

# From annular cavity to rotor-stator flow: nonlinear dynamics of axisymmetric rolls

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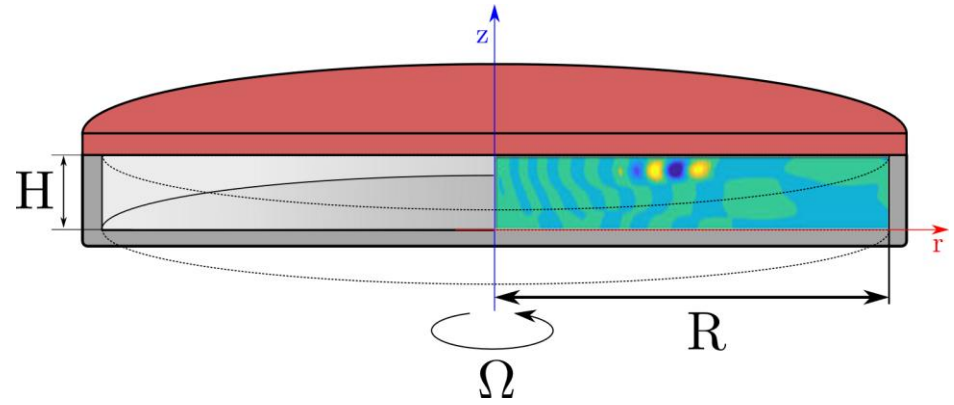
<sup>2</sup> Université Paris-Saclay, LISN-CNRS

<sup>3</sup> Université Claude Bernard Lyon 1, LMFA

\* now at EPFL, group HEAD with Eunok Yim

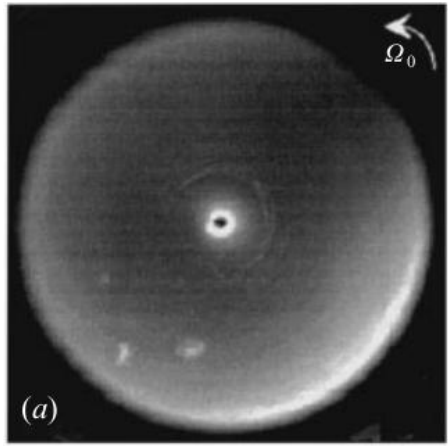
$$Re = H^2 \Omega / \nu$$

$$\text{Aspect ratio : } \Gamma = R / H$$

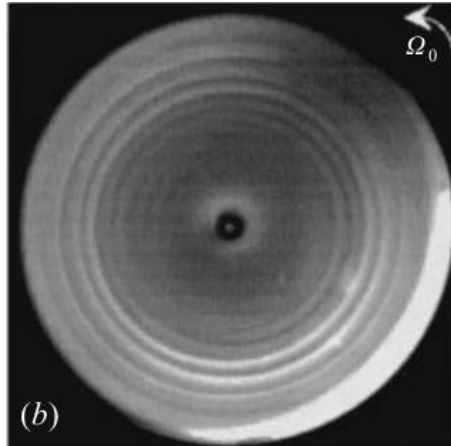


Circular rolls

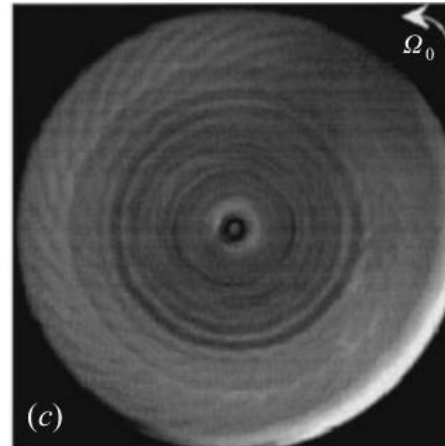
Circular + spiral rolls



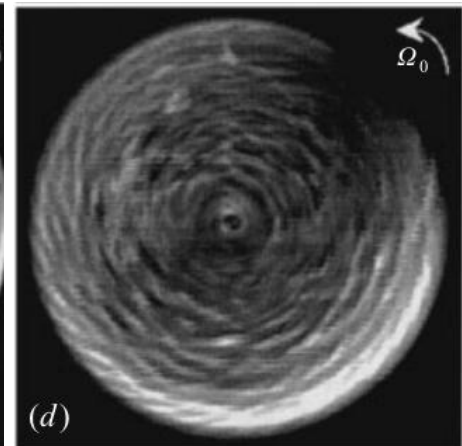
Re=150



Re=500



Re=700



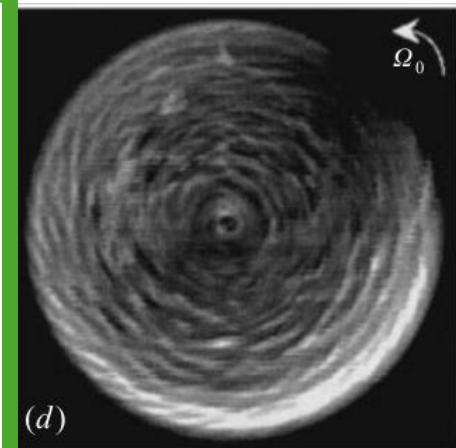
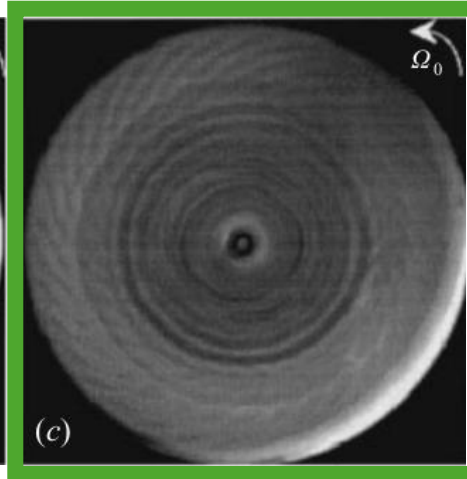
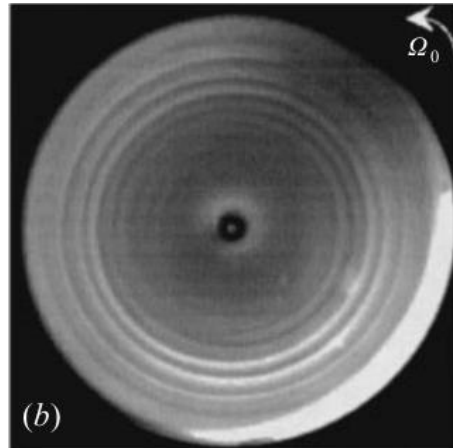
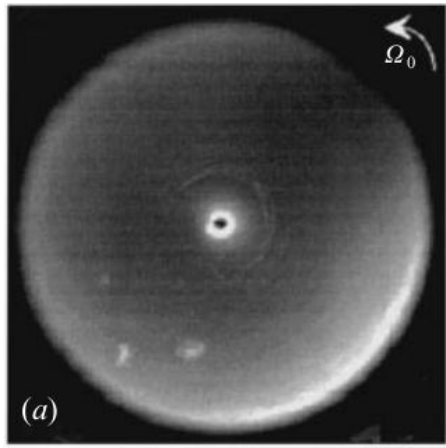
Re=1200

Re

Gauthier *et al.*, 1999

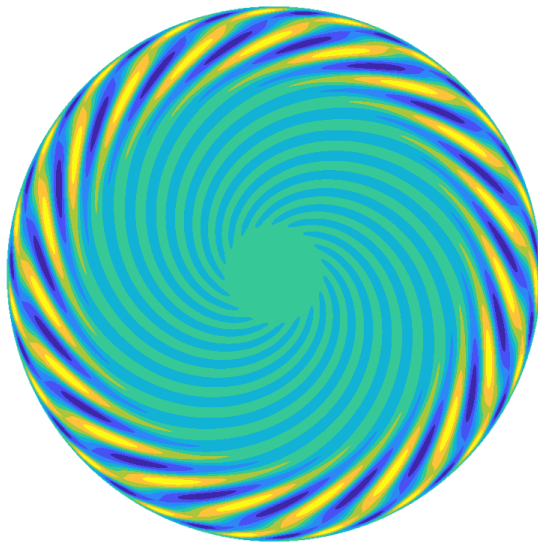
Circular rolls

Circular + spiral rolls



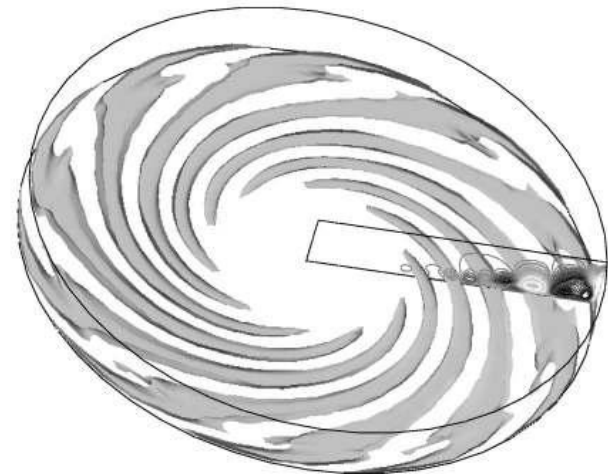
Gauthier *et al.*, 1999

First linearly unstable mode,  $R/H=10$



Gelfgat, Fluid Dyn. Res. 2015

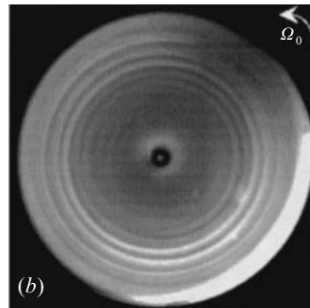
Saturated state, DNS,  $R/H=5$



Serre *et al.*, PoF 2004

Challenge: No clear explanation of the circular rolls.

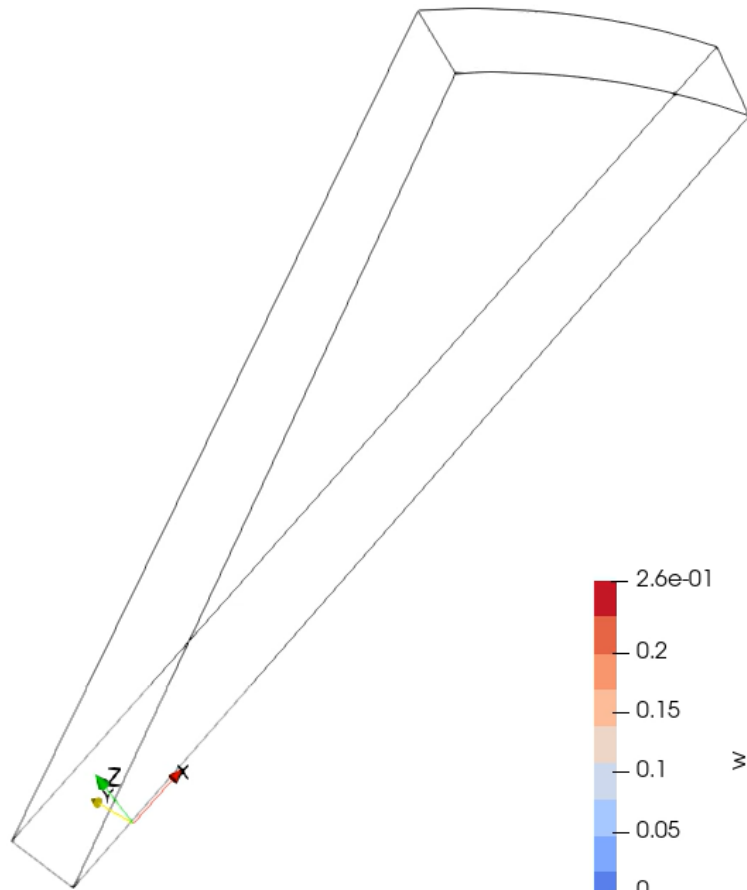
Experiment



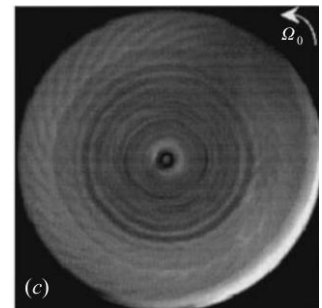
DNS

Re=500

Time: 0.000000



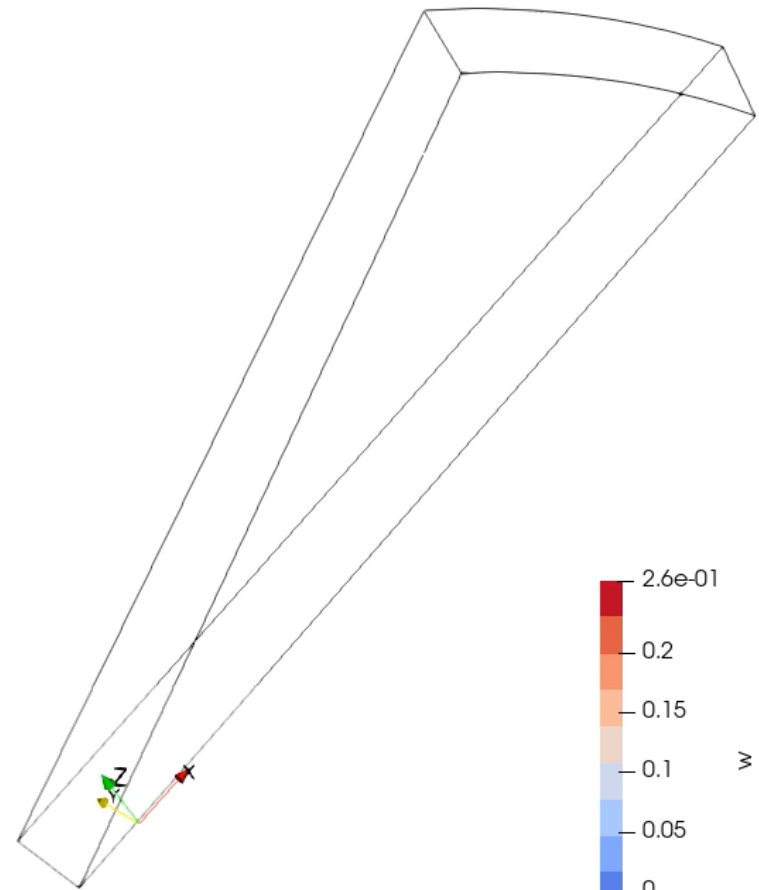
Experiment



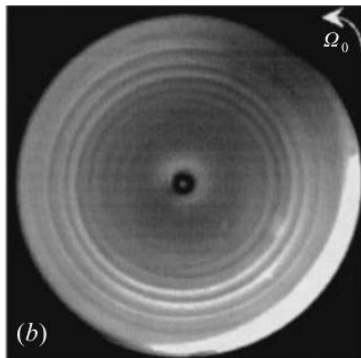
DNS

Re=700

Time: 0.000000



## Experiment

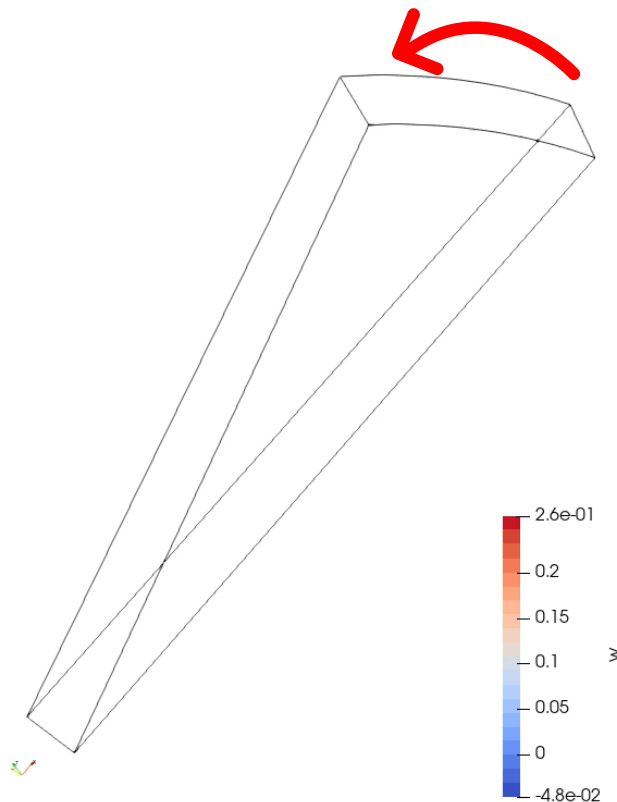


Re=500 - forcing

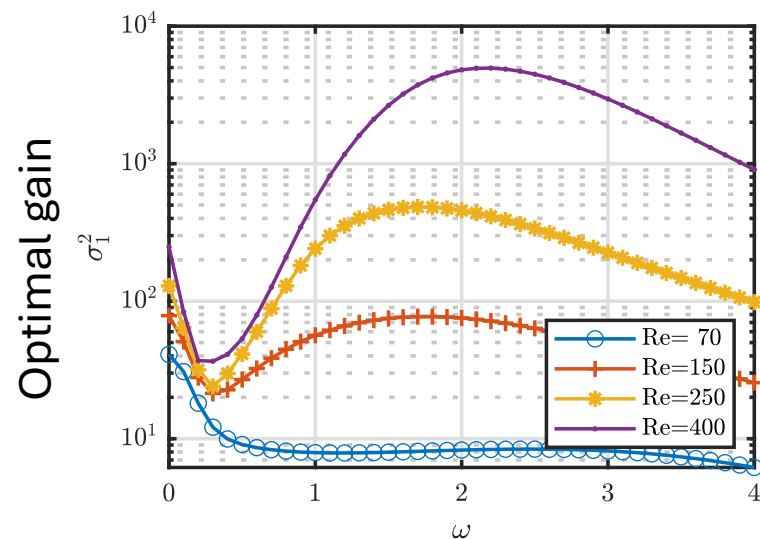
Time: 0.000000

Experimentally seen  
circular rolls are a  
forcing response.

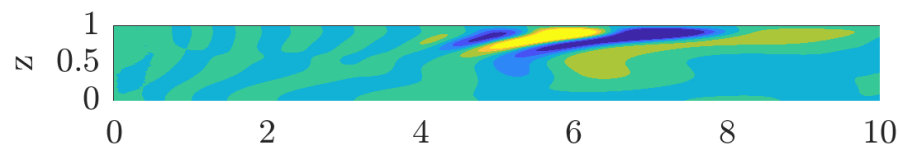
$$\Omega(t) = 1 + \text{white noise}$$



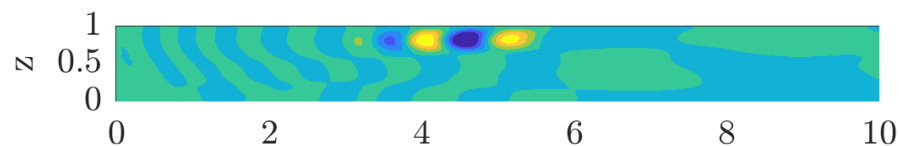
## Resolvent analysis (axisymmetric)



$f_\theta$  optimal forcing



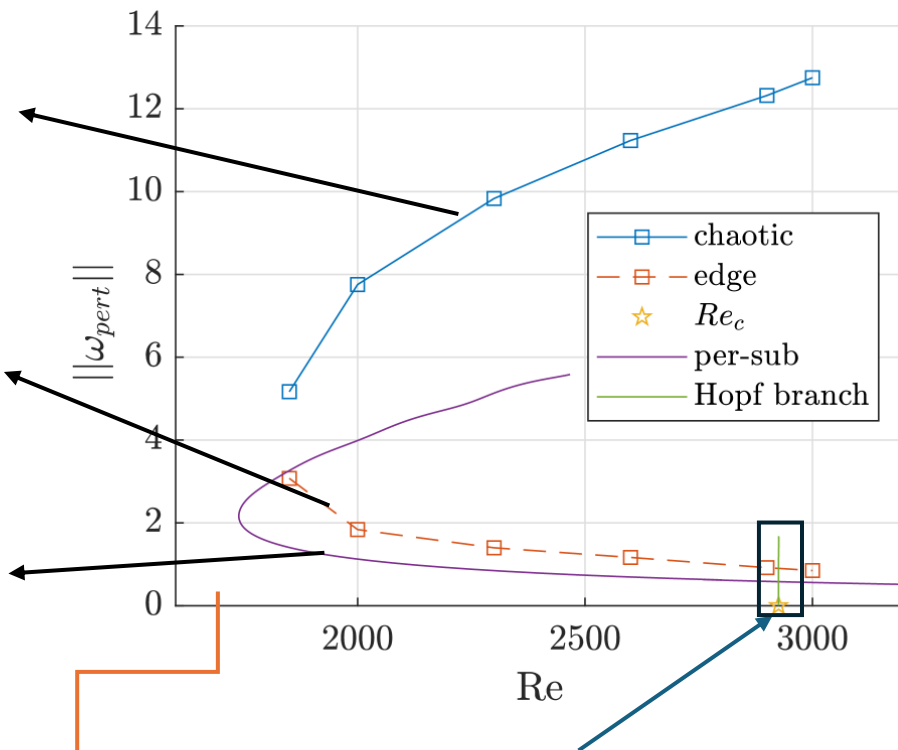
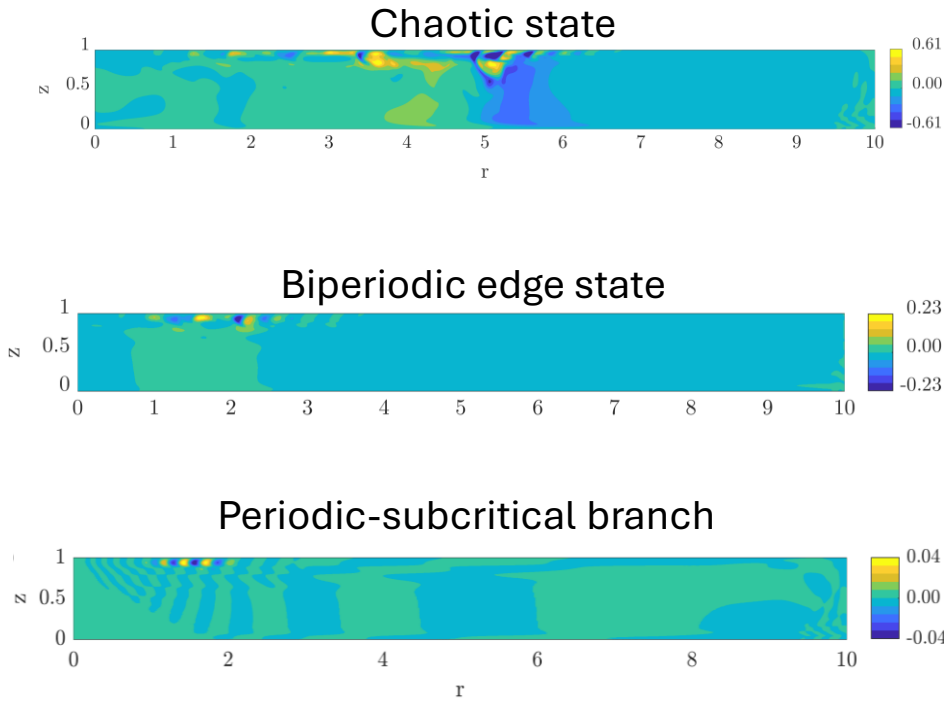
$u_\theta$  optimal response



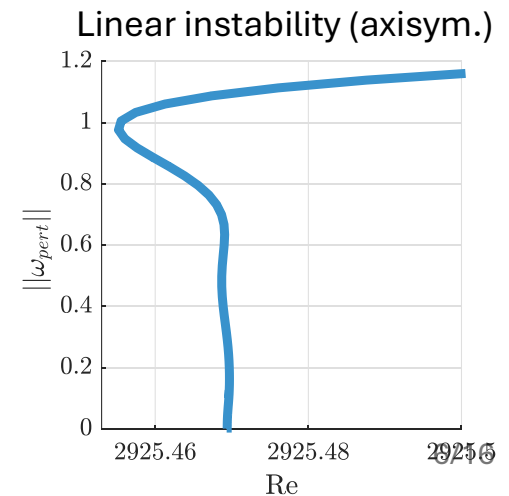
Gesla *et al.* On the origin of circular rolls in rotor-stator flow, JFM 2024.

# Self-sustained axisymmetric solutions

Gesla *et al.* Subcritical axisymmetric solutions in rotor-stator flow, PRF 2024.

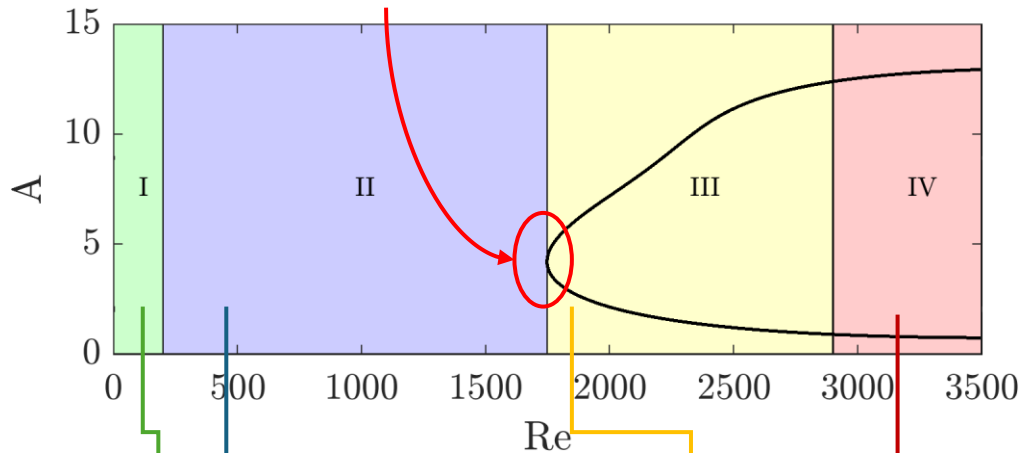


Experimental rolls –  $Re \approx 500$



# Dynamics of circular rolls

Finite lifetimes



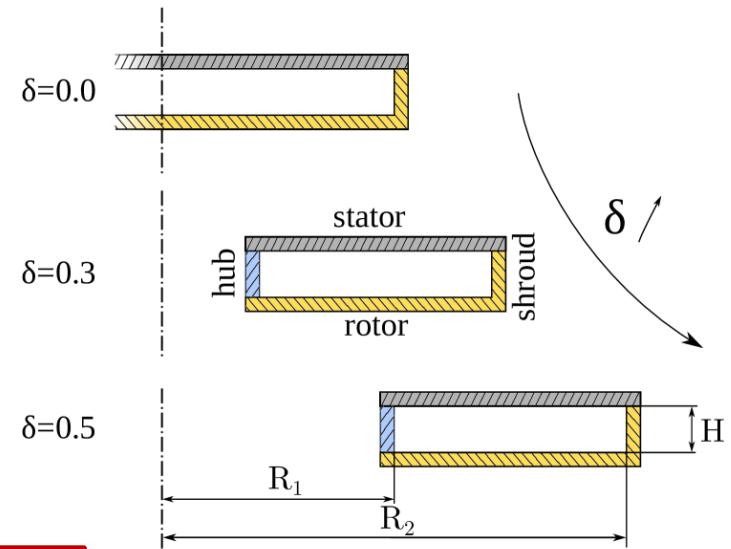
Stable flow,  
no rolls

Lin. stable flow, self-  
sustained solutions exist

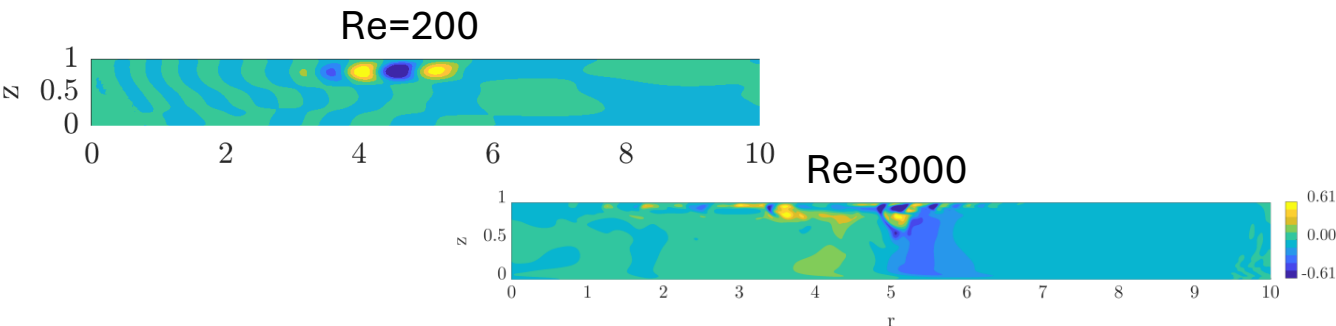
Stable flow, circular rolls  
appear as a forcing response

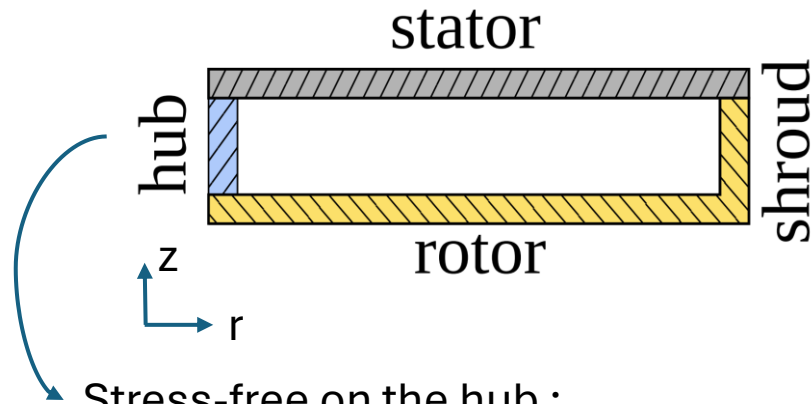
Unstable flow, evolution  
to chaotic dynamics

Use **homotopy** to explain the  
origin of the self-sustained states.



$$\delta = R_1/R_2$$

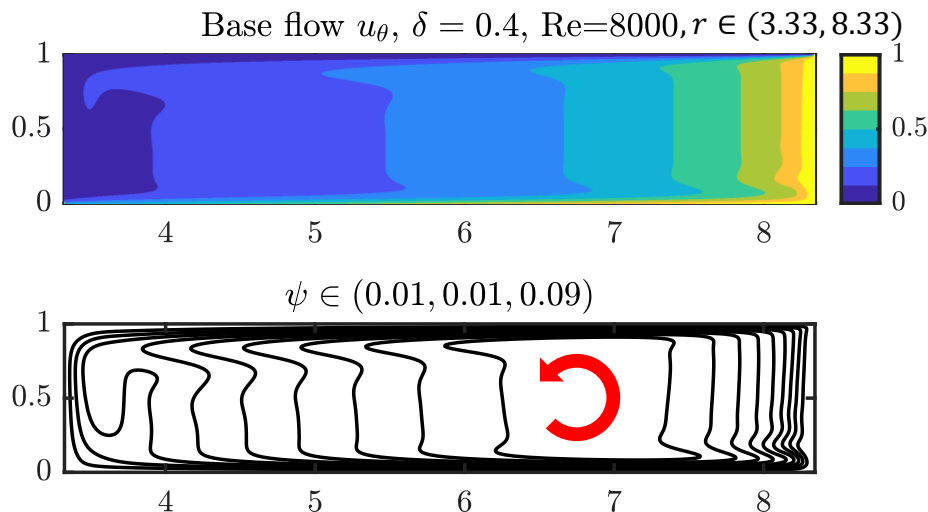




$$(u_r, \frac{\partial u_\theta}{\partial r} - \frac{u_\theta}{r}, \frac{\partial u_z}{\partial r}) = (0, 0, 0) \text{ at } r_1 = R_1/H.$$

If hub at  $r=0$  (rotor-stator) :

$$(u_r, u_\theta, \frac{\partial u_z}{\partial r}) = (0, 0, 0).$$



$$\frac{\partial \mathbf{u}}{\partial t} + \nabla(\mathbf{u} \otimes \mathbf{u}) = -\nabla p + \frac{1}{Re} \nabla^2 \mathbf{u}$$

$$\nabla \cdot \mathbf{u} = 0$$

Aspect ratio  $R/H=5$

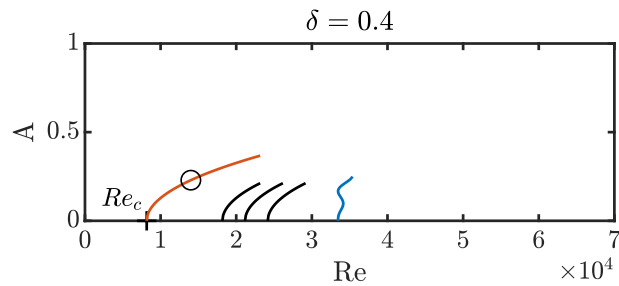
Boundary conditions :

$$\begin{aligned} \mathbf{u} &= \mathbf{0} \text{ at } z = 1 && \text{stator,} \\ \mathbf{u} &= H r / R_2 \mathbf{e}_\theta \text{ at } z = 0 && \text{rotor,} \\ \mathbf{u} &= \mathbf{e}_\theta \text{ at } r_2 = R_2/H, && \text{shroud.} \end{aligned}$$

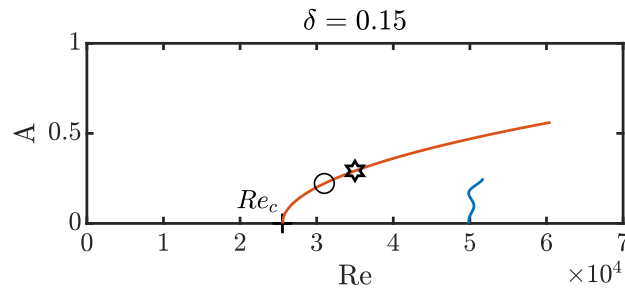
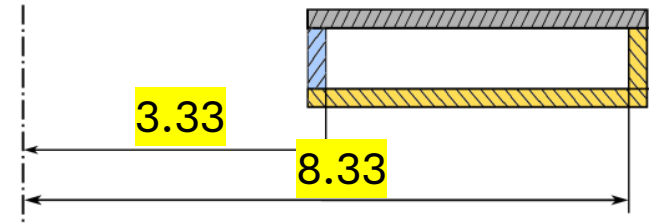
Numerical methods:

- Finite Volume discretization r-z
- Steady state - Newton method
- Instability - ARPACK
- Time integration with BDF2 scheme

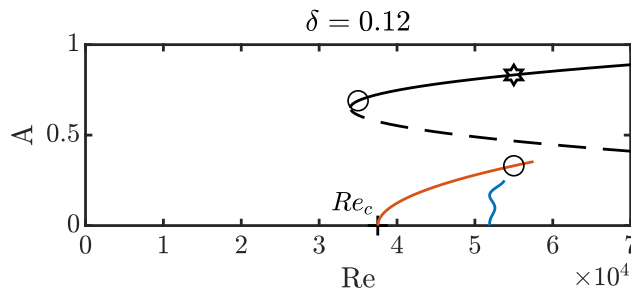
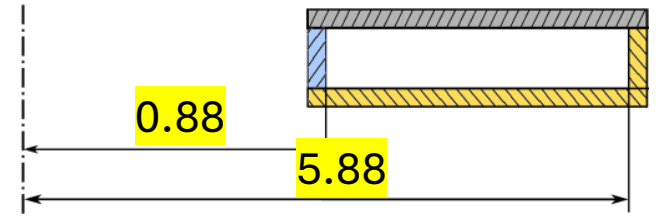




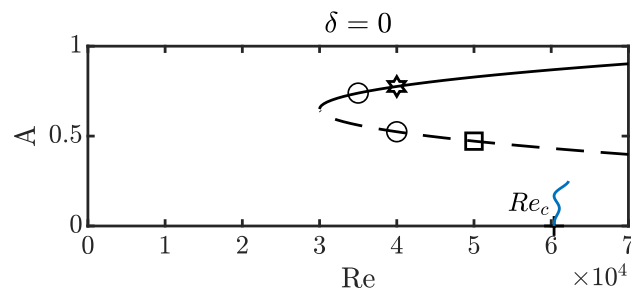
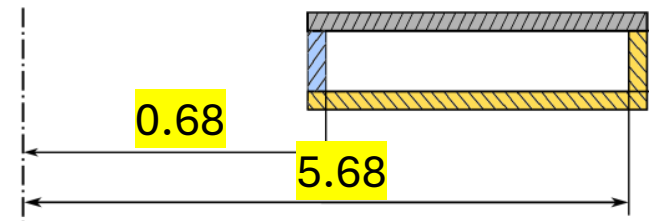
Supercritical transition  
 ○ – periodic state



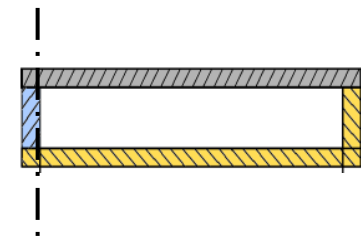
Supercritical transition  
 ○ – periodic state  
 ★ - chaotic state

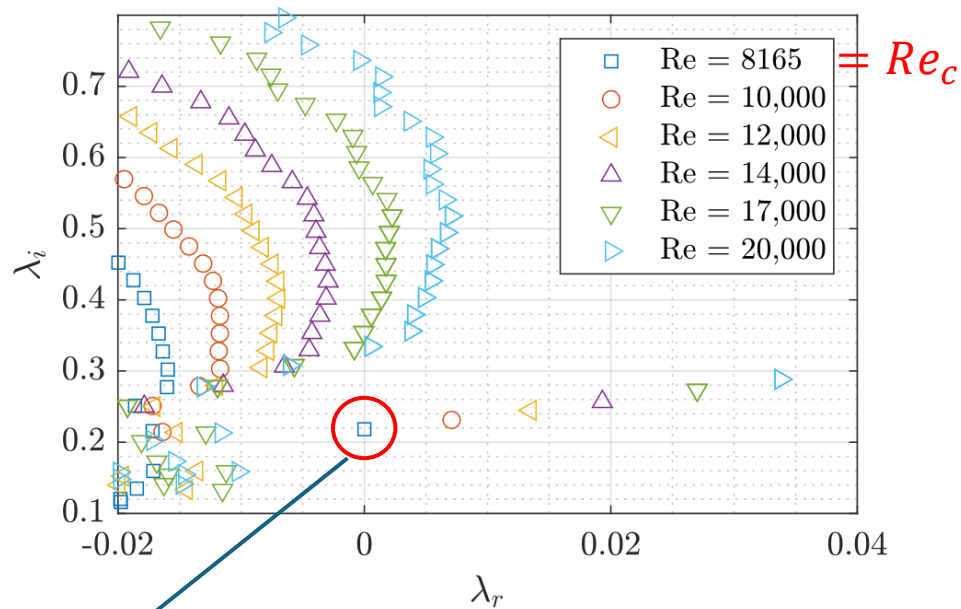
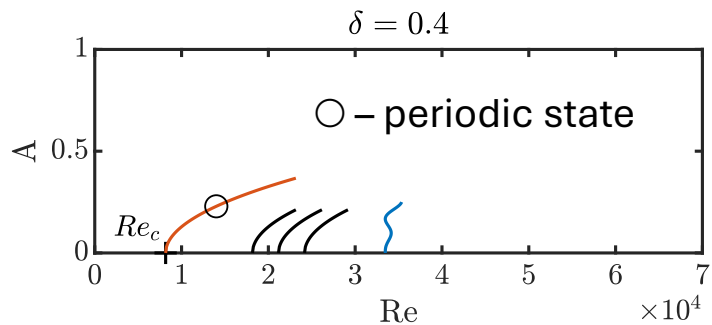


Competition between branches  
 ○ – periodic state  
 ★ - chaotic state

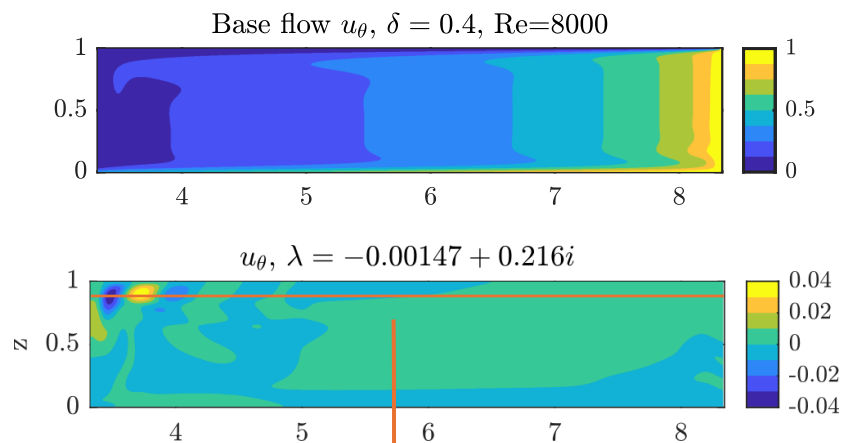


Subcritical transition  
 ○ – periodic state  
 □ - biperiodic state  
 ★ - chaotic state

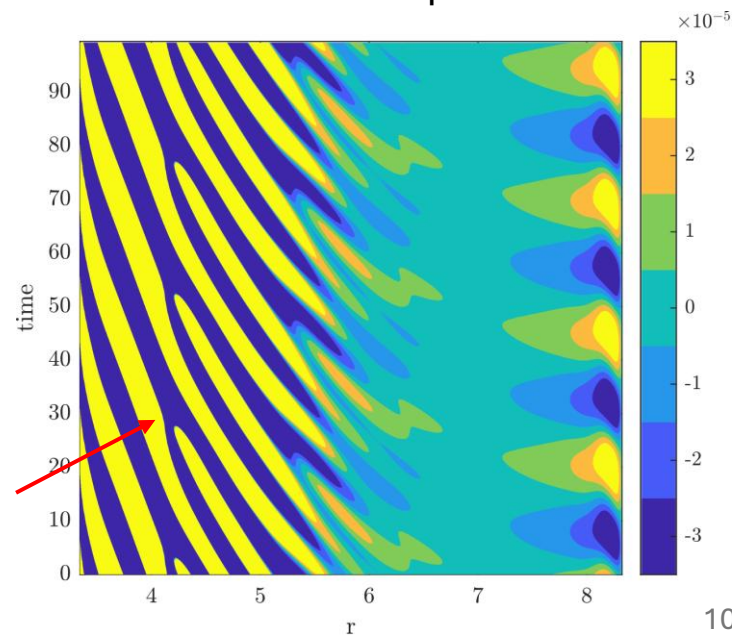




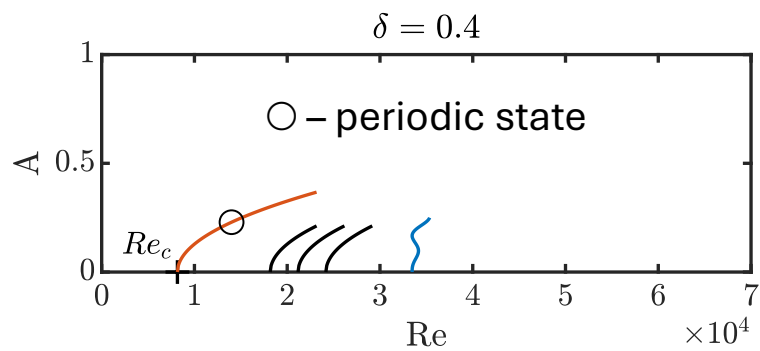
Periodic state on a supercritical branch



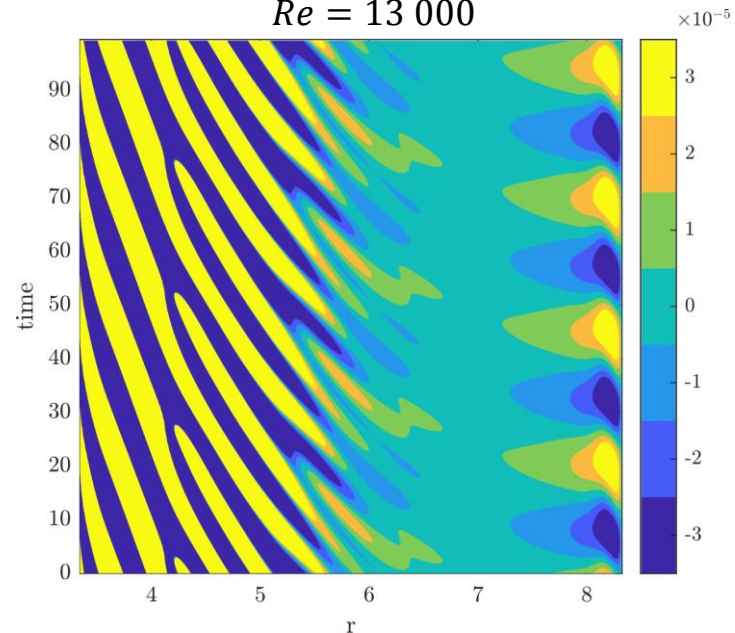
Roll pairing



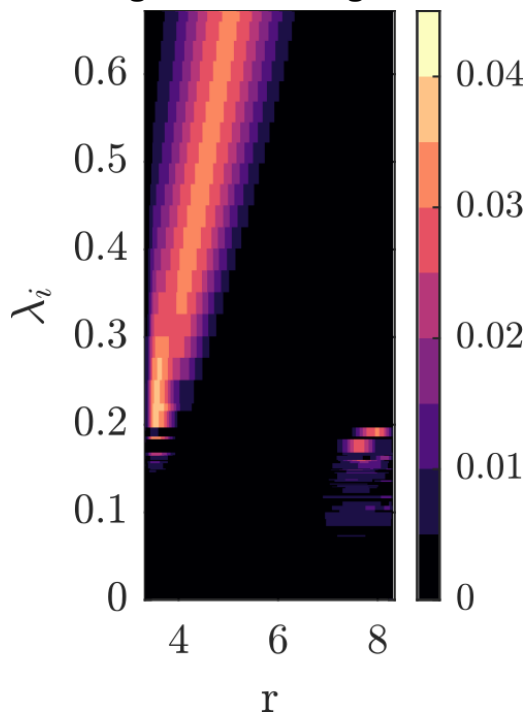
# Pairing scenario



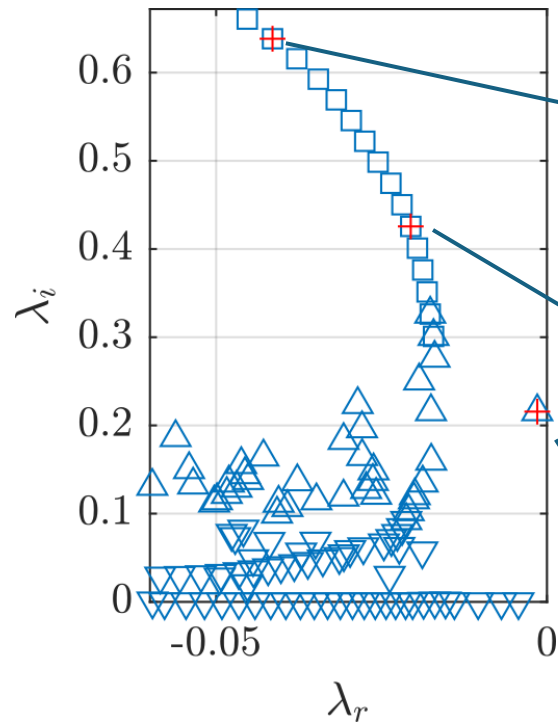
$Re = 13\,000$



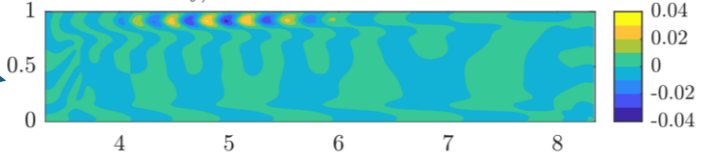
## Eigenvector magnitude



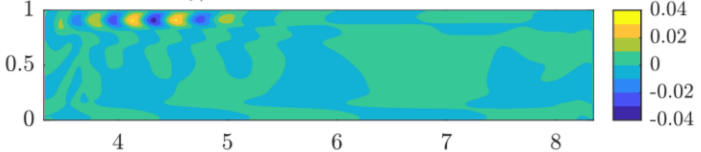
$Re = 8\,000$



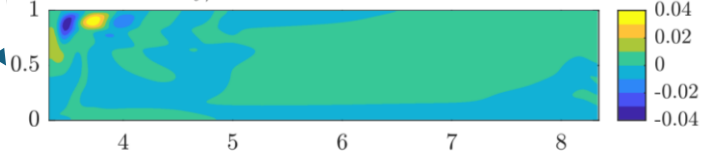
$u_\theta, \lambda = -0.04147 + 0.638i$



$u_\theta, \lambda = -0.02059 + 0.426i$

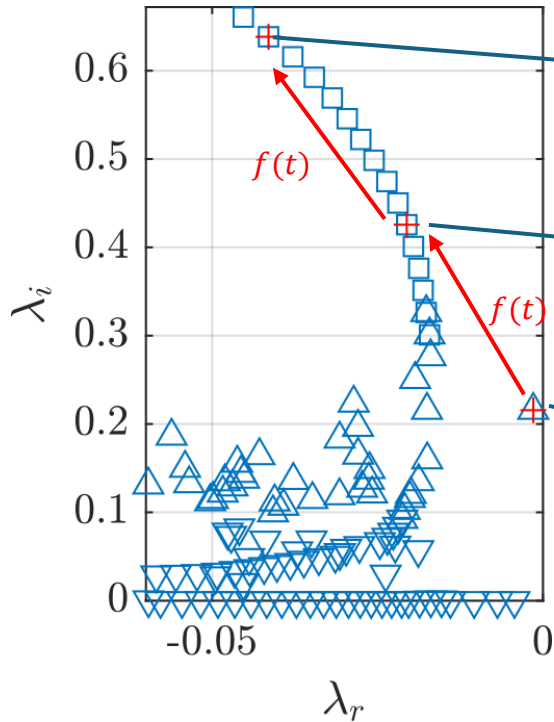


$u_\theta, \lambda = -0.00147 + 0.216i$

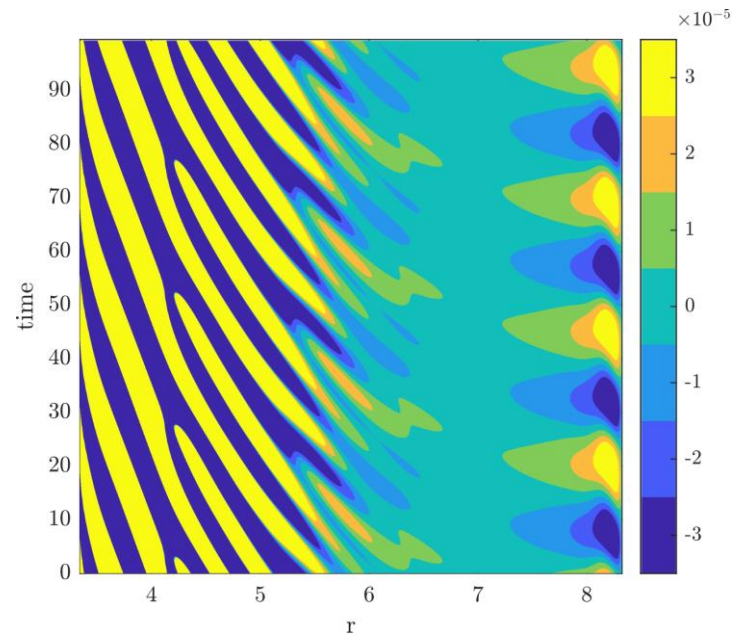
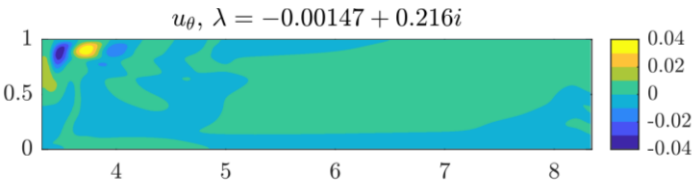
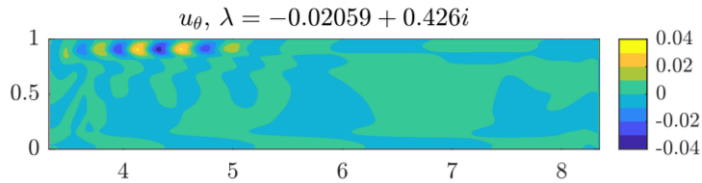
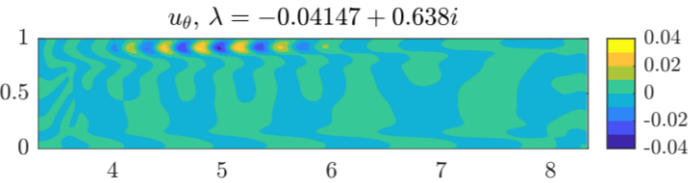


## Pairing scenario

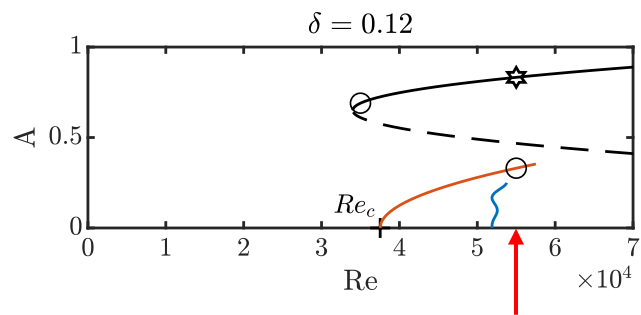
$Re = 8\,000$



$$u \propto e^{\lambda t}$$



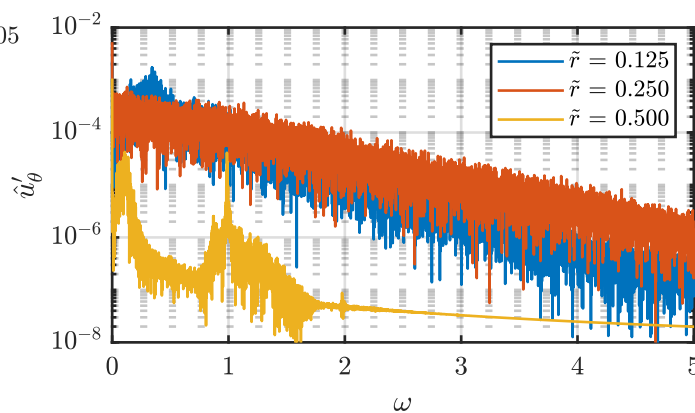
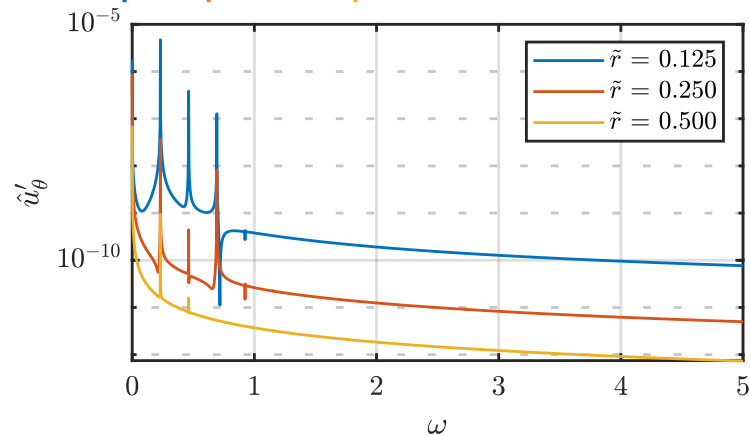
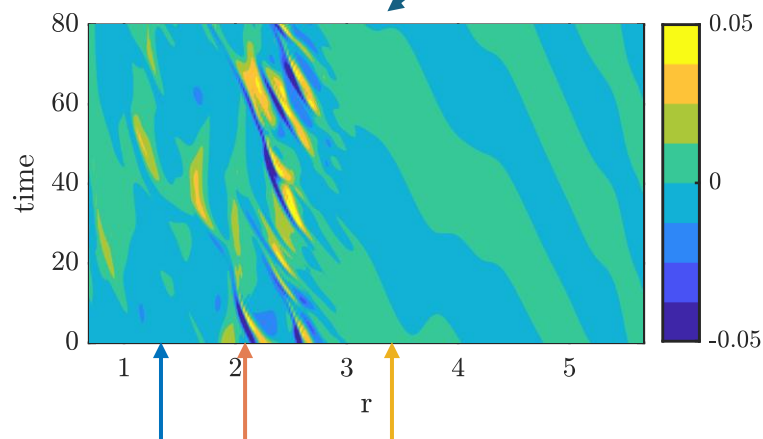
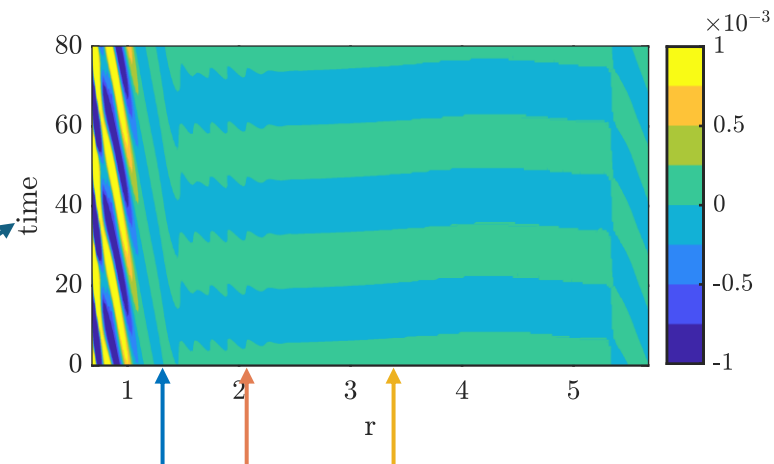
- Interaction of  $\lambda_i = 0.216i$  with itself forces the  $\lambda_i = 0.426i$
- Interaction of these two forces the  $\lambda_i = 0.638i$
- Different spatial support and number of rolls give an illusion of pairing

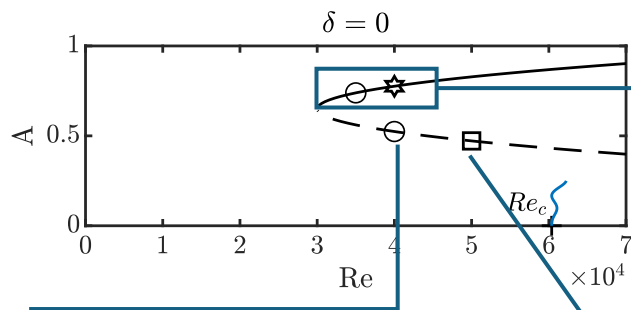


**Re = 55 000**

○ – periodic state

★ - chaotic state

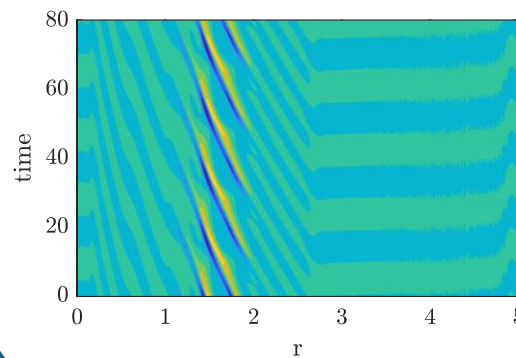




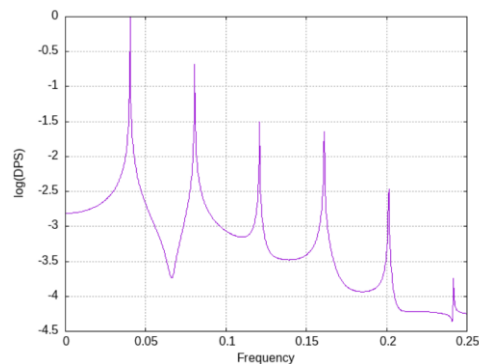
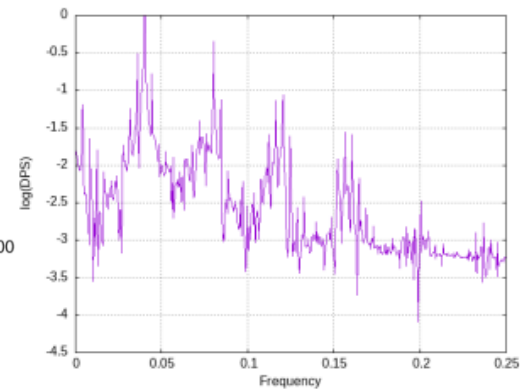
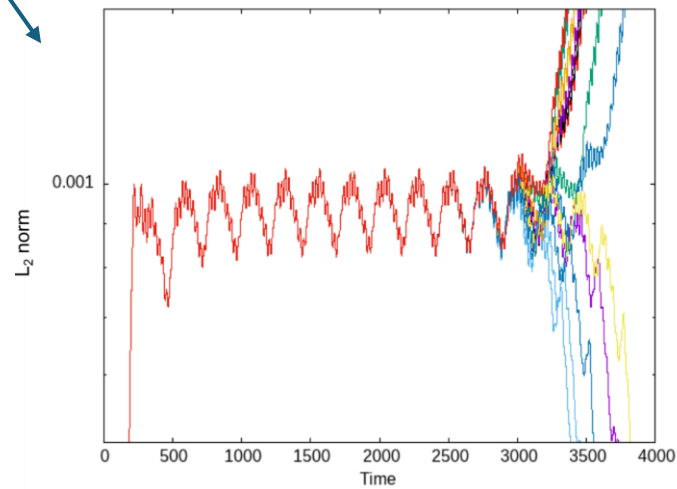
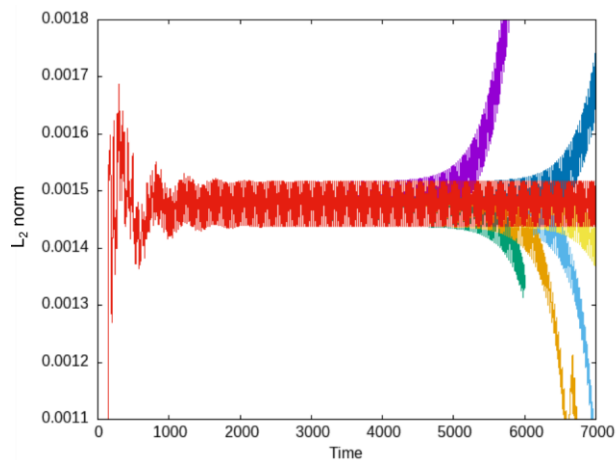
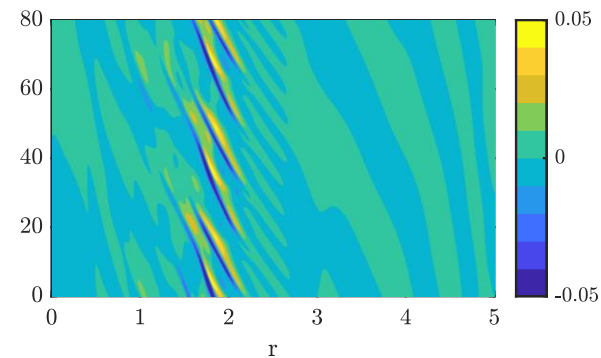
Subcritical transition

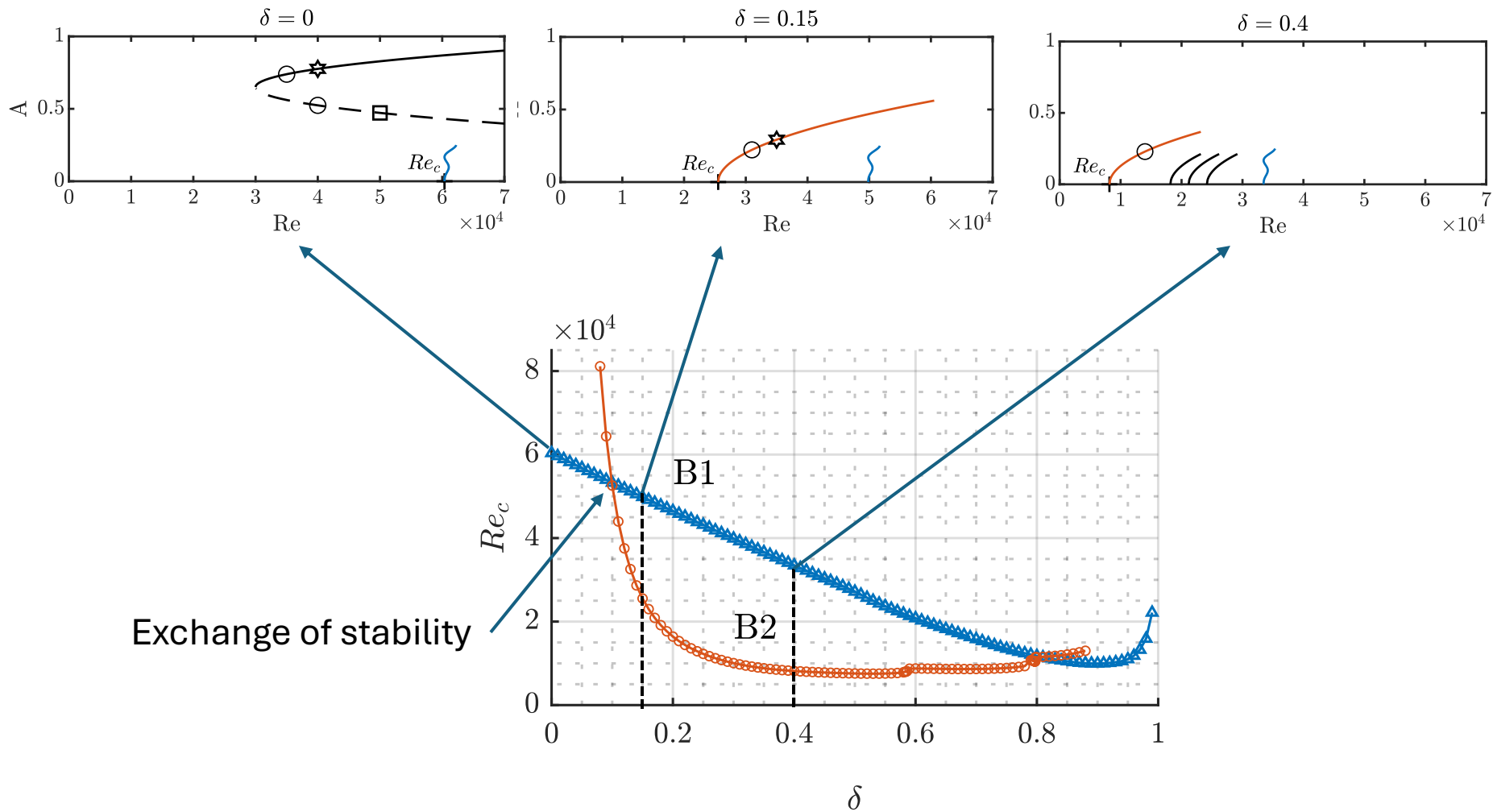
- – periodic state
- – biperiodic state
- ★ – chaotic state

Periodic top branch



Chaotic top branch



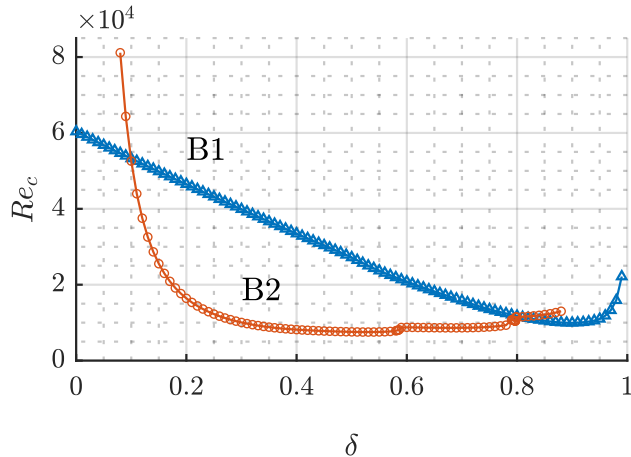


A supercritical branch is replaced by a steep Hopf branch for decreasing  $\delta$ .  
Laminar-turbulent transition becomes subcritical.

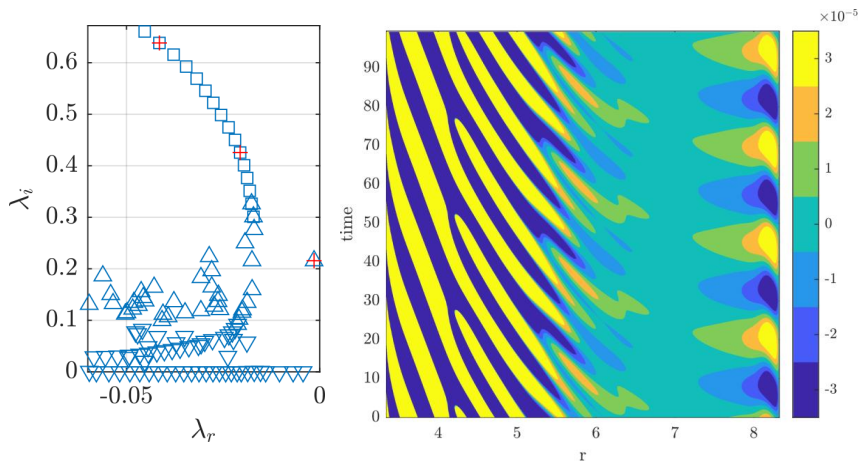




## Summary:



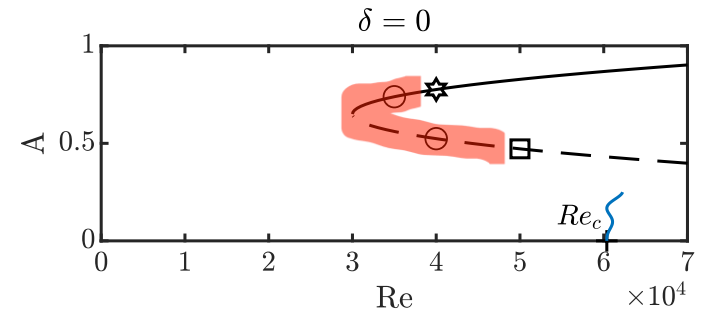
Stability exchange for a radially displaced cavity.



New perspective on the roll pairing.

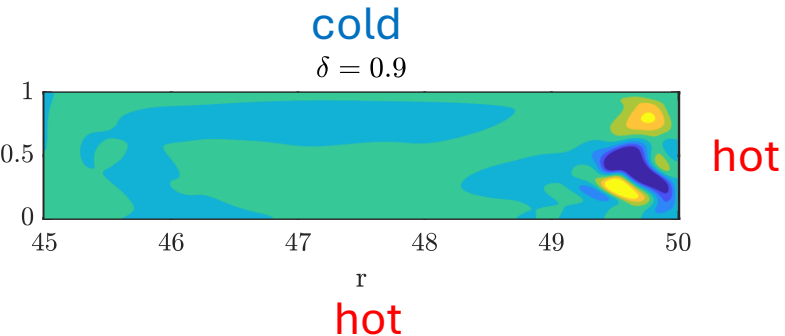
Gesla *et al.* From annular cavity to rotor-stator flow: nonlinear dynamics of axisymmetric rolls, accepted PRF 2025.

## Outlooks:



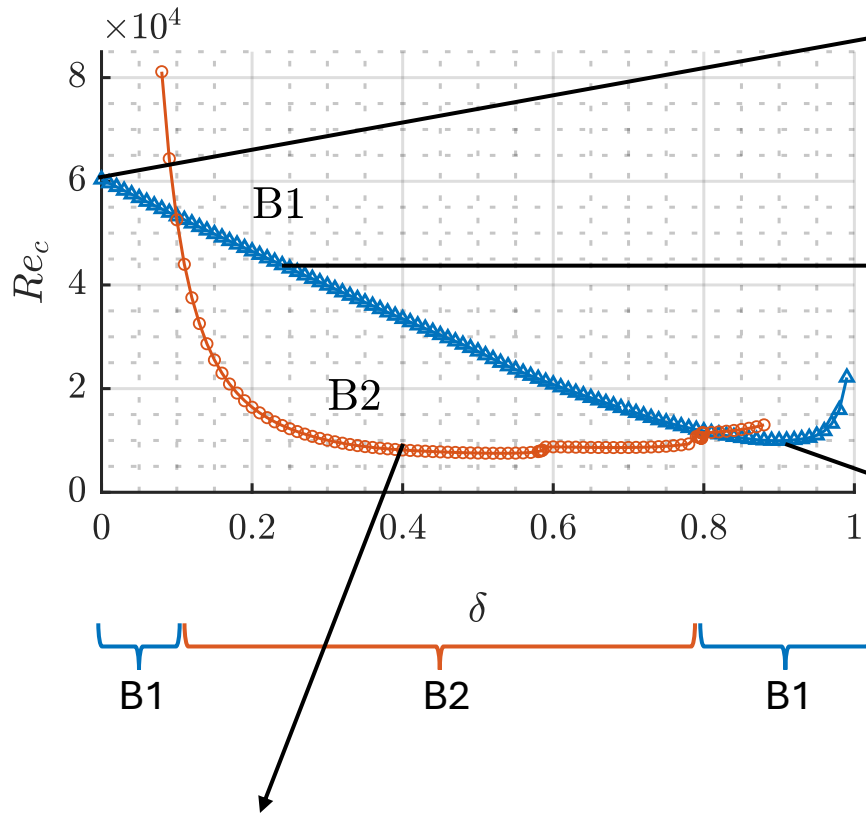
Self-Consistent Model, Harmonic Balance Method, Floquet Analysis on the periodic states identified

adiabatic

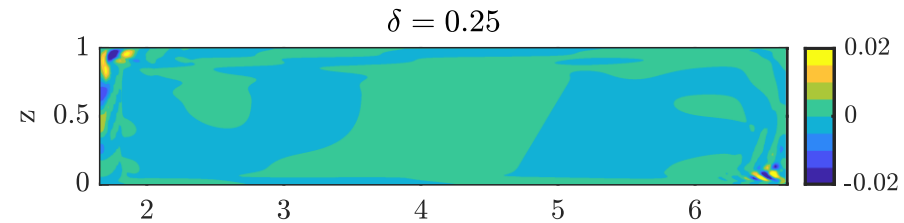
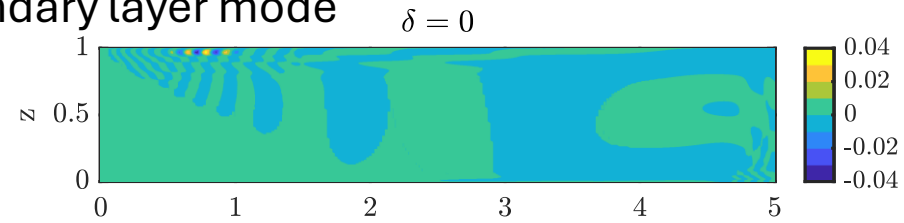


$\delta \rightarrow 1$  analysis, similarity to a Differentially Heated Cavity for  $Pr = 1$  and  $\Delta\Omega \ll \Omega$ , comparison of thresholds in  $Ra$  and  $Re$

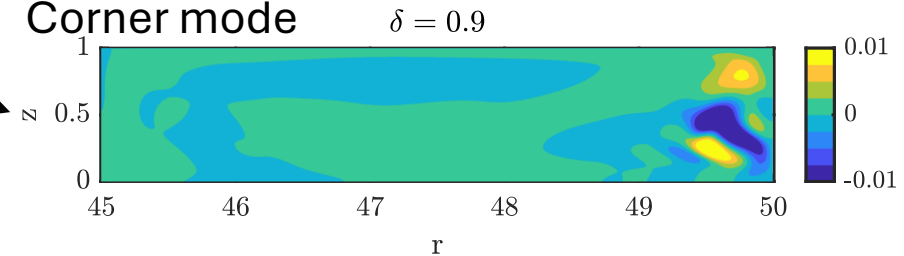
# Critical Re



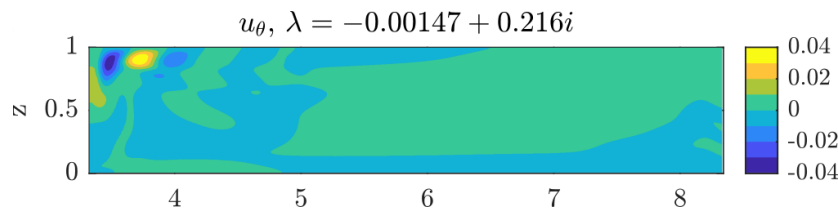
## Boundary layer mode



## Corner mode



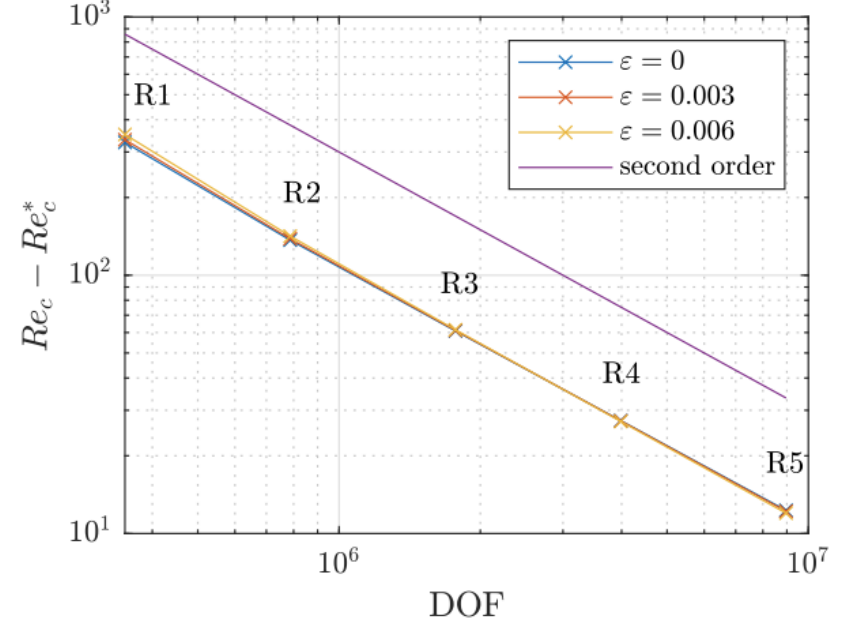
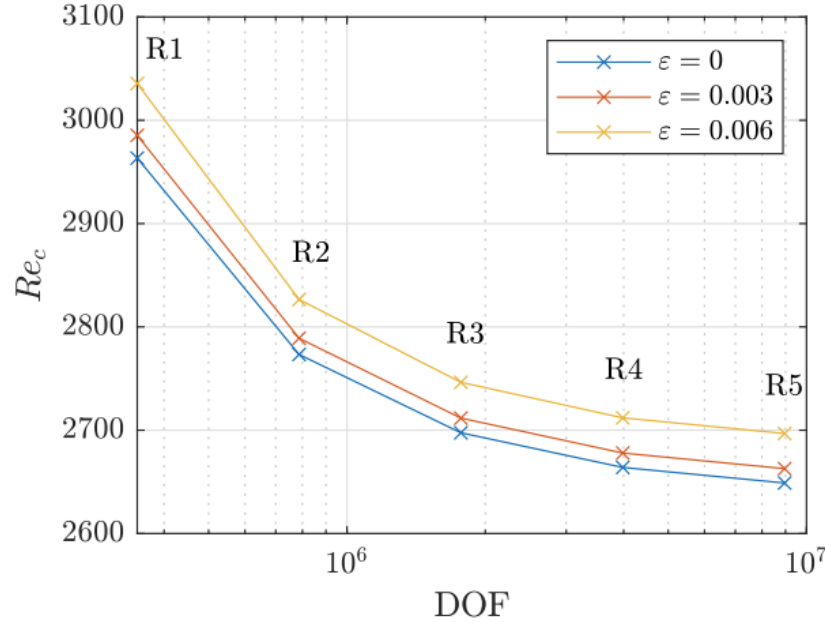
- Branch B1 is responsible for the linear instability at  $\delta < 0.1$  and  $\delta > 0.8$ .
- Eigenmode changes continuously from a boundary layer structure to a corner structure.





# Corner singularity

the corner singularity is also considered. This is achieved here by smoothing out the boundary condition at the Bödewadt corner, imposing an exponential velocity profile of the form  $u_\theta = r \exp\left(\frac{r-\Gamma}{\varepsilon}\right)$ . Two regularisations have been considered :  $\varepsilon = 0.003$  and  $\varepsilon = 0.006$ . The case without any regularization ( $u_\theta = 0$ ) is referred to as  $\varepsilon = 0$  for ease of notation.



resolution	$N_r$	$N_z$	type	DOF	$\varepsilon = 0$	$\varepsilon = 0.003$	$\varepsilon = 0.006$
<b>R0</b>	<b>600</b>	<b>160</b>	uniform	<b>390 k</b>	<b>2925.47</b>		
R1	683	128	non-uniform	356 k	2963.41	2985.43	3035.61
<b>R2</b>	<b>1024</b>	<b>192</b>	non-uniform	<b>796 k</b>	<b>2773.3</b>	2789.01	2826.55
R3	1536	288	non-uniform	1.7 m	2697.48	2711.71	2746.33
R4	2304	432	non-uniform	4.0 m	2663.96	2677.9	2711.97
R5	3456	648	non-uniform	9.0 m	2648.9	2662.8	2696.8
				extrapolation	2636.61	2650.59	2684.81
				order	2.01	2.04	2.09

TABLE V. Critical Reynolds number  $Re_c$  depending on the the spatial discretisation. From R1 to R5 the ratio between two consecutive grid resolutions is 1.5 in each direction.