

Modern Quantum Chemistry, Szabo & Ostlund

HW

WSF

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7 The 1-Particle Many-body Green's Function

7.1 Green's Function in Single-Particle Systems

Ex 7.1

$$\mathbf{V} = \mathbf{G}_0(E)^{-1} - \mathbf{G}(E)^{-1} \quad (7.1.1)$$

thus

$$\begin{aligned} \mathbf{G}_0(E)\mathbf{V}\mathbf{G}(E) &= \mathbf{G}_0(E)[\mathbf{G}_0(E)^{-1} - \mathbf{G}(E)^{-1}]\mathbf{G}(E) \\ &= \mathbf{G}(E) - \mathbf{G}_0(E) \end{aligned} \quad (7.1.2)$$

i.e.

$$\mathbf{G}(E) = \mathbf{G}_0(E) + \mathbf{G}_0(E)\mathbf{V}\mathbf{G}(E) \quad (7.1.3)$$

Ex 7.2

a. When $x = 0$,

$$\begin{aligned} \left. \frac{d^2}{dx^2} |x| \right|_{x=0} &= \lim_{\epsilon \rightarrow 0} \frac{\left. \frac{d|x|}{dx} \right|_{x=\epsilon} - \left. \frac{d|x|}{dx} \right|_{x=-\epsilon}}{2\epsilon} \quad (\epsilon > 0) \\ &= \lim_{\epsilon \rightarrow 0} \frac{1 - (-1)}{2\epsilon} \\ &= \infty \end{aligned} \quad (7.1.4)$$

otherwise,

$$\begin{aligned} \frac{d^2}{dx^2} |x| &= \frac{d^2}{dx^2} [x \operatorname{sgn}(x)] \\ &= \frac{d}{dx} [1 \times \operatorname{sgn}(x) + x \times 0] \\ &= 0 \end{aligned} \quad (7.1.5)$$

b.

$$\begin{aligned} \int_{-\infty}^{\infty} \frac{d^2}{dx^2} |x| dx &= \int_{-\infty}^{\infty} d \left(\frac{d}{dx} |x| \right) \\ &= \left. \frac{d}{dx} |x| \right|_{-\infty}^{\infty} \\ &= 1 - (-1) \\ &= 2 \end{aligned} \quad (7.1.6)$$

thus

$$\frac{d^2}{dx^2} |x| = 2\delta(x) \quad (7.1.7)$$

c.

$$\begin{aligned} \frac{d^2}{dx^2} a(x) &= \frac{d^2}{dx^2} \frac{1}{2} \int_{\alpha}^{\beta} dx' |x - x'| b(x') \\ &= \frac{d^2}{dx^2} \frac{1}{2} \int_{\alpha}^x dx' (x - x') b(x') + \frac{d^2}{dx^2} \frac{1}{2} \int_x^{\beta} dx' [-(x - x')] b(x') \\ &= \frac{d}{dx} \frac{1}{2} \int_{\alpha}^x dx' b(x') - \frac{d}{dx} \frac{1}{2} \int_x^{\beta} dx' b(x') \\ &= \frac{1}{2} b(x) - \frac{1}{2} [-b(x)] \\ &= b(x) \end{aligned} \quad (7.1.8)$$

Ex 7.3

$$\begin{aligned}
\left(E + \frac{1}{2} \frac{d^2}{dx^2}\right) G_0(x, x', E) &= \left(E + \frac{1}{2} \frac{d^2}{dx^2}\right) \frac{1}{i(2E)^{1/2}} e^{i(2E)^{1/2}|x-x'|} \\
&= \frac{E}{i(2E)^{1/2}} e^{i(2E)^{1/2}|x-x'|} + \frac{1}{2} \frac{1}{i(2E)^{1/2}} \frac{d^2}{dx^2} e^{i(2E)^{1/2}|x-x'|} \\
&= \frac{E}{i(2E)^{1/2}} e^{i(2E)^{1/2}|x-x'|} + \frac{1}{2} \frac{1}{i(2E)^{1/2}} \frac{d}{dx} \left[e^{i(2E)^{1/2}|x-x'|} i(2E)^{1/2} \frac{d}{dx} |x-x'| \right] \\
&= \frac{E}{i(2E)^{1/2}} e^{i(2E)^{1/2}|x-x'|} + \frac{1}{2} \left[e^{i(2E)^{1/2}|x-x'|} i(2E)^{1/2} \left(\frac{d}{dx} |x-x'| \right)^2 + e^{i(2E)^{1/2}|x-x'|} \frac{d^2}{dx^2} |x-x'| \right] \\
&= \frac{E}{i(2E)^{1/2}} e^{i(2E)^{1/2}|x-x'|} + \frac{1}{2} e^{i(2E)^{1/2}|x-x'|} \left[i(2E)^{1/2} \times 1 + 2\delta(x-x') \right] \\
&= e^{i(2E)^{1/2}|x-x'|} \left[\frac{E}{i(2E)^{1/2}} + \frac{-E}{i(2E)^{1/2}} + \delta(x-x') \right] \\
&= e^{i(2E)^{1/2}|x-x'|} \delta(x-x') \\
&= \delta(x-x')
\end{aligned} \tag{7.1.9}$$

Ex 7.4

$$\begin{aligned}
\phi_n(x) \phi_n^*(x') &= \lim_{E \rightarrow E_n} (E - E_n) \frac{1}{i(2E)^{1/2}} \left[e^{i(2E)^{1/2}|x-x'|} - \frac{e^{i(2E)^{1/2}(|x|+|x'|)}}{1 + i(2E)^{1/2}} \right] \\
&= \lim_{E \rightarrow -1/2} (E + 1/2) \frac{1}{-1} \left[e^{-|x-x'|} - \frac{e^{-(|x|+|x'|)}}{1 + i(2E)^{1/2}} \right] \\
&= - \lim_{E \rightarrow -1/2} (E + 1/2) e^{-|x-x'|} + \lim_{E \rightarrow -1/2} (E + 1/2) \frac{e^{-(|x|+|x'|)}}{1 + i(2E)^{1/2}} \\
&= 0 + \lim_{E \rightarrow -1/2} (E + 1/2) \frac{e^{-(|x|+|x'|)} (1 - i(2E)^{1/2})}{(1 + i(2E)^{1/2})(1 - i(2E)^{1/2})} \\
&= \lim_{E \rightarrow -1/2} (E + 1/2) \frac{e^{-(|x|+|x'|)} (1 - i(2E)^{1/2})}{1 + 2E} \\
&= \frac{1}{2} e^{-(|x|+|x'|)} (1 - (-1)) \\
&= e^{-(|x|+|x'|)}
\end{aligned} \tag{7.1.10}$$

Let $x = x'$,

$$\phi_n^2(x) = e^{-2|x|} \tag{7.1.11}$$

thus

$$\phi_n(x) = e^{-|x|} \tag{7.1.12}$$

Ex 7.5

$$\begin{aligned}
\mathcal{H} \phi &= \left[-\frac{1}{2} \frac{d^2}{dx^2} - \delta(x) \right] e^{-|x|} \\
&= -\frac{1}{2} \frac{d}{dx} \left[e^{-|x|} \left(-\frac{d}{dx} |x| \right) \right] - \delta(x) e^{-|x|} \\
&= \frac{1}{2} \left[-e^{-|x|} \left(\frac{d}{dx} |x| \right)^2 + e^{-|x|} \frac{d^2}{dx^2} |x| \right] - \delta(x) e^{-|x|} \\
&= \frac{1}{2} \left[-e^{-|x|} + e^{-|x|} \times 2\delta(x) \right] - \delta(x) e^{-|x|} \\
&= -\frac{1}{2} e^{-|x|}
\end{aligned} \tag{7.1.13}$$

thus the eigenvalue is $-\frac{1}{2}$.

Ex 7.6

a.

$$\begin{aligned} \mathrm{i} \frac{\partial}{\partial t} \phi(x, t) &= \mathrm{i} \int dx' \frac{\partial G(x, x', t)}{\partial t} \psi(x') \\ &= \int dx' \mathcal{H} G(x, x', t) \psi(x') \\ &= \mathcal{H} \phi(x, t) \end{aligned} \quad (7.1.14)$$

b. From

$$\mathrm{i} \frac{\partial G(x, x', t)}{\partial t} = \mathcal{H} G(x, x', t) \quad (7.1.15)$$

we get

$$\lim_{\varepsilon \rightarrow 0} \int_0^\infty dt \mathrm{i} \frac{\partial G(x, x', t)}{\partial t} [-\mathrm{i} e^{(\mathrm{i} E - \varepsilon)t}] = \lim_{\varepsilon \rightarrow 0} \int_0^\infty dt \mathcal{H} G(x, x', t) [-\mathrm{i} e^{(\mathrm{i} E - \varepsilon)t}] \quad (7.1.16)$$

$$\begin{aligned} \lim_{\varepsilon \rightarrow 0} \int_0^\infty dt \frac{\partial G(x, x', t)}{\partial t} e^{(\mathrm{i} E - \varepsilon)t} &= \int_0^\infty dt \mathcal{H} G(x, x', t) [-\mathrm{i} e^{\mathrm{i} Et}] \\ &= \mathcal{H} G(x, x', E) \end{aligned} \quad (7.1.17)$$

thus

$$\lim_{\varepsilon \rightarrow 0} \left[G(x, x', t) e^{(\mathrm{i} E - \varepsilon)t} \right]_{t=0}^\infty - \int_0^\infty dt G(x, x', t) e^{(\mathrm{i} E - \varepsilon)t} (\mathrm{i} E - \varepsilon) = \mathcal{H} G(x, x', E) \quad (7.1.18)$$

$$\begin{aligned} \mathcal{H} G(x, x', E) &= -G(x, x', 0) - \mathrm{i} E \int_0^\infty dt G(x, x', t) e^{\mathrm{i} Et} \\ &= -G(x, x', 0) - \mathrm{i} E G(x, x', E) / (-\mathrm{i}) \\ &= -\delta(x - x') + E G(x, x', E) \end{aligned} \quad (7.1.19)$$

\therefore

$$(E - \mathcal{H})G(x, x', E) = \delta(x - x') \quad (7.1.20)$$

c.

$$\begin{aligned} \mathrm{i} \frac{\partial}{\partial t} \mathcal{G}(t) &= \mathrm{i} \frac{\partial}{\partial t} e^{-\mathrm{i} \mathcal{H} t} \\ &= \mathrm{i} e^{-\mathrm{i} \mathcal{H} t} (-\mathrm{i} \mathcal{H}) \\ &= \mathcal{H} \mathcal{G}(t) \end{aligned} \quad (7.1.21)$$

$$\lim_{\varepsilon \rightarrow 0} \int_0^\infty dt e^{(\mathrm{i} E - \varepsilon)t} \mathrm{i} \frac{\partial}{\partial t} \mathcal{G}(t) = \lim_{\varepsilon \rightarrow 0} \int_0^\infty dt e^{(\mathrm{i} E - \varepsilon)t} \mathcal{H} \mathcal{G}(t) \quad (7.1.22)$$

$$\lim_{\varepsilon \rightarrow 0} \left[e^{(\mathrm{i} E - \varepsilon)t} \mathcal{G}(t) \right]_0^\infty - (\mathrm{i} E - \varepsilon) \int_0^\infty dt e^{(\mathrm{i} E - \varepsilon)t} \mathcal{G}(t) = \mathcal{H} \mathcal{G}(E) \quad (7.1.23)$$

\therefore

$$\begin{aligned} \mathcal{H} \mathcal{G}(E) &= \lim_{\varepsilon \rightarrow 0} \left[-\mathcal{G}(0) - (\mathrm{i} E - \varepsilon) \int_0^\infty dt e^{(\mathrm{i} E - \varepsilon)t} \mathcal{G}(t) \right] \\ &= -\mathcal{G}(0) + E \mathcal{G}(E) \\ &= -1 + E \mathcal{G}(E) \end{aligned} \quad (7.1.24)$$

thus

$$\mathcal{G}(E) = \frac{1}{E - \mathcal{H}} \quad (7.1.25)$$

7.2 The 1-Particle Many-body Green's Function

7.2.1 The Self-Energy

Ex 7.7

$$\begin{aligned}\Sigma_{ij}^{(2)}(E) &= \frac{1}{2} \sum_{ars} \frac{\langle rs || ia \rangle \langle ja || rs \rangle}{E + \varepsilon_a - \varepsilon_r - \varepsilon_s} + \frac{1}{2} \sum_{abr} \frac{\langle ab || ir \rangle \langle jr || ab \rangle}{E + \varepsilon_r - \varepsilon_a - \varepsilon_b} \\ &= \frac{1}{2} \sum_{ars} \frac{(\langle rs | ia \rangle - \langle rs | ai \rangle)(\langle ja | rs \rangle - \langle ja | sr \rangle)}{E + \varepsilon_a - \varepsilon_r - \varepsilon_s} + \frac{1}{2} \sum_{abr} \frac{(\langle ab | ir \rangle - \langle ab | ri \rangle)(\langle jr | ab \rangle - \langle jr | ba \rangle)}{E + \varepsilon_r - \varepsilon_a - \varepsilon_b}\end{aligned}\quad (7.2.1)$$

In the 1st summation:

To make the terms non-zero, the spin of r is fixed in the first and last term, and r, s, a are all fixed in the second and third term, thus

$$\begin{aligned}\text{the 1st term} &= \frac{1}{2} \sum_{ars}^{N/2} \frac{1}{E + \varepsilon_a - \varepsilon_r - \varepsilon_s} [2 \langle rs | ia \rangle \langle ja | rs \rangle - \langle rs | ai \rangle \langle ja | rs \rangle - \langle rs | ia \rangle \langle ja | sr \rangle + 2 \langle rs | ai \rangle \langle ja | sr \rangle] \\ &= \sum_{ars}^{N/2} \frac{1}{E + \varepsilon_a - \varepsilon_r - \varepsilon_s} [2 \langle rs | ia \rangle \langle ja | rs \rangle - \langle rs | ia \rangle \langle ja | sr \rangle] \\ &= \sum_{ars}^{N/2} \frac{\langle rs | ia \rangle [2 \langle ja | rs \rangle - \langle aj | rs \rangle]}{E + \varepsilon_a - \varepsilon_r - \varepsilon_s}\end{aligned}\quad (7.2.2)$$

Similarly,

$$\Sigma_{ij}^{(2)}(E) = \sum_{ars}^{N/2} \frac{\langle rs | ia \rangle [2 \langle ja | rs \rangle - \langle aj | rs \rangle]}{E + \varepsilon_a - \varepsilon_r - \varepsilon_s} + \sum_{abr}^{N/2} \frac{\langle ab | ir \rangle [2 \langle jr | ab \rangle - \langle rj | ab \rangle]}{E + \varepsilon_r - \varepsilon_a - \varepsilon_b}\quad (7.2.3)$$

Ex 7.8

$$\begin{aligned}[\mathbf{G}_0(E)]_{ij} &= \sum_m \frac{\langle {}^N\Psi_0 | a_i^\dagger a_m | {}^N\Psi_0 \rangle \langle a_m {}^N\Psi_0 | a_j | {}^N\Psi_0 \rangle}{E - (\langle {}^N\Psi_0 | \mathcal{H} | {}^N\Psi_0 \rangle - \langle a_m {}^N\Psi_0 | \mathcal{H} | a_m {}^N\Psi_0 \rangle)} + \sum_p \frac{\langle {}^N\Psi_0 | a_j a_p^\dagger | {}^N\Psi_0 \rangle \langle a_p^\dagger {}^N\Psi_0 | a_i^\dagger | {}^N\Psi_0 \rangle}{E + (\langle {}^N\Psi_0 | \mathcal{H} | {}^N\Psi_0 \rangle - \langle a_p^\dagger {}^N\Psi_0 | \mathcal{H} | a_p^\dagger {}^N\Psi_0 \rangle)} \\ &= \sum_m \frac{\delta_{im} \delta_{mj}}{E - \varepsilon_m} + 0 \\ &= \sum_m \frac{\delta_{ij}}{E - \varepsilon_m}\end{aligned}\quad (7.2.4)$$

7.2.2 The Solution of the Dyson Equation

7.3 Application of the Formalism to H_2 and HeH^+

Ex 7.9

a.