Notes of Modern Quantum Chemistry, Szabo & Ostlund

hebrewsnabla

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0

spatial mol orb – ψ – i,j,k,...spatial basis fxn – ϕ – $\mu,\nu,\lambda,...$ spin orb – χ occ mol orb – a,b,c,...vir mol orb – r,s,t,...exact many-elec wfn – Φ approx many-elec wfn – Ψ exact energy – $\mathscr E$ approx energy – E

1

- 1.1
- 1.2
- 1.3

1.4 N-D Complex Vector Spaces

Suppose

$$\mathcal{O}|a\rangle = |b\rangle \tag{1.1}$$

$$\langle i \mid \mathcal{O} \mid j \rangle = O_{ij} \tag{1.2}$$

def the **adjoint** of \mathcal{O} as \mathcal{O}^{\dagger}

$$\langle a | \mathcal{O}^{\dagger} = \langle b | \tag{1.3}$$

$$\langle i \mid \mathcal{O}^{\dagger} \mid j \rangle = O_{ii}^{*} \tag{1.4}$$

1.4.1 Change of Basis

$$|\alpha\rangle = \sum_{i} |i\rangle \langle i|\alpha\rangle = \sum_{i} |i\rangle U_{i\alpha}$$
 (1.5)

$$|i\rangle = \sum_{\alpha} |\alpha\rangle \langle i|\alpha\rangle = \sum_{\alpha} |\alpha\rangle U_{i\alpha}^*$$
 (1.6)

If i, α are all orthonormal, **U** must be unitary.

$$\Omega_{\alpha\beta} = \langle \alpha \, | \, \mathcal{O} \, | \, \beta \rangle = \dots \sum_{ij} U_{\alpha i}^* O_{ij} U_{j\beta}$$
(1.7)

or

$$\mathbf{\Omega} = \mathbf{U}^{\dagger} \mathbf{O} \mathbf{U} \tag{1.8}$$

2.1 The Electronic Problem

- 2.1.1 Atomic Units
- 2.1.2 The B-O Approximation
- 2.1.3 The Antisymmetry or Pauli Exclusion Principle
- 2.2 Orbitals, Slater Determinants, and Basis Functions
- 2.2.1 Spin Orbitals and Spatial Orbitals
- 2.2.2 Hartree Products
- 2.2.3 Slater Determinants

def

$$|\chi_{i}(\mathbf{x}_{1})\chi_{j}(\mathbf{x}_{2})\cdots\chi_{k}(\mathbf{x}_{N})\rangle \equiv \frac{1}{\sqrt{N!}}\begin{vmatrix} \chi_{i}(\mathbf{x}_{1}) & \chi_{j}(\mathbf{x}_{1}) & \cdots & \chi_{k}(\mathbf{x}_{1}) \\ \chi_{i}(\mathbf{x}_{2}) & \chi_{j}(\mathbf{x}_{2}) & \cdots & \chi_{k}(\mathbf{x}_{2}) \\ \vdots & \vdots & & \vdots \\ \chi_{i}(\mathbf{x}_{N}) & \chi_{j}(\mathbf{x}_{N}) & \cdots & \chi_{k}(\mathbf{x}_{N}) \end{vmatrix}$$
(2.1)

It can be further shortened to

$$|\chi_i \chi_j \cdots \chi_k\rangle$$
 (2.2)

2.2.4 The Hartree-Fock Approximation

2.2.5 The Minimal Basis H_2 Model

gerade, ungerade

2.2.6 Excited Determinants

Suppose the ground state det

$$|\psi_0\rangle = |\chi_1 \cdots \chi_a \cdots \chi_b \cdots \chi_N\rangle \tag{2.3}$$

thus, singly excited det

$$|\psi_a^r\rangle = |\chi_1 \cdots \chi_r \cdots \chi_b \cdots \chi_N\rangle \tag{2.4}$$

$$|\psi_{ab}^{rs}\rangle = |\chi_1 \cdots \chi_r \cdots \chi_s \cdots \chi_N\rangle \tag{2.5}$$

How does program determine what dets can exist? by gerade/ungerade?

2.3 Operators and Matrix Elements

2.3.1 Minimal Basis H₂ Matrix Elements

2.3.2 Notations for 1- and 2-Electron Integrals

For spin orb, physicists'

$$\langle ij | kl \rangle = \left\langle \chi_i(1)\chi_j(2) \left| \frac{1}{r_{12}} \right| \chi_k(1)\chi_l(2) \right\rangle$$
 (2.6)

$$\langle ij \parallel kl \rangle = \langle ij \mid kl \rangle - \langle ij \mid lk \rangle \tag{2.7}$$

chemists?

$$[ij|kl] = \left\langle \chi_i(1)\chi_j(1) \left| \frac{1}{r_{12}} \right| \chi_k(2)\chi_l(2) \right\rangle$$
 (2.8)

For spatial orb

$$(ij|kl) = \left\langle \psi_i(1)\psi_j(1) \left| \frac{1}{r_{12}} \right| \psi_k(2)\psi_l(2) \right\rangle$$
 (2.9)

2.3.3 General Rules for Matrix Elements

1. $|K\rangle = |\cdots mn \cdots \rangle$

$$\langle K | \mathcal{H} | K \rangle = \sum_{m}^{N} [m|h|m] + \frac{1}{2} \sum_{m}^{N} \sum_{n}^{N} ([mm|nn] - [mn|nm])$$
 (2.10)

or (Since [mm|mm] - [mm|mm] = 0)

$$\langle K | \mathcal{H} | K \rangle = \sum_{m}^{N} [m|h|m] + \sum_{m}^{N} \sum_{n>m}^{N} ([mm|nn] - [mn|nm])$$
 (2.11)

2. $|K\rangle = |\cdots mn \cdots \rangle, |L\rangle = |\cdots pn \cdots \rangle$

$$\langle K \mid \mathcal{H} \mid L \rangle = [m|h|p] + \sum_{n=1}^{N} ([mp|nn] - [mn|np])$$
 (2.12)

3.
$$|K\rangle = |\cdots mn \cdots \rangle, |L\rangle = |\cdots pq \cdots \rangle$$

$$\langle K \mid \mathcal{H} \mid L \rangle = [mp|nq] - [mq|np] \tag{2.13}$$

2.4 Second Quantization

2.4.1 Creation and Annihilation Operators and Their Anticommutation Relations

$$a_i^{\dagger} a_i^{\dagger} + a_i^{\dagger} a_i^{\dagger} = 0 \quad a_i a_j + a_j a_i^{\dagger} 0 \tag{2.14}$$

$$\{a_i, a_j^{\dagger}\} \equiv a_i a_j^{\dagger} + a_j^{\dagger} a_i = \delta_{ij} \tag{2.15}$$

2.5 Spin-Adapted Configurations

2.5.1 Spin Operators

total spin

$$\hat{\mathscr{P}} = \sum_{i}^{N} \hat{\mathbf{s}}(i) \tag{2.16}$$

$$\hat{\mathscr{S}}_x = \sum_{i}^{N} \hat{\mathbf{s}}_x(i) \tag{2.17}$$

$$\hat{\mathscr{S}}_{+} = \sum_{i}^{N} \hat{\mathbf{s}}_{+}(i) \tag{2.18}$$

$$\hat{\mathcal{S}}^2 = \hat{\mathcal{S}}_+ \hat{\mathcal{S}}_- - \hat{\mathcal{S}}_z + \hat{\mathcal{S}}_z^2 \tag{2.19}$$

$$\hat{\mathscr{S}}^2 |\Phi\rangle = S(S+1) |\Phi\rangle \tag{2.20}$$

$$\hat{\mathscr{S}}_z |\Phi\rangle = M_S |\Phi\rangle \tag{2.21}$$

$$\hat{\mathscr{S}}_z |ij\cdots k\rangle = \frac{1}{2} (N^\alpha - N^\beta) |ij\cdots k\rangle$$
 (2.22)

2.5.2 Restricted Determinants and Spin-Adapted Configurations

$$|^{1}\Psi_{1}^{2}\rangle = \frac{1}{\sqrt{2}}(|1\bar{2}\rangle + |\bar{1}2\rangle) = \frac{1}{\sqrt{2}}(\psi_{1}(1)\psi_{2}(2)\alpha(1)\beta(2) - \psi_{2}(1)...)$$
 (2.23)

3 The Hartree-Fock Approximation

3.1 The HF Equations

- 3.1.1 The Coulomb and Exchange Operators
- 3.1.2 The Fock Operator

3.2 Derivation of the HF Equations

- 3.2.1 Functional Variation
- 3.2.2 Minimization of the Energy of a Single Determinant
- 3.2.3 The Canonical HF Equations

3.3 Interpretation of Solutions to the HF Equations

3.3.1 Orbital Energies and Koopmans' Theorem

$$\varepsilon_{i} = \langle i \mid h \mid i \rangle + \sum_{b} \langle ib \parallel ib \rangle \tag{3.1}$$

$$= \langle i \mid h \mid i \rangle + \sum_{b} (\langle ib \mid ib \rangle - \langle ib \mid bi \rangle)$$
 (3.2)

Koopmans' Theorem

$$IP = -\varepsilon_a \quad EA = -\varepsilon_r \tag{3.3}$$

Koopmans' EA is often bad.

3.3.2 Brillouin's Theorem

$$\langle \Psi_0 \, | \, \mathcal{H} \, | \, \Psi_a^r \rangle = 0 \tag{3.4}$$

3.3.3 The HF Hamiltonian

$$\mathcal{H}_0 = \sum_{i}^{N} f(i) \tag{3.5}$$

3.4 Restricted Closed-shell HF: The Roothaan Equations

3.4.1 Closed-shell HF: Restricted Spin Orbitals

$$E_0 = 2\sum_a h_{aa} + \sum_a \sum_b (2J_{ab} - K_{ab})$$
(3.6)

3.4.2 Introduction of a Basis: The Roothaan Equations

3.4.3 The Charge Density

3.4.4 Expression for the Fock Matrix

3.4.5 Orthogonalization of the Basis

$$\mathbf{X}^{\dagger}\mathbf{S}\mathbf{X} = \mathbf{1} \tag{3.7}$$

 ${f S}$ can be diagonalized by unitary matrix ${f U}$:

$$\mathbf{U}^{\dagger}\mathbf{S}\mathbf{U} = \mathbf{s} \tag{3.8}$$

Symmetric Orthogonalization

$$\mathbf{X} = \mathbf{S}^{-1/2} = \mathbf{U}\mathbf{s}^{-1/2}\mathbf{U}^{\dagger} \tag{3.9}$$

(linear dependence must be removed)

Canonical Orthogonalization

$$\mathbf{X} = \mathbf{U}\mathbf{S}^{-1/2} = \mathbf{U}\mathbf{s}^{-1/2}\mathbf{U}^{\dagger} \tag{3.10}$$

Suppose s has m small values, we make a truncated $K \times (K-m)$ matrix

$$\tilde{\mathbf{X}} = \dots \tag{3.11}$$

thus

$$\phi'_{\mu} = \sum_{\nu}^{K} \phi_{\nu} \tilde{X}_{\nu\mu} \qquad \mu = 1, 2, \cdots, K - m$$
 (3.12)

However, calculate 2e integrals in transformed matrix is very time-consuming. Since

$$\phi' = \phi \mathbf{X} \tag{3.13}$$

$$\psi = \phi \mathbf{C} \tag{3.14}$$

we have

$$\psi = \phi' \mathbf{X}^{-1} \mathbf{C} \tag{3.15}$$

Let

$$\mathbf{C}' = \mathbf{X}^{-1}\mathbf{C} \qquad or \ \mathbf{C} = \mathbf{X}\mathbf{C}' \tag{3.16}$$

thus

$$\mathbf{FXC'} = \mathbf{SXC'}\varepsilon \tag{3.17}$$

$$(\mathbf{X}^{\dagger}\mathbf{F}\mathbf{X})\mathbf{C}' = (\mathbf{X}^{\dagger}\mathbf{S}\mathbf{X})\mathbf{C}'\boldsymbol{\varepsilon} = \mathbf{C}'\boldsymbol{\varepsilon}$$
(3.18)

def:

$$\mathbf{F}' = \mathbf{X}^{\dagger} \mathbf{F} \mathbf{X} \tag{3.19}$$

3.4.6 The SCF Procedure

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I think \mathbf{C}' should be $(K-m) \times K$.

3.4.7 Expectation Values and Population Analysis