

Geometry

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§1 Topology

This is (mostly!) a pure course, so we build up our object of study from definitions, and then eventually get to prove interesting things about those object. The following definition is the most important and foundational in this course.

Definition 1.1 (Surface). A **surface** is a topological space Σ such that every $p \in \Sigma$ has a neighbourhood homeomorphic to \mathbb{R}^2 .

In this course, we also impose the conditions that Σ is Hausdorff and second countable.

Remark 1.2. Generalising the above to homeomorphism to \mathbb{R}^n gives rise to a **manifold**, a more general object.

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¹Available here: <https://github.com/vEnhance/dotfiles/blob/master/texmf/tex/latex/evan/evan.sty>.

Remark 1.3. Forgotten what Hausdorff means again? Never forget again: a topological space X is Hausdorff iff we can ‘house off’ every pair of points: that is to say $\forall p \neq q$, there exist *disjoint* open sets $U, V \subset X$ such that $p \in U$ and $q \in V$.

(1.1) is a *local* condition on our topological space. We will want to work with more global properties of our surfaces, so we define **atlases**.

Definition 1.4. An **atlas** for a surface Σ is a set of open sets called **charts** $\{U_i\}$ (indexed by some index set \mathcal{I} , say), such that

$$\bigcup_{i \in \mathcal{I}} U_i = \Sigma. \quad (1)$$

Usually, we associate with each U_i a homeomorphism $\phi_i : U_i \rightarrow V_i \subset \mathbb{R}^2$ and call (U_i, ϕ_i) a **chart**.

What’s the point of this definition? We know that all $p \in \Sigma$ have local ngbds homeomorphic to \mathbb{R}^2 , so don’t we essentially already have a bunch of open sets that cover our surface? The elegance of atlases is that they allow us to describe surfaces we do not have a clean parametrisation for *with a single atlas*.

Example 1.5 (Single charts do not suffice)

In 1A Vector Calculus, we commonly parametrised S^2 by

$$\sigma(u, v) = \begin{pmatrix} \cos u \sin v \\ \sin u \sin v \\ \cos v \end{pmatrix} \quad (2)$$

where $u \in U := [0, 2\pi]$ and $v \in V := [0, \pi]$. But we can’t just choose $U \times V$ as our ngbd to all points on S^2 ; to flesh this point out: if we *tried* to make a homeomorphism $\phi : S^2 \rightarrow U \times V$ we would fail, since $v = 0$ (or π) leads to the u coordinate being arbitrary and injectivity breaking down ^a.

Atlases fix this problem in a clean way. For example can just pick $U = (0, 2\pi)$ and $V = (0, \pi)$, which then only misses out the north and south poles, and half a great circle that joins them. Another open set that’s a rotation of $U \times V$ will then allow us to cover all S^2 in two charts.

^ain addition, there are issues with U not being open

In the following section, we specialise to surfaces that are subspaces of \mathbb{R}^3 . Note that note this is not possible for all surfaces; the classic example is the Klein bottle, which self-intersects when we try and embed it into \mathbb{R}^3 , and hence (considering the subspace topology on \mathbb{R}^3) at these points of intersection points do not have local neighbourhoods homeomorphic to \mathbb{R}^3 .

Definition 1.6. **Smooth** means infinitely differentiable.

Definition 1.7. A **diffeomorphism** is a homeomorphism that is smooth, and has a smooth inverse.

Definition 1.8 (Transition map). The **transition maps** between charts are intuitively defined (from a restricted domain that’s a subset of V_α , \tilde{V} say)

$$\phi_\beta \circ \phi_\alpha^{-1} : \tilde{V} \rightarrow V_\beta \quad (3)$$

this maps between open sets in \mathbb{R}^2 , and has the explicit, suitably restricted domain

$$\tilde{V} = \phi_\alpha(U_\alpha \cap U_\beta) \quad (4)$$

(the notation gets dense when discussing transition maps, but all we're doing is always making sure things are well-defined).

Definition 1.9 (Smooth Surface). A smooth surface in \mathbb{R}^3 is a surface Σ given as the union of several U_i such that each transition map is a diffeomorphism.

Remark 1.10. Note that infinite differentiability is a strong condition; we're trying to set things up so we can prove nice things about shapes in space, worrying about wild real analytic pathologies as little as we possibly can.

To study smooth surfaces in \mathbb{R}^3 , actually we can specialise to certain restricted local parametrisations called *allowable* parametrisations:

Theorem 1.11

The following are equivalent:

- Σ is a smooth surface in \mathbb{R}^3 .
- Σ is locally the graph of a smooth function over one of the coordinate planes.
- Σ is locally cut out by the vanishing set of a smooth function with nonzero derivative. That is, $\forall p \in \Sigma$, there is an open $U \subset \mathbb{R}^3$ such that $\Sigma \cap U = f^{-1}(0)$ where $f : U \rightarrow \mathbb{R}$ is smooth and $Df|_p \neq 0$.
- Σ is locally the image of an **allowable** parametrisation, i.e. some $\sigma : V \rightarrow \Sigma$ where $V \subset \mathbb{R}^2$ is open and $D\sigma$ has full rank throughout V .

Proof. Some implications are easy.

The harder implications are those that we'd like to use the inverse function theorem to show, but since our maps are those of the form $\mathbb{R}^m \rightarrow \mathbb{R}^n$ where $m \neq n$, this is harder. The submersion theorem and the implicit function theorem are the corollaries of the inverse function theorem needed to show these implications.

@todo; flesh out details in implicit function theorem proof. \square

§2 Linear Algebra

Continuing to work with smooth surfaces in \mathbb{R}^3 , we can now consider how to measure familiar quantities such as length, area, angle given our σ parametrization. A central tool is the *first fundamental form* ... @todo just write these formulae out and leave as exercises.

The second fundamental form is another bilinear form that is motivated by the divergence of Σ from its tangent space locally near a point: by Taylor's, for h, j small,

$$\sigma(u+h, v+j) \approx \sigma(u, v) + h\sigma_u + j\sigma_v + \frac{1}{2}(\sigma_{uu}h^2 + 2\sigma_{uv}hj + \sigma_{vv}j^2) \quad (5)$$

to second order. Hence the divergence from the tangent plane $T_p\Sigma$ locally at p is (to second order)

$$[\sigma(u+h, v+j) - \sigma(u, v)] \cdot n = \frac{1}{2} \begin{pmatrix} h & j \end{pmatrix} \begin{pmatrix} n \cdot \sigma_{uu} & n \cdot \sigma_{uv} \\ n \cdot \sigma_{uv} & n \cdot \sigma_{vv} \end{pmatrix} \begin{pmatrix} h \\ j \end{pmatrix} \quad (6)$$

So we define the second fundamental form as the bilinear form

@todo finish this and other relation for bilinear form.

Theorem 2.1 (Alternative characterisation of SFF)

$$SFF = -(Dn)^T D\sigma \quad (7)$$

□

Since we have an allowable parametrisation everywhere on Σ , we can define a unit normal

$$\frac{\sigma_u \times \sigma_v}{\|\sigma_u \times \sigma_v\|} \quad (8)$$

everywhere⁴ on Σ . We define the **Gauss map** as exactly this map, $N : \Sigma \rightarrow S^2$. It is an intuitive result that this is independent of paramterisation.

Remark 2.2. Important: keep track of dimensions. The Gauss map is a map from (a subset of) \mathbb{R}^3 to (a subset of) \mathbb{R}^3 . In what follows however, we will specialise to two dimensional bilinear forms on the tangent space.

Now also consider n as the map $U \rightarrow S^2$ (defined locally on charts) such that

$$n = N \circ \sigma. \quad (9)$$

Now since $n \cdot n = 1$ always, in fact n_u and n_v are perpendicular to n , so $Dn|_p$ (as a map) always sends things to the tangent plane at $n(p)$ on S^2 , which of course coincides with the tangent plane at p on Σ (denoted $T_p\Sigma$).

So, having begun with the Gauss map between (subsets of) \mathbb{R}^3 , we now have a map between (subsets of) \mathbb{R}^2 . In fact, since

$$Dn|_p = DN|_{\sigma(p)} \circ D\sigma|_p \quad (10)$$

by the chain rule, defining

Definition 2.3. The shape operator \mathbb{S} is the negative derivative of the Gauss map

$$\mathbb{S} = -DN|_{\sigma(p)}. \quad (11)$$

Considered as a map from the tangent space (in the σ parametrization) at p to itself, we have, in the evaluation of the first fundamental form below $\mathbb{S}v = -DN|_{\sigma(p)} D\sigma|_p v = -Dn|_p v$ so⁵

⁴Subscripts generally denote partial differentiation.

⁵I am not entirely happy with my symbol pushing here. To elaborate, I am interpreting $I(v, w)$ as $(D\sigma v)^T D\sigma w$, i.e v and w contain coordinates (two numbers in \mathbb{R}) that $D\sigma$ turn into the familiar first fundamental form expression.

Theorem 2.4

$$I(\mathbb{S}v, w) = II(v, w) \quad (12)$$

Follows immediately from the chain rule and characterisations of I and II given:

$$I(\mathbb{S}v, w) = -v^T (D\sigma)^T (DN)^T (D\sigma)w = -v^T (Dn)^T D\sigma w = II(v, w). \quad (13)$$

We can now use these tools to build further upon the abstract surface formalisms began earlier. □

Definition 2.5 (Riemannian metrics on \mathbb{R}^2 subsets.). We can define a **Riemannian metric** on a subset $V \subset \mathbb{R}^2$ by a smooth map

$$z \mapsto \begin{pmatrix} E(z) & F(z) \\ F(z) & G(z) \end{pmatrix} \quad (14)$$

where the map is to positive definite bilinear forms (since a map to matrices would be too general).

The next definition is nothing surprising

Definition 2.6 (Riemannian metrics on abstract smooth surfaces). A **Riemannian metric** on a surface given as an atlas is a map defined for each such atlas as in (2.5) with the following compatibility condition on the transition maps $f_{\alpha\beta} := \phi_\beta \circ \phi_\alpha^{-1}$: we unsurprisingly want the first fundamental forms to agree, which amounts to needing the following to hold, and so pretending for a moment we have embedded surfaces in \mathbb{R}^3 with associated parametrizations σ_α and σ_β , we would want

$$(D\sigma_\alpha)^T D\sigma_\alpha = (D(\sigma_\beta \circ f))^T D(\sigma_\beta \circ f) \quad (15)$$

but of course we're being abstract, so actually we expand via the chain rule and insist

$$\begin{pmatrix} E & F \\ F & G \end{pmatrix}_\alpha = (Df)^T \begin{pmatrix} E & F \\ F & G \end{pmatrix}_\beta Df. \quad (16)$$

§3 Groups

§3.1 Möbius transformations and the Riemann sphere

You probably remember the following from IA Groups:

Definition 3.1. A Möbius transformation \mathcal{M} is a bijection from \mathbb{C}_∞ to itself given by

$$z \mapsto \frac{az + b}{cz + d} \quad (17)$$

where $a, b, c, d \in \mathbb{C}$ and $ad - bc \neq 0$.

These turn out to have special significance to **stereographic projection**:

Definition 3.2. **Stereographic projection** is the map from the unit sphere S^2 to the plane $\mathbb{R}^3 \cap \{z = 0\}$ given by

$$(x, y, z) \mapsto \left(\frac{x}{1-z}, \frac{y}{1-z}, 0 \right) \quad (18)$$

which is interpreted as projection from the north pole N :

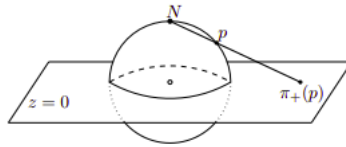


Figure 1: Stereographic projection (given by the map π_+).

Exercise 3.3. Check the formula works!

Remark 3.4. Where possible, avoid using the formula, and use nice geometric arguments (brevity is great for clarity, and exams!). This isn't always possible, however (see next claim).

The important link between Möbius transformations and stereographic projection is the following:

Definition 3.5. The subgroup $\text{PSU}(2) \leq \text{PSL}(2, \mathbb{C})$ is given by matrices of the form

$$\begin{pmatrix} a & b \\ -\bar{b} & \bar{a} \end{pmatrix} \quad (19)$$

where $a\bar{a} + b\bar{b} = 1$.

Proposition 3.6

Rotations of S^2 (namely, the action of $\text{SO}(3)$ on it) correspond 1-to-1 with Möbius transformations (of the plane) in the group $\text{PSU}(2)$ of unitary matrices in $\text{PSL}(2, \mathbb{C})$.

§4 Analysis and VP

§4.1 Geodesics

This part is mostly applying our tool of the first fundamental form in the Euler-Lagrange setting.

Proposition 4.1

Lengths of curves are given by

$$L(\gamma) := \int_{\gamma} \|\gamma'(t)\| dt = \int_{\gamma} \sqrt{E\dot{u}^2 + 2F\dot{u}\dot{v} + G\dot{v}^2} dt. \quad (20)$$

Exercise 4.2. Show this (hint⁸).

⁸write $\|\gamma'(t)\| = \sqrt{\gamma'(t) \cdot \gamma'(t)}$

We don't actually want to apply Euler-Lagrange to $L(\gamma)$ directly, since the square root is messy and it turns out we can avoid it, and still get equations that tell us when curves are length minimisers.

Definition 4.3 (Energy). The energy functional is given by

$$E(\gamma) := \int_{\gamma} \|\gamma'(t)\|^2 dt \quad (21)$$

Remark 4.4. Unsurprisingly, L is parametrisation independent (exercise ...). However, WARNING: E is not.

From here, we can apply Euler-Lagrange to study energy minimisers over curves between two points, since E is an expression in the form

$$\int f(t, u(t), u'(t), v(t), v'(t)) dt. \quad (22)$$

Proposition 4.5

Doing this, we get the fairly gross geodesic equations

$$\frac{d}{dt} (E\dot{u} + F\dot{v}) = \frac{1}{2} (E_u \dot{u}^2 + 2F_u \dot{u}\dot{v} + G_u \dot{v}^2) \quad (23)$$

and

$$\frac{d}{dt} (F\dot{u} + G\dot{v}) = \frac{1}{2} (E_v \dot{u}^2 + 2F_v \dot{u}\dot{v} + G_v \dot{v}^2). \quad (24)$$

Remark 4.6. As with most Euler-Lagrange things, these need not be committed to memory; the Euler-Lagrange idea of considering small perturbations $f \mapsto \delta f$ (see any VP notes) is quick to rederive even in timed conditions.

As promised, we now show that studying geodesics, i.e length minimisers, amounts to studying E minimisers:

Proposition 4.7 (Integral Cauchy-Schwarz)

Define an inner product on continuous functions $f : [a, b] \rightarrow \mathbb{R}_{ge0}$ as follows:

$$\langle f, g \rangle = \int_a^b f(t)g(t) dt. \quad (25)$$

Then

$$\langle f, g \rangle^2 \leq \langle f, f \rangle \langle g, g \rangle \quad (26)$$

With equality iff f and g are scalar multiples.

Proof. Any Cauchy-Schwarz proof generalises to this infinite dimensional setting with no issues. \square

Applying C-S with $||\gamma'(t)||$ and the identically 1 function,

$$L(\gamma)^2 \leq kE(\gamma) \quad (27)$$

where k is the length of the interval used to parametrise γ .

Proposition 4.8

Energy minimisers are length minimisers.

Proof. Suppose that for contradiction's sake γ minimised energy but wasn't a length minimiser. Then pick some γ' that was a length minimiser, parametrised with constant speed. Then (27) gives us that

$$E(\gamma) > \frac{1}{k}L(\gamma)^2 \geq \frac{1}{k}L(\gamma')^2 = E(\gamma') \quad (28)$$

i.e γ isn't in fact an energy minimiser, contradiction. \square

To avoid using the messy geodesic equations, we might be able to use the following alternative characterisation:

Proposition 4.9

Let Σ be a smooth surface in \mathbb{R}^3 and $\gamma : [a, b] \rightarrow \Sigma$ be a smooth curve on Σ . Then γ is a geodesic if and only if the derivative of the tangent vector to γ is everywhere normal to Σ .

Proof. Exercise! Hint^a. \square

^aThere are two geodesic equations, and normal to the tangent plane is equivalent to perpendicular to σ_u and σ_v ... you guess what you need to do!

§4.2 Geodesic normal form

I assume that the rigour won't be covered in exams (and it is pretty technical analysis anyway), so the following will be handwaved:

Proposition 4.10 (Geodesic existence)

$\forall p \in \Sigma$ embedded in \mathbb{R}^3 and all tangent vectors v at p , there's a geodesic $\gamma : [0, \varepsilon) \rightarrow \Sigma$ locally, with $\gamma(0) = p$ and $\gamma'(0) = v$.

Proof. Working with the gross geodesic equations, this is a question about existence of solutions to differential equations, specifically some of the form

$$\begin{bmatrix} E & F \\ F & G \end{bmatrix} \begin{bmatrix} u'' \\ v'' \end{bmatrix} = \dots \quad (29)$$

Where the RHS is some gross (but smooth!) expression in terms of u , v , u' and v' .

Now the first fundamental form is invertible (wth smooth inverse), so after inverting it we're indeed in the Picard-Lindelöf setting. Strengthening what we proved in Analysis and Topology, it turns out that this smooth setup means solutions are also smooth (proof omitted). Hand wave over. \square

Given this existence result, we now can make a coordinate system by travelling along the (locally existent) geodesic $\gamma(v)$ and then constructing for every $\gamma(v)$ a geodesic γ_v that is perpendicular to γ and is tangent to γ . This can be thought of as a bunch of orthogonal ‘offshoot’ branches from a common stem.

Here’s the picture to have in mind:

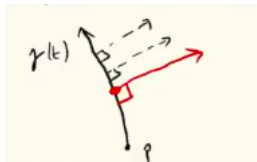


Figure 2: γ is parametrised by t , not v , but the ‘offshoot’ intuition is shown.

Remark 4.11. The curves $u \mapsto \gamma_v(u)$ for fixed v are of course geodesics, but WARNING there’s no guarantee that $v \mapsto \gamma_v(u)$ is a geodesic for fixed u , except when $u = 0$. This is the reason the geodesic normal form doesn’t just collapse to something isometric to the plane below!

Proposition 4.12

Embedded surfaces $\Sigma \subseteq \mathbb{R}^3$ have parametrisations of the form

$$du^2 + G(u, v)dv^2 \quad (30)$$

locally.

Proof. Use (via the handwaving...) the geodesic normal coordinates above:

$$\sigma(u, v) = \gamma_v(u). \quad (31)$$

There’s reasonable intuition that this works; though I remain uncertain as to why (even working locally) this map is injective.

$E = 1$ follows from the geodesics being parametrised by arc length, but it’s not obvious that F is identically zero: we have to use the geodesic equation (24) on a curve with v constant, i.e an offshooting geodesic. Then the geodesic equation gives that

$$\dot{F} = E_v = 0 \quad (32)$$

since we saw $E = 1$. But on the other hand by the MV chain rule

$$\dot{F} = F_u \dot{u} + F_v \dot{v} = F_u = 0 \quad (33)$$

so in fact F is u independent. When $u = 0$, $F = 0$ so $F \equiv 0$ everywhere. \square

@todo write Gaussian curvature part.

§5 Hyperbolic geometry

To motivate this section (and it really needs motivation!), we need some heavy machinery:

Given some surface Σ with *constant* curvature, we can rescale all the atlases by a factor of $k \neq 0$, namely $\phi \mapsto k\phi$ and WLOG assume that Σ has curvature $\kappa \in \{-1, 0, 1\}$. We've extensively studied the two non-negative cases here, and while we've seen on the example sheet the *tractoid* which has constant curvature -1, we will now see that there are much easier (2D!) 'models' to work with than that surface.

Proposition 5.1

Let $\Sigma = \mathfrak{h}$ the upper half plane model, equipped with the following abstract Riemannian metric:

$$\frac{dx^2 + dy^2}{y^2}. \quad (34)$$

Then Σ is hyperbolic (i.e has constant curvature -1 throughout Σ).

Proof. Sadly, despite the disk model being probably the most common hyperbolic model we will work with, this comes down a clever trick:

Under the change of coordinates $x = e^{-u} \tanh(v)$ and $y = e^{-u} \operatorname{sech}(v)$ (check that these indeed parametrise in the way we need) we can explicitly calculate dx and dy via the multivariable chain rule from IA^a. Then the metric is

$$du^2 + \cosh^2(u)dv^2 \quad (35)$$

and hence the Gaussian curvature of \mathfrak{h} is

$$-\frac{\sqrt{G_{uu}}}{\sqrt{G}} = -1 \quad (36)$$

@todo write the geodesic normal form section to reference the formula. □

^aI don't think that formalising this applied math hand-waving with differential forms (i.e dx etc.) is actually too far out of reach of this course, and it is a nice application of the multivariable calculus developed in Analysis and Topology. See <https://venhance.github.io/napkin/Napkin.pdf> section XII or even <http://pi.math.cornell.edu/~sjamaar/manifolds/manifold.pdf> for a complete exposition. However, I currently do not have a good intuition for why squaring these differential forms is a legitimate manipulation, and this is needed to derive the result.

Important intuition: if we're given a point to work with in a hyperbolic question, use the disk model with this point as the origin. Alternatively, if we're given a line, use the half-plane model with the line as the (positive) imaginary axis.

§5.1 Möbius transformations

At this point in the course, we return to working with inversions and Möbius transformations, this time in the hyperbolic models. It's easy to get lost in wondering why we're doing all this. To set the record straight

Definition 5.2. In \mathbb{H}^2 , a **flag** is a triple consisting of

- A (global) geodesic γ .
- A point $p \in \gamma$.
- A particular side of γ .

We care about flags because of the following result, essentially the KEY theorem of the disk models!

Proposition 5.3

Let \mathcal{F} be the set of isometries that act transitively on the set of flags in \mathbb{H}^2 . Then \mathcal{F} is generated by the Möbius transformations that preserve the particular model of \mathbb{H}^2 that we're working with, and a reflection in a geodesic.

Remark 5.4. We only need one particular reflection (e.g. $z \mapsto -\bar{z}$, reflection in the imaginary axis^a) to generate \mathcal{F} due to being able to obtain all such geodesics with MTs.

^athis works in either model.

Remark 5.5. If we insist such isometries preserve orientation, we in fact only get Möbius transformations, without the reflections.

Remark 5.6. In fact we have a uniqueness statement too: we have a unique hyperbolic isometry that maps a flag to another flag.

Proof. One direction here is not too difficult: to show Möbius transformations are isometries, it is not too difficult to check that their generators are isometries.

The proof that all orientation-preserving isometries are Möbius transformations is more difficult, and the claims made should be verified! Work in the disk model. Given some orientation-preserving isometry T , we can find a Möbius transformation M such that $M \circ T$ fixes 0 and 1. It's clear that isometries must also preserve geodesics, so in fact $M \circ T$ preserves $\mathbb{R} \cap D$. In addition, isometries in this model must preserve angles, because if locally an isometry sheared space, then curve lengths would not be preserved¹⁴. So, once again applying geodesic preservation, we have that $i\mathbb{R} \cap D$ is preserved by $M \circ T$ (it can't be reflected since we're assuming we're in the orientation-preserving case). From this point, the isometry must be the identity since points are determined uniquely by their distance (in hyperbolic, NOT Euclidean metric) to the real and imaginary axis, along with whether they are closest to the positive or negative parts of both of these axes. I made a picture!

¹⁴I would appreciate a clearer reason for this statement!

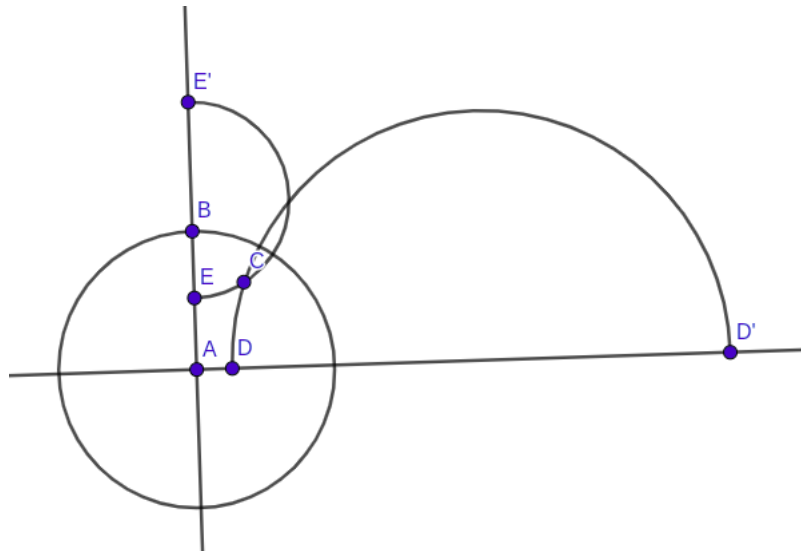


Figure 3: Given some point C in the disk, the distance of C from the real and imaginary axes are the lengths of the geodesic segments CE and CD shown. This gives rise to the fact that an isometry that fixes \mathbb{R} and $i\mathbb{R}$ is in fact the identity.

@todo: write orientation preserving section. □

So what actually are the Möbius transformations that preserve our various models?

Proposition 5.7

In the disk model, the Möbius transformations that preserve the disk (for the last time, equivalently the orientation preserving isometries!) are those of the form

$$z \mapsto \lambda \frac{z - a}{\bar{a}z - 1} \quad (37)$$

where $|\lambda| = 1$ and $|a| < 1$.

Proof. Complex analysis example sheet 1. □

Proposition 5.8

In the half-plane model, the relevant Möbius transformations are those of the form

$$z \mapsto \frac{az + b}{cz + d} \quad (38)$$

where $\begin{bmatrix} a & b \\ c & d \end{bmatrix} \in \mathrm{SL}(2, \mathbb{R})$.

Proof. Sketch: this isn't actually too hard: we need to fix the real line, which quickly yields that all the coefficients need be real, and then we can just isolate the cases $c = 0$ and $c \neq 0$. In the former the result falls out quickly, and in the latter case, WLOG $c = 1$ so write out

$$\frac{az + b}{z + d} = \frac{-1}{z + d} + a \quad (39)$$

which is the composition of a bunch of two translations and also $z \mapsto -\frac{1}{z}$. This final map is an isometry since we have metric

$$\frac{|dz|^2}{4|\mathrm{Im}z|^2} \quad (40)$$

and changing variables to $w = -\frac{1}{z}$ preserves this^a. \square

^aNote that again the differential form manipulation should make you a little uneasy...

Example 5.9 (The hyperbolic plane is *weird*)

This can definitely be skipped for exams.

Proposition 5.10 (Pairs of pants (!))

Important: the right angles lead to smooth gluing of geodesics.

References