

# Modern Quantum Chemistry, Szabo & Ostlund

## HW

WSF

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## 7 The 1-Particle Many-body Green's Function

### 7.1 Green's Function in Single-Particle Systems

**Ex 7.1**

$$\mathbf{V} = \mathbf{G}_0(E)^{-1} - \mathbf{G}(E)^{-1} \quad (7.1.1)$$

thus

$$\begin{aligned} \mathbf{G}_0(E)\mathbf{V}\mathbf{G}(E) &= \mathbf{G}_0(E)[\mathbf{G}_0(E)^{-1} - \mathbf{G}(E)^{-1}]\mathbf{G}(E) \\ &= \mathbf{G}(E) - \mathbf{G}_0(E) \end{aligned} \quad (7.1.2)$$

i.e.

$$\mathbf{G}(E) = \mathbf{G}_0(E) + \mathbf{G}_0(E)\mathbf{V}\mathbf{G}(E) \quad (7.1.3)$$

**Ex 7.2**

**a.** When  $x = 0$ ,

$$\begin{aligned} \left. \frac{d^2}{dx^2} |x| \right|_{x=0} &= \lim_{\epsilon \rightarrow 0} \frac{\left. \frac{d|x|}{dx} \right|_{x=\epsilon} - \left. \frac{d|x|}{dx} \right|_{x=-\epsilon}}{2\epsilon} \quad (\epsilon > 0) \\ &= \lim_{\epsilon \rightarrow 0} \frac{1 - (-1)}{2\epsilon} \\ &= \infty \end{aligned} \quad (7.1.4)$$

otherwise,

$$\begin{aligned} \frac{d^2}{dx^2} |x| &= \frac{d^2}{dx^2} [x \operatorname{sgn}(x)] \\ &= \frac{d}{dx} [1 \times \operatorname{sgn}(x) + x \times 0] \\ &= 0 \end{aligned} \quad (7.1.5)$$

**b.**

$$\begin{aligned} \int_{-\infty}^{\infty} \frac{d^2}{dx^2} |x| dx &= \int_{-\infty}^{\infty} d \left( \frac{d}{dx} |x| \right) \\ &= \left. \frac{d}{dx} |x| \right|_{-\infty}^{\infty} \\ &= 1 - (-1) \\ &= 2 \end{aligned} \quad (7.1.6)$$

thus

$$\frac{d^2}{dx^2} |x| = 2\delta(x) \quad (7.1.7)$$

**c.**

$$\begin{aligned} \frac{d^2}{dx^2} a(x) &= \frac{d^2}{dx^2} \frac{1}{2} \int_{\alpha}^{\beta} dx' |x - x'| b(x') \\ &= \frac{d^2}{dx^2} \frac{1}{2} \int_{\alpha}^x dx' (x - x') b(x') + \frac{d^2}{dx^2} \frac{1}{2} \int_x^{\beta} dx' [-(x - x')] b(x') \\ &= \frac{d}{dx} \frac{1}{2} \int_{\alpha}^x dx' b(x') - \frac{d}{dx} \frac{1}{2} \int_x^{\beta} dx' b(x') \\ &= \frac{1}{2} b(x) - \frac{1}{2} [-b(x)] \\ &= b(x) \end{aligned} \quad (7.1.8)$$

**Ex 7.3**

$$\begin{aligned}
\left(E + \frac{1}{2} \frac{d^2}{dx^2}\right) G_0(x, x', E) &= \left(E + \frac{1}{2} \frac{d^2}{dx^2}\right) \frac{1}{i(2E)^{1/2}} e^{i(2E)^{1/2}|x-x'|} \\
&= \frac{E}{i(2E)^{1/2}} e^{i(2E)^{1/2}|x-x'|} + \frac{1}{2} \frac{1}{i(2E)^{1/2}} \frac{d^2}{dx^2} e^{i(2E)^{1/2}|x-x'|} \\
&= \frac{E}{i(2E)^{1/2}} e^{i(2E)^{1/2}|x-x'|} + \frac{1}{2} \frac{1}{i(2E)^{1/2}} \frac{d}{dx} \left[ e^{i(2E)^{1/2}|x-x'|} i(2E)^{1/2} \frac{d}{dx} |x-x'| \right] \\
&= \frac{E}{i(2E)^{1/2}} e^{i(2E)^{1/2}|x-x'|} + \frac{1}{2} \left[ e^{i(2E)^{1/2}|x-x'|} i(2E)^{1/2} \left( \frac{d}{dx} |x-x'| \right)^2 + e^{i(2E)^{1/2}|x-x'|} \frac{d^2}{dx^2} |x-x'| \right] \\
&= \frac{E}{i(2E)^{1/2}} e^{i(2E)^{1/2}|x-x'|} + \frac{1}{2} e^{i(2E)^{1/2}|x-x'|} \left[ i(2E)^{1/2} \times 1 + 2\delta(x-x') \right] \\
&= e^{i(2E)^{1/2}|x-x'|} \left[ \frac{E}{i(2E)^{1/2}} + \frac{-E}{i(2E)^{1/2}} + \delta(x-x') \right] \\
&= e^{i(2E)^{1/2}|x-x'|} \delta(x-x') \\
&= \delta(x-x')
\end{aligned} \tag{7.1.9}$$

**Ex 7.4**

$$\begin{aligned}
\phi_n(x) \phi_n^*(x') &= \lim_{E \rightarrow E_n} (E - E_n) \frac{1}{i(2E)^{1/2}} \left[ e^{i(2E)^{1/2}|x-x'|} - \frac{e^{i(2E)^{1/2}(|x|+|x'|)}}{1 + i(2E)^{1/2}} \right] \\
&= \lim_{E \rightarrow -1/2} (E + 1/2) \frac{1}{-1} \left[ e^{-|x-x'|} - \frac{e^{-(|x|+|x'|)}}{1 + i(2E)^{1/2}} \right] \\
&= - \lim_{E \rightarrow -1/2} (E + 1/2) e^{-|x-x'|} + \lim_{E \rightarrow -1/2} (E + 1/2) \frac{e^{-(|x|+|x'|)}}{1 + i(2E)^{1/2}} \\
&= 0 + \lim_{E \rightarrow -1/2} (E + 1/2) \frac{e^{-(|x|+|x'|)} (1 - i(2E)^{1/2})}{(1 + i(2E)^{1/2})(1 - i(2E)^{1/2})} \\
&= \lim_{E \rightarrow -1/2} (E + 1/2) \frac{e^{-(|x|+|x'|)} (1 - i(2E)^{1/2})}{1 + 2E} \\
&= \frac{1}{2} e^{-(|x|+|x'|)} (1 - (-1)) \\
&= e^{-(|x|+|x'|)}
\end{aligned} \tag{7.1.10}$$

Let  $x = x'$ ,

$$\phi_n^2(x) = e^{-2|x|} \tag{7.1.11}$$

thus

$$\phi_n(x) = e^{-|x|} \tag{7.1.12}$$

**Ex 7.5**

$$\begin{aligned}
\mathcal{H} \phi &= \left[ -\frac{1}{2} \frac{d^2}{dx^2} - \delta(x) \right] e^{-|x|} \\
&= -\frac{1}{2} \frac{d}{dx} \left[ e^{-|x|} \left( -\frac{d}{dx} |x| \right) \right] - \delta(x) e^{-|x|} \\
&= \frac{1}{2} \left[ -e^{-|x|} \left( \frac{d}{dx} |x| \right)^2 + e^{-|x|} \frac{d^2}{dx^2} |x| \right] - \delta(x) e^{-|x|} \\
&= \frac{1}{2} \left[ -e^{-|x|} + e^{-|x|} \times 2\delta(x) \right] - \delta(x) e^{-|x|} \\
&= -\frac{1}{2} e^{-|x|}
\end{aligned} \tag{7.1.13}$$

thus the eigenvalue is  $-\frac{1}{2}$ .

**Ex 7.6**

**a.**

$$\begin{aligned}
 i \frac{\partial}{\partial t} \phi(x, t) &= i \int dx' \frac{\partial G(x, x', t)}{\partial t} \psi(x') \\
 &= \int dx' \mathcal{H} G(x, x', t) \psi(x') \\
 &= \mathcal{H} \phi(x, t)
 \end{aligned} \tag{7.1.14}$$

**b.** From

$$i \frac{\partial G(x, x', t)}{\partial t} = \mathcal{H} G(x, x', t) \tag{7.1.15}$$

we get

$$\lim_{\varepsilon \rightarrow 0} \int_0^\infty dt i \frac{\partial G(x, x', t)}{\partial t} [-i e^{(iE-\varepsilon)t}] = \lim_{\varepsilon \rightarrow 0} \int_0^\infty dt \mathcal{H} G(x, x', t) [-i e^{(iE-\varepsilon)t}] \tag{7.1.16}$$

$$\begin{aligned}
 \lim_{\varepsilon \rightarrow 0} \int_0^\infty dt \frac{\partial G(x, x', t)}{\partial t} e^{(iE-\varepsilon)t} &= \int_0^\infty dt \mathcal{H} G(x, x', t) [-i e^{iEt}] \\
 &= \mathcal{H} G(x, x', E)
 \end{aligned} \tag{7.1.17}$$

thus

$$\lim_{\varepsilon \rightarrow 0} \left[ G(x, x', t) e^{(iE-\varepsilon)t} \right]_{t=0}^\infty - \int_0^\infty dt G(x, x', t) e^{(iE-\varepsilon)t} (iE - \varepsilon) = \mathcal{H} G(x, x', E) \tag{7.1.18}$$

$$\begin{aligned}
 \mathcal{H} G(x, x', E) &= -G(x, x', 0) - iE \int_0^\infty dt G(x, x', t) e^{iEt} \\
 &= -G(x, x', 0) - iEG(x, x', E)/(-i) \\
 &= -\delta(x - x') + EG(x, x', E)
 \end{aligned} \tag{7.1.19}$$

$\therefore$

$$(E - \mathcal{H})G(x, x', E) = \delta(x - x') \tag{7.1.20}$$

**c.**

$$\begin{aligned}
 i \frac{\partial}{\partial t} \mathcal{G}(t) &= i \frac{\partial}{\partial t} e^{-i\mathcal{H}t} \\
 &= i e^{-i\mathcal{H}t} (-i\mathcal{H}) \\
 &= \mathcal{H} \mathcal{G}(t)
 \end{aligned} \tag{7.1.21}$$

$$\lim_{\varepsilon \rightarrow 0} \int_0^\infty dt e^{(iE-\varepsilon)t} i \frac{\partial}{\partial t} \mathcal{G}(t) = \lim_{\varepsilon \rightarrow 0} \int_0^\infty dt e^{(iE-\varepsilon)t} \mathcal{H} \mathcal{G}(t) \tag{7.1.22}$$

$$\lim_{\varepsilon \rightarrow 0} \left[ e^{(iE-\varepsilon)t} \mathcal{G}(t) \right]_0^\infty - (iE - \varepsilon) \int_0^\infty dt e^{(iE-\varepsilon)t} \mathcal{G}(t) = \mathcal{H} \mathcal{G}(E) \tag{7.1.23}$$

$\therefore$

$$\begin{aligned}
 \mathcal{H} \mathcal{G}(E) &= \lim_{\varepsilon \rightarrow 0} \left[ -\mathcal{G}(0) - (iE - \varepsilon) \int_0^\infty dt e^{(iE-\varepsilon)t} \mathcal{G}(t) \right] \\
 &= -\mathcal{G}(0) + E\mathcal{G}(E) \\
 &= -1 + E\mathcal{G}(E)
 \end{aligned} \tag{7.1.24}$$

thus

$$\mathcal{G}(E) = \frac{1}{E - \mathcal{H}} \tag{7.1.25}$$

## 7.2 The 1-Particle Many-body Green's Function

### 7.2.1 The Self-Energy

**Ex 7.7**

$$\begin{aligned}\Sigma_{ij}^{(2)}(E) &= \frac{1}{2} \sum_{ars} \frac{\langle rs || ia \rangle \langle ja || rs \rangle}{E + \varepsilon_a - \varepsilon_r - \varepsilon_s} + \frac{1}{2} \sum_{abr} \frac{\langle ab || ir \rangle \langle jr || ab \rangle}{E + \varepsilon_r - \varepsilon_a - \varepsilon_b} \\ &= \frac{1}{2} \sum_{ars} \frac{(\langle rs | ia \rangle - \langle rs | ai \rangle)(\langle ja | rs \rangle - \langle ja | sr \rangle)}{E + \varepsilon_a - \varepsilon_r - \varepsilon_s} + \frac{1}{2} \sum_{abr} \frac{(\langle ab | ir \rangle - \langle ab | ri \rangle)(\langle jr | ab \rangle - \langle jr | ba \rangle)}{E + \varepsilon_r - \varepsilon_a - \varepsilon_b}\end{aligned}\quad (7.2.1)$$

In the 1st summation:

To make the terms non-zero, the spin of  $r$  is fixed in the first and last term, and  $r, s, a$  are all fixed in the second and third term, thus

$$\begin{aligned}\text{the 1st term} &= \frac{1}{2} \sum_{ars}^{N/2} \frac{1}{E + \varepsilon_a - \varepsilon_r - \varepsilon_s} [2 \langle rs | ia \rangle \langle ja | rs \rangle - \langle rs | ai \rangle \langle ja | rs \rangle - \langle rs | ia \rangle \langle ja | sr \rangle + 2 \langle rs | ai \rangle \langle ja | sr \rangle] \\ &= \sum_{ars}^{N/2} \frac{1}{E + \varepsilon_a - \varepsilon_r - \varepsilon_s} [2 \langle rs | ia \rangle \langle ja | rs \rangle - \langle rs | ia \rangle \langle ja | sr \rangle] \\ &= \sum_{ars}^{N/2} \frac{\langle rs | ia \rangle [2 \langle ja | rs \rangle - \langle aj | rs \rangle]}{E + \varepsilon_a - \varepsilon_r - \varepsilon_s}\end{aligned}\quad (7.2.2)$$

Similarly,

$$\Sigma_{ij}^{(2)}(E) = \sum_{ars}^{N/2} \frac{\langle rs | ia \rangle [2 \langle ja | rs \rangle - \langle aj | rs \rangle]}{E + \varepsilon_a - \varepsilon_r - \varepsilon_s} + \sum_{abr}^{N/2} \frac{\langle ab | ir \rangle [2 \langle jr | ab \rangle - \langle rj | ab \rangle]}{E + \varepsilon_r - \varepsilon_a - \varepsilon_b}\quad (7.2.3)$$

**Ex 7.8**

$$\begin{aligned}[\mathbf{G}_0(E)]_{ij} &= \sum_m \frac{\langle {}^N\Psi_0 | a_i^\dagger a_m | {}^N\Psi_0 \rangle \langle a_m {}^N\Psi_0 | a_j | {}^N\Psi_0 \rangle}{E - (\langle {}^N\Psi_0 | \mathcal{H} | {}^N\Psi_0 \rangle - \langle a_m {}^N\Psi_0 | \mathcal{H} | a_m {}^N\Psi_0 \rangle)} + \sum_p \frac{\langle {}^N\Psi_0 | a_j a_p^\dagger | {}^N\Psi_0 \rangle \langle a_p^\dagger {}^N\Psi_0 | a_i^\dagger | {}^N\Psi_0 \rangle}{E + (\langle {}^N\Psi_0 | \mathcal{H} | {}^N\Psi_0 \rangle - \langle a_p^\dagger {}^N\Psi_0 | \mathcal{H} | a_p^\dagger {}^N\Psi_0 \rangle)} \\ &= \sum_m \frac{\delta_{im} \delta_{mj}}{E - \varepsilon_m} + 0 \\ &= \sum_m \frac{\delta_{ij}}{E - \varepsilon_m}\end{aligned}\quad (7.2.4)$$

### 7.2.2 The Solution of the Dyson Equation

## 7.3 Application of the Formalism to $\text{H}_2$ and $\text{HeH}^+$

**Ex 7.9**

**a.**

$${}^{N+1}\mathcal{O}_0 = \quad (7.3.1)$$

**Ex 7.10** Since

$$\begin{aligned}\Sigma_{11}^{(2)}(\varepsilon_1) &= \frac{K_{12}}{\varepsilon_1 + \varepsilon_1 - 2\varepsilon_2} \\ &= \frac{K_{12}}{2(\varepsilon_1 - \varepsilon_2)}\end{aligned}\quad (7.3.2)$$

$$\begin{aligned}
\Sigma_{11}^{(3)}(\varepsilon_1) &= \frac{K_{12}^2(J_{22} - 2J_{12} + K_{12})}{(\varepsilon_1 - 2\varepsilon_2 + \varepsilon_1)^2} + \frac{K_{12}^2(J_{11} - 2J_{12} + K_{12})}{(\varepsilon_1 - 2\varepsilon_2 + \varepsilon_1)(\varepsilon_1 - \varepsilon_2)} + \frac{K_{12}^2(2J_{12} - K_{12} - J_{11})}{4(\varepsilon_1 - \varepsilon_2)^2} \\
&= \frac{K_{12}^2(J_{22} - 2J_{12} + K_{12})}{4(\varepsilon_1 - \varepsilon_2)^2} + \frac{K_{12}^2(J_{11} - 2J_{12} + K_{12})}{2(\varepsilon_1 - \varepsilon_2)^2} + \frac{K_{12}^2(2J_{12} - K_{12} - J_{11})}{4(\varepsilon_1 - \varepsilon_2)^2} \\
&= \frac{K_{12}^2(J_{22} + J_{11} - 4J_{12} + 2K_{12})}{4(\varepsilon_1 - \varepsilon_2)^2}
\end{aligned} \tag{7.3.3}$$

thus

$$\Sigma_{11}^{(2)}(\varepsilon_1) = E_0^{(2)} \tag{7.3.4}$$

$$\Sigma_{11}^{(3)}(\varepsilon_1) = E_0^{(3)} \tag{7.3.5}$$

Similarly,

$$\begin{aligned}
\Sigma_{22}^{(2)}(\varepsilon_2) &= \frac{K_{12}}{\varepsilon_2 + \varepsilon_2 - 2\varepsilon_1} \\
&= \frac{K_{12}}{2(\varepsilon_2 - \varepsilon_1)}
\end{aligned} \tag{7.3.6}$$

$$\begin{aligned}
\Sigma_{22}^{(3)}(\varepsilon_2) &= \frac{K_{12}^2(2J_{12} - K_{12} - J_{11})}{(\varepsilon_2 - 2\varepsilon_1 + \varepsilon_2)^2} + \frac{K_{12}^2(J_{22} - 2J_{12} + K_{12})}{(\varepsilon_2 - 2\varepsilon_1 + \varepsilon_2)(\varepsilon_1 - \varepsilon_2)} + \frac{K_{12}^2(J_{22} + K_{12} - 2J_{12})}{4(\varepsilon_1 - \varepsilon_2)^2} \\
&= \frac{K_{12}^2(2J_{12} - K_{12} - J_{11})}{4(\varepsilon_1 - \varepsilon_2)^2} - \frac{K_{12}^2(J_{22} - 2J_{12} + K_{12})}{2(\varepsilon_1 - \varepsilon_2)^2} + \frac{K_{12}^2(J_{22} + K_{12} - 2J_{12})}{4(\varepsilon_1 - \varepsilon_2)^2} \\
&= \frac{K_{12}^2(-J_{11} - J_{22} + 4J_{12} - 2K_{12})}{4(\varepsilon_1 - \varepsilon_2)^2}
\end{aligned} \tag{7.3.7}$$

thus

$$\Sigma_{22}^{(2)}(\varepsilon_2) = -E_0^{(2)} \tag{7.3.8}$$

$$\Sigma_{22}^{(3)}(\varepsilon_2) = -E_0^{(3)} \tag{7.3.9}$$

**Ex 7.11**

**Ex 7.12**

## 7.4 Perturbation Theory and the Green's Function Method

**Ex 7.13**

**Ex 7.14**

**Ex 7.15**

## 7.5 Some Illustrative Calculations

**Ex 7.16**