# CPEN 400Q / EECE 571Q Lecture 11 The quantum Fourier transform and quantum phase estimation

Tuesday 15 February 2022

### Announcements

- Project group / topic selection due today
- Please upgrade to PennyLane v0.21; new requirements.txt file will be included later with Quiz 5 and with Assignment 3.

Quiz 5 after class today.

### Last time

We introduced the quantum Fourier transform, and saw how it is the analog of the classical inverse discrete Fourier transform.

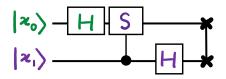
$$QFT|x\rangle = rac{1}{\sqrt{N}} \sum_{k=0}^{N-1} \omega^{xk} |k
angle$$

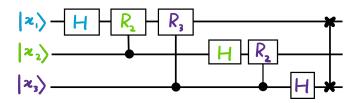
$$QFT = \frac{1}{\sqrt{N}} \begin{pmatrix} 1 & 1 & 1 & \cdots & 1\\ 1 & \omega & \omega^{2} & \cdots & \omega^{N-1}\\ 1 & \omega^{2} & \omega^{4} & \cdots & \omega^{2(N-1)}\\ \vdots & \vdots & \vdots & \ddots & \vdots\\ 1 & \omega^{N-1} & \omega^{2(N-1)} & \cdots & \omega^{(N-1)(N-1)} \end{pmatrix}$$

where for *n* qubits,  $N=2^n$ , and  $\omega=e^{2\pi i/N}$ 

### Last time

We saw the circuits for some special cases. For 1 qubit, it is just the Hadamard. For 2 and 3 qubits:





## Quantum Fourier transform

I showed you what the general form of the circuit looked like:

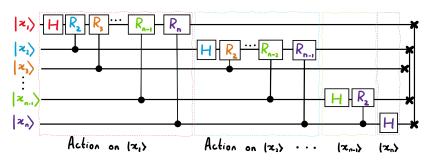


Image credit: Xanadu Quantum Codebook node F.3

## Learning outcomes

- Derive the QFT circuit and implement it in PennyLane
- Describe the steps of the quantum phase estimation (QPE) subroutine
- Use the QFT to implement QPE

## Review: fractional binary notation

## Example

Let  $k = k_1 k_2 k_3 k_4 = 0.1001$ . The numerical value of this is:

$$0.1001 = \frac{1}{2} + \frac{0}{2^2} + \frac{0}{2^3} + \frac{1}{2^4}$$
$$= \frac{1}{2} + \frac{1}{16}$$
$$= 0.5625$$

We need this for the QFT because in the exponent, we have

$$\frac{1}{\sqrt{N}}\sum_{k=0}^{N-1}\omega^{\times k}|k\rangle = \frac{1}{\sqrt{N}}\sum_{k=0}^{N-1}e^{2\pi i \times (k/N)}|k\rangle$$

and k/N is a fractional value.

What we are going to show is that

$$\frac{1}{\sqrt{N}}\sum_{k=0}^{N-1}\omega^{xk}|k\rangle$$

can be rewritten in the following factorized form:

$$\frac{\left(|0\rangle+e^{2\pi i 0.x_n}|1\rangle\right)\left(|0\rangle+e^{2\pi i 0.x_{n-1}x_n}|1\rangle\right)\cdots\left(|0\rangle+e^{2\pi i 0.x_1\cdots x_n}|1\rangle\right)}{\sqrt{N}}$$

Then, we will see how this form reveals to us the circuit that creates this state!

Start by rewriting k/N using fractional binary.

$$|x\rangle \rightarrow \frac{1}{\sqrt{N}} \sum_{k=0}^{N-1} \omega^{xk} |k\rangle$$

$$= \frac{1}{\sqrt{N}} \sum_{k=0}^{N-1} e^{2\pi i x (k/N)} |k\rangle$$

$$= \frac{1}{\sqrt{N}} \sum_{k_1=0}^{1} \cdots \sum_{k_n=0}^{1} e^{2\pi i x (\sum_{\ell=1}^{n} k_{\ell} 2^{-\ell})} |k_1 \cdots k_n\rangle$$

(keeping the last equation from the previous slide)

$$= \frac{1}{\sqrt{N}} \sum_{k_{1}=0}^{1} \cdots \sum_{k_{n}=0}^{1} e^{2\pi i x \left(\sum_{\ell=1}^{n} k_{\ell} 2^{-\ell}\right)} |k_{1} \cdots k_{n}\rangle$$

$$= \frac{1}{\sqrt{N}} \sum_{k_{1}=0}^{1} \cdots \sum_{k_{n}=0}^{1} \bigotimes_{\ell=1}^{n} e^{2\pi i x k_{\ell} 2^{-\ell}} |k_{\ell}\rangle$$

$$= \frac{1}{\sqrt{N}} \bigotimes_{\ell=1}^{n} \left(\sum_{k_{\ell}=0}^{1} e^{2\pi i x k_{\ell} 2^{-\ell}} |k_{\ell}\rangle\right)$$

(keeping the last equation from the previous slide)

$$= \frac{1}{\sqrt{N}} \bigotimes_{\ell=1}^{n} \left( \sum_{k_{\ell}=0}^{1} e^{2\pi i x k_{\ell} 2^{-\ell}} |k_{\ell}\rangle \right)$$

$$= \frac{1}{\sqrt{N}} \bigotimes_{\ell=1}^{n} \left( |0\rangle + e^{2\pi i x 2^{-\ell}} |1\rangle \right)$$

$$= \frac{\left( |0\rangle + e^{2\pi i 0.x_{n}} |1\rangle \right) \left( |0\rangle + e^{2\pi i 0.x_{n-1} x_{n}} |1\rangle \right) \cdots \left( |0\rangle + e^{2\pi i 0.x_{1} \cdots x_{n}} |1\rangle \right)}{\sqrt{N}}$$

Starting with the state

$$|x\rangle = |x_1 \cdots x_n\rangle,$$

apply a Hadamard to qubit 1:

$$\frac{1}{\sqrt{2}}\left(|0\rangle + e^{2\pi i 0.x_1}|1\rangle\right)|x_2\cdots x_n\rangle$$

$$|x_1\rangle$$
 — $H$ —

$$x_3\rangle$$
 ———

$$\frac{1}{\sqrt{2}}\left(|0\rangle+e^{2\pi i0.x_1}|1\rangle\right)|x_2\cdots x_n\rangle$$

If  $x_1=0$ ,  $e^0=1$  and we get the  $|+\rangle$  state.

If 
$$x_1 = 1$$
,  $e^{2\pi i(1/2)} = e^{\pi i} = -1$  and we get the  $|-\rangle$  state.

$$|x_{1}\rangle - H - |x_{2}\rangle - |x_{3}\rangle - |x_{3}\rangle - |x_{n-1}\rangle - |x_{n}\rangle - |x_{n}\rangle$$

We are trying to make a state that looks like this:

$$|x\rangle \rightarrow \frac{\left(|0\rangle + e^{2\pi i 0.x_n}|1\rangle\right) \left(|0\rangle + e^{2\pi i 0.x_{n-1}x_n}|1\rangle\right) \cdots \left(|0\rangle + e^{2\pi i 0.x_1 \cdots x_n}|1\rangle\right)}{\sqrt{N}}$$

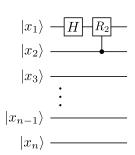
Every qubit has a different *phase* on the  $|1\rangle$  state. We are going to need some way of creating this.

We define the gate:

$$R_k = \begin{pmatrix} 1 & 0 \\ 0 & e^{2\pi i/2^k} \end{pmatrix}$$

Now let's apply a controlled  $R_2$  gate from qubit 2 to qubit 1

$$R_2 = \begin{pmatrix} 1 & 0 \\ 0 & e^{2\pi i/2^2} \end{pmatrix}$$



The first qubit picks up a phase:

$$\begin{split} \frac{1}{\sqrt{2}} \left( |0\rangle + e^{2\pi i 0.x_1} |1\rangle \right) |x_2 \cdots x_n\rangle &\rightarrow \frac{1}{\sqrt{2}} \left( |0\rangle + e^{2\pi i 0.x_1} e^{\frac{2\pi i}{2^2} x_2} |1\rangle \right) |x_2 \cdots x_n\rangle \\ &= \frac{1}{\sqrt{2}} \left( |0\rangle + e^{2\pi i 0.x_1 x_2} |1\rangle \right) |x_2 \cdots x_n\rangle \end{split}$$

Now let's apply a controlled  $R_3$  gate from qubit 3 to qubit 1

$$R_3 = \begin{pmatrix} 1 & 0 \\ 0 & e^{2\pi i/2^3} \end{pmatrix}$$

$$|x_{1}\rangle - H - R_{2} - R_{3}$$

$$|x_{2}\rangle - \vdots$$

$$|x_{n-1}\rangle - \vdots$$

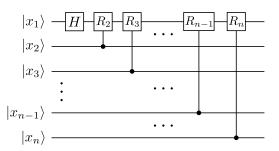
$$|x_{n}\rangle - \vdots$$

The first qubit picks up another phase:

$$\frac{1}{\sqrt{2}}\left(|0\rangle + e^{2\pi i 0.x_1 x_2}|1\rangle\right)|x_2 \cdots x_n\rangle \rightarrow \frac{1}{\sqrt{2}}\left(|0\rangle + e^{2\pi i 0.x_1 x_2 x_3}|1\rangle\right)|x_2 \cdots x_n\rangle$$

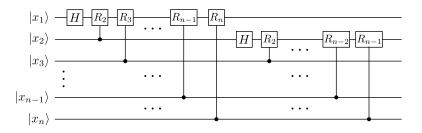
We can apply a controlled  $R_4$  from the fourth qubit, etc. up to the n-th qubit to get

$$\frac{1}{\sqrt{2}}\left(|0\rangle + e^{2\pi i 0.x_1 x_2 \cdots x_n}|1\rangle\right)|x_2 \cdots x_n\rangle$$



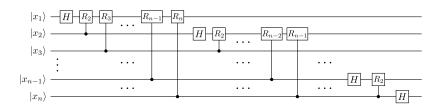
Next, ignore the first qubit and do the same thing with the second qubit: apply H, and then controlled rotations from every qubit from 3 to n to get

$$\frac{1}{\sqrt{2}^2} \left( |0\rangle + e^{2\pi i 0.x_1 x_2 \cdots x_n} |1\rangle \right) \left( |0\rangle + e^{2\pi i 0.x_2 \cdots x_n} |1\rangle \right) |x_3 \cdots x_n\rangle$$



If we do this for all qubits, we get something similar to that big ugly state from earlier:

$$|x\rangle \rightarrow \frac{\left(|0\rangle + e^{2\pi i 0.x_1 \cdots x_n}|1\rangle\right) \cdots \left(|0\rangle + e^{2\pi i 0.x_{n-1}x_n}|1\rangle\right) \left(|0\rangle + e^{2\pi i 0.x_n}|1\rangle\right)}{\sqrt{N}}$$

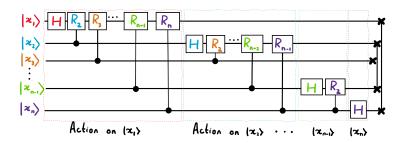


This is almost what we want: the order of the qubits is backwards. This is easily fixed with some SWAP gates.

## Quantum Fourier transform

So the QFT can be implemented using:

- n Hadamard gates
- n(n-1)/2 controlled rotations
- |n/2| SWAP gates if you care about the order



The number of gates is *polynomial in n*, so this can be implemented efficiently on a quantum computer! Let's try it...

# Eigenvalues of unitary matrices

Fun fact: eigenvalues of unitary matrices are complex numbers with magnitude 1.

Proof: Let  $\lambda_k$  be the eigenvalue associated with eigenvector  $|k\rangle$  of a unitary U:

$$U|k\rangle = \lambda_k |k\rangle$$

We can take the conjugate transpose of this equation:

$$\langle k | U^{\dagger} = \langle k | \lambda_k^*$$

Multiply the two sides together:

$$\langle k | U^{\dagger} U | k \rangle = \langle k | \lambda_k^* \lambda_k | k \rangle$$
  
 $\langle k | k \rangle = |\lambda_k|^2 \langle k | k \rangle$   
 $1 = |\lambda_k|^2$ 

# Eigenvalues of unitary matrices

So we can write

$$\lambda_k = e^{2\pi i \theta_k}$$

where  $\theta_k$  is some phase angle such that  $|\theta_k| \leq 1$ .

What if we want to *learn* an unknown  $\theta_k$ ?

## Eigenvalues of unitary matrices

Idea: apply U to the relevant eigenvector, because that's "what makes the phase come out".

$$U|k\rangle = e^{2\pi i\theta_k}|k\rangle$$

...but this is an unobservable global phase!

We have to do something different: eigenvalue estimation, or quantum phase estimation (QPE).

Given a unitary U and one of its eigenvectors  $|k\rangle$ , estimate the value of  $\theta_k$  such that

$$U|k\rangle = e^{2\pi i\theta_k}|k\rangle$$

### Must determine:

- How to design a circuit that extracts the  $\theta_k$
- To what precision can we estimate it
- What to do if we don't know a  $|k\rangle$  in advance

(You will explore the last two in your homework!)

Let *U* be an *n*-qubit unitary; therefore  $|k\rangle$  is an *n*-qubit state.

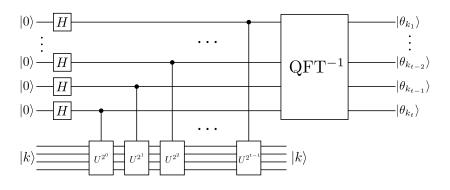
Assume for now that  $\theta_k$  can be represented *exactly* using t bits in *fractional binary*:

$$\theta_k = 0.\theta_{k_1} \cdots \theta_{k_t}$$

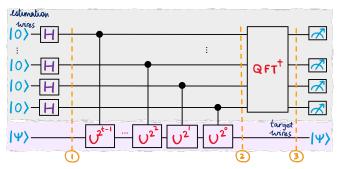
Fact: We can construct a circuit with n + t qubits that recover the value of  $\theta_k$  exactly by:

- 1. Preparing n qubits in state  $|k\rangle$
- 2. Applying controlled applications of U to those qubits in a special way
- 3. Applying the inverse QFT

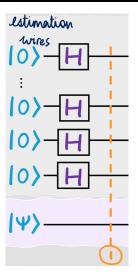
This is one version of the circuit:



The order of the controlled operations is irrelevant though, so you may see this too:



Why does this work? Let's analyze the state at points 1, 2, and 3 above.



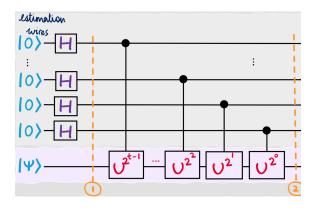
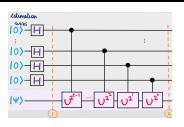


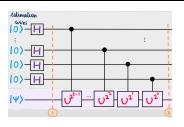
Image credit: Xanadu Quantum Codebook node P.2



### Consider the top-most qubit:

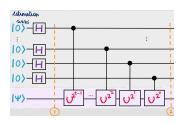
$$\begin{split} (CU)^{2^{t-1}} \left( \frac{1}{\sqrt{2}} (|0\rangle + |1\rangle) |+\rangle^{\otimes t-1} |k\rangle \right) &= (CU)^{2^{t-1}} \left( \frac{1}{\sqrt{2}} |0\rangle |+\rangle^{\otimes t-1} |k\rangle \right) \\ &+ (CU)^{2^{t-1}} \left( \frac{1}{\sqrt{2}} |1\rangle |+\rangle^{\otimes t-1} |k\rangle \right) \\ &= \left( \frac{1}{\sqrt{2}} |0\rangle |+\rangle^{\otimes t-1} |k\rangle \right) \\ &+ \left( \frac{1}{\sqrt{2}} |1\rangle |+\rangle^{\otimes t-1} (e^{2\pi i \theta_k})^{2^{t-1}} |k\rangle \right) \end{split}$$

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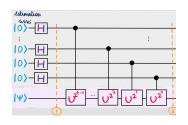
### Use phase kickback

$$\begin{split} &\left(\frac{1}{\sqrt{2}}|0\rangle|+\rangle^{\otimes t-1}|k\rangle\right)+\left(\frac{1}{\sqrt{2}}|1\rangle|+\rangle^{\otimes t-1}\left(e^{2\pi i\theta_{k}}\right)^{2^{t-1}}|k\rangle\right)\\ &=\left(\frac{1}{\sqrt{2}}|0\rangle|+\rangle^{\otimes t-1}|k\rangle\right)+\left(\frac{1}{\sqrt{2}}(e^{2\pi i\theta_{k}})^{2^{t-1}}|1\rangle|+\rangle^{\otimes t-1}|k\rangle\right)\\ &=\frac{1}{\sqrt{2}}(|0\rangle+\left(e^{2\pi i\theta_{k}}\right)^{2^{t-1}}|1\rangle)|+\rangle^{\otimes t-1}|k\rangle \end{split}$$



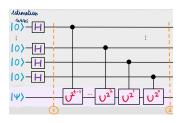
### What is happening in the exponent?

$$\begin{split} \left(e^{2\pi i\theta_{k}}\right)^{2^{t-1}} &= e^{2\pi i\theta_{k} \cdot 2^{t-1}} \\ &= e^{2\pi i\left(\frac{\theta_{k_{1}}}{2^{1}} + \frac{\theta_{k_{2}}}{2^{2}} + \cdots \frac{\theta_{k_{t}}}{2^{t}}\right) \cdot 2^{t-1}} \\ &= e^{2\pi i\left(2^{t-2}\theta_{k_{1}} + 2^{t-3}\theta_{k_{2}} + \cdots \frac{\theta_{k_{t}}}{2}\right)} \\ &= e^{2\pi i\frac{\theta_{k_{t}}}{2}} \\ &= e^{2\pi i0 \cdot \theta_{k_{t}}} \end{split}$$



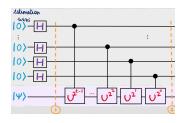
So we have the combined state:

$$\frac{1}{\sqrt{2}}(|0\rangle + (e^{2\pi i\theta_k})^{2^{t-1}}|1\rangle)|+\rangle^{\otimes t-1}|k\rangle = \frac{1}{\sqrt{2}}(|0\rangle + e^{2\pi i0.\theta_{k_t}}|1\rangle)|+\rangle^{\otimes t-1}|k\rangle$$



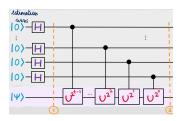
Let's do the second-last qubit (ignore what happens to others for now):

$$(CU)^2\left(|+\rangle^{\otimes t-2}\frac{1}{\sqrt{2}}(|0\rangle+|1\rangle)|+\rangle|k\rangle\right)=|+\rangle^{\otimes t-2}\frac{1}{\sqrt{2}}(|0\rangle+e^{2\pi i\theta_k\cdot 2}|1\rangle)|+\rangle|k\rangle$$



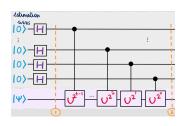
Again check the exponent...

$$\begin{split} (e^{2\pi i\theta_k})^2 &= e^{2\pi i\theta_k \cdot 2} \\ &= e^{2\pi i (\frac{\theta_{k_1}}{2^1} + \frac{\theta_{k_2}}{2^2} + \cdots \frac{\theta_{k_t}}{2^t}) \cdot 2} \\ &= e^{2\pi i (\theta_{k_1} + \frac{\theta_{k_2}}{2} + \cdots \frac{\theta_{k_t}}{2^{t-1}})} \\ &= e^{2\pi i 0 \cdot \theta_{k_2} \cdots \theta_{k_t}} \end{split}$$



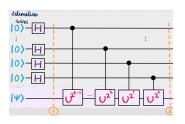
So we have the combined state:

$$|+
angle^{\otimes t-2}rac{1}{\sqrt{2}}(|0
angle+(e^{2\pi i heta_k})^2|1
angle)|+
angle|k
angle=|+
angle^{\otimes t-2}rac{1}{\sqrt{2}}(|0
angle+e^{2\pi i0. heta_{k_2}\cdots heta_{k_t}}|1
angle)|+
angle|k
angle$$



Can show in the same way that for the last qubit

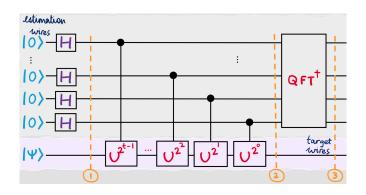
$$|+\rangle^{\otimes t-1}\frac{1}{\sqrt{2}}(|0\rangle+(e^{2\pi i\theta_k})|1\rangle)|k\rangle=|+\rangle^{\otimes t-1}\frac{1}{\sqrt{2}}(|0\rangle+e^{2\pi i0.\theta_{k_1}\cdots\theta_{k_t}}|1\rangle)|k\rangle$$



### After step 2, we have the state

$$\frac{1}{\sqrt{2}}(|0\rangle + e^{2\pi i 0.\theta_{k_t}}|1\rangle) \cdots \frac{1}{\sqrt{2}}(|0\rangle + e^{2\pi i 0.\theta_{k_2} \cdots \theta_{k_t}}|1\rangle) \frac{1}{\sqrt{2}}(|0\rangle + e^{2\pi i 0.\theta_{k_1} \cdots \theta_{k_t}}|1\rangle)|k\rangle$$

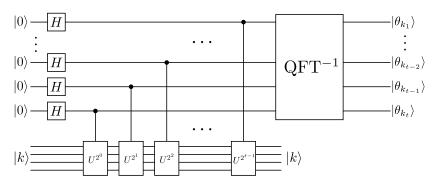
Should look familiar!



Last step is to apply the inverse QFT to recover the state...

Image credit: Xanadu Quantum Codebook node P.2

We can then measure to learn the numerical value of  $\theta_k$ .



Let's implement it.

### Next time

### Content:

■ Starting with Shor's algorithm

### Action items:

1. E-mail me your project team and paper selection by end of day

### Recommended reading:

- Codebook nodes F.1-F.3, P.1-P.4
- Nielsen & Chuang 5.1, 5.2