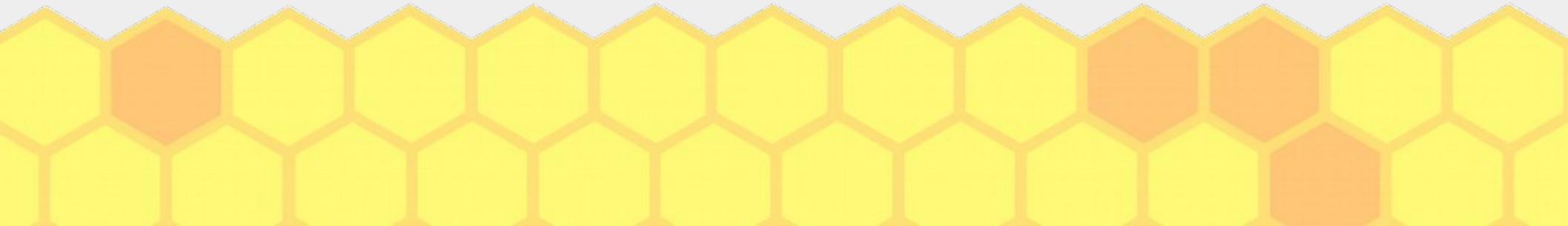


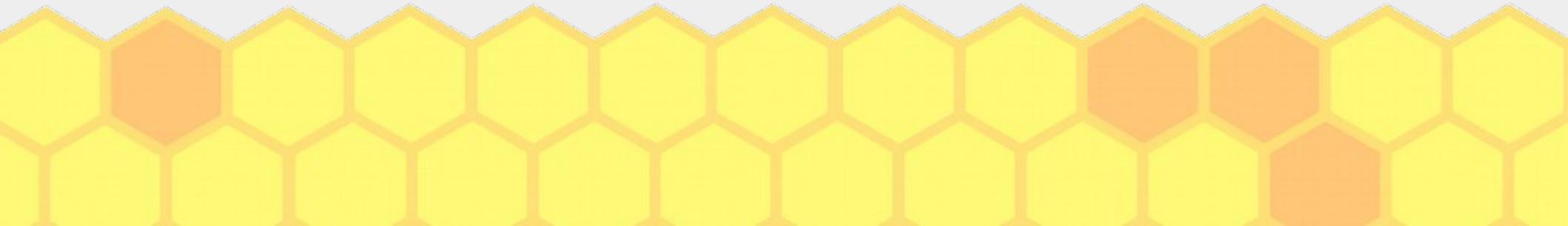
Stochastic methods in Mathematical Modelling

Lecture 13. Anomalous random processes 1. Part 1.

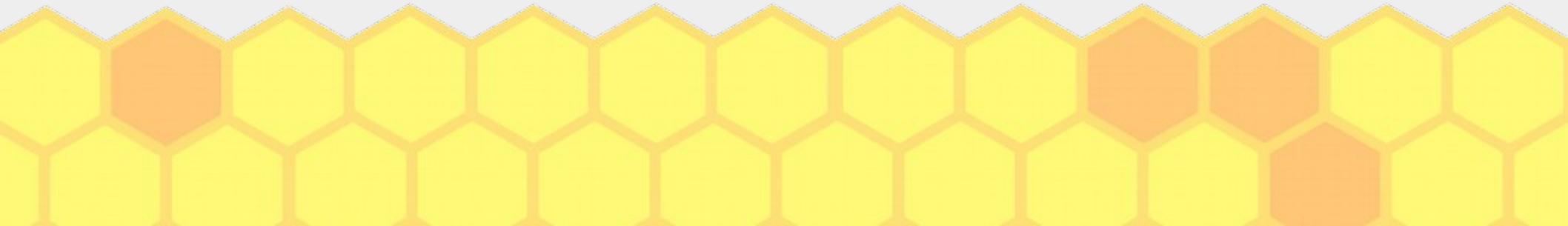
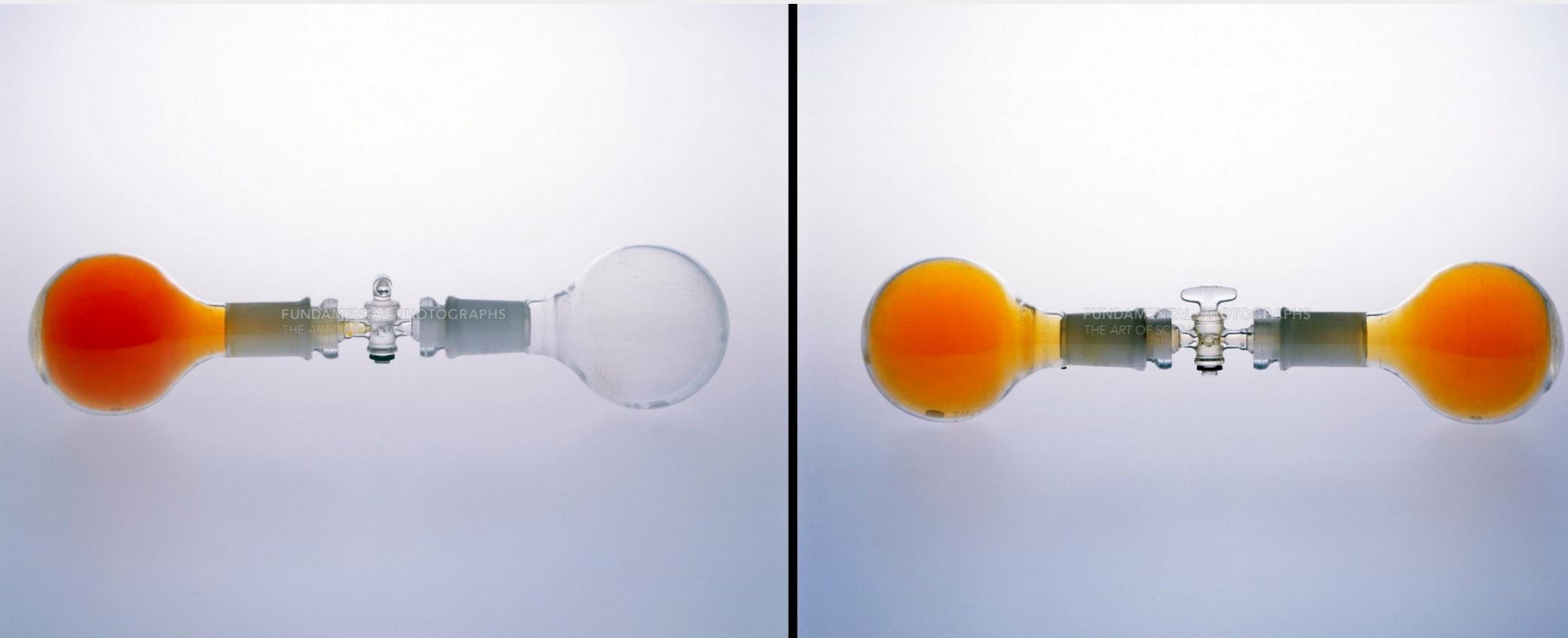


The plan of the lecture

1. Examples of Brownian and anomalous phenomena
Ergodicity, sub-, superdiffusion
2. CTRW model. Solution.
3. Fractional derivatives
4. Solution of subdiffusive fractional Fokker-Planck equation.
Mittag-Leffler and Fox H-functions



Diffusion



Diffusion

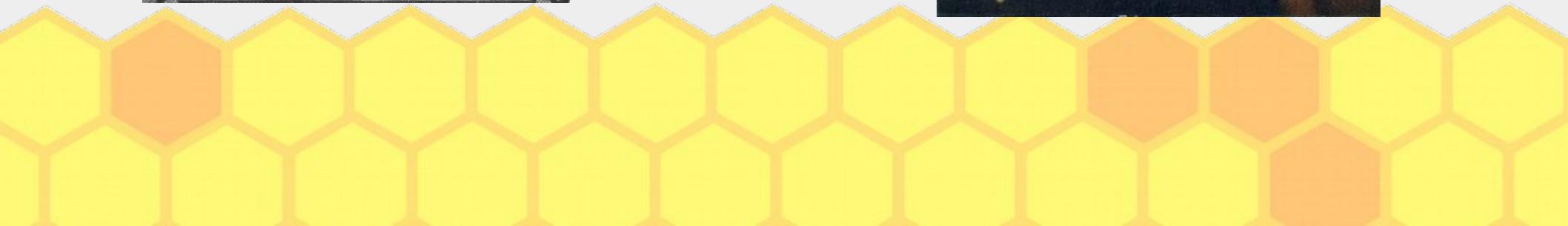
Jan Ingenhousz (1730-1799)

1785. Ingenhousz found that under a microscope the coal dust on the surface of alcohol moves erratically



Robert Brown (1773-1858)

1827. Brown sees that small lipid and starch organelles performing jittery motion known now as Brownian motion



Perrin experiment. Brownian motion

Fig. 6.

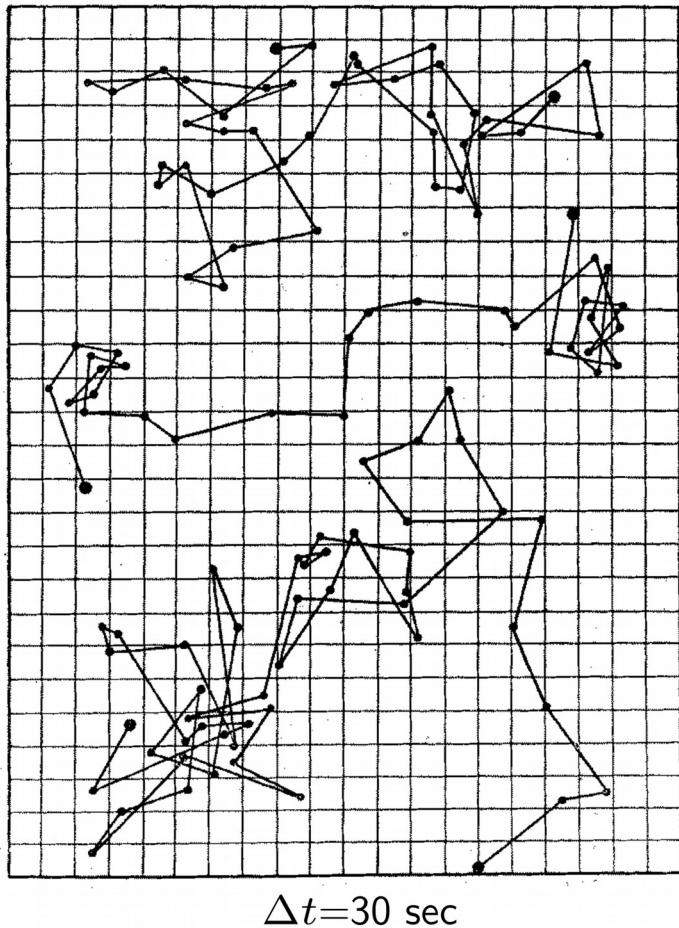
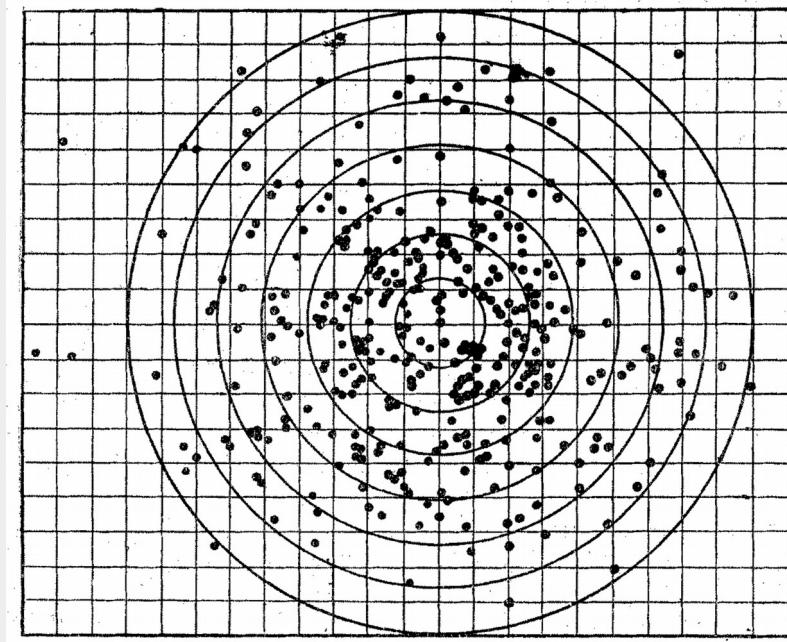


Fig. 7.



$$P(r, \Delta t) = \frac{1}{(4\pi D \Delta t)^{d/2}} \exp\left(-\frac{r^2}{4D\Delta t}\right)$$

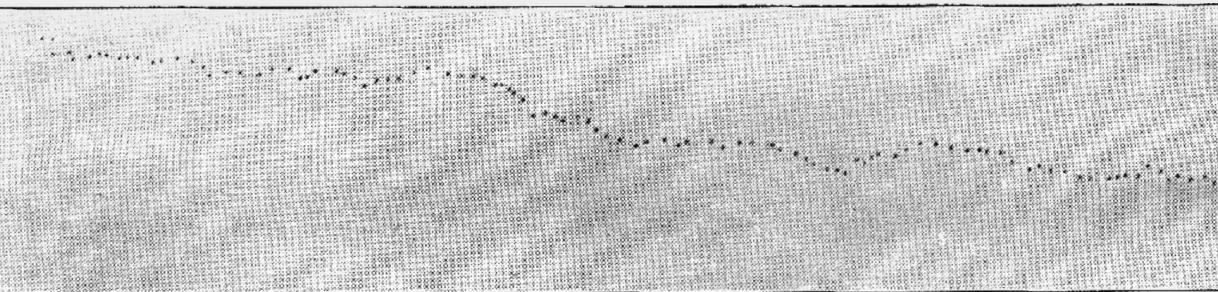
Einstein-(Smoluchowski) relation

$$D = \frac{k_B T}{m \eta} = \frac{R/N_A T}{m \eta}$$

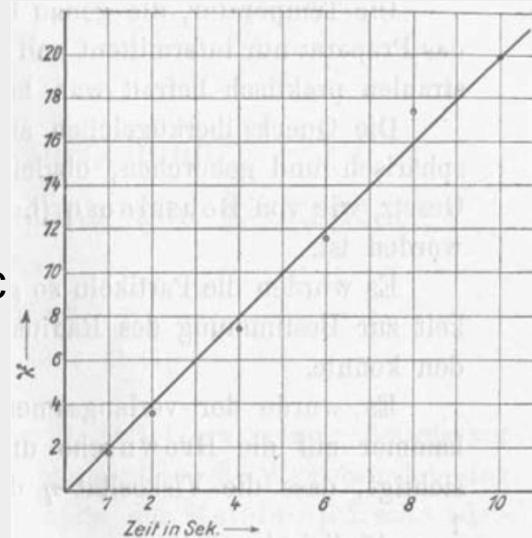
$$N_A = 7.05 \times 10^{23}$$

J Perrin, Comptes Rendus (Paris) 146 (1908) 967

Brownian motion II: single trajectory time series analysis



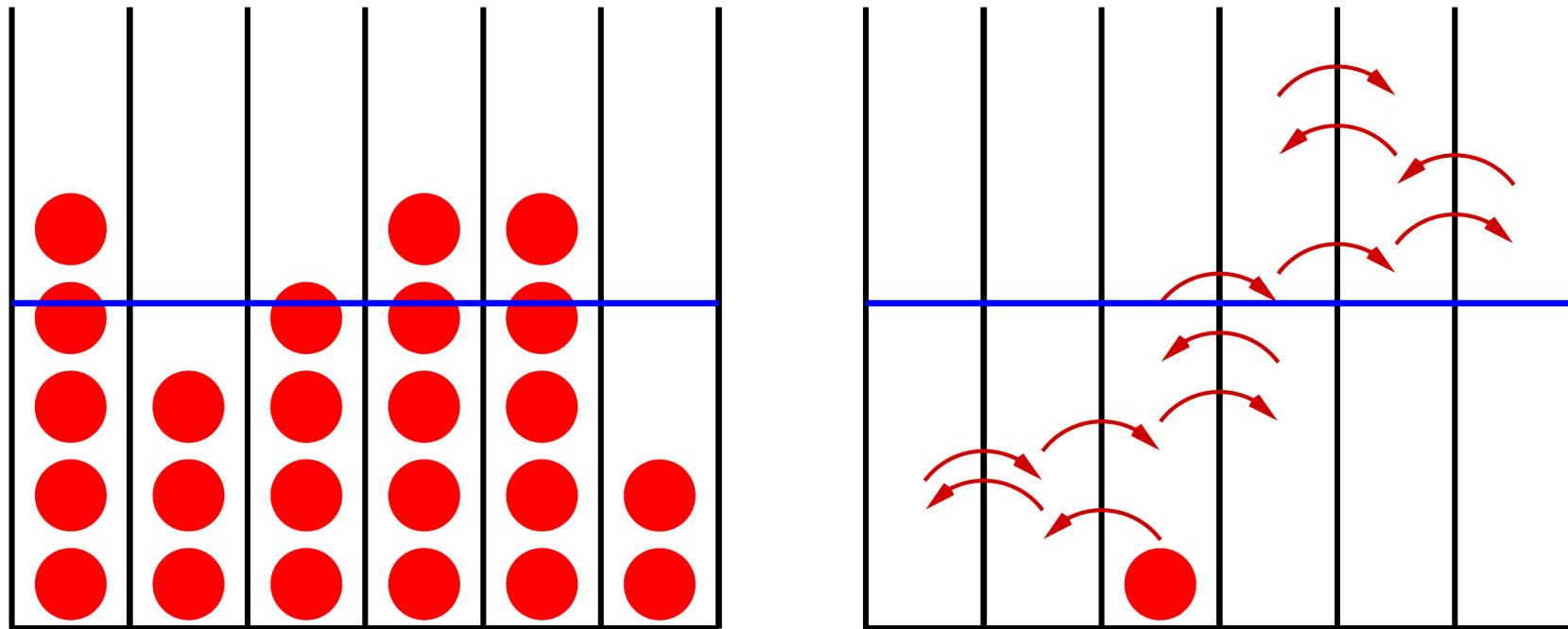
The stochastic, Brownian motion around the deterministic sedimentation with constant velocity can be clearly distinguished



example for the time averaged mean squared displacement versus time (in seconds) from a single recorded falling mercury droplet

I Nordlund, Z Physik (1914): $N_A = 5.91 \times 10^{23}$

Time versus ensemble averages: ergodic hypothesis



Ergodicity. Ensemble average = Time average:

$$\langle p_i \rangle = \lim_{N \rightarrow \infty} \frac{N_i}{N} \equiv \bar{p}_i = \lim_{t \rightarrow \infty} \frac{t_i}{t}$$

Skoltech Extracting information from single trajectories

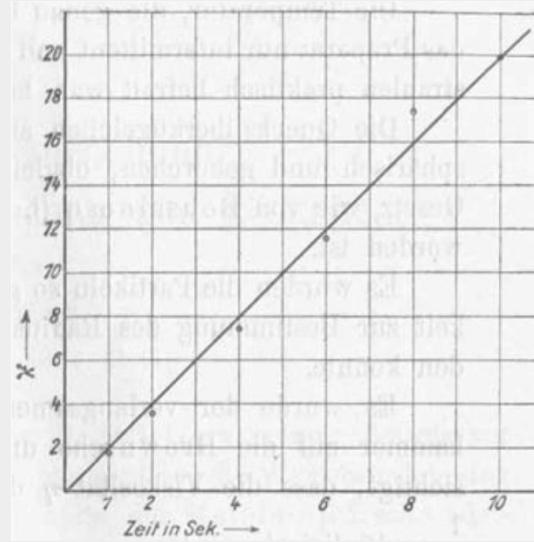
Skolkovo Institute of Science and Technology

Ensemble average for normal diffusion

$$\langle x^2(t) \rangle = \int x^2 P(x, t) dx = 2 D_1 t$$

Single particle trajectory time series $x(t)$

$$\overline{\delta^2(\Delta, T)} = \frac{\int_0^{T-\Delta} [x(t+\Delta) - x(t)]^2 dt'}{T - \Delta}$$



Normal diffusion: On average the number of jumps is proportional to the elapsed time

$$\langle [x(t+\Delta) - x(t)]^2 \rangle = \langle \delta x^2 \rangle n(t+\Delta, t) \simeq \frac{\langle \delta x^2 \rangle}{\tau} \Delta$$

Single trajectory information equals ensemble information

$$\overline{\delta^2(\Delta, T)} \sim 2 D_1 \Delta = \lim_{T \rightarrow \infty} \overline{\delta^2(\Delta, T)}, \text{ where } D_1 = \frac{\langle \delta x^2 \rangle}{2 \tau}$$

Normal vs Anomalous

Normal

$$\langle x^2(t) \rangle = 2D_1 t$$

Anomalous

$$\langle x^2(t) \rangle \neq 2Kt$$

Mostly

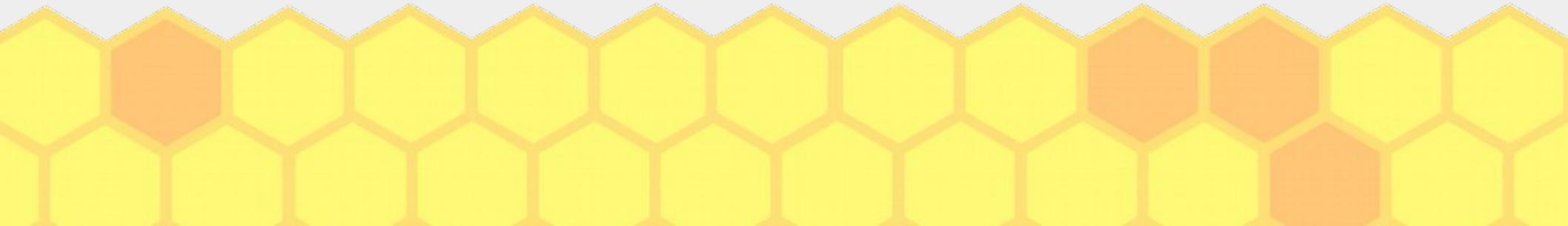
$$\langle x^2(t) \rangle \sim K_\alpha t^\alpha$$

Superdiffusion $\alpha > 1$

Subdiffusion $\alpha < 1$

Could be more complicated (Sinai diffusion, for instance)

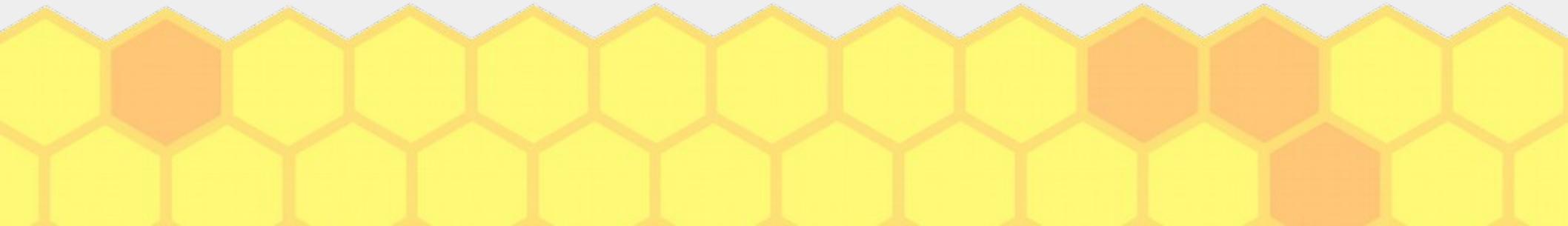
$$\langle x^2(t) \rangle \sim \log^\beta t$$



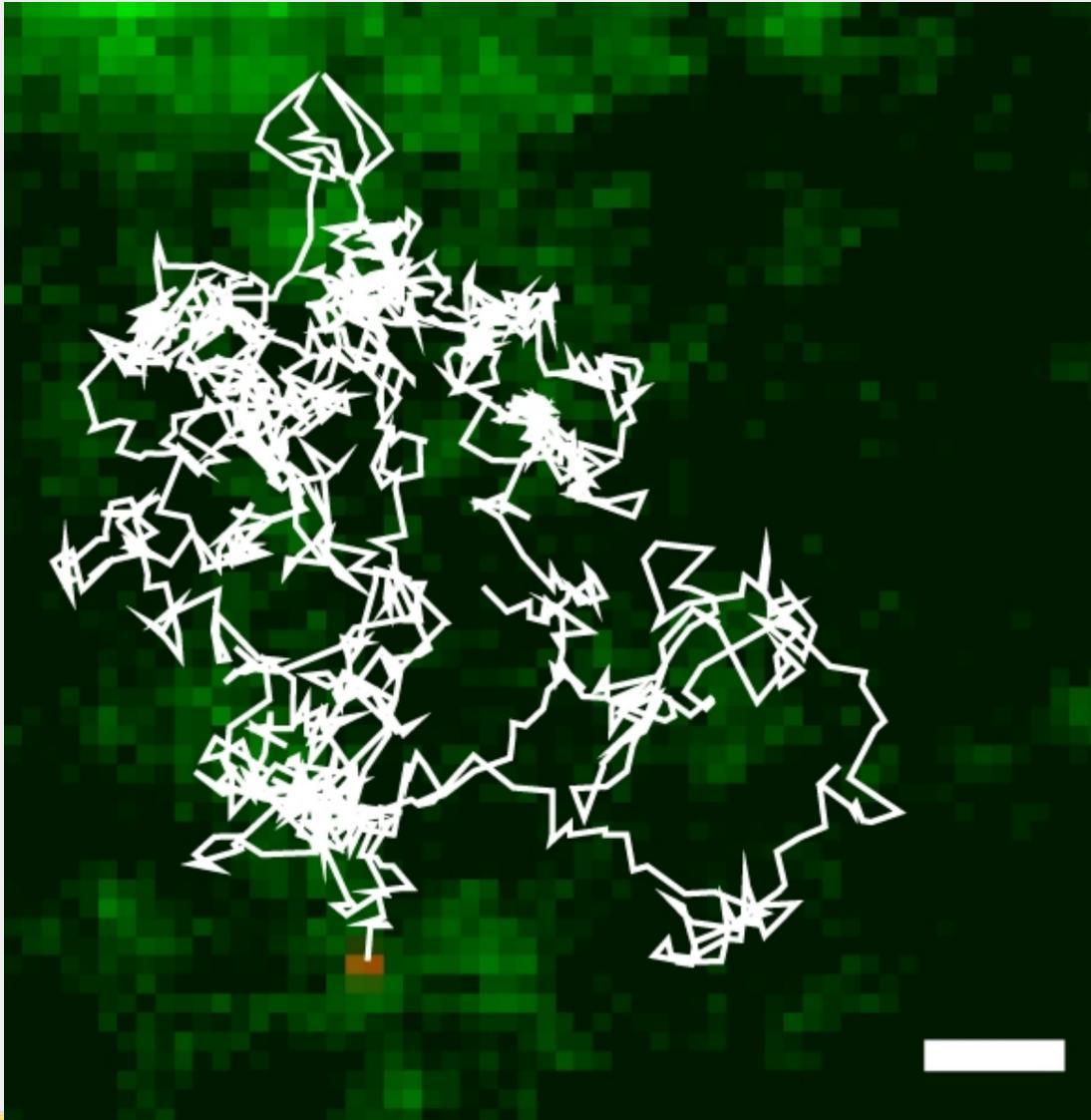
Reasons for anomaly

$$\langle x^2(t) \rangle \sim K_\alpha t^\alpha$$

1. Specific power-law distributions for jump lengths/waiting times (diverging second or first moments)
2. Correlations in jump directions
3. Specific geometry of the system (percolation cluster) or time evolution of parameters

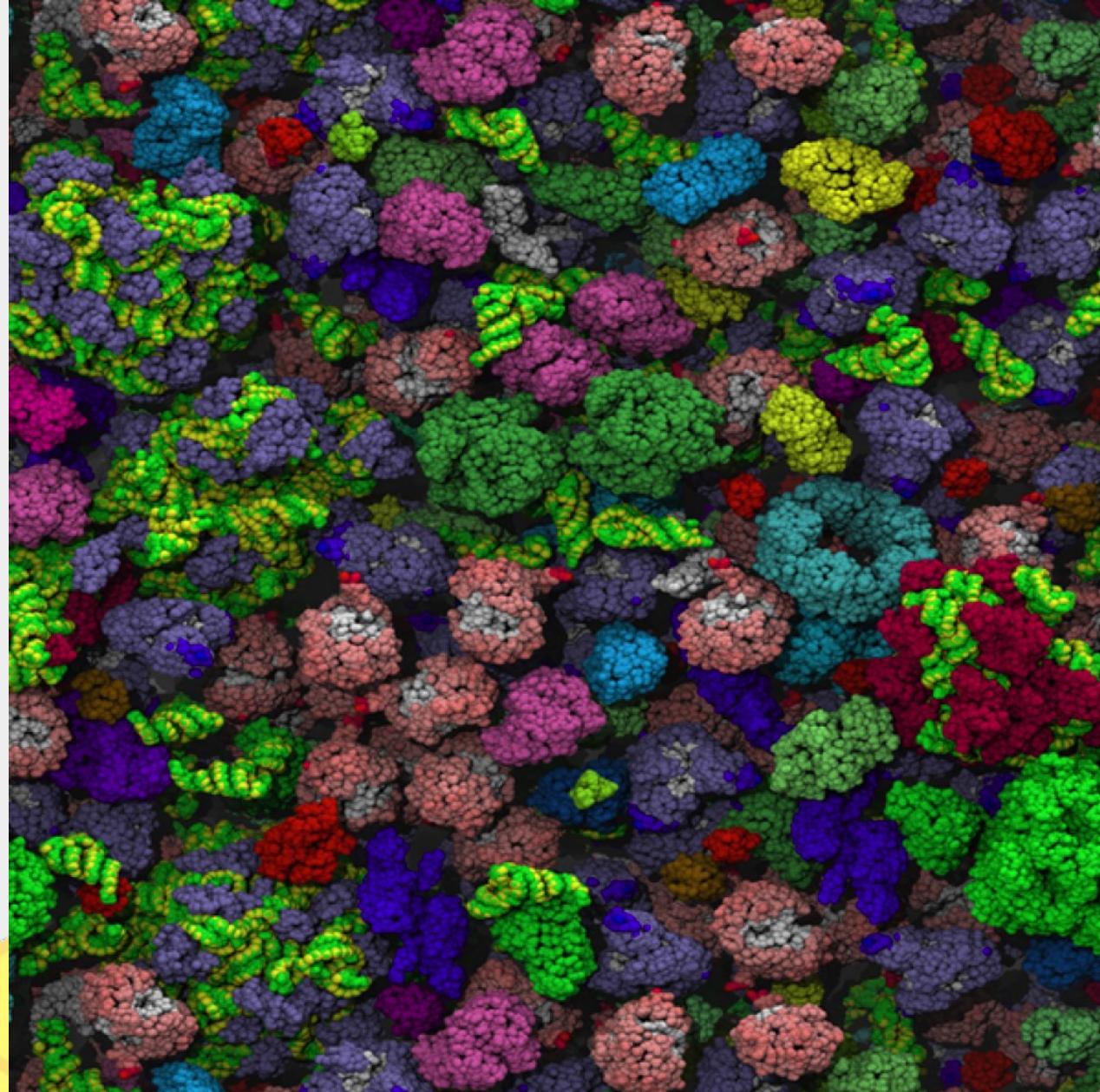


Single particle tracking I



In living human embryonic kidney cells (scale bar: 1 μ m)
D. Krapf et al.

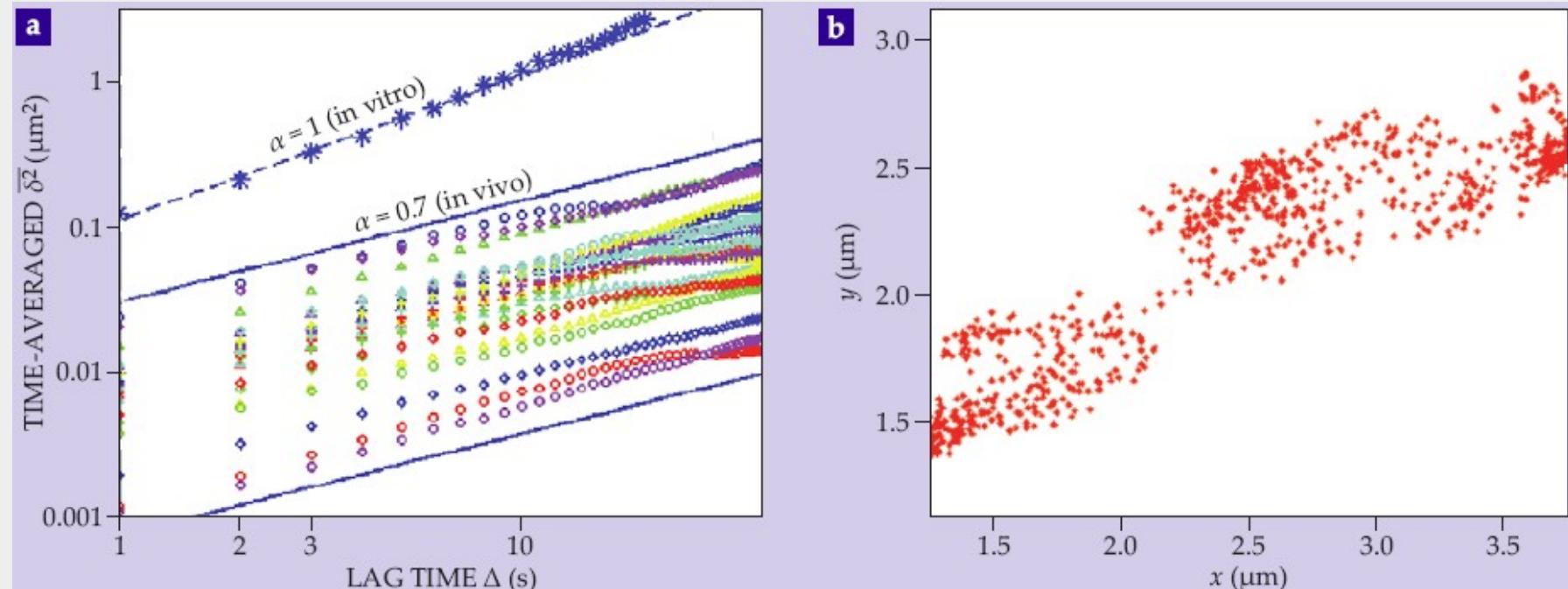
Inside of a cell



McGuee & AH Elcock, PLoS Comp Biol (2010)

Skoltech Single particle tracking II

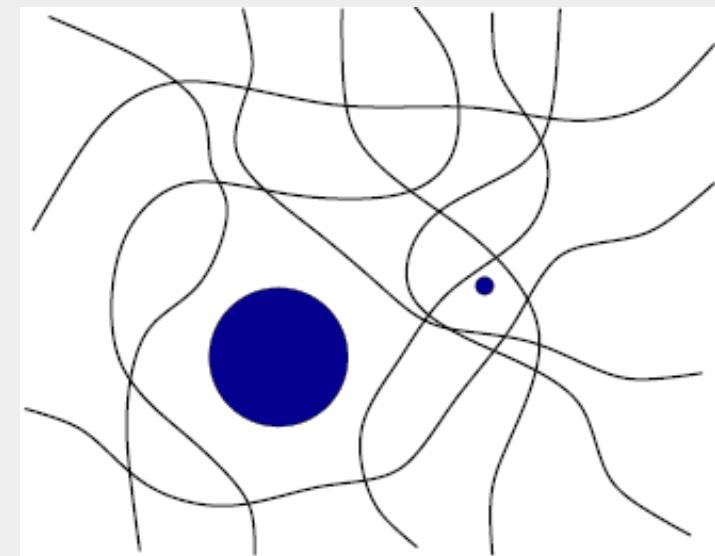
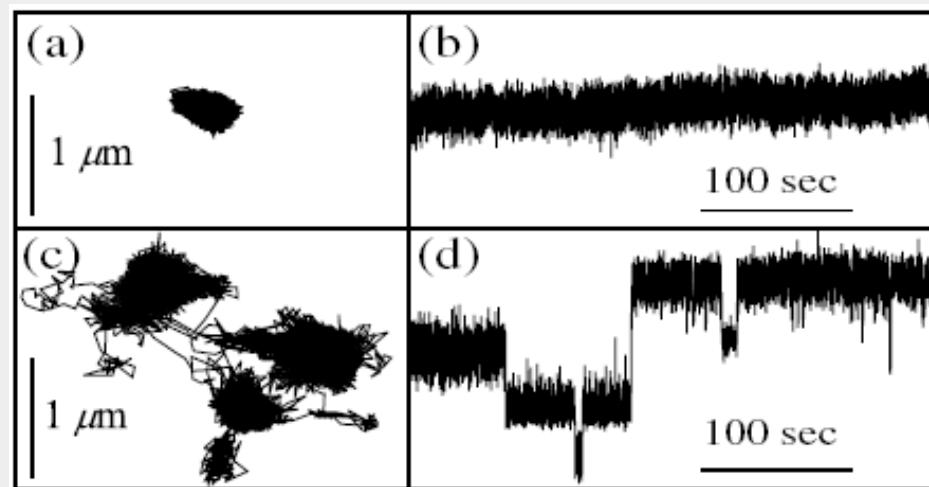
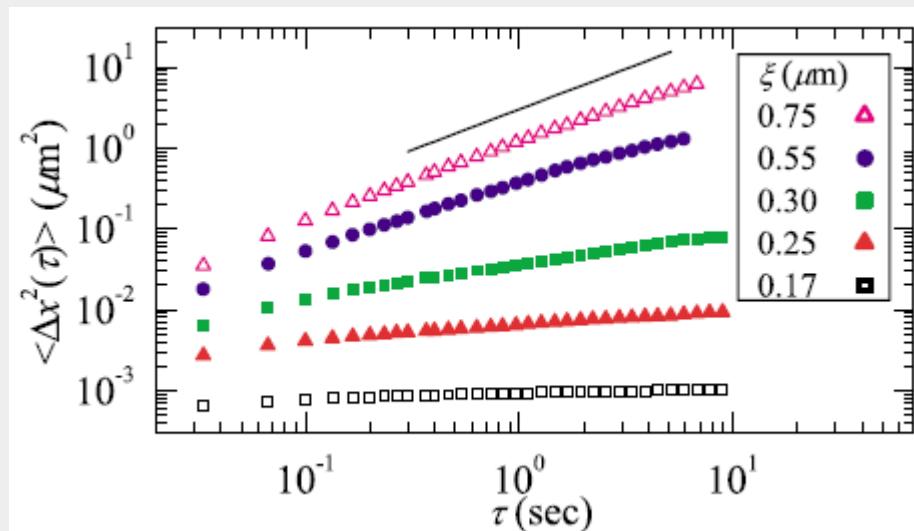
Skolkovo Institute of Science and Technology



J.-H. Jeon et al. Phys. Rev. Lett. 106, 048103 (2011)

Skoltech Actin net trapping experiment

Skolkovo Institute of Science and Technology



ξ is a typical cell size

The ball size is 0.25 micrometers

Jump probability $P \sim \tau^{1-\gamma}$

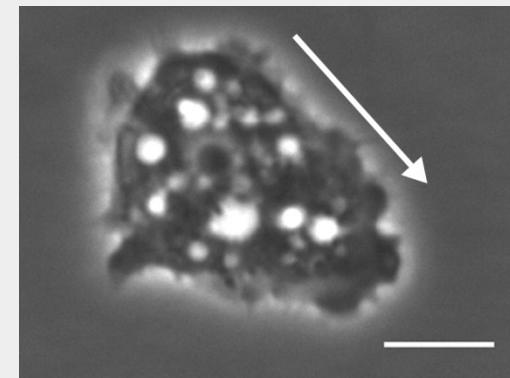
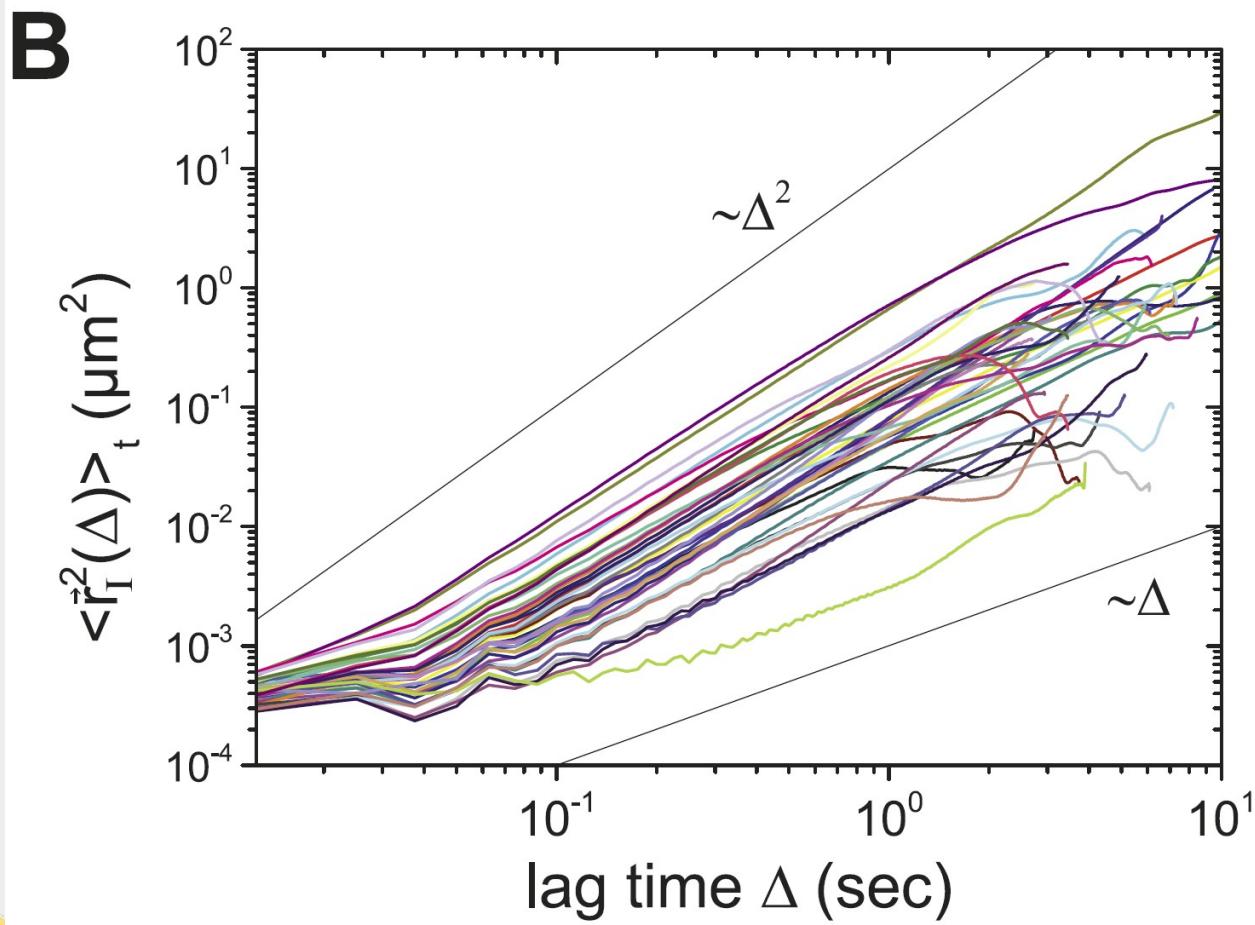
I.Y.Wong et al., Phys. Rev. Lett. 92, 178101 (2004)

Travel adventures of a dollar bill



D Brockmann, Physics World (2010) <https://www.wheresgeorge.com>

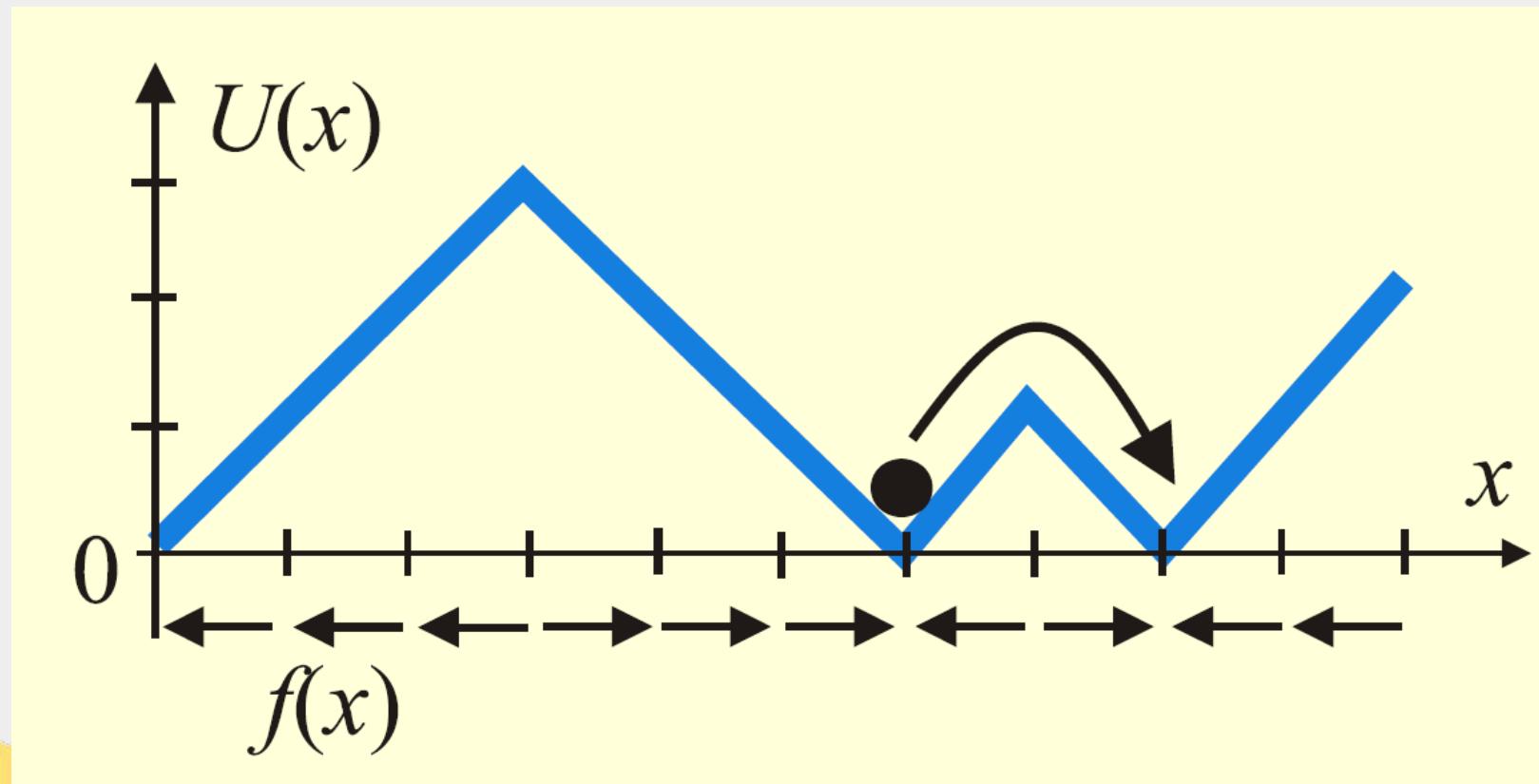
Superdiffusion in amoebae cells



Sinai diffusion

A simple random walk in a potential landscape which is a simple random walk itself

$$\langle x^2(t) \rangle \sim \log^4 t$$



Continuous-Time Random Walk model (CTRW)

$$X(t) = X_0 + \sum_{i=1}^{N(t)} \Delta X_i$$

The lengths and waiting times (durations) of the jumps are drawn from $\varphi(x,t)$

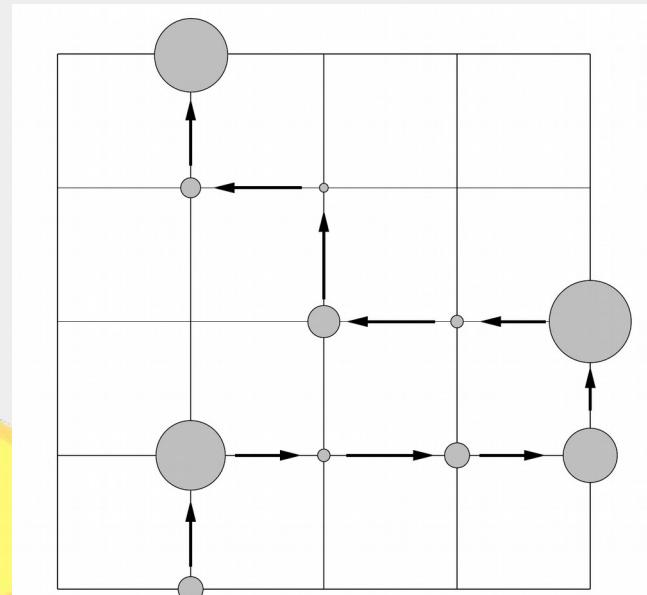
Jump pdf

$$\lambda(x) = \int dt \varphi(x, t)$$

Waiting time pdf

$$w(t) = \int dx \varphi(x, t)$$

The concept of waiting times



Continuous-Time Random Walk model (CTRW)

$$X(t) = X_0 + \sum_{i=1}^{N(t)} \Delta X_i$$

Jump pdf

general jump-WT pdf $\varphi(x, t)$

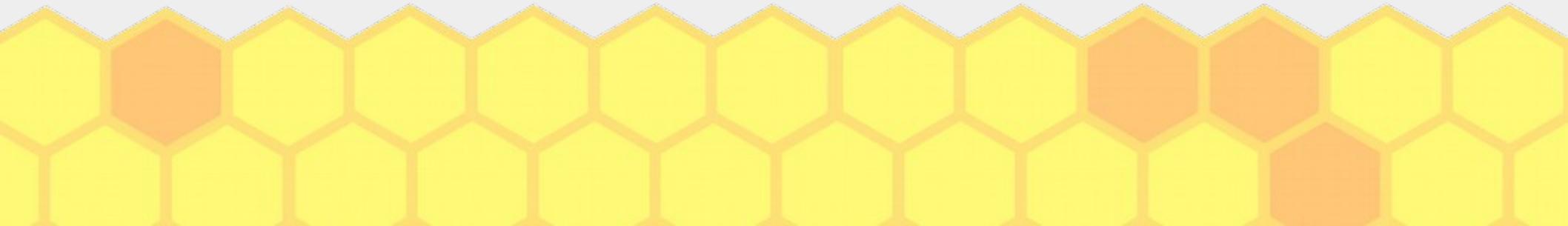
$$\lambda(x) = \int dt \varphi(x, t)$$

Waiting time pdf

$$w(t) = \int dx \varphi(x, t)$$

General solution for the CTRW PDF

$$P(k, s) = \frac{1 - w(s)}{s} \frac{P_0(k)}{1 - \varphi(k, s)}$$



$\lambda(x)$ is jump distribution

$w(t)$ are waiting times

$$T = \int_0^t w(t) dt$$
$$\sum = \int_{-\infty}^{+\infty} x^2 \lambda(x) dx$$

could diverge

CTRW

$$\lambda(x) = \int_0^t dt' \psi(x, t')$$
$$w(t) = \frac{1}{\int_{-\infty}^{+\infty} dx \psi(x, t)}$$

PDF of just arriving to (x, t) :

$$1) \gamma(x, t) = \int_{-\infty}^{+\infty} dx' \int_0^t \psi(x', t') \psi(x - x', t - t') + \underbrace{\delta(x) \delta(t)}_{P_0(x, t=0)}$$

$$2) P(x, t) = \int_0^t \gamma(x, t') \psi(t - t')$$

Prob that we never moved in $t - t'$

$$3) \Psi(t) = 1 - \int_0^t dt' w(t') \rightarrow$$

PDF that we left

Fourier + Laplace (1),(2),(3)

$$\left\{ \begin{array}{l} \Psi(s) = \frac{1}{s} - \frac{\omega(s)}{s} \\ P(k,s) = \eta(k,s) \Psi(s) \\ \eta(k,s) = \eta(k,s) \vartheta(k,s) + P_0(k) \end{array} \right. \Rightarrow P(k,s) = \frac{1-\omega(s)}{s} \frac{P_0(k)}{1-\vartheta(k,s)}$$

if $\vartheta(x,t) = \omega(t)\lambda(x)$

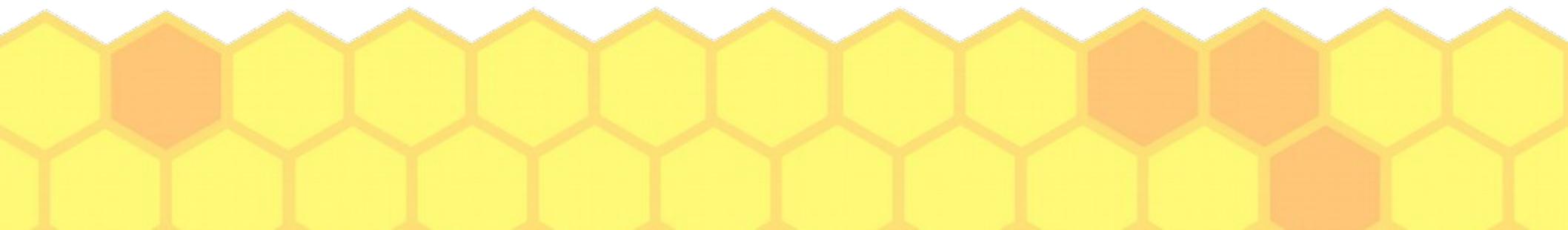
If Ψ & \sum^2 are finite \Rightarrow Brownian motion

$$\left. \begin{array}{l} \omega(t) = t^{-1} e^{-t/k} \\ \lambda(x) = \frac{1}{\sqrt{4\pi\sigma^2}} e^{-\frac{x^2}{4\sigma^2}} \end{array} \right\} \rightarrow$$

$$\omega(s) \sim 1 - s^2 + o(s^2)$$

$$\lambda(k) \sim 1 - k^2 + o(k^2)$$

$$P(k,s) = \frac{P_0}{s + Dk^2}$$



$$P(k, s) = \frac{P_0}{s + Dk^2}$$

$$s P(k, s) - P_0 = -Dk^2 P(k, s)$$
$$\left[\frac{\partial P(k, t)}{\partial t} \right] = D \left[\frac{\partial^2 P(x, s)}{\partial x^2} \right]$$

Long rests

$$\varphi(x, t) = w(t) \star \psi(x)$$

$$w(t) \sim A \left(\frac{t}{\tau}\right)^{1+\alpha}, t \gg \tau$$

Tauberian theorem

$$w(s) \sim 1 - (s\tau)^\alpha, s \rightarrow 0$$

$$\tilde{\chi}(k) \sim 1 - \sigma^2 k^2$$

$$P(k, s) = \frac{(s\tau)^\alpha}{s} \cdot \frac{P_0(k)}{1 - (1 - (s\tau)^\alpha)(1 - \sigma^2 k^2)} \underset{s \rightarrow 0}{\simeq} \frac{P_0(k)/s}{1 + K_\alpha s^{-\alpha} k^2}$$

$$\underbrace{s P(k, s) - P_0(k)}_{L\left\{ \frac{\partial P(x, t)}{\partial t} \right\}} = -K_\alpha s^{1-\alpha} k^2 P(k, s) \quad \nabla \left\{ -\frac{\partial^2 P(x, s)}{\partial x^2} \right\}$$

$$h^{-1} \{ S^P P(x, s) \} = \underline{\mathcal{D}_t^P} P(x, t), \quad P \geq 0$$

$$P(x, s) - \frac{P_0(x)}{s} = + K_2 s^{-\alpha} \frac{\partial^2 P(x, t)}{\partial x^2}$$

$S P - P_0$

$$\boxed{\frac{\partial P(x, t)}{\partial t} = \mathcal{D}_t^{1-\alpha} \frac{\partial^2 P(x, t)}{\partial x^2}}$$

$$\mathcal{D}_t^{1-\alpha} P(x, t) = \frac{\partial}{\partial t} \mathcal{D}_t^{-\alpha} P(x, t)$$

Riemann-Liouville derivative (operator)

fractional derivative

$$\mathcal{D}_t^{1-\alpha} P(x, t) = \frac{1}{\Gamma(\alpha)} \frac{\partial}{\partial t} \int_0^t \frac{P(x, t')}{(t-t')^{1-\alpha}} dt, \quad 0 < \alpha < 1$$



Long Rests: Fractional Diffusion Equation

$$P(k,s) = \frac{1-w(s)}{s} \frac{P_0(k)}{1-\varphi(k,s)}$$

Assume for $t \rightarrow \infty$ $w(t) \sim A \left(\frac{\tau}{t}\right)^{1+\alpha}$

Assume that $\lambda(x)$ has
a finite variance

$$w(s) \sim 1 - (s \tau)^\alpha$$

$$P(k,s) = \frac{P_0(k)/s}{1 + K_\alpha k^2 s^{-\alpha}}$$



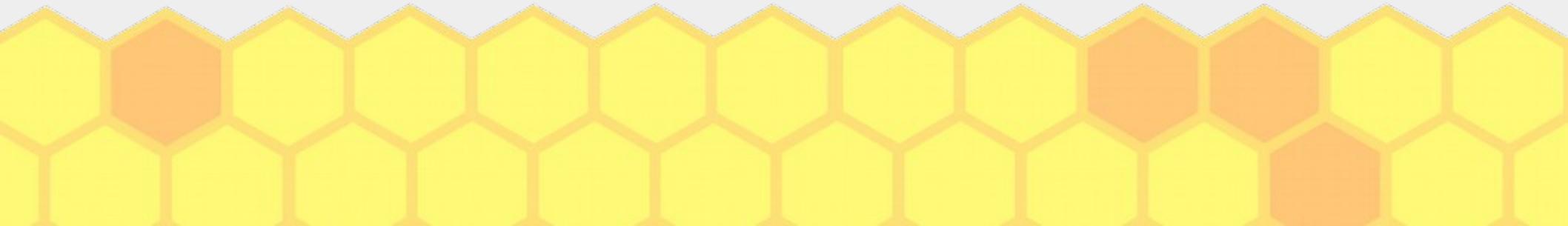
Long Rests: Fractional Diffusion Equation

$$P(k,s) = \frac{P_0(k)/s}{1 + K_\alpha k^2 s^{-\alpha}}$$

$$w(t) \sim A \left(\frac{\tau}{t} \right)^{1+\alpha}$$

$\lambda(x)$ has
a finite variance

$$\frac{\partial P(x,t)}{\partial t} = {}_0D_t^{1-\alpha} K_\alpha \frac{\partial^2}{\partial x^2} P(x,t), K_\alpha = \sigma^2 / \tau^\alpha$$



Fractional derivatives

Riemann-Liouville fractional derivative

$${}_0D_t^q P(x, t) = \frac{1}{\Gamma(1-q)} \frac{d}{dt} \int_0^t \frac{P(x, t')}{(t-t')^q} dt' \quad q < 1$$

$${}_aD_t^p P(t) = \frac{1}{\Gamma(n-p)} \frac{d^n}{dt^n} \int_a^t P(t') (t-t')^{n-p-1} dt', \quad n-1 \leq p < n$$

Examples

$${}_0D_t^\nu t^\mu = \frac{\Gamma(1+\mu)}{\Gamma(1+\mu-\nu)} t^{\mu-\nu}$$

$${}_0D_t^\nu 1 = \frac{1}{\Gamma(1-\nu)} t^{-\nu}$$

$${}_0D_t^\nu e^t = \frac{t^{-\nu}}{\Gamma(1-\nu)} {}_1F_1(1, 1-\nu, t)$$

$$-\infty D_t^\nu e^t = e^t$$

Weyl fractional operator

Fractional derivatives

$$Df(x) = \frac{d f(x)}{dx} = I^{-1} f(x)$$

$$I f(x) = \int_0^x f(x) dx = D^{-1} f(x)$$

$$D(D f(x)) = \frac{d^2 f(x)}{dx^2}$$

$$I(I f(x)) = \int_0^x \int_0^{x'} f(x') dx''$$

$$(I^n f)(x) = \frac{1}{(n-1)!} \underbrace{\Gamma(n)}_{\Gamma(n)} \int_0^x (x-x')^{n-1} f(x') dx'$$

$$I^\alpha f(x) = \frac{1}{\Gamma(\alpha)} \int_0^x (x-x')^{\alpha-1} f(x') dx \quad \alpha > 0$$

Riemann-Liouville integral

$$I^\alpha(I^\beta(x)) = I^{\alpha+\beta}(x)$$

$${}_{a+}^P D_t^p f(t) = \left(\frac{d}{dt} \right)^{m+1} \int_a^t P(t-\tau) {}^{m-p} f(\tau) d\tau$$

$m \leq p < m+1$

Fractional derivatives

Riemann-Liouville fractional derivative

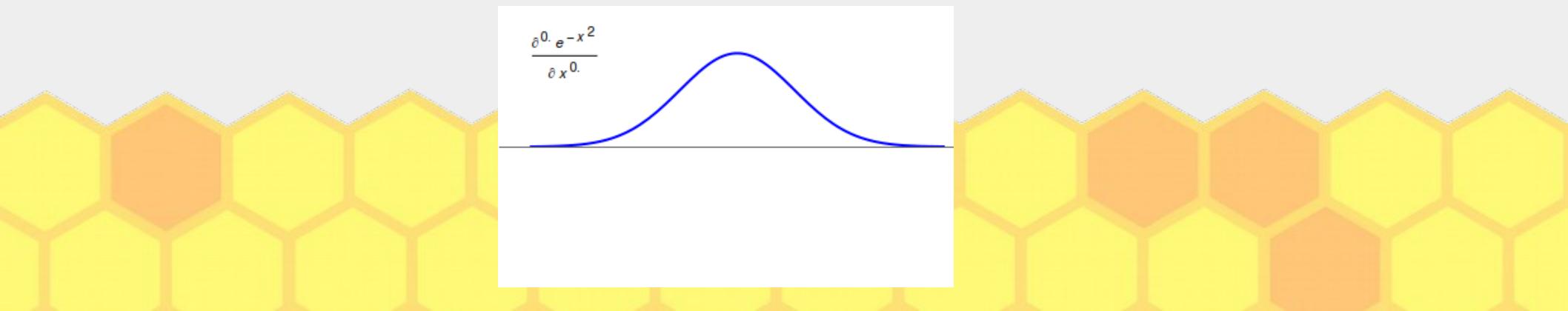
$${}_0D_t^q P(x, t) = \frac{1}{\Gamma(1-q)} \frac{d}{dt} \int_0^t \frac{P(x, t')}{(t-t')^q} dt' , \quad q < 1$$

$$\mathcal{L}\left\{{}_0D_t^{-q}f(t)\right\} = u^{-q}f(u) \quad q > 0$$

$$\mathcal{L}\left\{{}_0D_t^p f(t)\right\} = u^p f(u) - \sum_{j=0}^{n-1} u^j c_j$$

$$c_j = \lim_{t \rightarrow 0} {}_0D_t^{p-1-j} f(t)$$

Example: a fractional derivative from a Gaussian changing from 0 to 1



Fractional derivatives

Weyl fractional derivative

$$\mathcal{F}\left\{-_{-\infty} D_x^\mu f(x)\right\} = (ik)^\mu f(k)$$

q<1

Grünwald-Letnikov fractional derivative: Finite element representation

$$f'(t) = \frac{df}{dt} = \lim_{h \rightarrow 0} \frac{f(t) - f(t-h)}{h}$$

$$f'''(t) = \frac{d^3 f}{dt^3} = \lim_{h \rightarrow 0} \frac{f(t) - 3f(t-h) + 3f(t-2h) - f(t-3h)}{h^3}$$

$$f^{(n)}(t) = \frac{d^n f}{dt^n} = \lim_{h \rightarrow 0} \frac{1}{h^n} \sum_{r=0}^n (-1)^r \binom{n}{r} f(t - rh)$$

$$\binom{n}{r} = \frac{n(n-1)(n-2)\dots(n-r+1)}{r!}$$

Fractional derivatives

Grünwald-Letnikov fractional derivative: Finite element representation

$$f^{(n)}(t) = \frac{d^n f}{dt^n} = \lim_{h \rightarrow 0} \frac{1}{h^n} \sum_{r=0}^n (-1)^r \binom{n}{r} f(t - rh)$$

$$\binom{n}{r} = \frac{n(n-1)(n-2)\dots(n-r+1)}{r!}$$

$$f_h^{(p)}(t) = \frac{1}{h^p} \sum_{r=0}^n (-1)^r \binom{p}{r} f(t - rh)$$

$${}_a D_t^p f(t) = \lim_{\substack{h \rightarrow 0 \\ nh=t-a}} h^{-p} \sum_{r=0}^n (-1)^r \binom{p}{r} f(t - rh) = \lim_{\substack{h \rightarrow 0 \\ nh=t-a}} f_h^{(p)}(t)$$

If we consider a class of functions $f(t)$ having $m + 1$ continuous derivatives for t , then the Grünwald-Letnikov definition is equivalent to the Riemann-Liouville definition

Solution of subdiffusive FFPE

$$\frac{\partial P(x,t)}{\partial t} = {}_0D_t^{1-\alpha} K_\alpha \frac{\partial^2}{\partial x^2} P(x,t), K_\alpha = \sigma^2 / \tau^\alpha$$

$$\langle x^2(t) \rangle = \frac{2K_\alpha}{\Gamma(1 + \alpha)} t^\alpha$$

$$P(x,t) = \frac{1}{\sqrt{4K_\alpha t^\alpha}} H_{1,1}^{1,0} \left[\frac{|x|}{\sqrt{K_\alpha t^\alpha}} \middle| (1 - \alpha/2, \alpha/2) \right] \left. \middle| (0, 1) \right]$$



$$\int_{-\infty}^{+\infty} x^2 dx$$

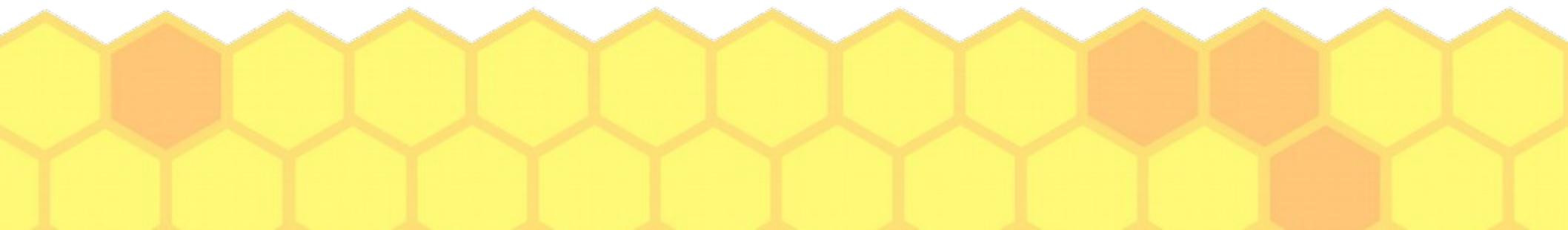
$$\frac{\partial P}{\partial t} = {}_0D_t^{1-\alpha} k_2 \frac{\partial^2 p(x,t)}{\partial x^2} \quad \text{Second moment}$$

$-\infty < x < \infty$

$$\frac{d \langle x^2(t) \rangle}{dt} = {}_0D_t^{1-\alpha} k_2 \int_{-\infty}^{+\infty} \frac{d^2 p(x,t)}{dx^2} x^2 dx = {}_0D_t^{1-\alpha} (2k_2) = \frac{2k_2}{\Gamma(\alpha+1)} t^{\alpha-1}$$

$$0 < \alpha \leq 1$$

$$\langle x^2(t) \rangle = \frac{2k_2}{\Gamma(\alpha+1)} t^\alpha$$



Fox H-functions

$$H_{pq}^{mn}(z) = H_{pq}^{mn} \left[z \begin{matrix} |(a_p, A_p)| \\ |(b_q, B_q)| \end{matrix} \right] = H_{pq}^{mn} \left[z \begin{matrix} |(a_1, A_1), (a_2, A_2), \dots, (a_p, A_p)| \\ |(b_1, B_1), (b_2, B_2), \dots, (b_q, B_q)| \end{matrix} \right] = \frac{1}{2\pi i} \int_L ds \chi(s) z^s$$

$$\chi(s) = \frac{\prod_1^m \Gamma(b_j - B_j s) \prod_1^n \Gamma(1 - a_j + A_j s)}{\prod_{m+1}^q \Gamma(1 - b_j + B_j s) \prod_{n+1}^p \Gamma(a_j - A_j s)}$$

Examples:

$$e^{-z} = H_{01}^{10} \left[z \left| \begin{array}{c} \hline \\ (0, 1) \end{array} \right. \right]$$

$$\frac{1}{p + k^{2-\alpha}} = \frac{1}{p} \frac{1}{2 - \alpha} H_{11}^{11} \left[\frac{k}{p^{2-\alpha}} \left| \begin{array}{c} (0, \frac{1}{2-\alpha}) \\ (0, \frac{1}{2-\alpha}) \end{array} \right. \right]$$

Way to compute is to use the expansion for $z \rightarrow 0$ under the set of conditions

$$H_{pq}^{mn}(z) = \sum_{h=1}^m \sum_{v=0}^{\infty} \frac{\prod_{j=1, j \neq h}^m \Gamma(b_j - B_j(b_h + v)/B_h)}{\prod_{j=m+1}^q \Gamma(1 - b_j + B_j(b_h + v)/B_h)} \frac{\prod_{j=1}^n \Gamma(1 - a_j + A_j(b_h + v)/B_h)}{\prod_{j=n+1}^p \Gamma(a_j - A_j(b_h + v)/B_h)} \\ \times \frac{(-1)^v z^{(b_h + v)/B_h}}{v! B_h}$$



Fox H-functions

$$H_{pq}^{mn}(z) = H_{pq}^{mn} \left[z \begin{matrix} |(a_p, A_p) \\ (b_q, B_q) \end{matrix} \right] = H_{pq}^{mn} \left[z \begin{matrix} |(a_1, A_1), (a_2, A_2), \dots, (a_p, A_p) \\ (b_1, B_1), (b_2, B_2), \dots, (b_q, B_q) \end{matrix} \right] = \frac{1}{2\pi i} \int_L ds \chi(s) z^s$$

Pros and contras

Pros:

1. Perform integrations one cannot do otherwise
2. Analytical inverse Laplace transform
3. Always either $z \rightarrow 0$ or $z \rightarrow \infty$ expansion exists
4. Very helpful for equations with fractional derivatives

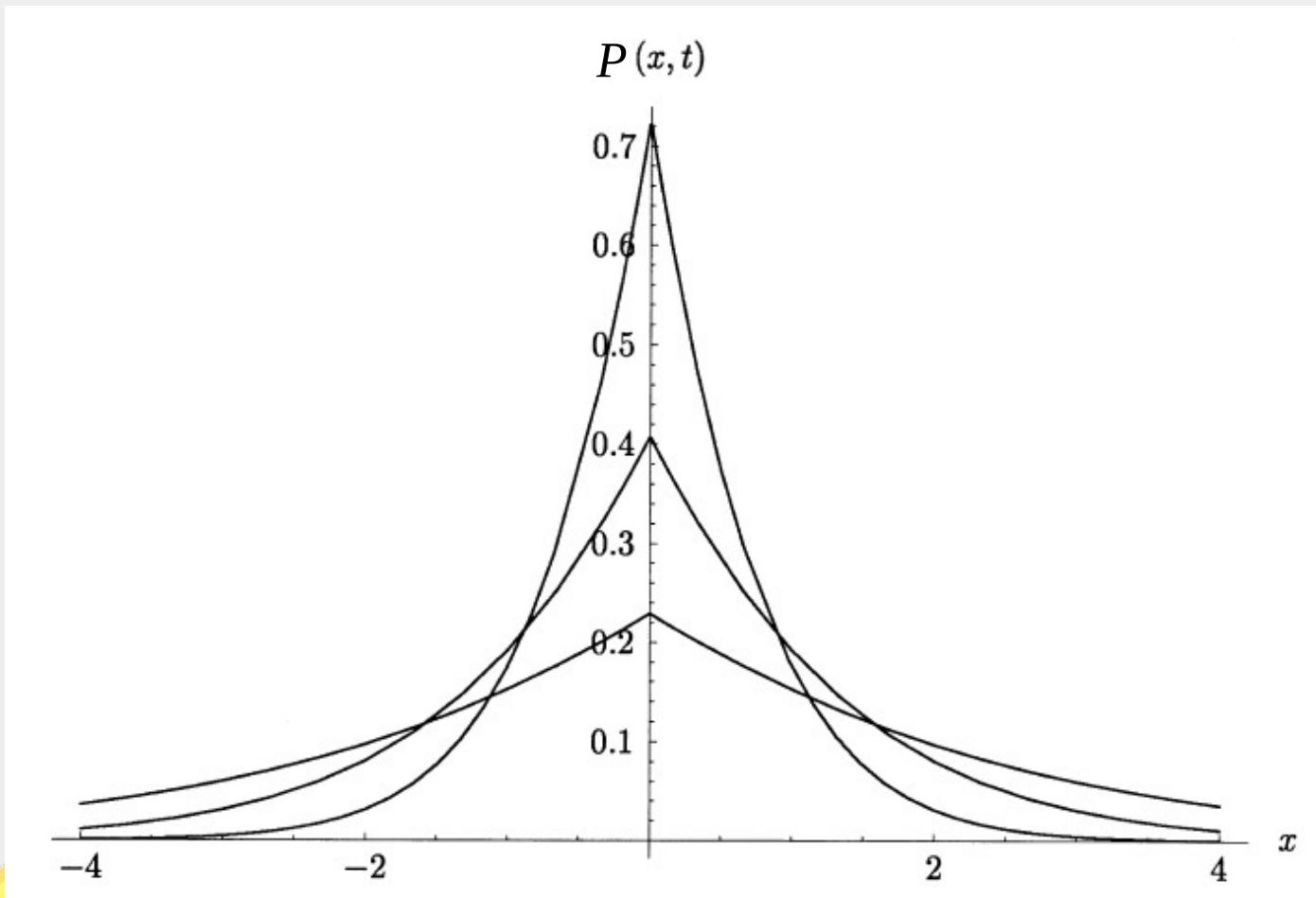
Contras

1. Not always possible to express a needed function as an H -function
2. Just by looking at coefficients we do not have a clue what the function is



Numerical results for subdiffusive FFPE

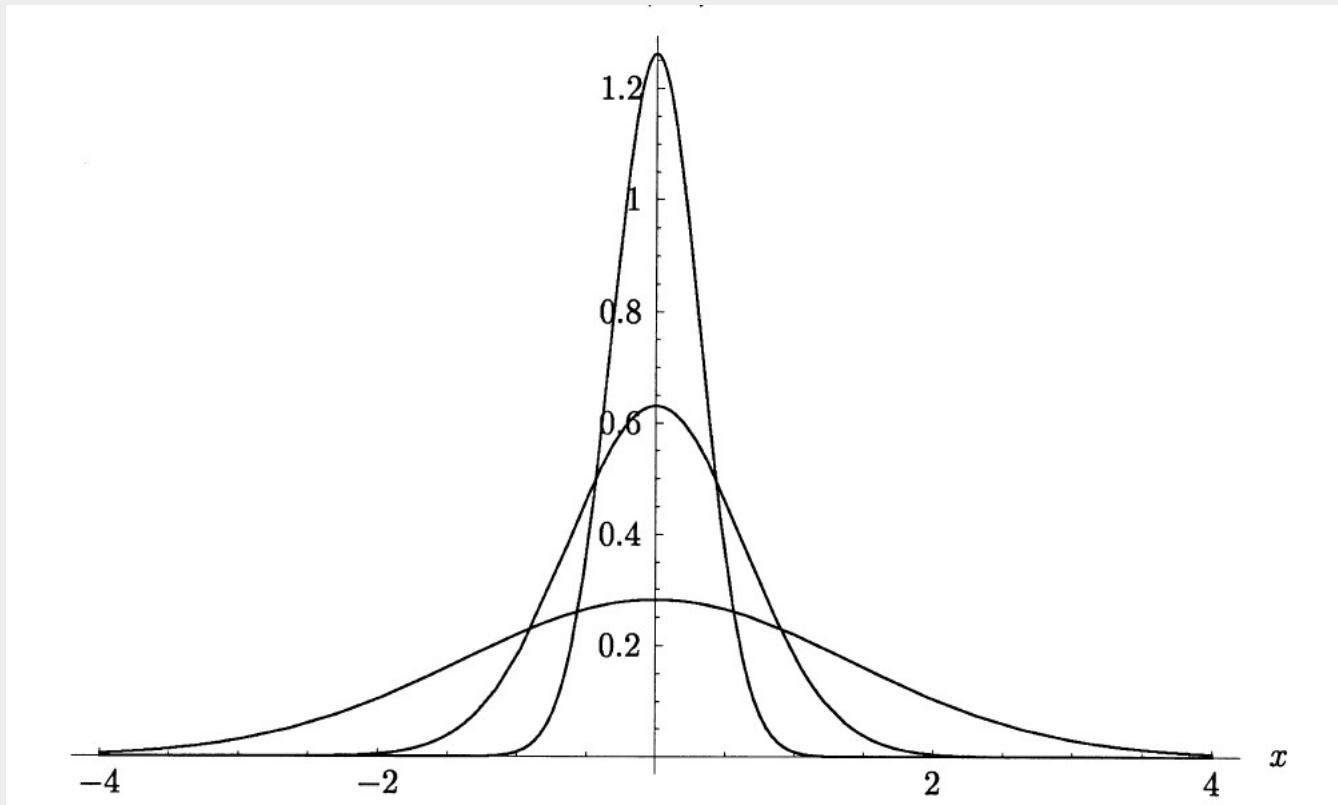
$$\frac{\partial P(x,t)}{\partial t} = {}_0D_t^{1-\alpha} K_\alpha \frac{\partial^2}{\partial x^2} P(x,t), K_\alpha = \sigma^2 / \tau^\alpha \quad \alpha = 0.5; t = 0.1, 1, 10$$



R. Metzler, J. Klafter, Phys. Rep. 2000

Reminder: Numerical results for normal FPE

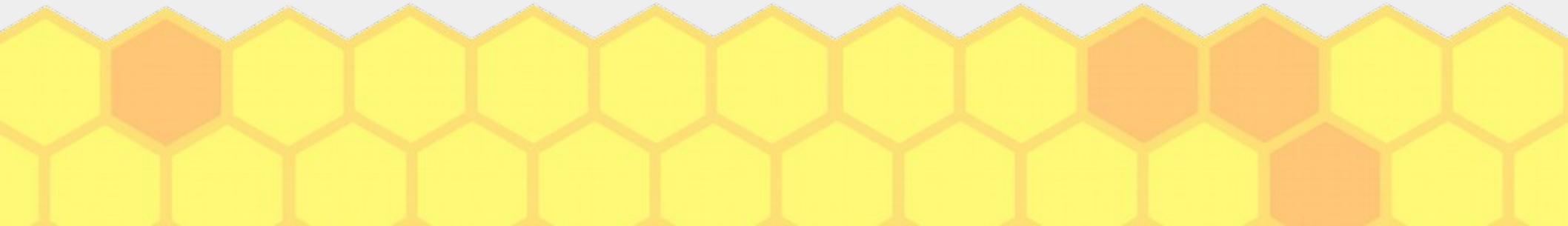
$$\frac{\partial P(x,t)}{\partial t} = K_2 \frac{\partial^2}{\partial x^2} P(x,t) \quad t = 0.1, 1, 10$$



R. Metzler, J. Klafter, Phys. Rep. 2000

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