(dev)
def rotate_with_rz(theta):
    qml.Hadamard(wires=0)
    qml.RZ(theta, wires=0)
    return qml.probs()
```

## Extract the state

Since we are running on a simulator...

```
# Note that we did NOT specify shots: analytic mode
dev = qml.device('default.qubit', wires=1)

@qml.qnode(dev)
def rotate_with_rz(theta):
    qml.Hadamard(wires=0)
    qml.RZ(theta, wires=0)
    return qml.state()
```

(Can analytically compute probabilities too. But of course we cannot do this with a real device!)

# Observables

Generally, we are interested in measuring real, physical quantities. In physics, these are called **observables**. They are represented by Hermitian matrices. An operator (matrix)  $H$  is Hermitian if

$$H = H^\dagger$$

**Why Hermitian?** The possible measurement outcomes are given by the eigenvalues of the operator, and eigenvalues of Hermitian operators are **real**.

# Observables

Example:

$$Z = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$$

$Z$  is Hermitian:

Its eigensystem is

Example:

$$X = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}$$

$X$  is Hermitian and its (normalized) eigensystem is



Example:

$$Y = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix}$$

$Y$  is Hermitian and its (normalized) eigensystem is

# Expectation values

When we measure  $X$ ,  $Y$ , or  $Z$  on a state, for each shot we will get one of the eigenstates (/eigenvalues). If we take multiple shots, what do we expect to see *on average*?

Analytically, the **expectation value** of measuring the observable  $M$  given the state  $|\psi\rangle$  is

$$\langle M \rangle = \langle \psi | M | \psi \rangle.$$

## Expectation values: analytical

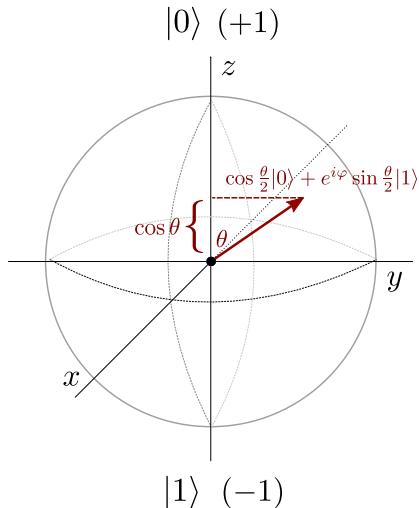
Example: consider the quantum state

$$|\psi\rangle = \frac{1}{2}|0\rangle - i\frac{\sqrt{3}}{2}|1\rangle.$$

Let's compute the expectation value of  $Y$ :

# Expectation values and the Bloch sphere

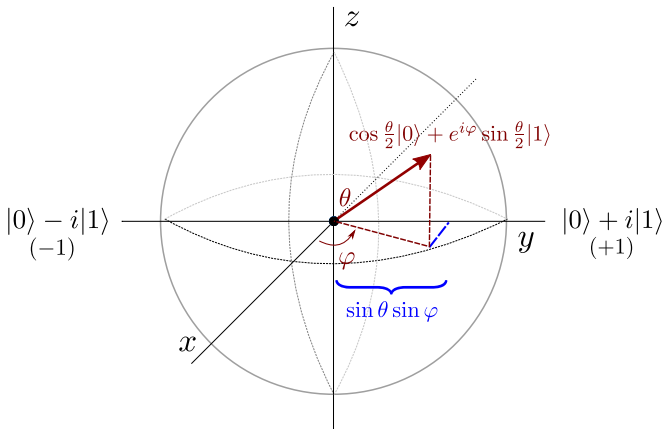
The Bloch sphere offers us some more insight into what a projective measurement is.



$$\begin{aligned} |\psi\rangle &= \cos \frac{\theta}{2} |0\rangle + e^{i\varphi} \sin \frac{\theta}{2} |1\rangle \\ Z|\psi\rangle &= \cos \frac{\theta}{2} |0\rangle - e^{i\varphi} \sin \frac{\theta}{2} |1\rangle \\ \langle\psi| Z|\psi\rangle &= \cos^2 \frac{\theta}{2} - \sin^2 \frac{\theta}{2} \\ &= \cos \theta \end{aligned}$$

## Expectation values and the Bloch sphere

In this picture, we can visualize measurement in different bases by projecting onto different axes.



Exercise: derive this by computing  $\langle \psi | Y | \psi \rangle$ .

## Expectation values: from measurement data

Let's compute the expectation value of  $Z$  for the following circuit using 10 samples:

```
dev = qml.device('default.qubit', wires=1, shots=10)

@qml.qnode(dev)
def circuit():
    qml.RX(2*np.pi/3, wires=0)
    return qml.sample()
```

Results might look something like this:

```
[1, 1, 1, 0, 1, 1, 1, 0, 1, 1]
```

## Expectation values: from measurement data

The expectation value pertains to the measured eigenvalue; recall  $Z$  eigenstates are

$$\begin{aligned}\lambda_1 &= +1, & |\psi_1\rangle &= \begin{pmatrix} 1 \\ 0 \end{pmatrix} \\ \lambda_2 &= -1, & |\psi_2\rangle &= \begin{pmatrix} 0 \\ 1 \end{pmatrix}\end{aligned}$$

So when we observe  $|0\rangle$ , this is eigenvalue  $+1$  (and if  $|1\rangle$ ,  $-1$ ).  
Our samples shift from

$$[1, 1, 1, 0, 1, 1, 1, 0, 1, 1]$$

to

$$[-1, -1, -1, 1, -1, -1, -1, 1, -1, -1]$$

## Expectation values: from measurement data

The expectation value is the weighted average of this, where the weights are the eigenvalues:

$$\langle Z \rangle = \frac{1 \cdot n_1 + (-1) \cdot n_{-1}}{N}$$

where

- $n_1$  is the number of +1 eigenvalues
- $n_{-1}$  is the number of -1 eigenvalues
- $N$  is the total number of shots

For our example,  $\langle Z \rangle = -0.6$ .



# Expectation values

Let's do this in PennyLane instead:

```
dev = qml.device('default.qubit', wires=1)

@qml.qnode(dev)
def measure_z():
    qml.RX(2*np.pi/3, wires=0)
    return qml.expval(qml.PauliZ(0))
```

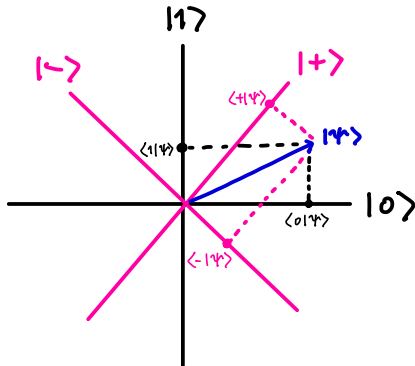
So far we've seen 4 ways of extracting information out of a QNode:

1. `qml.state()`
2. `qml.probs(wires=x)`
3. `qml.sample()`
4. `qml.expval(observable)`

The first three all return results of measurements taken with respect to the computational basis; and most hardware only allows for computational basis measurements. How can we measure with respect to *different bases* with that restriction? (and what does that mean?)

# Basis rotations

What does it mean to measure in a different bases? Projective measurement with respect to a different set of orthonormal states. For example,  $\{|+\rangle, |-\rangle\}$  are an orthonormal basis.



## Basis rotations

Use a basis rotation to “trick” the quantum computer into measuring in a different basis.

Suppose we want to measure in the  $Y$  basis:

$$|i\rangle = \frac{1}{\sqrt{2}} (|0\rangle + i|1\rangle), \quad |-i\rangle = \frac{1}{\sqrt{2}} (|0\rangle - i|1\rangle).$$

Unitary operations preserve length *and* angles between normalized quantum state vectors.

There exists some unitary transformation that will convert between these eigenvectors, and the eigenvectors of  $Z$  (the basis in which we will take the measurement).

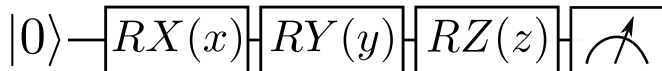
Let's try to turn

$$\begin{aligned}|i\rangle &= \frac{1}{\sqrt{2}}(|0\rangle + i|1\rangle) \rightarrow |0\rangle \\ | - i\rangle &= \frac{1}{\sqrt{2}}(|0\rangle - i|1\rangle) \rightarrow |1\rangle\end{aligned}$$

That way, if we measure and observe  $|0\rangle$ , we know that this was previously  $|i\rangle$  in the  $Y$  basis (and similarly for  $|1\rangle$ ).

## Basis rotations: hands-on

Let's run the following circuit, and measure in the  $Y$  basis



# Recap

- Implement single-qubit quantum algorithms in PennyLane
- Describe the behaviour of common single-qubit gates
- Calculate the expectation value of an observable
- Perform measurements in other bases

What topics did you find unclear today?

# Next time

## Content:

- Multi-qubit states, operations, and measurements
- Entanglement

## Action items:

1. Finish Assignment 0 (due before class Tuesday)
2. Start on Assignment 1 once posted (you can do problem 1)
3. Quiz next class

## Recommended reading:

- Codebook nodes I.5-I.10
- Nielsen & Chuang 4.2