

SC Quantathon v1 (Team 1)

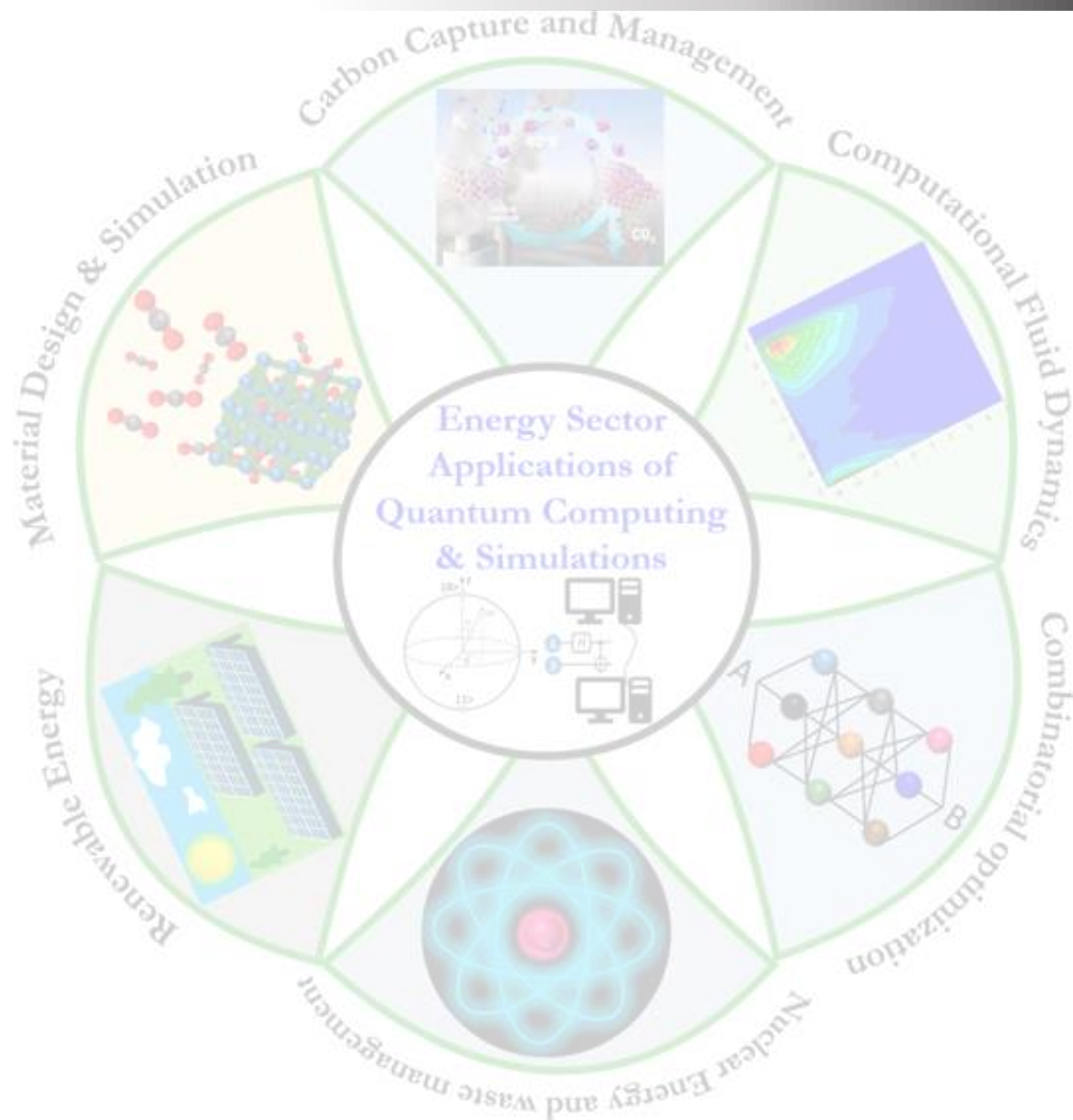
Xanadu Challenge: Spectral gap for molecules

Anjali A. Agrawal, Omar Alsheikh, Norman Hogan, Arvin Kushwaha and Heba Labib

North Carolina State University

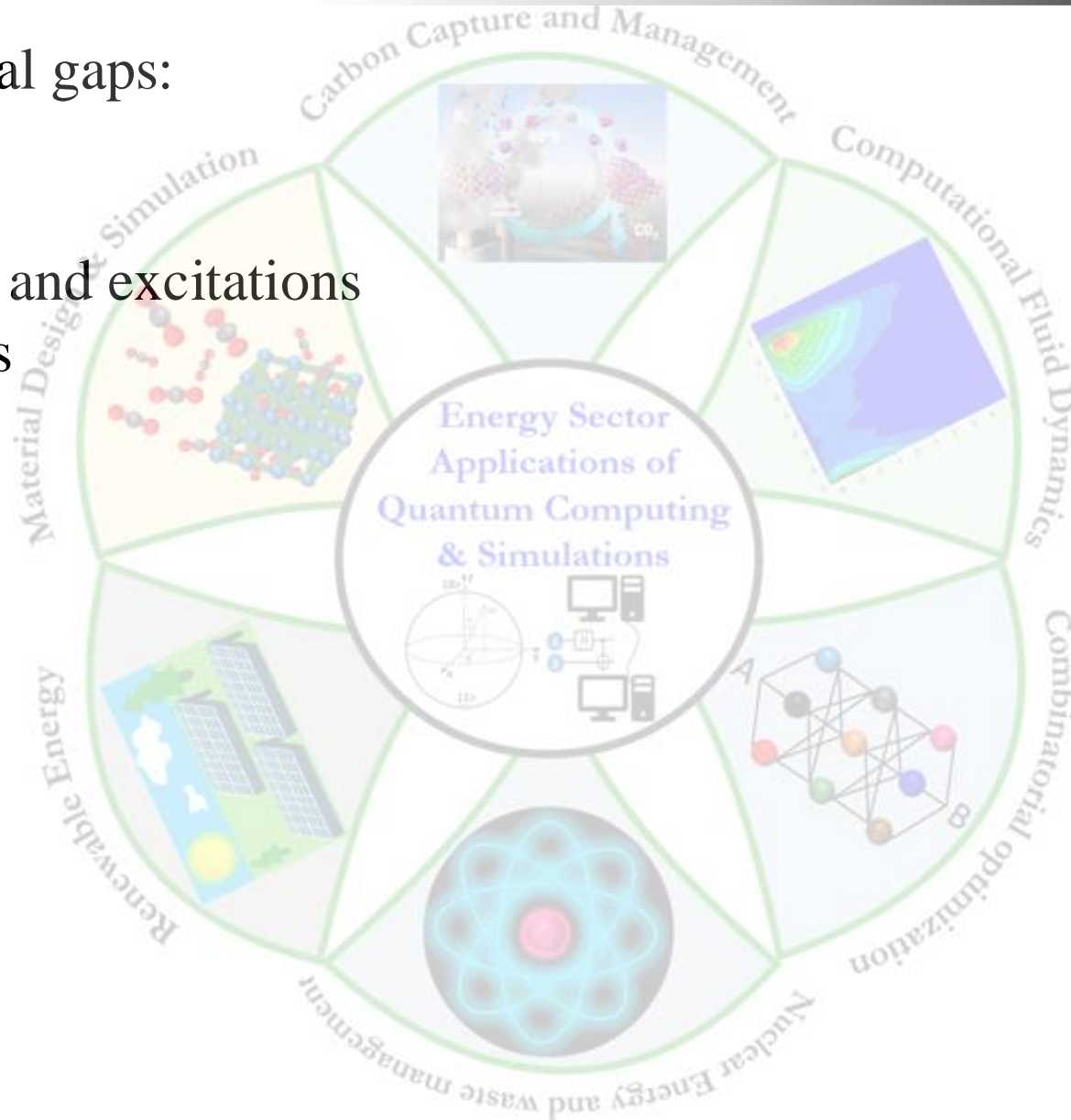


Motivation



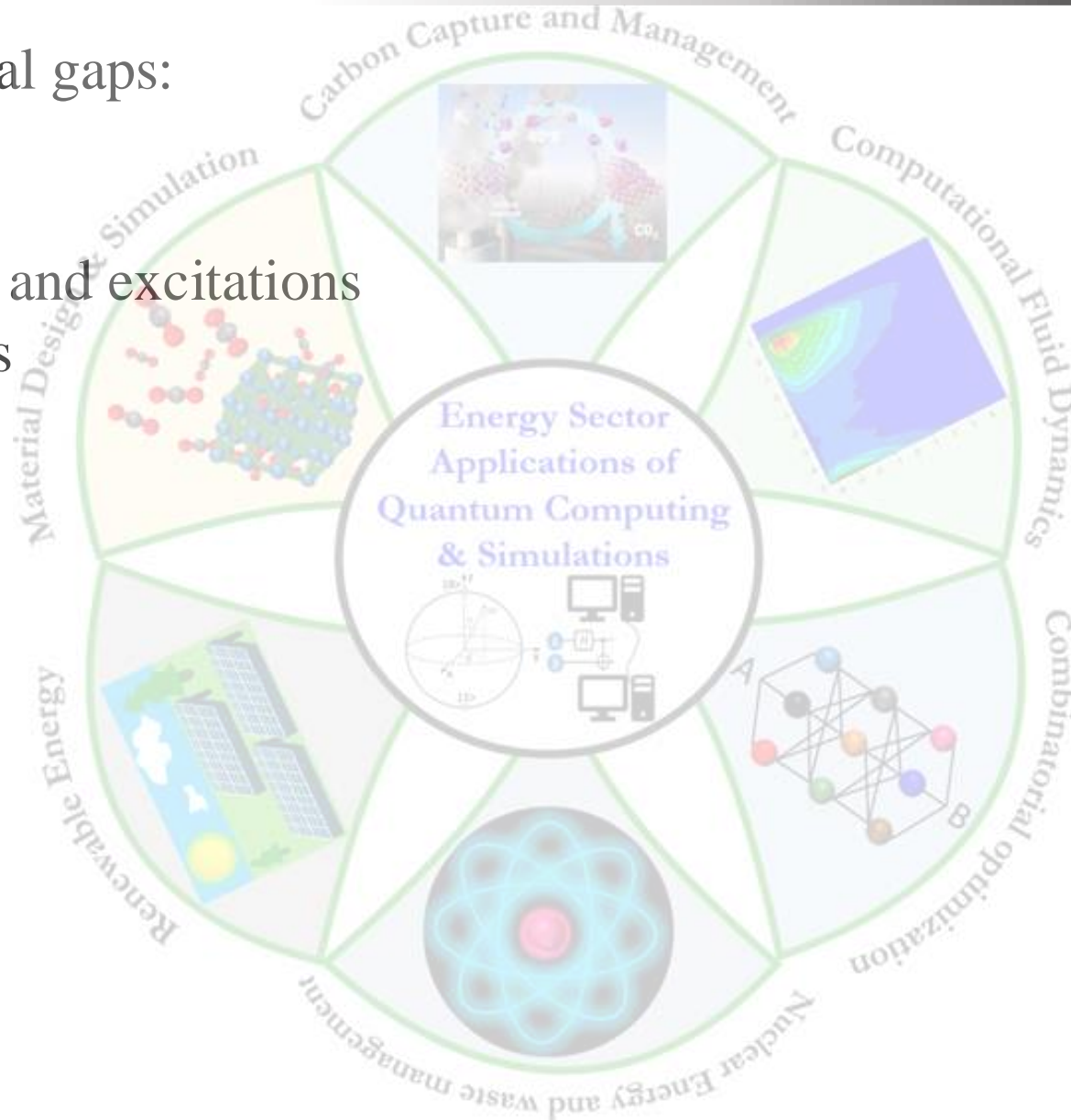
Importance of spectral gaps:

- Stability
- Reactivity
- Photo-interactions and excitations
- Material properties
- Algorithm designs



Importance of spectral gaps:

- Stability
- Reactivity
- Photo-interactions and excitations
- Material properties
- Algorithm designs

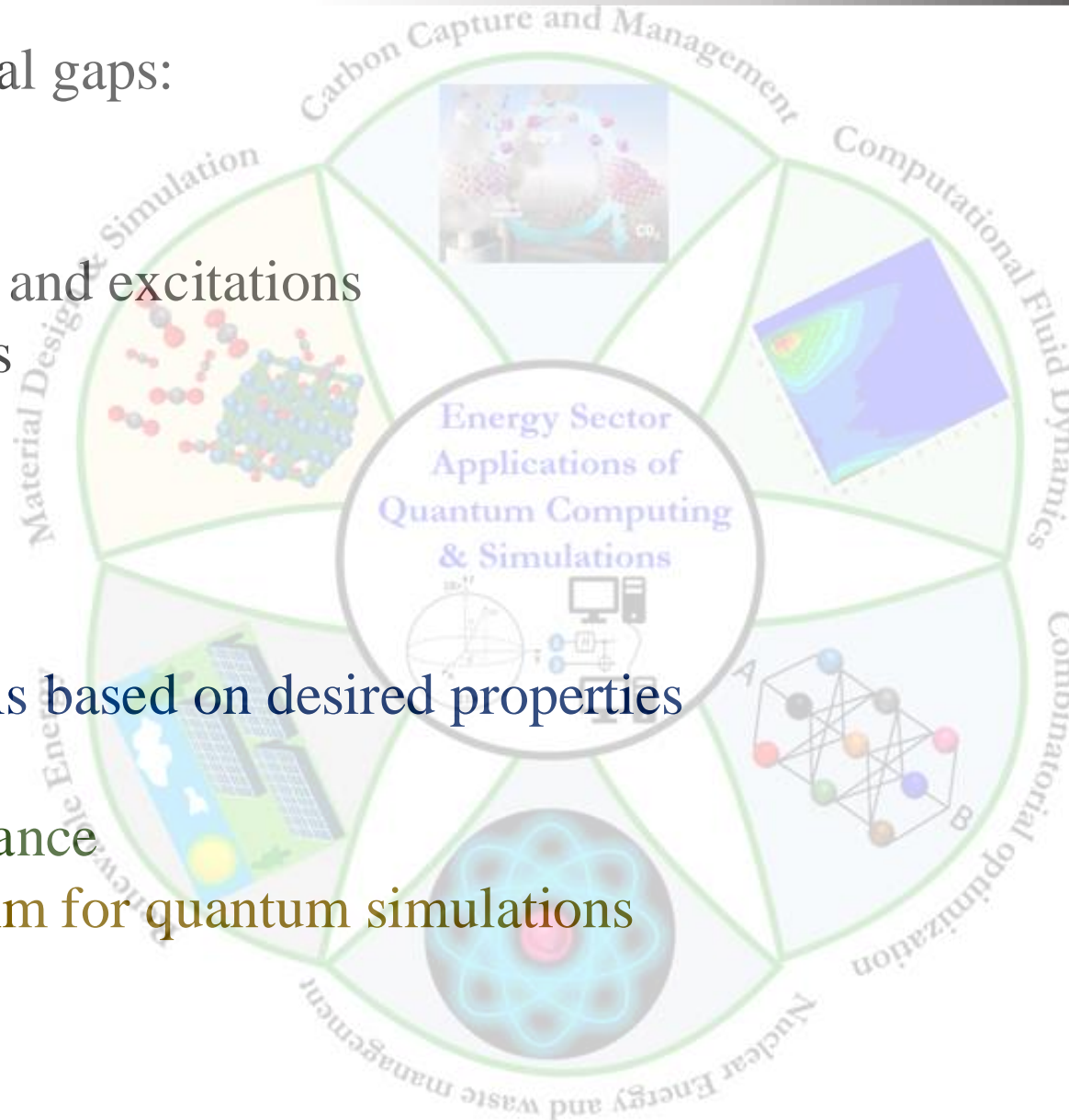


Importance of spectral gaps:

- Stability
- Reactivity
- Photo-interactions and excitations
- Material properties
- Algorithm designs

Applications:

1. Choosing materials based on desired properties
(Material search)
2. Experiment assistance
3. Choice of algorithm for quantum simulations
4. Material design



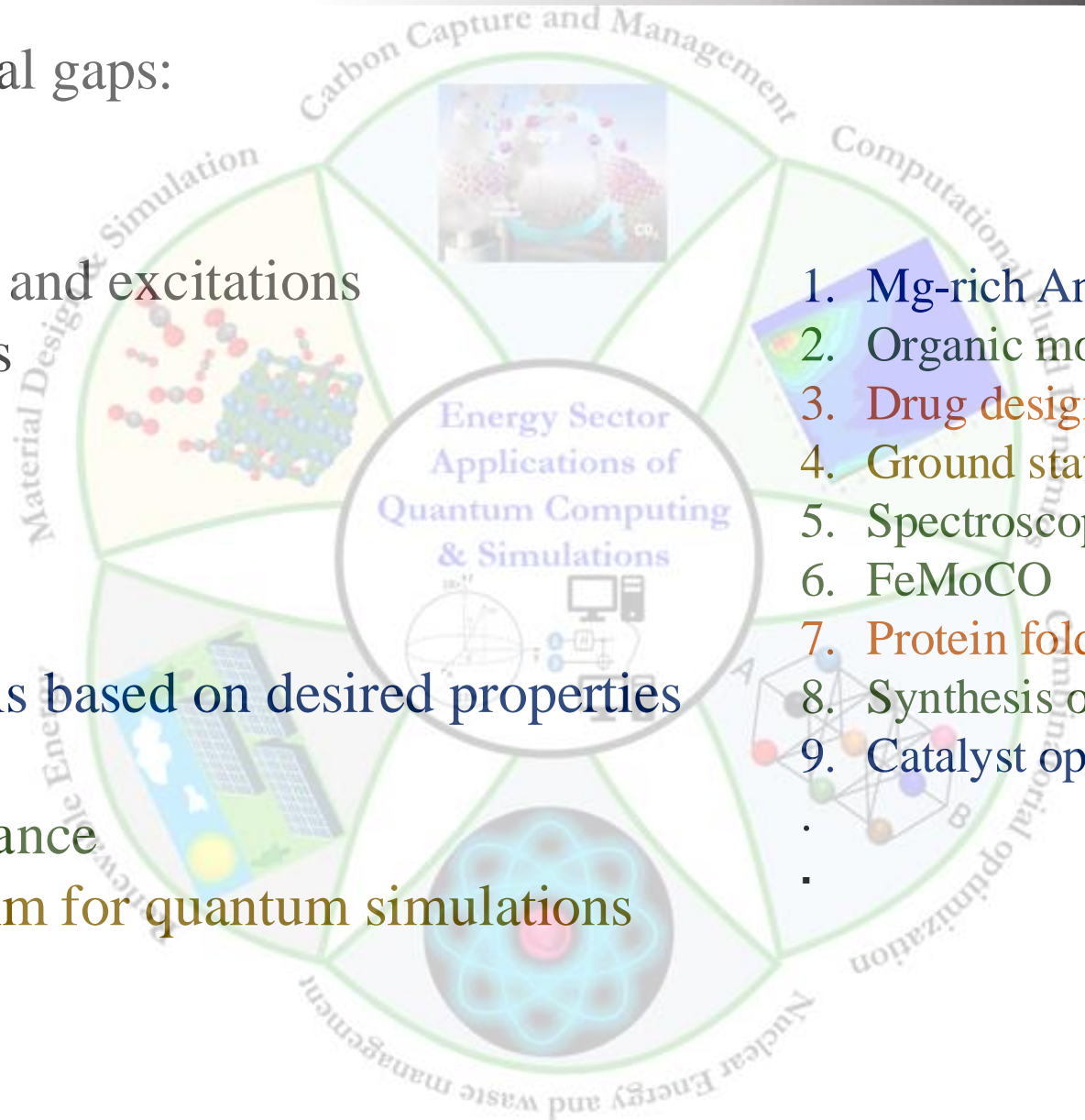
Importance of spectral gaps:

- Stability
- Reactivity
- Photo-interactions and excitations
- Material properties
- Algorithm designs

Applications:

1. Choosing materials based on desired properties (Material search)
2. Experiment assistance
3. Choice of algorithm for quantum simulations
4. Material design

1. Mg-rich Anticorrosive Coating Design
2. Organic molecules combustion energy
3. Drug design
4. Ground state energy estimation
5. Spectroscopy experiments
6. FeMoCO
7. Protein folding
8. Synthesis of cyclic ozone (Fuel tech)
9. Catalyst optimization and fabrication



Application choice: Metallification of H₂ molecule

- Variation of bond length impacts the properties of the molecule (Insulator-semiconductor-metal)
- Experiments are designed to get different phases by changing the bond length
- Bond length changes the spectral gaps and hence the reaction properties

Application choice: Metallification of H₂ molecule

- Variation of bond length impacts the properties of the molecule (Insulator-semiconductor-metal)
- Experiments are designed to get different phases by changing the bond length
- Bond length changes the spectral gaps and hence the reaction properties

Application choice: Metallification of H₂ molecule

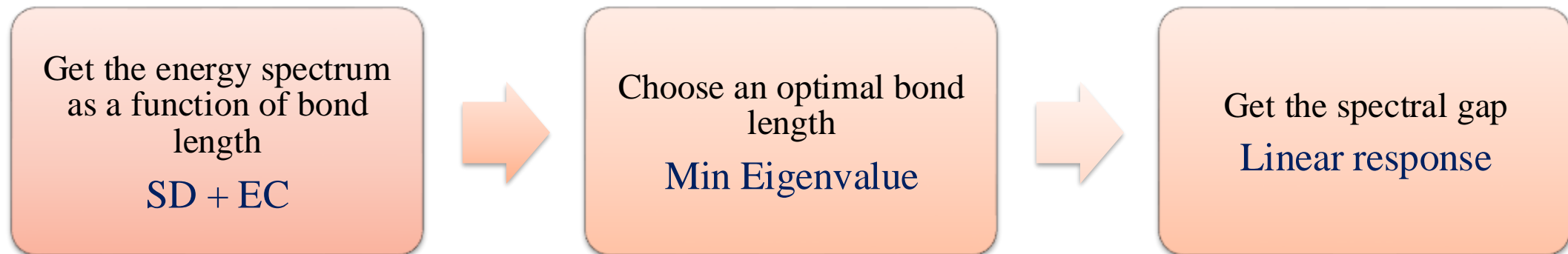
- Variation of bond length impacts the properties of the molecule (Insulator-semiconductor-metal)
- Experiments are designed to get different phases by changing the bond length
- Bond length changes the spectral gaps and hence the reaction properties

Different methods come with various challenges; we aimed to apply the strategies that suited the chosen Hamiltonian system.

Application choice: Metallification of H₂ molecule

- Variation of bond length impacts the properties of the molecule (Insulator-semiconductor-metal)
- Experiments are designed to get different phases by changing the bond length
- Bond length changes the spectral gaps and hence the reaction properties

Different methods come with various challenges; we aimed to apply the strategies that suited the chosen Hamiltonian system.



Eigenvector continuation as a subspace method

$$\mathcal{H}_{target}(\theta) = \{H(\theta_1), H(\theta_2), \dots, H(\theta_n)\}$$

Choose k_p points \downarrow in θ space

θ - Bond length

$$\{H(\theta_1), H(\theta_2), \dots, H(\theta_{k_p})\}$$

Use any (truncated) state preparation method to solve for (approximate) ground state

$$\{|\phi_1\rangle, |\phi_2\rangle, \dots, |\phi_{k_p}\rangle\}$$

Low-energy subspace

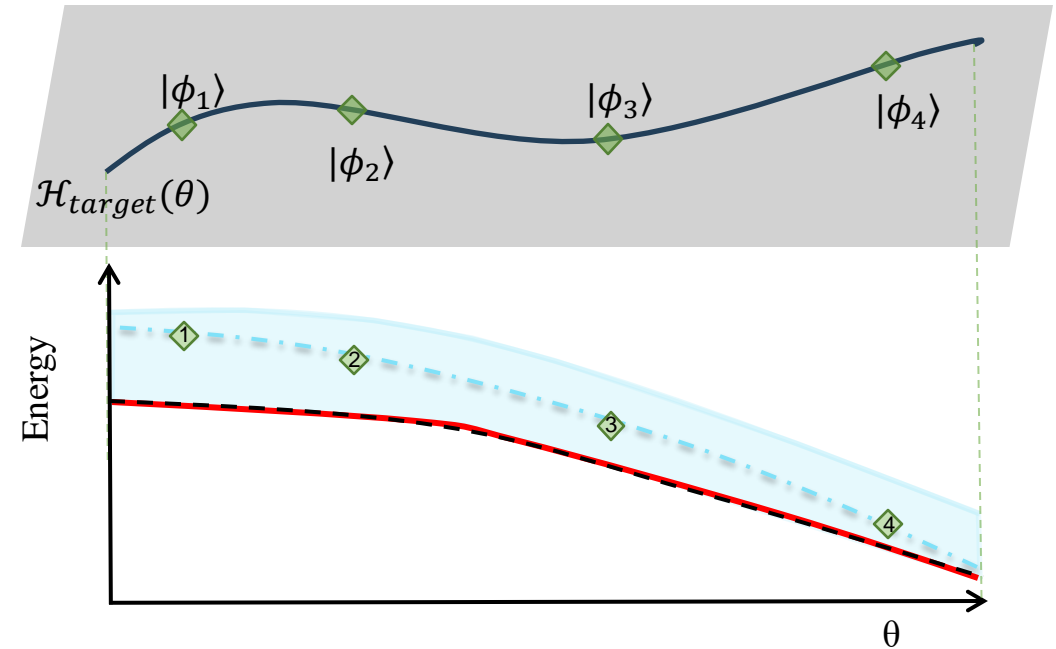
Subspace Diagonalization using Eigenvector continuation

Subspace = $\{|\phi_1\rangle, |\phi_2\rangle, |\phi_3\rangle, |\phi_4\rangle\}$; $k_p = 4$ (in fig.)

$$\mathbb{H}(\theta)_{ij} = \langle \phi_i | H(\theta) | \phi_j \rangle$$

$$S_{ij} = \langle \phi_i | \phi_j \rangle$$

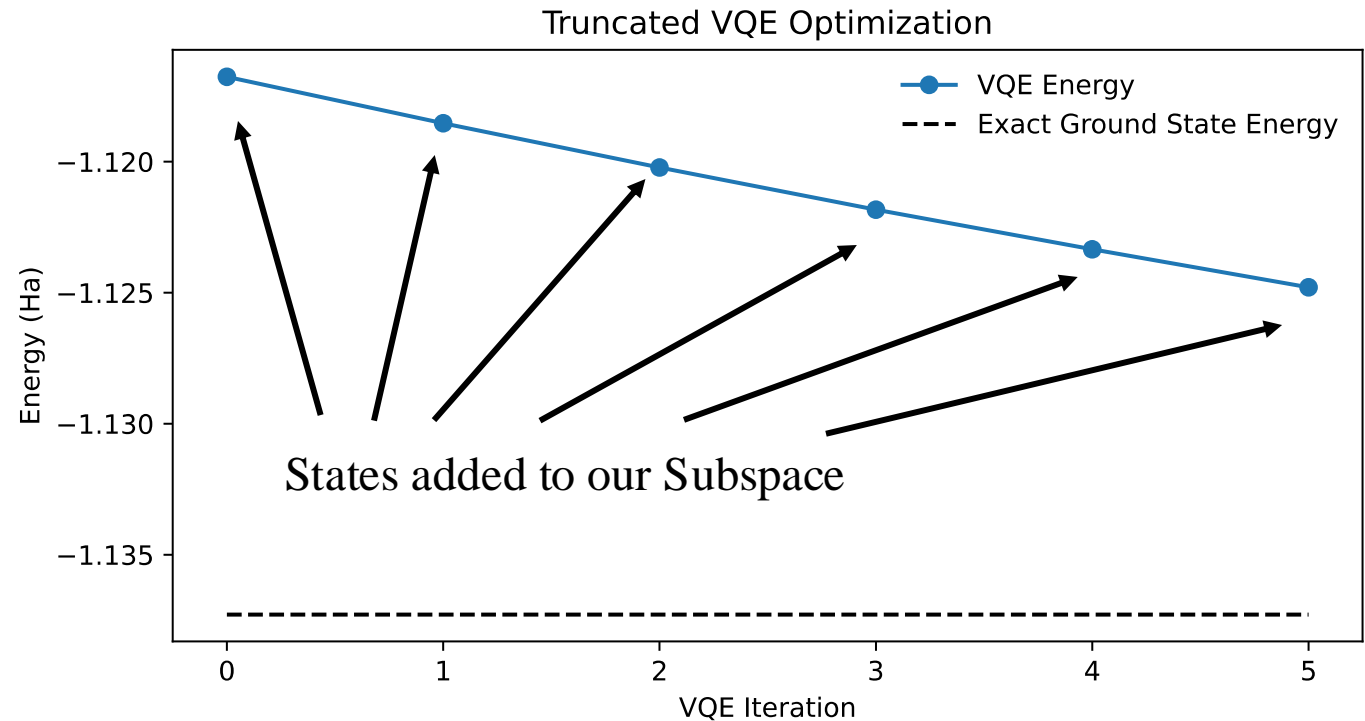
$$\mathbb{H}(\theta)|\psi\rangle = E S|\psi\rangle$$



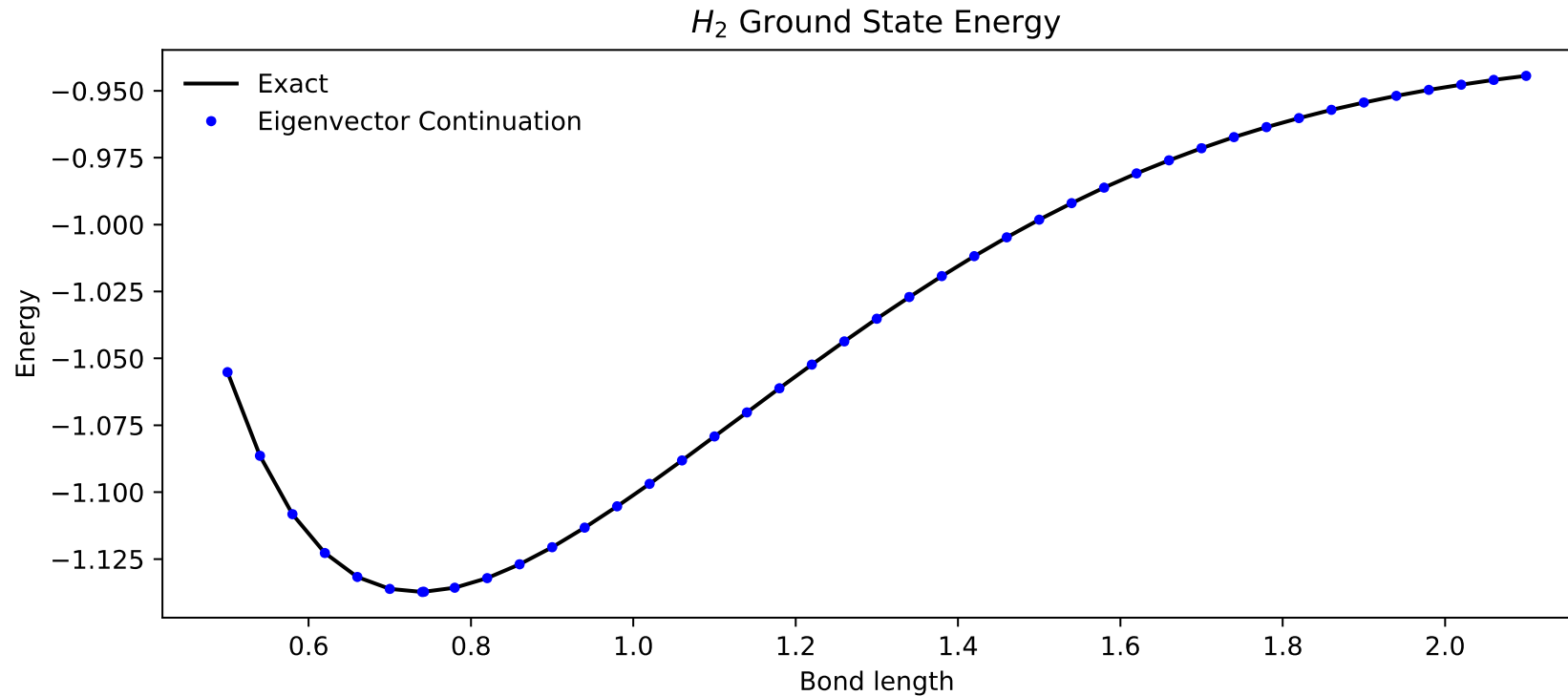
- / ◆ Truncated state preparation
- - - Exact ground state energy
- Subspace diagonalization using EC

Truncated VQE as our Basis

- Used UCCSD ansatz (60 CNOTS)
- Need only a few iterations of VQE for a subspace
- Employed a filtering procedure to ensure the subspace diagonalization is numerically stable
 1. Find the lowest energy state in subspace
 2. Add states with the least overlap to subspace – ensure the condition number is not too high
 3. Convergence is met when adding states doesn't change energy

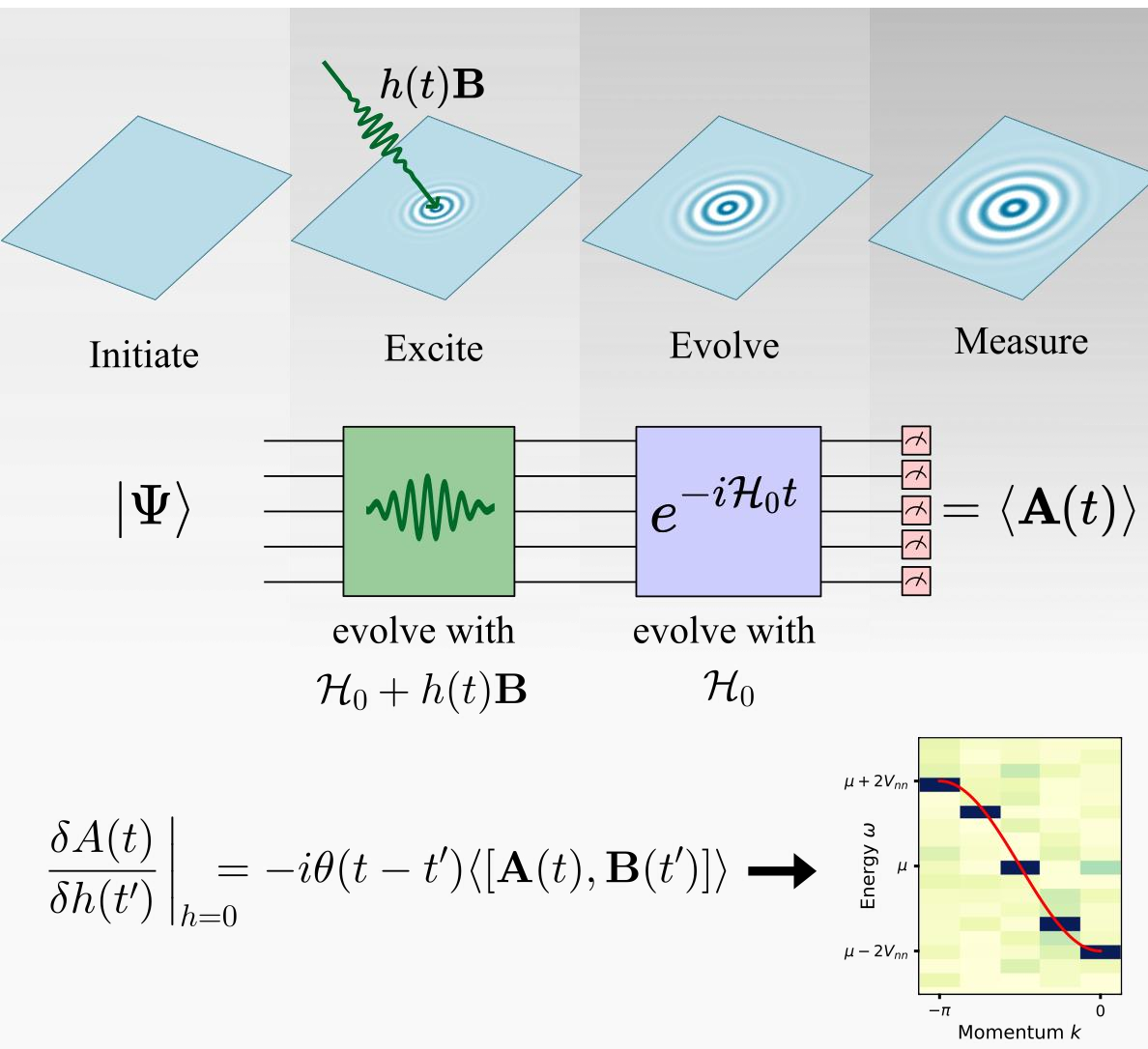


H₂ molecule spectrum using EC



- Only 2-4 states needed per data point
- VQE only done for 5 iterations with bond length of 0.74 Å, but we get the *entire* spectrum!

Linear Response and Quantum Simulations



Spectral Functions to measure energy differences ϵ_{nm} :

$$A_{jj}(T, \omega) = -\frac{1}{\pi} \text{Im} G_{jj}^R(\omega)$$

$$A_{jj}(T, \omega) = \bar{A}_{jj}(\omega) + R_{jj}(T, \omega),$$

$$\bar{A}_{jj}(\omega) = \sum_{n,m} |\psi_n|^2 \{ |C_{nm}^{(j)}|^2 \delta(\omega - \epsilon_{nm}) + n \leftrightarrow m \},$$

Starting with the ground state we get from VQE+EC: we confine ourselves with peaks ϵ_{n0} .

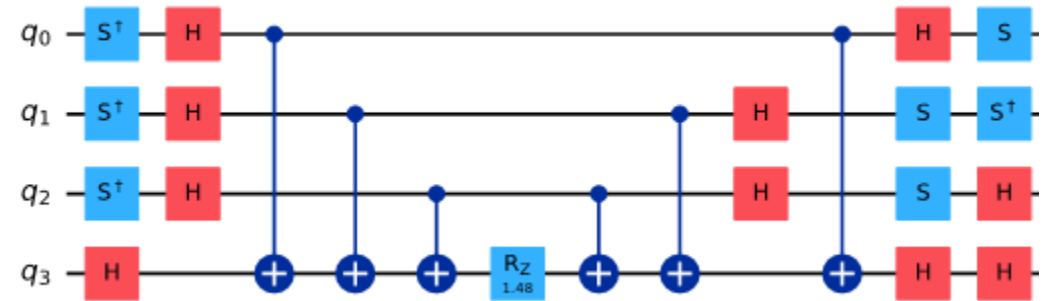
A way to get spectral gaps!

Time evolution using Cartan decomposition

We used a fixed depth technique to simulate the time evolution operator. $U(t) = e^{-it\mathcal{H}}$

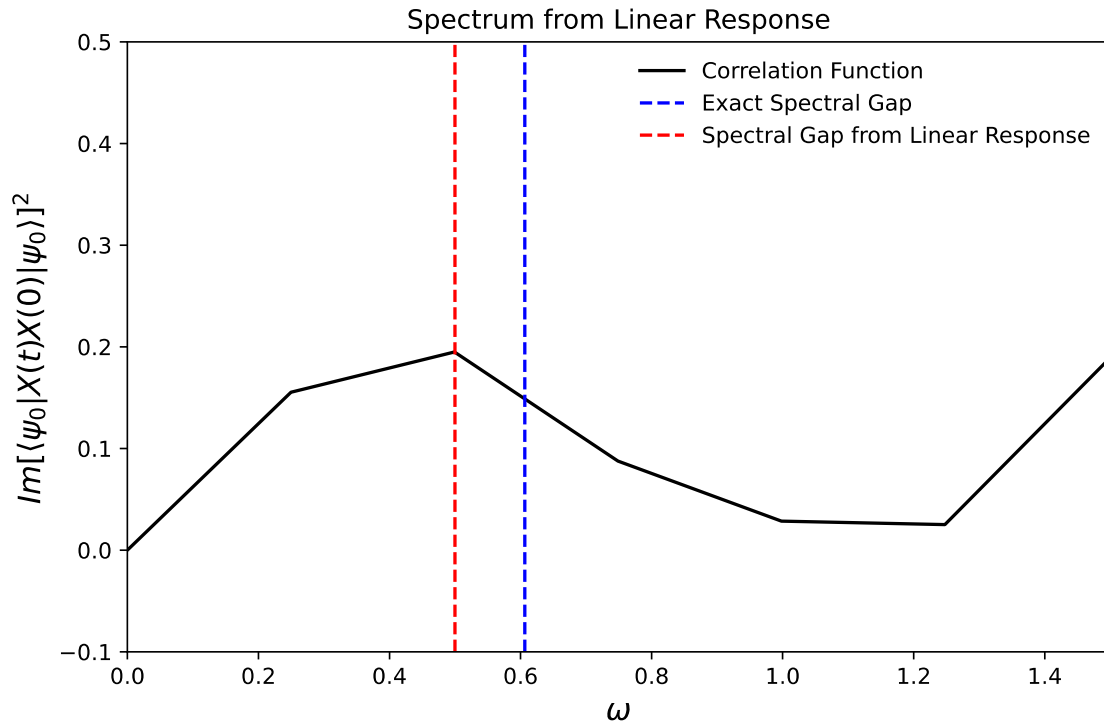
To compare:

- Direct circuit implementation (124 CNOTs)
- One first order Trotter step is (34 CNOTs)



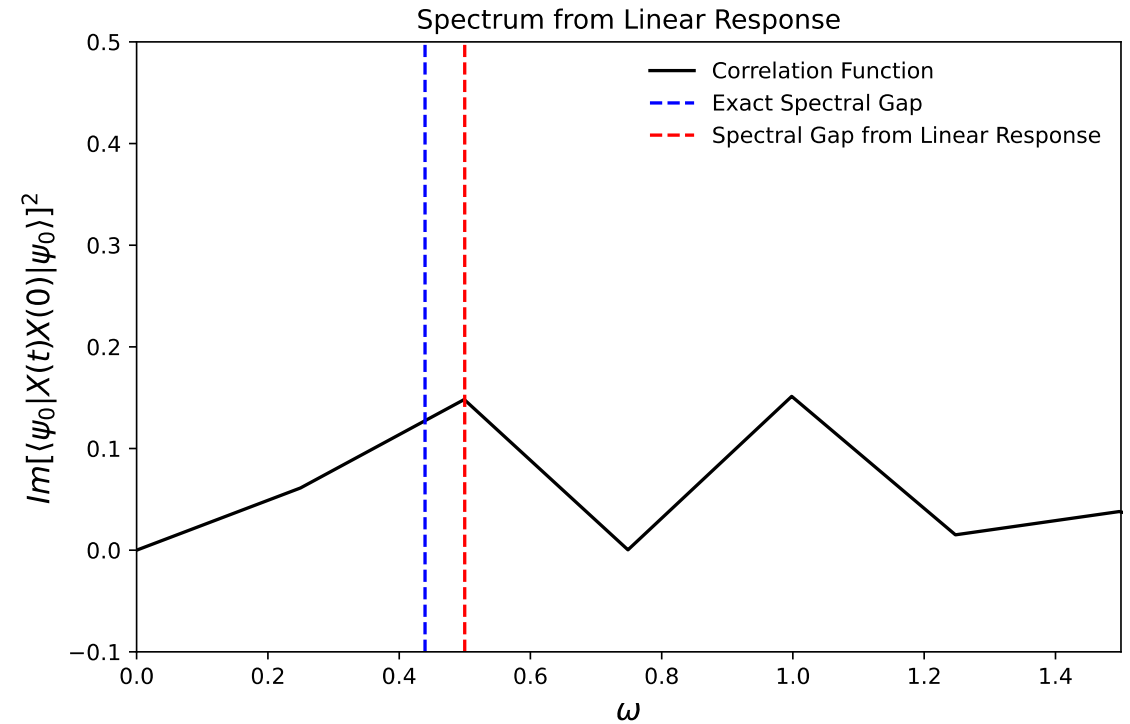
Spectral gaps using Linear Response

Bond length = 0.74 Å



Error: 17%

Bond length = 1.1 Å



Error: 13%

Right now:

VQE → Eigenvector Continuation → Cartan Fixed-Depth → Linear
Response

Next steps:

- Better VQE Ansatz
- Use Cartan for Imaginary Time Evolution to get excited states
- Post-processing signals
- Investigate other molecules

Eigenvector continuation:

1. Francis, Akhil, *et. al.*, 2209.10571
2. Agrawal, Anjali, *et. Al.*, 2406.17037
3. D. Frame *et. al.*, Phys. Rev. Lett. 121, 032501

Linear Response

1. Kökcü, E., *et. al.*, *Nat Commun* **15**, 3881 (2024)
2. Weidinger, Simon A., *et. Al.*, PhysRevB.98.224205

Cartan decomposition

1. Kökcü, E., *et. al.*, 2104.00728

Metallification of H₂ molecule

1. Tiwary, P., *et. al.*, 10.1073/PNAS.1600917113
2. Bing, Li, *et, al.*, Phys. Rev. Lett. 126.036402

Resources

1. PennyLane (& their Quantum Chemistry Datasets)
2. QBraid
3. DoraHacks



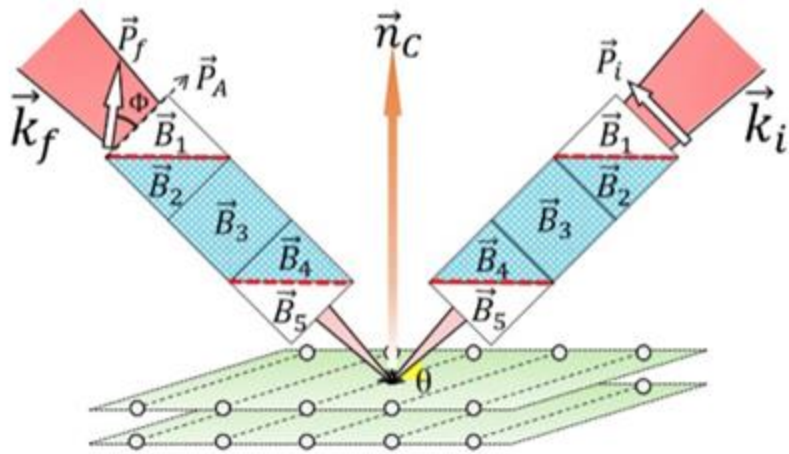
Thank you!



How do we know what electrons do inside matter?

How do we know what electrons do inside matter?

Neutron Scattering

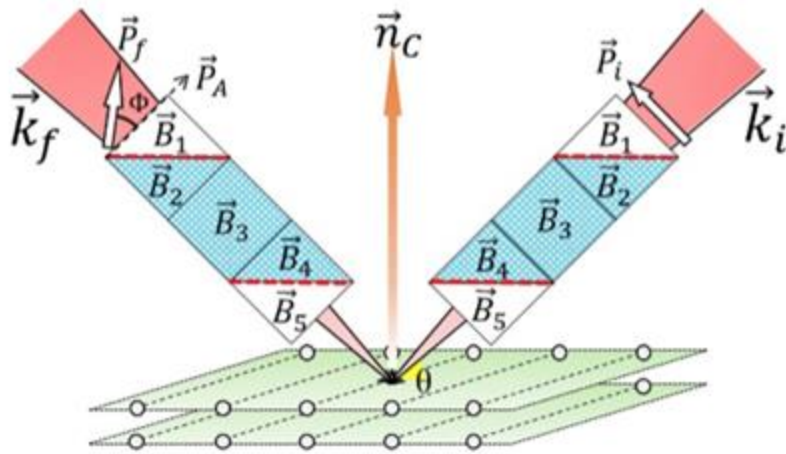


Spin-Spin Correlations

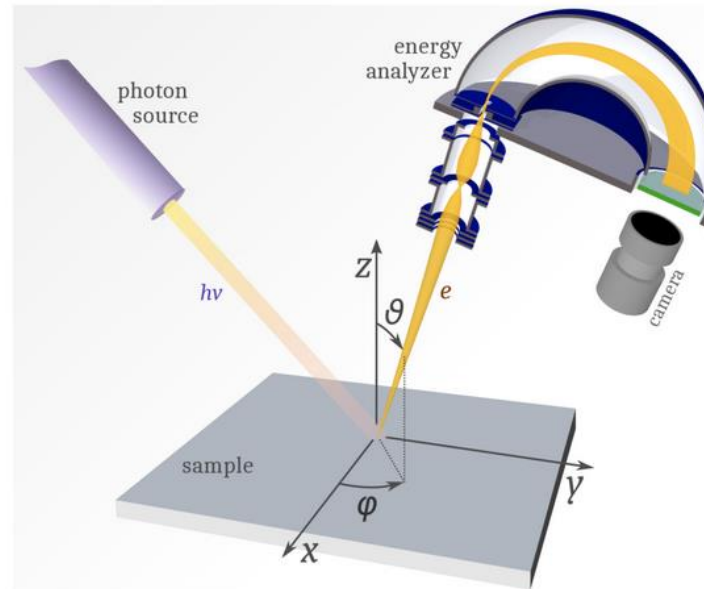
Figure Courtesy Shen Group

How do we know what electrons do inside matter?

Neutron Scattering



Angle-Resolved Photoemission Spectroscopy



Spin-Spin Correlations

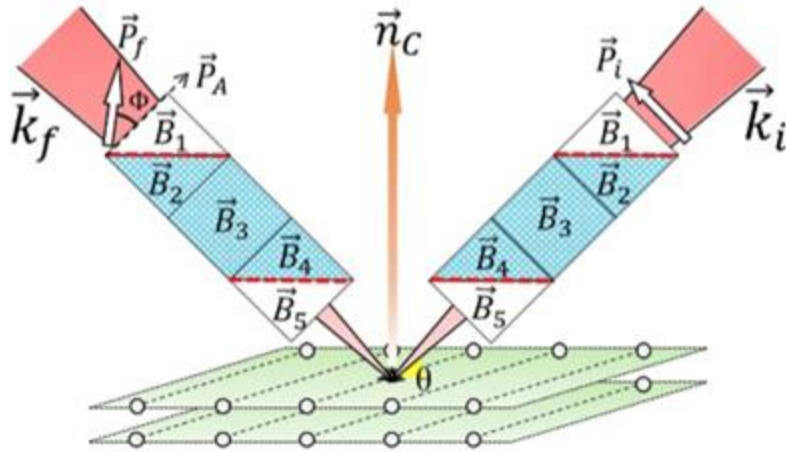
Figure Courtesy Shen Group

Single-Particle Green Functions

Public-domain figure from Wikimedia
<https://commons.wikimedia.org/wiki/File:ARPESgeneral.png>

How do we know what electrons do inside matter?

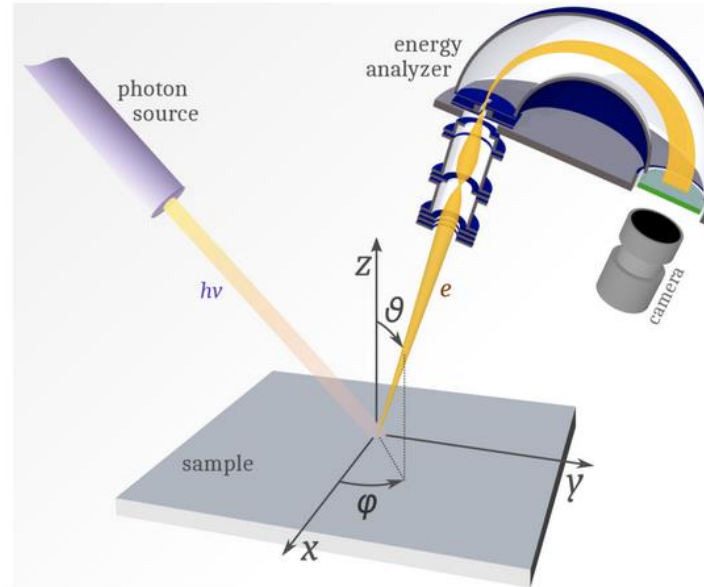
Neutron Scattering



Spin-Spin Correlations

Figure Courtesy Shen Group

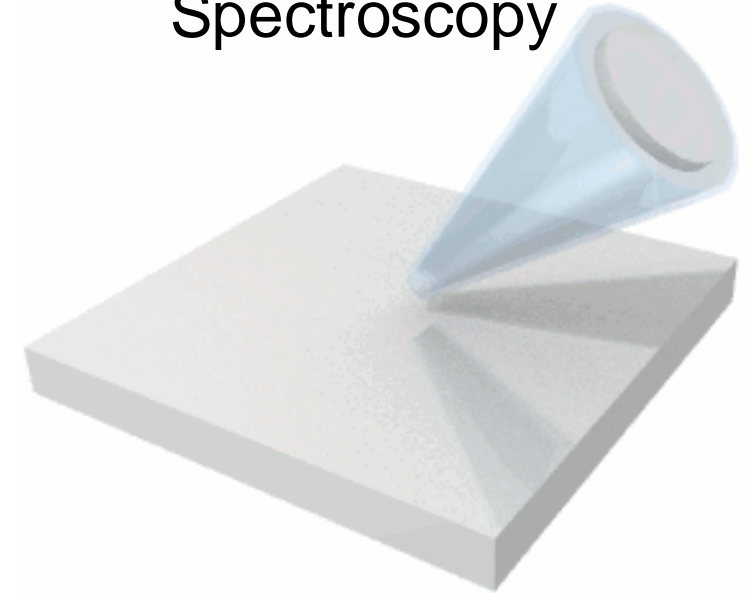
Angle-Resolved Photoemission Spectroscopy



Single-Particle Green Functions

Public-domain figure from Wikimedia
<https://commons.wikimedia.org/wiki/File:ARPESgeneral.png>

Time-Resolved Photoemission Spectroscopy



Determining Band Structure

Animation courtesy [F. Schmitt](#)

Consider: $H_0 = \sigma^z, A = B = \sigma^x$, A two-level system with $\Delta E = 2$:

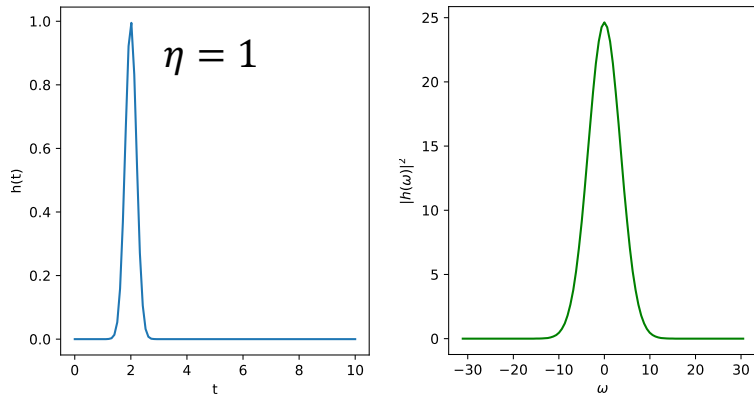
Consider: $H_0 = \sigma^z, A = B = \sigma^x$, A two-level system with $\Delta E = 2$:

$$A(\omega) = \chi(\omega)h(\omega)$$

Dictates Frequency Support!

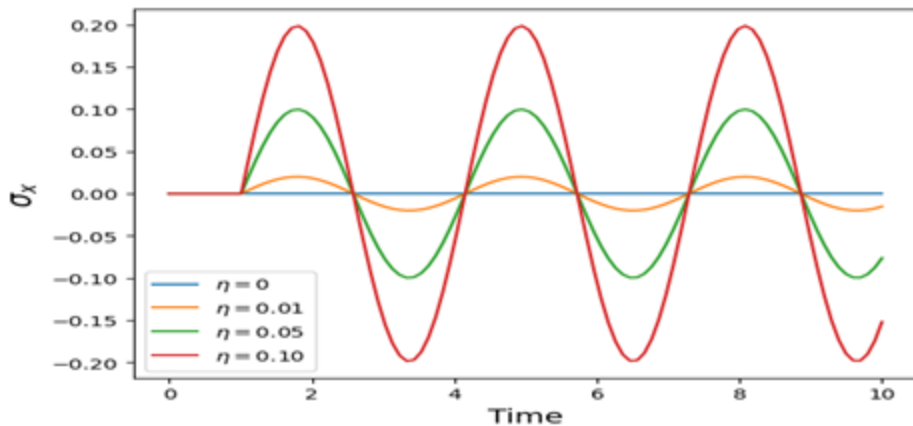
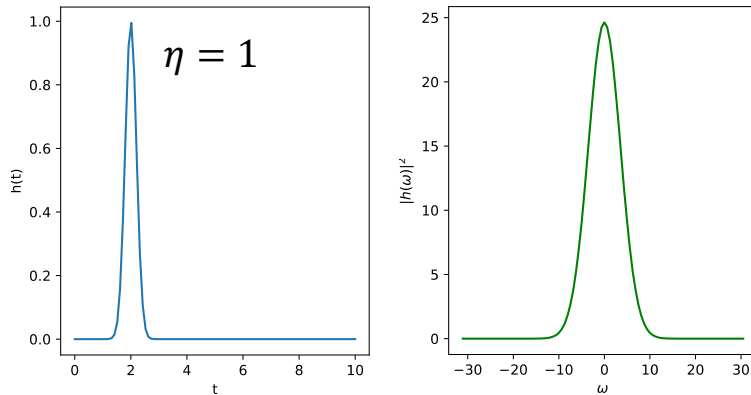
Consider: $H_0 = \sigma^z, A = B = \sigma^x$, A two-level system with $\Delta E = 2$:

1. Broad Coverage: $A(\omega) = \chi(\omega)h(\omega)$ Dictates Frequency Support!



Consider: $H_0 = \sigma^z, A = B = \sigma^x$, A two-level system with $\Delta E = 2$:

1. Broad Coverage: $A(\omega) = \chi(\omega)h(\omega)$ Dictates Frequency Support!



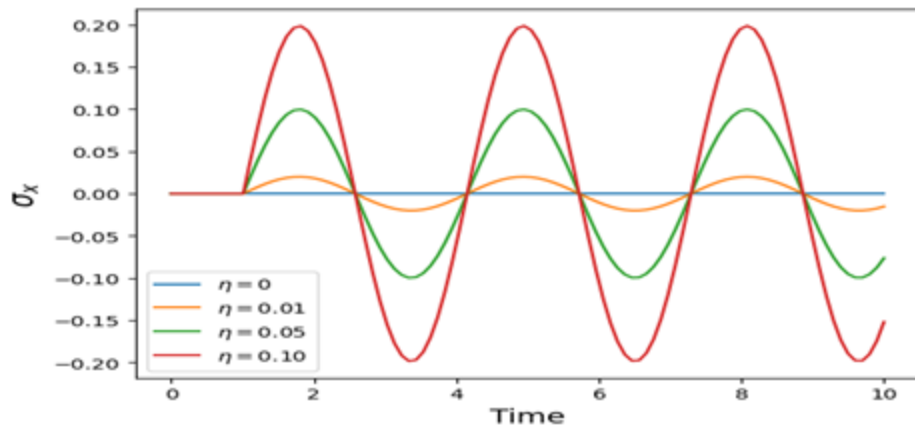
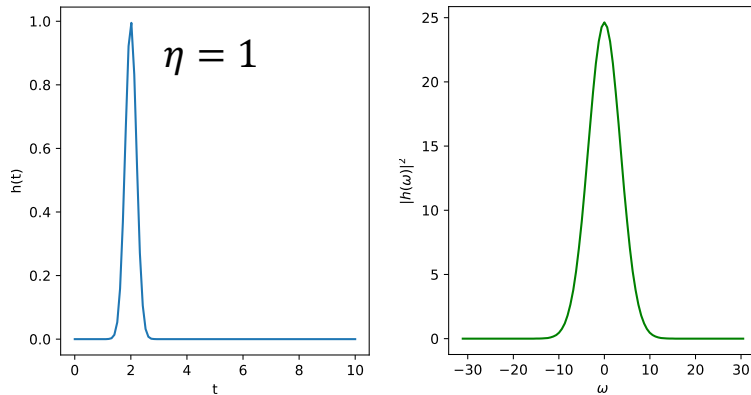
Consider: $H_0 = \sigma^z, A = B = \sigma^x$, A two-level system with $\Delta E = 2$:

$$A(\omega) = \chi(\omega)h(\omega)$$

Dictates Frequency Support!

1. Broad Coverage:

2. Selective Coverage:



Frequency Selectivity

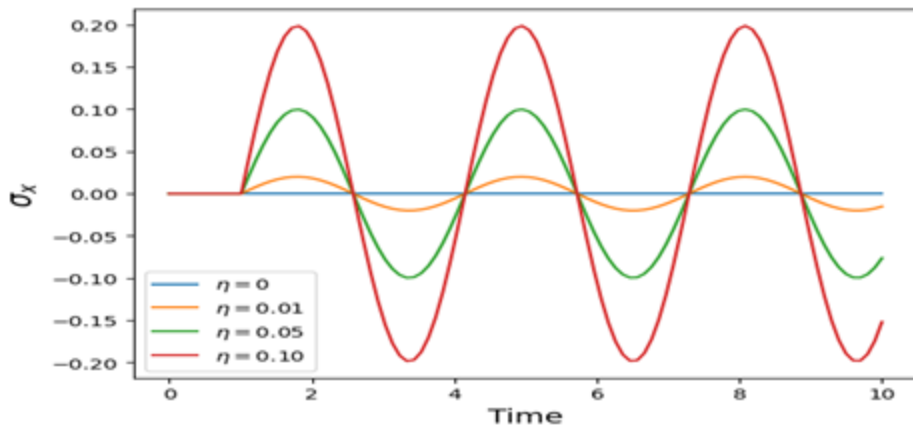
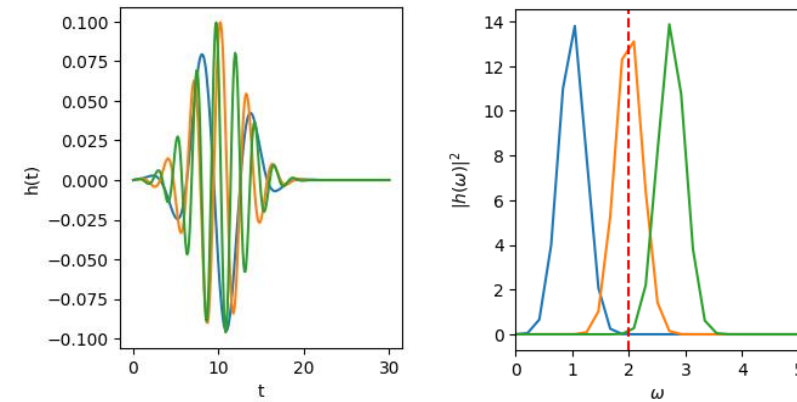
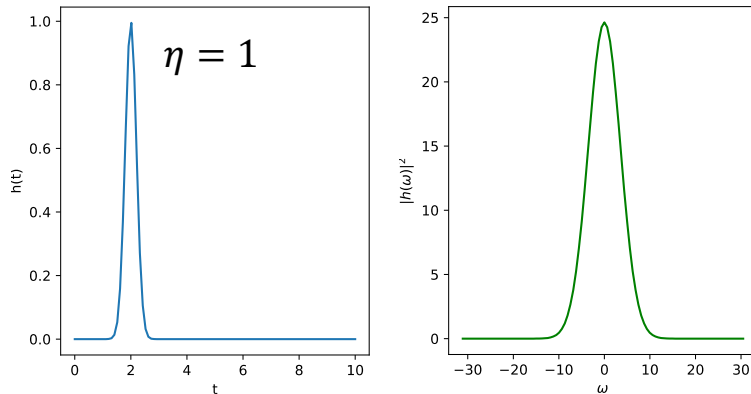
Consider: $H_0 = \sigma^z, A = B = \sigma^x$, A two-level system with $\Delta E = 2$:

$$A(\omega) = \chi(\omega)h(\omega)$$

Dictates Frequency Support!

1. Broad Coverage:

2. Selective Coverage:



Frequency Selectivity

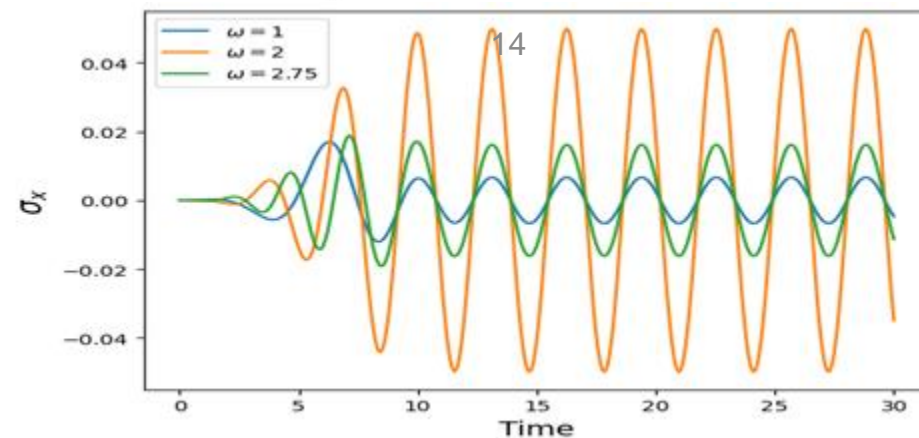
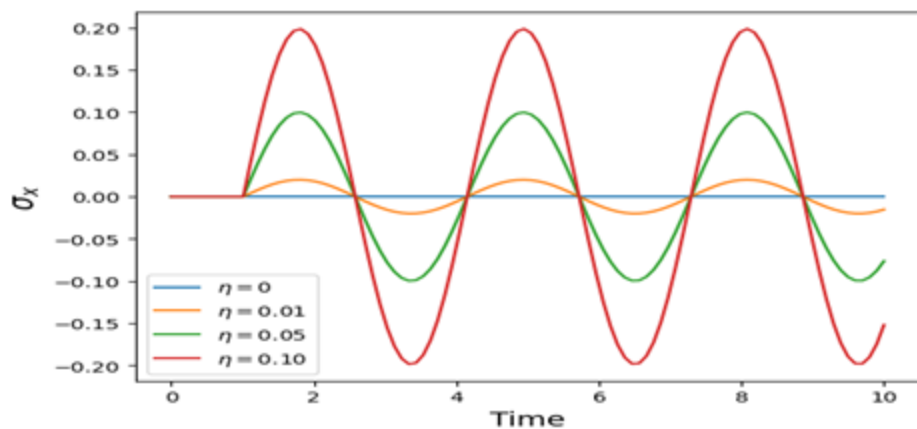
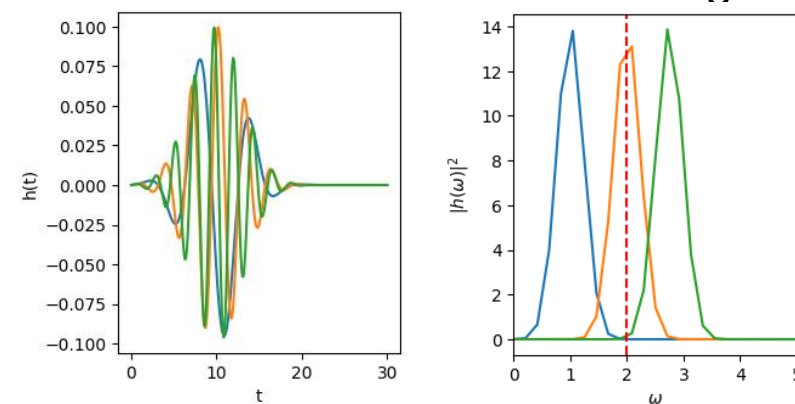
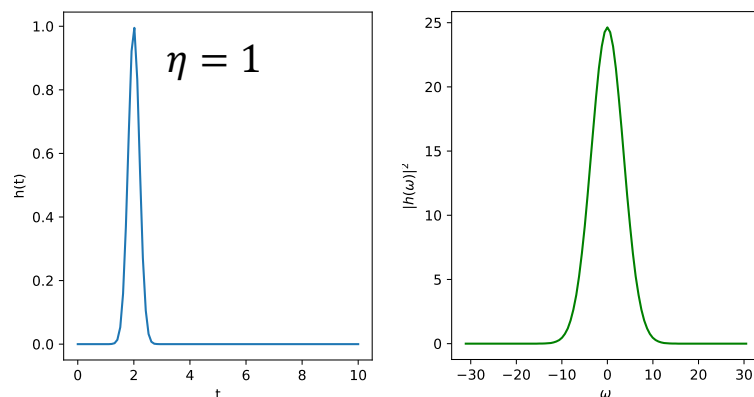
Consider: $H_0 = \sigma^z, A = B = \sigma^x$, A two-level system with $\Delta E = 2$:

$$A(\omega) = \chi(\omega)h(\omega)$$

Dictates Frequency Support!

1. Broad Coverage:

2. Selective Coverage:



Frequency Selectivity

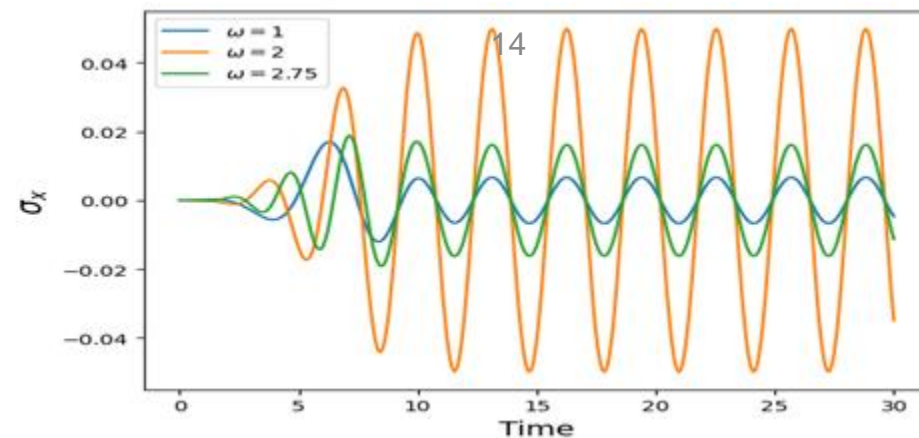
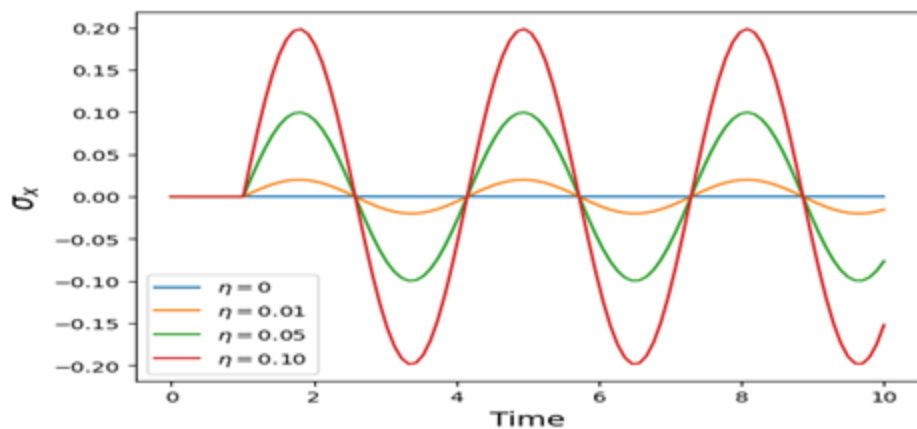
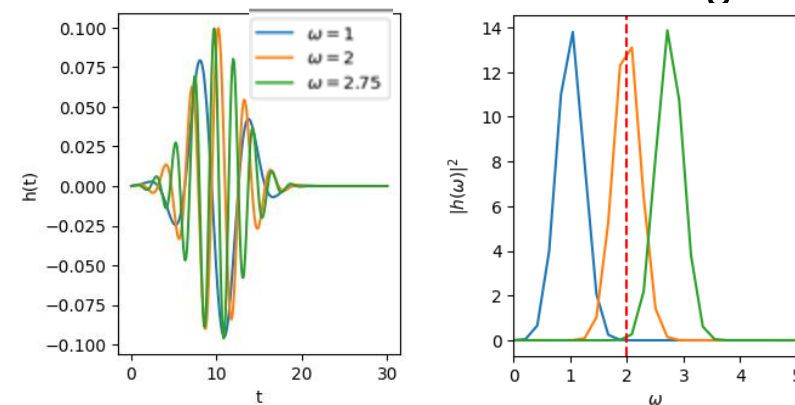
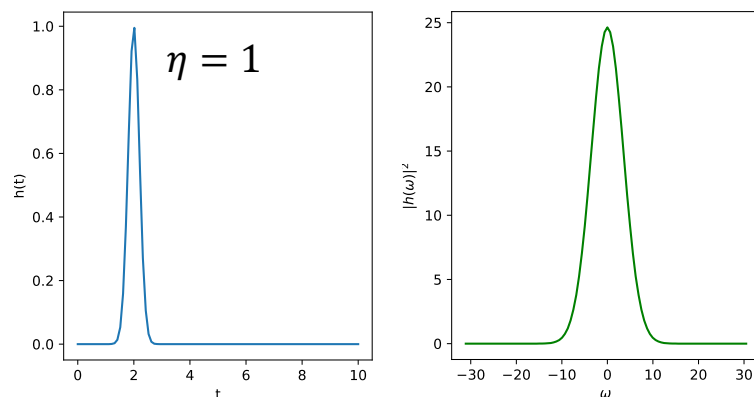
Consider: $H_0 = \sigma^z, A = B = \sigma^x$, A two-level system with $\Delta E = 2$:

$$A(\omega) = \chi(\omega)h(\omega)$$

Dictates Frequency Support!

1. Broad Coverage:

2. Selective Coverage:



Fermionic Correlation Functions

$\delta = 0$

$\delta = 0.4$

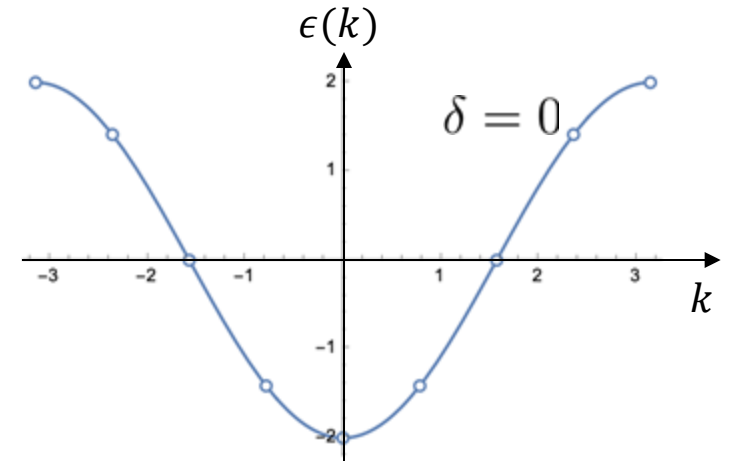
$\delta = 0.8$

Fermionic Correlation Functions

$\delta = 0$

$\delta = 0.4$

$\delta = 0.8$

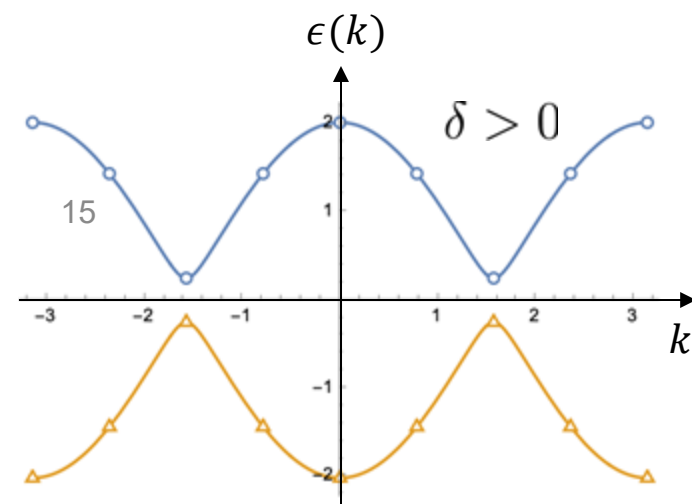
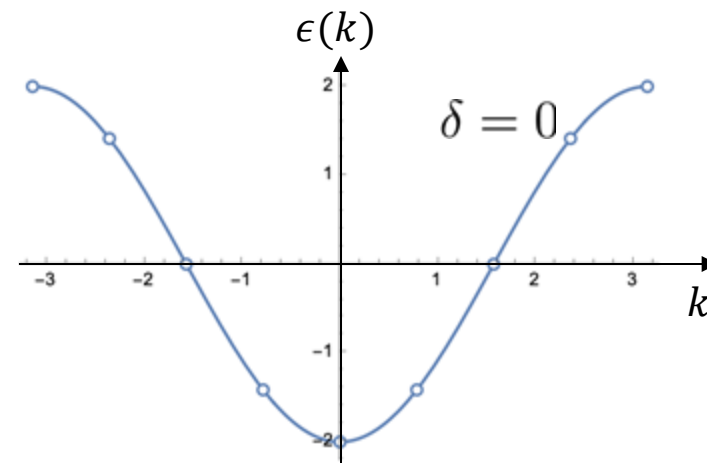


Fermionic Correlation Functions

$\delta = 0$

$\delta = 0.4$

$\delta = 0.8$



Fermionic Correlation Functions

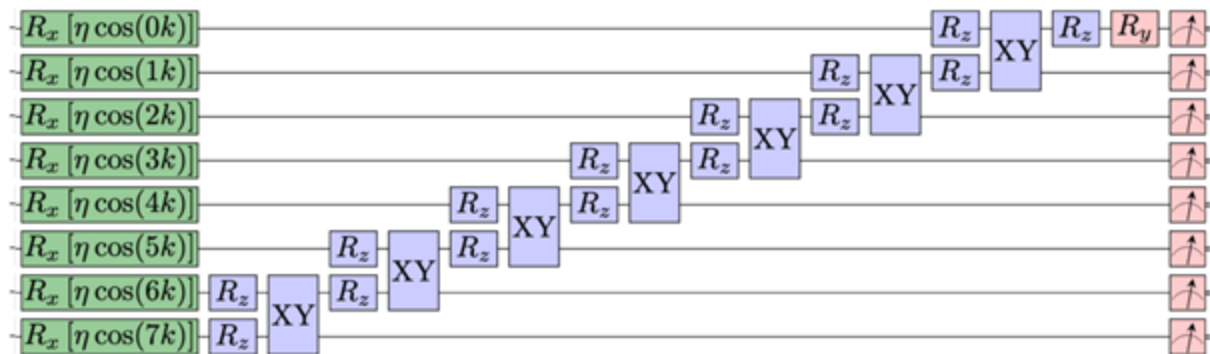
$$G^R(r_i, t; r_j, t') = -i\theta(t - t') \langle \psi_0 | \{c_i(t), c_j^\dagger(t')\} | \psi_0 \rangle$$

$$\mathbf{B} = \sum_i 2 \cos(kr_i) [c_i + c_i^\dagger]$$

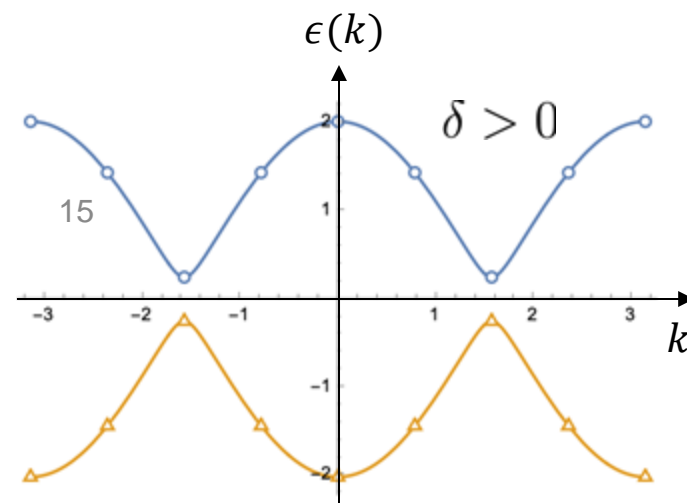
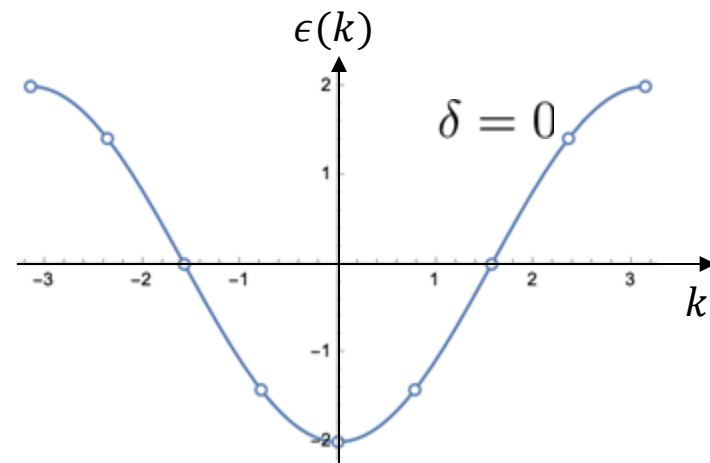
$\delta = 0$

$\delta = 0.4$

$\delta = 0.8$



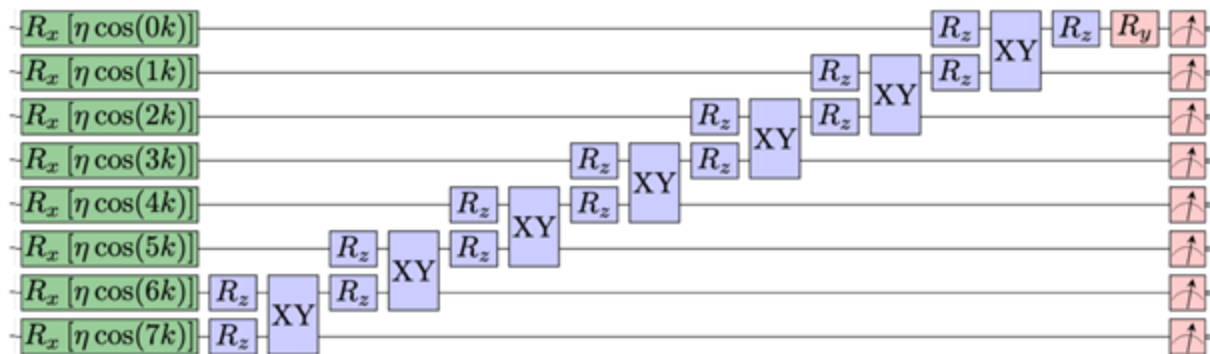
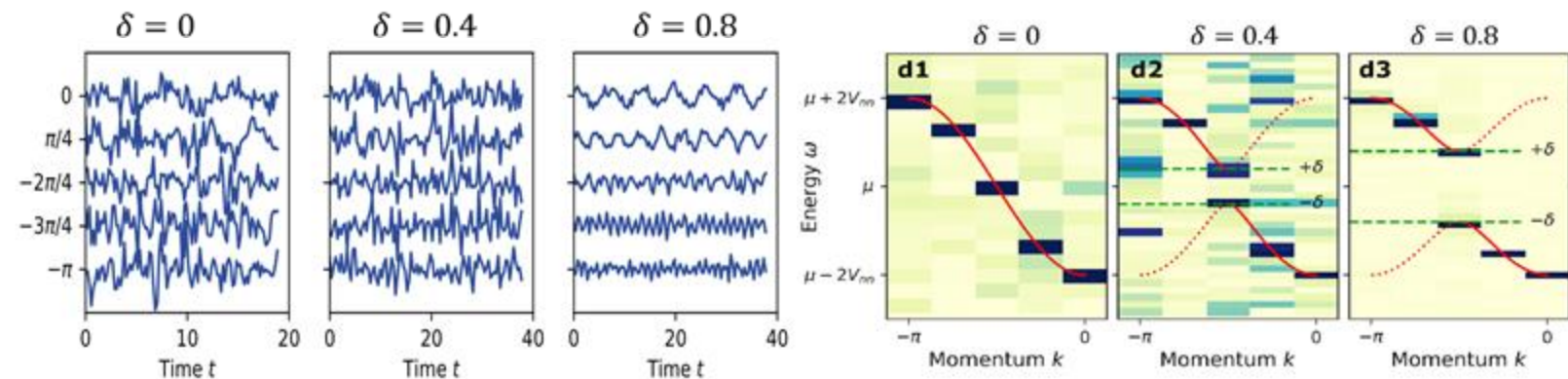
Compressed Circuit[] run on *ibm_auckland*



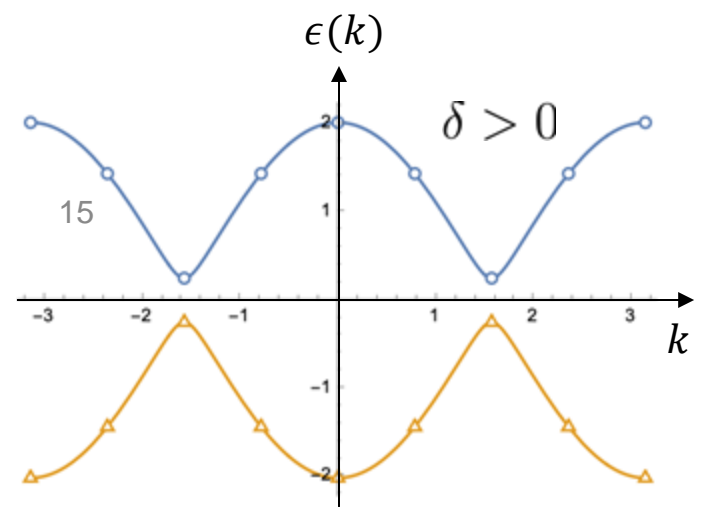
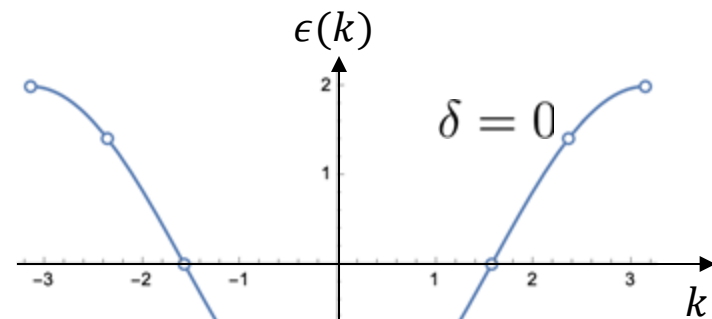
Fermionic Correlation Functions

$$G^R(r_i, t; r_j, t') = -i\theta(t - t') \langle \psi_0 | \{c_i(t), c_j^\dagger(t')\} | \psi_0 \rangle$$

$$\mathbf{B} = \sum_i 2 \cos(kr_i) [c_i + c_i^\dagger]$$



Compressed Circuit[] run on *ibm_auckland*



Cartan Decomposition and KHK Theorem

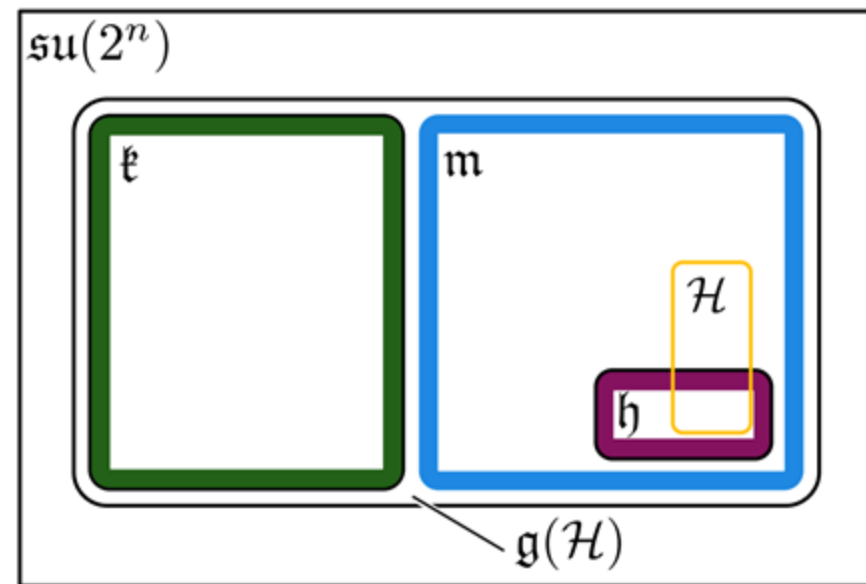
Definition 1 Consider a compact semi-simple Lie subgroup $G \subset SU(2^n)$, which has a corresponding Lie subalgebra \mathfrak{g} . A **Cartan decomposition** on \mathfrak{g} is defined as an orthogonal split $\mathfrak{g} = \mathfrak{k} \oplus \mathfrak{m}$ satisfying

$$[\mathfrak{k}, \mathfrak{k}] \subset \mathfrak{k} \quad [\mathfrak{m}, \mathfrak{m}] \subset \mathfrak{k} \quad [\mathfrak{k}, \mathfrak{m}] = \mathfrak{m} \quad (4)$$

and is referred as $(\mathfrak{g}, \mathfrak{k})$. **Cartan subalgebra** of this decomposition is defined as one of the maximal Abelian subalgebras of \mathfrak{m} , and denoted as \mathfrak{h} .

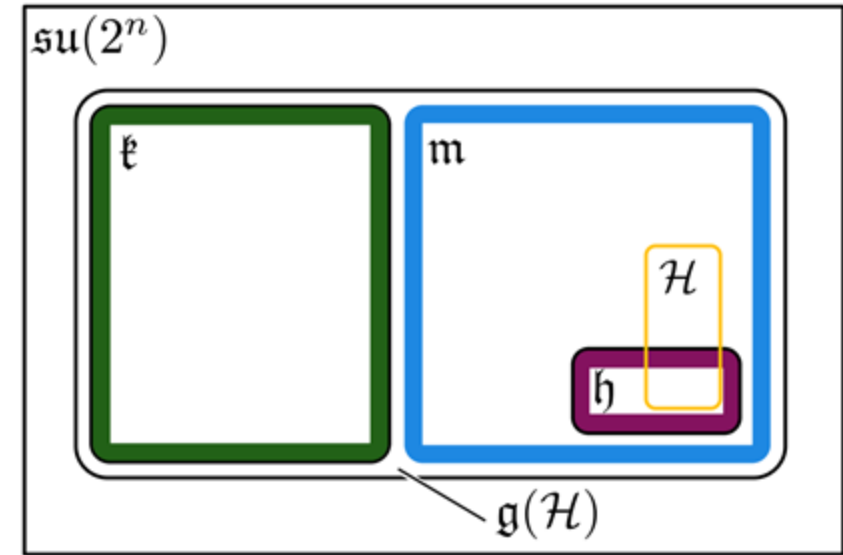
Theorem 1 Given a Cartan decomposition $\mathfrak{g} = \mathfrak{k} \oplus \mathfrak{m}$, for any element $\mathcal{H} \in \mathfrak{m}$ there exist a $K \in e^{\mathfrak{k}}$ and $h \in \mathfrak{h}$ such that

$$\mathcal{H} = KhK^\dagger \quad (5)$$



Algorithm

- 1) Generate Hamiltonian algebra $\mathfrak{g}(\mathcal{H})$
- 2) Find a Cartan decomposition such that \mathcal{H} is in \mathfrak{m}
- 3) Fit parameters via minimizing $f(K)$
- 4) Build the circuit using K and h
- 5) Then simulate for any time you want!



$$f(K) = \langle K v K^\dagger, \mathcal{H} \rangle$$

The diagram shows the decomposition of a quantum gate $e^{-it\mathcal{H}}$ into three gates: e^{-ik} , e^{-ith} , and e^{ik} .

Results

