# Notes of Modern Quantum Chemistry, Szabo & Ostlund

## hebrewsnabla

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#### 0

spatial mol orb –  $\psi$  – i,j,k,...spatial basis fxn –  $\phi$  –  $\mu,\nu,\lambda,...$ spin orb –  $\chi$ occ mol orb – a,b,c,...vir mol orb – r,s,t,...exact many-elec wfn –  $\Phi$ approx many-elec wfn –  $\Psi$ exact energy –  $\mathscr E$ approx energy – E

1

- 1.1
- 1.2
- 1.3

### 1.4 N-D Complex Vector Spaces

Suppose

$$\mathcal{O}|a\rangle = |b\rangle \tag{1.1}$$

$$\langle i \mid \mathcal{O} \mid j \rangle = O_{ij} \tag{1.2}$$

def the **adjoint** of  $\mathcal{O}$  as  $\mathcal{O}^{\dagger}$ 

$$\langle a|\,\mathcal{O}^{\dagger} = \langle b| \tag{1.3}$$

$$\langle i \mid \mathcal{O}^{\dagger} \mid j \rangle = O_{ji}^*$$
 (1.4)

#### 1.4.1 Change of Basis

$$|\alpha\rangle = \sum_{i} |i\rangle \langle i|\alpha\rangle = \sum_{i} |i\rangle U_{i\alpha}$$
 (1.5)

$$|i\rangle = \sum_{\alpha} |\alpha\rangle \langle i|\alpha\rangle = \sum_{\alpha} |\alpha\rangle U_{i\alpha}^*$$
 (1.6)

If  $i, \alpha$  are all orthonormal, **U** must be unitary.

$$\Omega_{\alpha\beta} = \langle \alpha \, | \, \mathcal{O} \, | \, \beta \rangle = \dots \sum_{ij} U_{\alpha i}^* O_{ij} U_{j\beta} \tag{1.7}$$

or

$$\mathbf{\Omega} = \mathbf{U}^{\dagger} \mathbf{O} \mathbf{U} \tag{1.8}$$

#### 2.1 The Electronic Problem

- 2.1.1 Atomic Units
- 2.1.2 The B-O Approximation
- 2.1.3 The Antisymmetry or Pauli Exclusion Principle
- 2.2 Orbitals, Slater Determinants, and Basis Functions
- 2.2.1 Spin Orbitals and Spatial Orbitals
- 2.2.2 Hartree Products
- 2.2.3 Slater Determinants

def

$$|\chi_{i}(\mathbf{x}_{1})\chi_{j}(\mathbf{x}_{2})\cdots\chi_{k}(\mathbf{x}_{N})\rangle \equiv \frac{1}{\sqrt{N!}}\begin{vmatrix} \chi_{i}(\mathbf{x}_{1}) & \chi_{j}(\mathbf{x}_{1}) & \cdots & \chi_{k}(\mathbf{x}_{1}) \\ \chi_{i}(\mathbf{x}_{2}) & \chi_{j}(\mathbf{x}_{2}) & \cdots & \chi_{k}(\mathbf{x}_{2}) \\ \vdots & \vdots & & \vdots \\ \chi_{i}(\mathbf{x}_{N}) & \chi_{j}(\mathbf{x}_{N}) & \cdots & \chi_{k}(\mathbf{x}_{N}) \end{vmatrix}$$
(2.1)

It can be further shortened to

$$|\chi_i\chi_j\cdots\chi_k\rangle$$
 (2.2)

- 2.2.4 The Hartree-Fock Approximation
- 2.2.5 The Minimal Basis  $\mathrm{H}_2$  Model

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#### 2.2.6 Excited Determinants

Suppose the ground state det

$$||psi_0\rangle = |\chi_1 \cdots \chi_a \cdots \chi_b \cdots \chi_N\rangle \tag{2.3}$$

thus, singly excited det

$$||psi_a^r\rangle = |\chi_1 \cdots \chi_r \cdots \chi_b \cdots \chi_N\rangle$$
 (2.4)

$$||psi_{ab}^{rs}\rangle = |\chi_1 \cdots \chi_r \cdots \chi_s \cdots \chi_N\rangle$$
 (2.5)

#### 2.3 Operators and Matrix Elements

- $2.3.1 \quad \text{Minimal Basis $H_2$ Matrix Elements} \\$
- 2.3.2 Notations for 1- and 2-Electron Integrals

For spin orb

$$\langle ij \mid kl \rangle = \left\langle \chi_i(1)\chi_j(2) \mid \frac{1}{r_{12}} \mid \chi_k(1)\chi_l(2) \right\rangle$$
 (2.6)

$$\langle ij \parallel kl \rangle = \langle ij \mid kl \rangle - \langle ij \mid lk \rangle \tag{2.7}$$

$$[ij|kl] = \left\langle \chi_i(1)\chi_j(1) \left| \frac{1}{r_{12}} \right| \chi_k(2)\chi_l(2) \right\rangle$$
 (2.8)

For spatial orb

$$(ij|kl) = \left\langle \psi_i(1)\psi_j(1) \middle| \frac{1}{r_{12}} \middle| \psi_k(2)\psi_l(2) \right\rangle$$
 (2.9)