

Lecture 6: Density Functional Theory

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3:22 PM

- Introduce a field $J(r)$
- Derive a variational principle to determine $\langle \hat{g}(r) \rangle$

$$Z_c[J] = z_0 \int dr^N e^{-\beta U(r) - \int dr J(r) \hat{g}(r)}$$
$$= e^{-\beta A[J]} \leftarrow$$

If $J(r) = 0$, this is the equil. ensemble

$$\begin{aligned} \Rightarrow - \frac{\delta \log Z_c[J]}{\delta J(r)} &= \frac{1}{Z_c} z_0 \int dr^N \hat{g}(r) e^{-\beta U(r) - \int dr J(r) \hat{g}(r)} \\ &= \langle \hat{g}(r) \rangle_J = \beta \frac{\delta A[J]}{\delta J(r)} \end{aligned}$$

$\beta = 1$ from here on

Assume the relation b/t $J(r)$, $\langle \hat{g}(r) \rangle$ is invertible

If we know $\langle \hat{g}(r) \rangle$, we can calculate $J(r)$

Higher derivatives give higher moments of the density:

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$$\frac{\delta^2 A}{\delta J(r) \delta J(r')} = \frac{\delta \langle \hat{\rho}(r) \rangle_J}{\delta J(r')} = - \langle \hat{\rho}(r) \hat{\rho}(r') \rangle$$

↓ b/c of invertible

$$\frac{\delta J(r)}{\delta \langle \hat{\rho}(r') \rangle_J} = - \langle \hat{\rho}(r) \hat{\rho}(r') \rangle^{-1}$$

Use Legendre transform to derive

$$F[\langle \hat{\rho} \rangle] \equiv A[J] - \int dr J(r) \underline{\frac{\delta A[J]}{\delta J(r)}}$$

$$= \underline{A[J]} - \int dr J(r) \langle \hat{\rho}(r) \rangle_J$$

$$\frac{\delta F}{\delta \langle \hat{\rho}(r) \rangle} = -J(r)$$

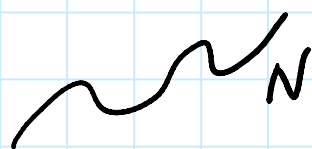
$$\frac{\delta^2 F}{\delta \langle \hat{\rho}(r) \rangle \delta \langle \hat{\rho}(r') \rangle} = - \frac{\delta J(r)}{\delta \langle \hat{\rho}(r') \rangle} = \langle \hat{\rho}(r) \hat{\rho}(r') \rangle^{-1}$$

↑

- If we have $F[J]$ as a power series in $\rho(r)$, the quadratic term is related to $S(k)$
- Interactions are reflected in the form of this coefficient

- Interactions are reflected in the form of this coefficient
- Stability of the homogeneous phase can be assessed
- Now we use FT to calculate this quadratic coefficient

Model A / Edwards model:



$$\beta U_0 = \frac{1}{2} \sum_i \sum_j \frac{1}{2b^2} |\mathbf{r}_{i,s} - \mathbf{r}_{j,s}|^2$$

$$\beta U_1 = \frac{u_0}{2} \int d\mathbf{r} [\hat{\phi}(\mathbf{r})]^2$$

Use H-S transform; w/ the field $J(\mathbf{r})$ in place:

$$Z[J] = z_0 \int d\mathbf{r}^N \underline{e^{-\beta U_0 - \beta U_1 - \int d\mathbf{r} J(\mathbf{r}) \hat{\phi}(\mathbf{r})}}$$

$$= z_1 \int D\omega \ e^{-\mathcal{H}[\omega, J]} = e^{-\mathcal{A}[J]}$$

$$\mathcal{H}[\omega, J] = \frac{1}{2u_0} \int d\mathbf{r} [\omega(\mathbf{r})]^2 - n \log Q[\mu]$$

→

$$\mu(\mathbf{r}) = \frac{1}{n} \omega(\mathbf{r}) + J(\mathbf{r})$$

$$\text{Recall } Q[\mu] = \frac{1}{V} \int d\mathbf{r}^N e^{-\mu(\mathbf{r}_N)} \phi_{N,N-1} e^{-\mu(\mathbf{r}_{N-1})} \dots \phi_{2,1} e^{-\mu(\mathbf{r}_1)}$$

Use the MF equation $\frac{\delta \mathcal{H}}{\delta \omega} = 0$

$$\frac{\delta \mathcal{H}}{\delta \omega} = \frac{1}{\omega_0} \omega(r) - i \tilde{f}(r; [\omega, J])$$

↑

Next time, use W.I.E. on \tilde{f} to relate $\omega(r)$ to $J(r)$