CPEN 400Q / EECE 571Q Lecture 11 The quantum Fourier transform and quantum phase estimation

Tuesday 15 February 2022

Announcements

- Project group / topic selection due today
- Please upgrade to PennyLane v0.21; new requirements.txt file will be included later with Quiz 5 and with Assignment 3.

Quiz 5 after class today.

Last time

We introduced the quantum Fourier transform, and saw how it is the analog of the classical inverse discrete Fourier transform.

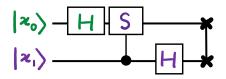
$$QFT|x\rangle = rac{1}{\sqrt{N}} \sum_{k=0}^{N-1} \omega^{xk} |k
angle$$

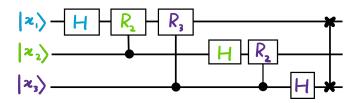
$$QFT = \frac{1}{\sqrt{N}} \begin{pmatrix} 1 & 1 & 1 & \cdots & 1\\ 1 & \omega & \omega^{2} & \cdots & \omega^{N-1}\\ 1 & \omega^{2} & \omega^{4} & \cdots & \omega^{2(N-1)}\\ \vdots & \vdots & \vdots & \ddots & \vdots\\ 1 & \omega^{N-1} & \omega^{2(N-1)} & \cdots & \omega^{(N-1)(N-1)} \end{pmatrix}$$

where for *n* qubits, $N=2^n$, and $\omega=e^{2\pi i/N}$

Last time

We saw the circuits for some special cases. For 1 qubit, it is just the Hadamard. For 2 and 3 qubits:





Quantum Fourier transform

I showed you what the general form of the circuit looked like:

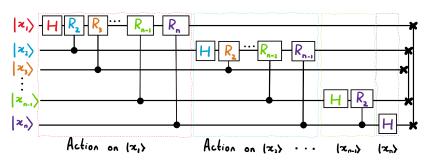


Image credit: Xanadu Quantum Codebook node F.3

Learning outcomes

- Derive the QFT circuit and implement it in PennyLane
- Describe the steps of the quantum phase estimation (QPE) subroutine
- Use the QFT to implement QPE

Review: fractional binary notation

Example

Let $k = k_1 k_2 k_3 k_4 = 0.1001$. The numerical value of this is:

$$0.1001 = \frac{1}{2^{1}} + \frac{0}{2^{2}} + \frac{0}{2^{3}} + \frac{1}{2^{4}}$$

$$= \frac{1}{2} + \frac{1}{16}$$

$$= 0.5625$$

$$= 0.5625$$

We need this for the QFT because in the exponent, we have

$$\frac{1}{\sqrt{N}}\sum_{k=0}^{N-1}\omega^{xk}|k\rangle = \frac{1}{\sqrt{N}}\sum_{k=0}^{N-1}e^{2\pi ix(k/N)}|k\rangle$$

and k/N is a fractional value.

What we are going to show is that

QFT
$$|\mathbf{x}\rangle$$
: $\frac{1}{\sqrt{N}}\sum_{k=0}^{N-1}\omega^{xk}|k\rangle$

can be rewritten in the following factorized form:

$$\frac{\left(|0\rangle+e^{2\pi i0.x_n}|1\rangle\right)\left(|0\rangle+e^{2\pi i0.x_{n-1}x_n}|1\rangle\right)\cdots\left(|0\rangle+e^{2\pi i0.x_1\cdots x_n}|1\rangle\right)}{\sqrt{N}}$$

Then, we will see how this form reveals to us the circuit that creates this state!

(keeping the last equation from the previous slide)

$$\frac{1}{N} \sum_{k_{1}=0}^{1} \sum_{k_{n}=0}^{1} \frac{2\pi i x \left(\sum_{k=1}^{\infty} \frac{k_{1}}{2^{k}}\right)}{2\pi i x \left(\sum_{k=1}^{\infty} \frac{k_{2}}{2^{k}}\right)} |k_{1} k_{2} \cdots k_{n}\rangle$$

$$= \frac{1}{N} \sum_{k_{1}=0}^{1} \sum_{k_{n}=0}^{\infty} \frac{2\pi i x \frac{k_{1}}{2^{1}}}{|k_{1}\rangle} \otimes \cdots \otimes \left(e^{2\pi i x \frac{k_{n}}{2^{n}}} |k_{n}\rangle\right)$$

$$= \frac{1}{N} \sum_{k_{1}=0}^{1} \sum_{k_{n}=0}^{\infty} \frac{2\pi i x \frac{k_{1}}{2^{1}}}{|k_{1}\rangle} \otimes \cdots \otimes \left(e^{2\pi i x \frac{k_{n}}{2^{n}}} |k_{n}\rangle\right)$$

$$= \frac{1}{N} \sum_{k_{1}=0}^{\infty} \sum_{k_{n}=0}^{\infty} \frac{2\pi i x \frac{k_{1}}{2^{1}}}{|k_{1}\rangle} |k_{2}\rangle$$

$$= \frac{1}{N} \sum_{k_{1}=0}^{\infty} \sum_{k_{n}=0}^{\infty} \frac{2\pi i x \frac{k_{1}}{2^{1}}}{|k_{2}\rangle} |k_{2}\rangle$$

(keeping the last equation from the previous slide)
$$= \sqrt{100} \left(\sum_{k=0}^{1} e^{2\pi i x_k} \frac{k!}{2!} |k| \right)$$

$$= \frac{1}{N} \left(\frac{10}{10} + \frac{2\pi i \times 1}{2} \right)$$

$$= \frac{1}{N} \left(\frac{10}{10} + \frac{2\pi i \times 1}{2} \times 1 + \frac{2\pi^{-2}}{2} \times 1 + \frac{2\pi^{-2}}{2$$

Starting with the state

$$|x\rangle = |x_1 \cdots x_n\rangle,$$

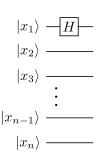
apply a Hadamard to qubit 1:

$$\frac{1}{\sqrt{2}}\left(|0\rangle + e^{-|1\rangle} |1\rangle\right)|x_2 \cdots x_n\rangle \qquad |x_{n-1}\rangle = -|x_n\rangle = -|x_n\rangle$$

$$|x_{n-1}\rangle$$
 ———

$$\frac{1}{\sqrt{2}} \left(|0\rangle + e^{2\pi i 0.x_1} |1\rangle \right) |x_2 \cdots x_n\rangle$$
If $x_1 = 0$, $e^0 = 1$ and we get the $|+\rangle$ state.

If $x_1=1$, $e^{2\pi i(1/2)}=e^{\pi i}=-1$ and we get the $|-\rangle$ state.



We are trying to make a state that looks like this:

$$|x\rangle \rightarrow \frac{\left(|0\rangle + e^{2\pi i 0.x_n}|1\rangle\right)\left(|0\rangle + e^{2\pi i 0.x_{n-1}x_n}|1\rangle\right)\cdots\left(|0\rangle + e^{2\pi i 0.x_1\cdots x_n}|1\rangle\right)}{\sqrt{N}}$$

Every qubit has a different *phase* on the $|1\rangle$ state. We are going to need some way of creating this.

We define the gate:

$$R_{k} = \begin{pmatrix} 1 & 0 \\ 0 & e^{2\pi i/2k} \end{pmatrix}$$

Now let's apply a controlled R_2 gate from qubit 2 to qubit 1

$$R_2 = \begin{pmatrix} 1 & 0 \\ 0 & e^{2\pi i}/2^2 \end{pmatrix}$$

$$|x_1\rangle$$
 H R_2 $|x_2\rangle$ $|x_3\rangle$ \vdots $|x_{n-1}\rangle$ $|x_n\rangle$ $|x_n\rangle$ $|x_n\rangle$ $|x_n\rangle$

The first qubit picks up a phase:

The first qubit picks up a phase:

$$\frac{1}{\sqrt{2}} \left(|0\rangle + e^{2\pi i \cdot 0 \cdot X_1} \right) |X_2 - X_n\rangle = \frac{1}{\sqrt{2}} \left(|0\rangle + e^{2\pi i \cdot 0 \cdot X_1} \right) |X_2 - X_n\rangle$$

$$= \frac{1}{\sqrt{2}} \left(|0\rangle + e^{2\pi i \cdot 0 \cdot X_1} \right) |X_2 - X_n\rangle$$

$$= \frac{1}{\sqrt{2}} \left(|0\rangle + e^{2\pi i \cdot 0 \cdot X_1} \right) |X_2 - X_n\rangle$$

Now let's apply a controlled R_3 gate from qubit 3 to qubit 1

$$R_3 = \begin{pmatrix} 1 & 0 \\ 0 & e^{2\pi i}/2^3 \end{pmatrix}$$

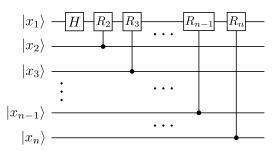
$$|x_1\rangle - H - R_2 - R_3 - |x_2\rangle - |x_3\rangle - |x_3\rangle - |x_1\rangle - |x_1$$

The first qubit picks up another phase:

$$\frac{1}{\sqrt{2}}\left(|0\rangle + e^{2\pi i 0.x_1 x_2}|1\rangle\right)|x_2 \cdots x_n\rangle \rightarrow \frac{1}{\sqrt{2}}\left(|0\rangle + e^{2\pi i 0.x_1 x_2 x_3}|1\rangle\right)|x_2 \cdots x_n\rangle$$

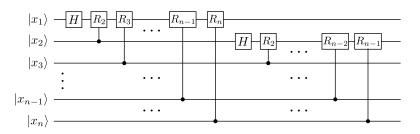
We can apply a controlled R_4 from the fourth qubit, etc. up to the n-th qubit to get

$$\frac{1}{\sqrt{2}}\left(|0\rangle + e^{2\pi i 0.x_1 x_2 \cdots x_n}|1\rangle\right)|x_2 \cdots x_n\rangle$$



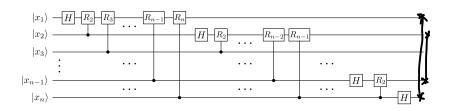
Next, ignore the first qubit and do the same thing with the second qubit: apply H, and then controlled rotations from every qubit from 3 to n to get

$$\frac{1}{\sqrt{2}^2} \left(|0\rangle + e^{2\pi i 0.x_1 x_2 \cdots x_n} |1\rangle \right) \left(|0\rangle + e^{2\pi i 0.x_2 \cdots x_n} |1\rangle \right) |x_3 \cdots x_n\rangle$$



If we do this for all qubits, we get something similar to that big ugly state from earlier:

$$|x\rangle \rightarrow \frac{\left(|0\rangle + e^{2\pi i 0.x_1 \cdots x_n}|1\rangle\right) \cdots \left(|0\rangle + e^{2\pi i 0.x_{n-1}x_n}|1\rangle\right) \left(|0\rangle + e^{2\pi i 0.x_n}|1\rangle\right)}{\sqrt{N}}$$

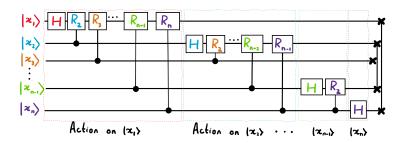


This is almost what we want: the order of the qubits is backwards. This is easily fixed with some SWAP gates.

Quantum Fourier transform

So the QFT can be implemented using:

- n Hadamard gates
- n(n-1)/2 controlled rotations
- | n/2 | SWAP gates if you care about the order



The number of gates is *polynomial in n*, so this can be implemented efficiently on a quantum computer! Let's try it...

Eigenvalues of unitary matrices

Fun fact: eigenvalues of unitary matrices are complex numbers with magnitude 1.

Proof: Let
$$\lambda_k$$
 be the eigen $|k\rangle$ associated with $|k\rangle = |\lambda_k| |k\rangle$ a unitary $|U|k\rangle = |\lambda_k| |k\rangle$. We can take the conjugate transpose of this equation $|k| |U|^{\frac{1}{2}} = |\langle k| |\lambda_k| |$

Multiply the two sides together

$$\langle k|U^{\dagger}U|k\rangle = \langle k|\lambda_{k}^{\dagger}\lambda_{k}|k\rangle$$

 $\langle k|k\rangle = |\lambda_{k}|^{2}\langle k|k\rangle$
 $1 = |\lambda_{k}|^{2}$

Eigenvalues of unitary matrices

So we can write
$$\lambda = \rho$$

where θ_k is some phase angle such that $|\theta_k| \leq 1$.

What if we want to *learn* an unknown θ_k ?

Eigenvalues of unitary matrices

Idea: apply U to the relevant eigenvector, because that's "what makes the phase come out".

$$U(k) = e^{2\pi i \theta_k}$$

...but this is an unobservable global phase!

We have to do something different: eigenvalue estimation, or quantum phase estimation (QPE).

Given a unitary U and one of its eigenvectors $|k\rangle$, estimate the value of θ_k such that

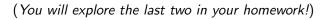
$$U|k\rangle = e^{2\pi i\theta_k}|k\rangle$$

Must determine:



 \blacksquare How to design a circuit that extracts the θ_k

- To what precision can we estimate it
- What to do if we don't know a $|k\rangle$ in advance



Let *U* be an *n*-qubit unitary; therefore $|k\rangle$ is an *n*-qubit state.

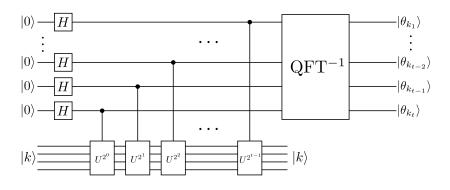
Assume for now that θ_k can be represented *exactly* using t bits in *fractional binary*:

$$\theta_k = 0.\theta_{k_1} \cdots \theta_{k_t}$$

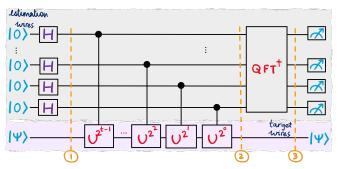
Fact: We can construct a circuit with n + t qubits that recover the value of θ_k exactly by:

- 1. Preparing n qubits in state $|k\rangle$
- 2. Applying controlled applications of U to those qubits in a special way
- 3. Applying the inverse QFT

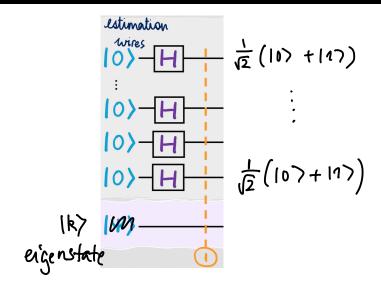
This is one version of the circuit:



The order of the controlled operations is irrelevant though, so you may see this too:



Why does this work? Let's analyze the state at points 1, 2, and 3 above.



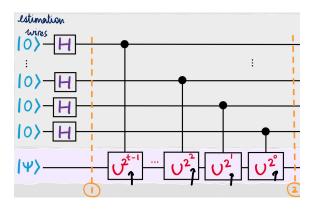
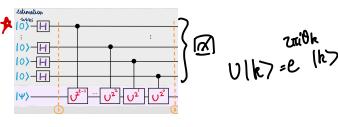
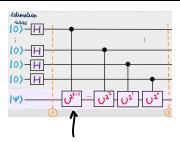


Image credit: Xanadu Quantum Codebook node P.2

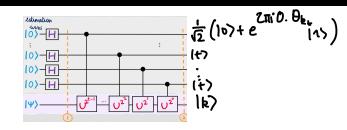


31 / 42



Use phase kickback
$$\frac{1}{\sqrt{2}} |0\rangle |+\rangle |k\rangle + \frac{1}{\sqrt{2}} \left(e^{2\pi i \theta_{k}}\right)^{2^{t-1}} |1\rangle |+\rangle |k\rangle$$

$$-\frac{1}{\sqrt{2}} (10) + \left(e^{2\pi i \theta_{k}}\right)^{2^{t-1}} |1\rangle |+\rangle |k\rangle$$



What is happening in the exponent?
$$\left(e^{2\pi i \theta k}\right)^{2^{k+1}} = e^{2\pi i \left(\frac{\theta k}{2} + \frac{\theta k}{2^{k}} + \dots + \frac{\theta k_{k}}{2^{k}}\right) \cdot 2^{k-1}}$$

$$= e^{2\pi i \left(2^{k-2} \cdot \theta k_{1} + 2^{k-2} \cdot \theta k + \dots + \frac{\theta k_{k}}{2^{k}}\right)}$$

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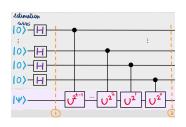
$$= e^{2\pi i \left(\frac{\theta k_{k}}{2} + \dots + \frac{\theta k_{k}}{2^{k}}\right)}$$

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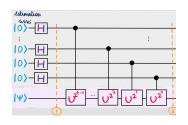
$$= e^{2\pi i \left(\frac{\theta k_{k}}{2} + \dots + \frac{\theta k_{k}}{2^{k}}\right)}$$

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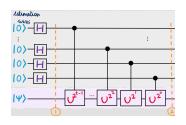
So we have the combined state:

$$\frac{1}{\sqrt{2}}(|0\rangle + (e^{2\pi i\theta_k})^{2^{t-1}}|1\rangle)|+\rangle^{\otimes t-1}|k\rangle = \frac{1}{\sqrt{2}}(|0\rangle + e^{2\pi i0.\theta_{k_t}}|1\rangle)|+\rangle^{\otimes t-1}|k\rangle$$



Let's do the second-last qubit (ignore what happens to others for now):

$$(CU)^{2}\left(|+\rangle^{\otimes t-2}\frac{1}{\sqrt{2}}(|0\rangle+|1\rangle)|+\rangle|k\rangle\right)=|+\rangle^{\otimes t-2}\frac{1}{\sqrt{2}}(|0\rangle+e^{2\pi i\theta_{k}\cdot2}|1\rangle)|+\rangle|k\rangle$$



Again check the exponent...

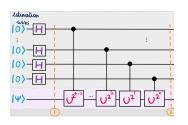
$$(e^{2\pi i\theta_k})^2 = e^{2\pi i\theta_k \cdot 2}$$

$$= e^{2\pi i(\frac{\theta_{k_1}}{2^1} + \frac{\theta_{k_2}}{2^2} + \cdots \frac{\theta_{k_t}}{2^t}) \cdot 2}$$

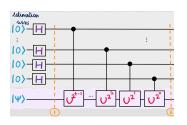
$$= e^{2\pi i(\theta_{k_1} + \frac{\theta_{k_2}}{2} + \cdots \frac{\theta_{k_t}}{2^{t-1}})}$$

$$= e^{2\pi i \theta_{k_1} + \frac{\theta_{k_2}}{2} + \cdots \frac{\theta_{k_t}}{2^{t-1}}}$$

$$= e^{2\pi i \theta_{k_1} + \frac{\theta_{k_2}}{2} + \cdots \frac{\theta_{k_t}}{2^{t-1}}}$$

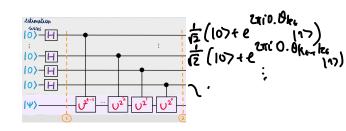


So we have the combined state:



Can show in the same way that for the last qubit

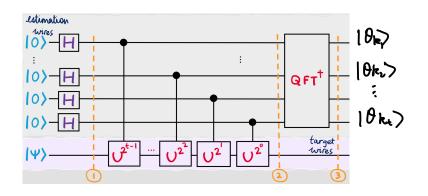
$$|+\rangle^{\otimes t-1}\frac{1}{\sqrt{2}}(|0\rangle+(e^{2\pi i\theta_k})|1\rangle)|k\rangle=|+\rangle^{\otimes t-1}\frac{1}{\sqrt{2}}(|0\rangle+e^{2\pi i0.\theta_{k_1}\cdots\theta_{k_t}}|1\rangle)|k\rangle$$



After step 2, we have the state

$$\frac{1}{\sqrt{2}}(|0\rangle+e^{2\pi i0.\theta_{k_t}}|1\rangle)\cdots\frac{1}{\sqrt{2}}(|0\rangle+e^{2\pi i0.\theta_{k_2}\cdots\theta_{k_t}}|1\rangle)\frac{1}{\sqrt{2}}(|0\rangle+e^{2\pi i0.\theta_{k_1}\cdots\theta_{k_t}}|1\rangle)|k\rangle$$

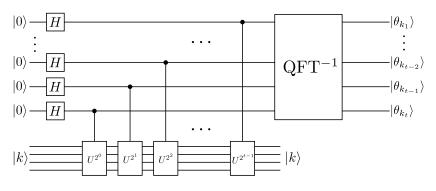
Should look familiar!



Last step is to apply the inverse QFT to recover the state...

Image credit: Xanadu Quantum Codebook node P.2

We can then measure to learn the numerical value of θ_k .



Let's implement it.

Next time

Content:

■ Starting with Shor's algorithm

Action items:

1. E-mail me your project team and paper selection by end of day

Recommended reading:

- Codebook nodes F.1-F.3, P.1-P.4
- Nielsen & Chuang 5.1, 5.2