M3M6: Methods of Mathematical Physics

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1 Lecture 11: Representing analytic functions by their behaviour near singularities

A key theme in complex analysis is representing functions by their behaviour near singularities. A simple example of this is partial fraction expansion: a rational function p(z)/q(z) can be expressed as a sum of its behaviour near poles and infinity. This is more complicated, but doable in a systematic manner for functions with branch cuts. In this lecture we:

- 1. Derive partial fraction expansion using Cauchy's integral formula
- 2. Recoverying functions like $\sqrt{(z-1)\sqrt{(z+1)}}$ from their behaviour on the branch cut

1.1 Partial fraction expansion

Corollary (Cauchy's integral representation around holes) Let $D \subset \mathbb{C}$ be a domain with g holes (i.e., genus g). Suppose f is holmorphic in and on the boundary of D. Given g simple closed negatively oriented contours surrounding the holes $\gamma_1, \ldots, \gamma_g$ and a simple closed positively oriented contour γ_{∞} surrounding the outer boundary of D, we have

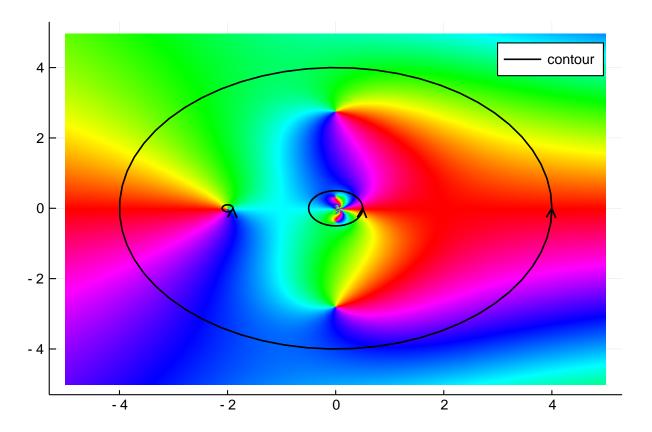
$$f(z) = \frac{1}{2\pi i} \left[\sum_{k=1}^{g} \oint_{\gamma_k} + \oint_{\gamma_\infty} \right] \frac{f(\zeta)}{\zeta - z} d\zeta$$

Here are is an example. Consider

$$f(z) = (e^{1/z} + e^z)/(z(z+2))$$

which has an essential singularity at 0 and ∞ and a simple pole at 2. We can recover f from contours around each singularity:

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using ApproxFun, ComplexPhasePortrait, Plots f = z \rightarrow (\exp(1/z) + \exp(z))/(z*(z+2)) \Gamma = \text{Circle}(0.0, 4.0) \cup \text{Circle}(0.0, 0.5, false) \cup \text{Circle}(-2.0, 0.1, false) phaseplot(-5..5, -5..5, f) plot!(\Gamma; color=:black, label=:contour, arrow=true, linewidth=1.5)
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Cauchy's integral formula is still valid:

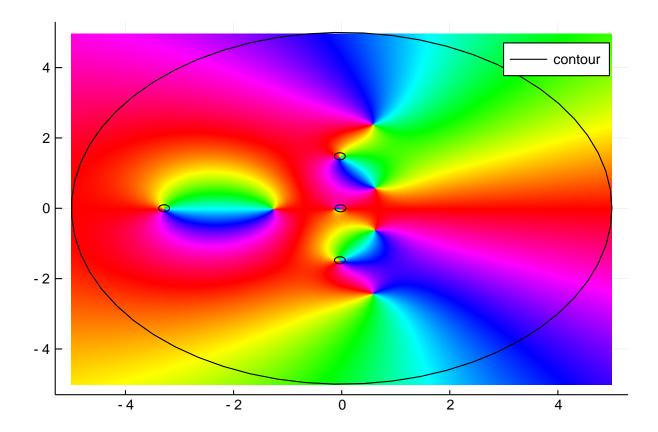
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\begin{split} & \zeta = \operatorname{Fun}(\Gamma) \\ & z = 2.0 + 1.0 \operatorname{im} \\ & \operatorname{sum}(\mathbf{f}.(\zeta)/(\zeta - z))/(2\pi * \operatorname{im}), \ \mathbf{f}(z) \end{split}
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(0.8671607060038513 + 0.10261889457156066im, 0.8671607060038514 + 0.10261889457156062im)

Now we specialise to the case where we have a rational function

$$r(z) = \frac{p(z)}{q(z)}$$

where p, q are both polynomials. This is analytic everywhere apart from the roots of q, which we enumerate $\lambda_1, \ldots, \lambda_g$. If we integrate over negatively oriented circles around each root:



we recover the function:

$$\zeta = \operatorname{Fun}(\Gamma)$$

$$z = 2.0+2.0\operatorname{im}$$

$$\operatorname{sum}(\mathbf{r}.(\zeta)/(\zeta - z))/(2\pi * \operatorname{im}) , \mathbf{r}(z)$$

(-0.43791652197566594 + 2.8021270643729217im, -0.43791652197567027 + 2.8021270643729177im)

But now we can use the residue theorem to simplify the integrals! Near the jth root we have the Laurent series

$$r(z) = r_{-N_i}^j (z - \lambda_j)^{-N_j} + \dots + r_{-1}^j (z - \lambda_j)^{-1} + r_0 + r_1 (z - \lambda_j) + \dots$$

where N_j is the order of the zero of q(z) at λ_j .

Then it follows that

$$\frac{1}{2\pi i} \oint_{\gamma_j} \frac{r(\zeta)}{z - \zeta} d\zeta = r_{-N_j}^j (z - \lambda_j)^{-N} + \dots + r_{-1}^j (z - \lambda_j)^{-1}$$

for z outside the contour γ_j .

Similarly, for the contour around infinity γ_{∞} , if we have the Laurent series

$$r(z) = \dots + r_{-1}^{\infty} z^{-1} + r_0^{\infty} + r_1^{\infty} z + \dots + r_{N_0}^{\infty} z^{N_0}$$

where N_{∞} is the degree of p(z) divided by the degree of q(z). Then we have

$$\frac{1}{2\pi \mathrm{i}} \oint_{\gamma_{\infty}} \frac{r(\zeta)}{z - \zeta} \mathrm{d}\zeta = r_0^{\infty} + r_1^{\infty} z + \dots + r_{N_0}^{\infty} z^{N_0}.$$

Thus we have the expansion summing over the behaviour near each singularity that holds for all z:

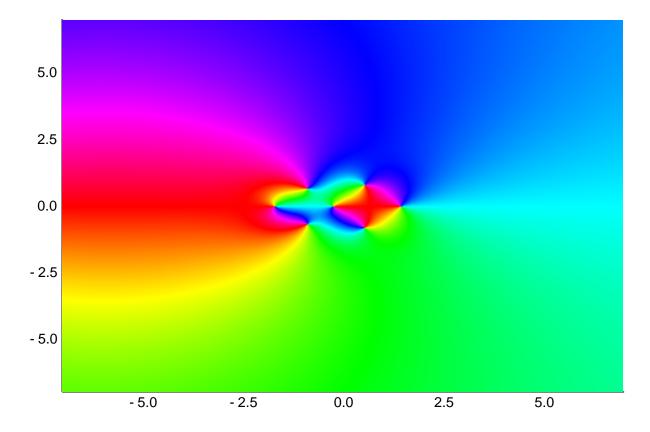
$$r(z) = \sum_{k=0}^{N_{\infty}} r_k^{\infty} z^k + \sum_{j=1}^{d} \sum_{k=-N_j}^{-1} r_k^j (z - \lambda_j)^k$$

Example When we only have simple poles and no polynomial growth at ∞ , this has a simple form in terms of residues:

$$r(z) = r(\infty) + \sum_{j=1}^{d} (z - \lambda_j)^{-1} \operatorname{Res}_{z = \lambda_j} r(z)$$

Here we demonstrate it on a random polynomial:

```
n = 5
m = 5
p = Fun(Taylor(), randn(n))
q = Fun(Taylor(), randn(m))
\[ \lambda = \text{complexroots(q)} \]
\[ r = z \to -> \text{extrapolate(p,z)/extrapolate(q,z)} \]
\[ phaseplot(-7..7, -7..7, r) \]
```



This constructs r_2 as the partial fraction expansion of r:

```
res = extrapolate.(p,\lambda)./extrapolate.(q',\lambda)

r\infty = p.coefficients[n]/q.coefficients[m]

r_2 = z \rightarrow r\infty + sum(res.*(z .- <math>\lambda).^(-1))
```

```
z = 0.1+0.2im

r(z) - r_2(z) # we match to high accuracy
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-1.1657341758564144e-15 + 1.6653345369377348e-16im

1.2 Recovering analytic functions

We now consider the above approach for 2 examples with branch cuts.

Example 1

Consider $\phi(z) = \log(z-1) - \log(z+1)$. For x < -1 the branch cuts cancel and we have

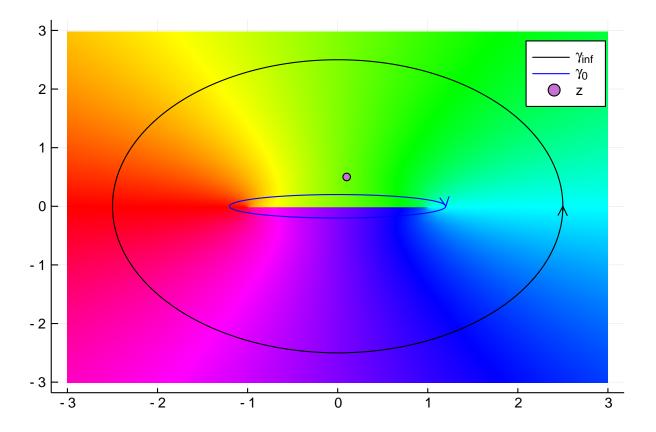
$$\phi_{+}(x) = \lim \phi(x + i\epsilon) = \log_{+}(x - 1) - \log_{+}(x + 1) = \log|x - 1| + i\pi - \log|x + 1| - i\pi = \log(1 - x) - \log(-1 - x).$$

Similarly

$$\phi_{-}(x) = \log(1-x) - \log(-1-x) = \phi_{+}(x)$$

i.e., we are continuous on the branch cut (with $\phi(x) := \phi_+(x)$) and therefore analytic. Thus $\phi(z)$ is analytic off [-1,1] which can be seen clearly from a phase portrait. Using the corollary above we can recover f from integrating over two contours: γ_{∞} surrounding ∞ and γ_0 surrounding the branch cut, with z in-between:

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\varphi = z \rightarrow \log(z-1) - \log(z+1) phaseplot(-3..3, -3..3, \varphi) \theta = range(0,2\pi; length=200) plot!(2.5cos.(\theta), 2.5sin.(\theta); color=:black, label="\gamma_inf", arrow=true) plot!(1.2cos.(\theta), 0.2sin.(-\theta); color=:blue, label="\gamma_0", arrow=true) scatter!([0.1],[0.5]; label="z")
```



That is, we have

$$\phi(z) = \frac{1}{2\pi i} \left[\oint_{\gamma_0} + \oint_{\gamma_\infty} \right] \frac{\phi(\zeta)}{\zeta - z} d\zeta$$

Note that $\phi(z)$ is analytic in a neighbourhood of ∞ and has weaker-than-pole growth: it can only grow at worse like $\log z$. Thus it must be analytic at ∞ . To determine the constant we use a simple Taylor series argument, evaluating on the real axis to simplify rules of logarithms: as $x \to \infty$ we have

$$\phi(x) = \log(x-1) - \log(x+1) = \log x + \log(1-1/x) - \log x - \log(1+1/x) \to 0,$$

that is $\phi(\infty) = 0$, and in fact $\phi(z) = O(z^{-1})$ as $\phi(z^{-1})$ is analytic at zero and therefore has a converging Taylor series. It follows from Cauchy's theorem (exterior) that

$$\oint_{\gamma_{\infty}} \frac{\phi(\zeta)}{\zeta - z} d\zeta = 0$$

as the integrand decays like $O(\zeta^{-2})$.

We are left with the integral on γ_0 . We can think of it as a rectangular contour with contours $[-1 - \epsilon - i\epsilon, -1 - \epsilon + i\epsilon, 1 + \epsilon + i\epsilon, 1 + \epsilon - i\epsilon]$. Letting $\epsilon \to 0$, on the contour above $\phi(z)$ tends to

$$\lim_{\epsilon \to 0} \phi(x + i\epsilon) = \phi_+(x)$$

and similar to the contour below. Since ϕ only has logarithmic singularities this limit can be done safely. Thus we end up with the expression

$$\phi(z) = \frac{1}{2\pi i} \oint_{\gamma_0} \frac{\phi(\zeta)}{\zeta - z} d\zeta = \frac{1}{2\pi i} \int_{-1}^1 \frac{\phi_+(x) - \phi_-(x)}{x - z} dx = \int_{-1}^1 \frac{1}{x - z} dx.$$

Example 2

We repeat the above procedure with $\phi(z) = \sqrt{z-1}\sqrt{z+1}$. Again this is analytic off [-1,1] and we can express it as integrals over γ_0 and γ_∞ . Now it grows like z at ∞ ,

$$\phi(z) = z + O(z^{-1}),$$

hence we have (as above)

$$\frac{1}{2\pi i} \oint_{\gamma_{\infty}} \frac{\phi(\zeta)}{\zeta - z} d\zeta = z.$$

The integral over the contour γ_0 can be collapsed. On the jump -1 < x < 1 we have

$$\phi_{+}(x) = \sqrt{x-1}_{+}\sqrt{x+1} = i\sqrt{|x-1|}\sqrt{x+1} = i\sqrt{1-x}\sqrt{x+1} = i\sqrt{1-x^2}$$
 while $\phi_{-}(x) = -\phi_{+}(x) = -i\sqrt{1-x^2}$. We thus have

$$\phi(z) = z + \frac{1}{2\pi i} \oint_{\gamma_0} \frac{\phi(\zeta)}{\zeta - z} d\zeta = z + \frac{1}{2\pi i} \int_{-1}^1 \frac{\phi_+(x) - \phi_-(x)}{x - z} dx$$
$$= z + \frac{1}{\pi} \int_{-1}^1 \frac{\sqrt{1 - x^2}}{x - z} dx.$$