

1 Lecture 11: Representing analytic functions by their behaviour near singularities

A *key* theme in complex analysis is representing functions by their behaviour near singularities. A simple example of this is partial fraction expansion: a rational function $p(z)/q(z)$ can be expressed as a sum of its behaviour near poles and infinity. This is more complicated, but doable in a systematic manner for functions with branch cuts. In this lecture we:

1. Derive partial fraction expansion using Cauchy's integral formula
2. Recovering functions like $\sqrt{(z-1)}\sqrt{(z+1)}$ from their behaviour on the branch cut

1.1 Partial fraction expansion

Corollary (Cauchy's integral representation around holes) Let $D \subset \mathbb{C}$ be a domain with g holes (i.e., genus g). Suppose f is holomorphic in and on the boundary of D . Given g simple closed negatively oriented contours surrounding the holes $\gamma_1, \dots, \gamma_g$ and a simple closed positively oriented contour γ_∞ surrounding the outer boundary of D , we have

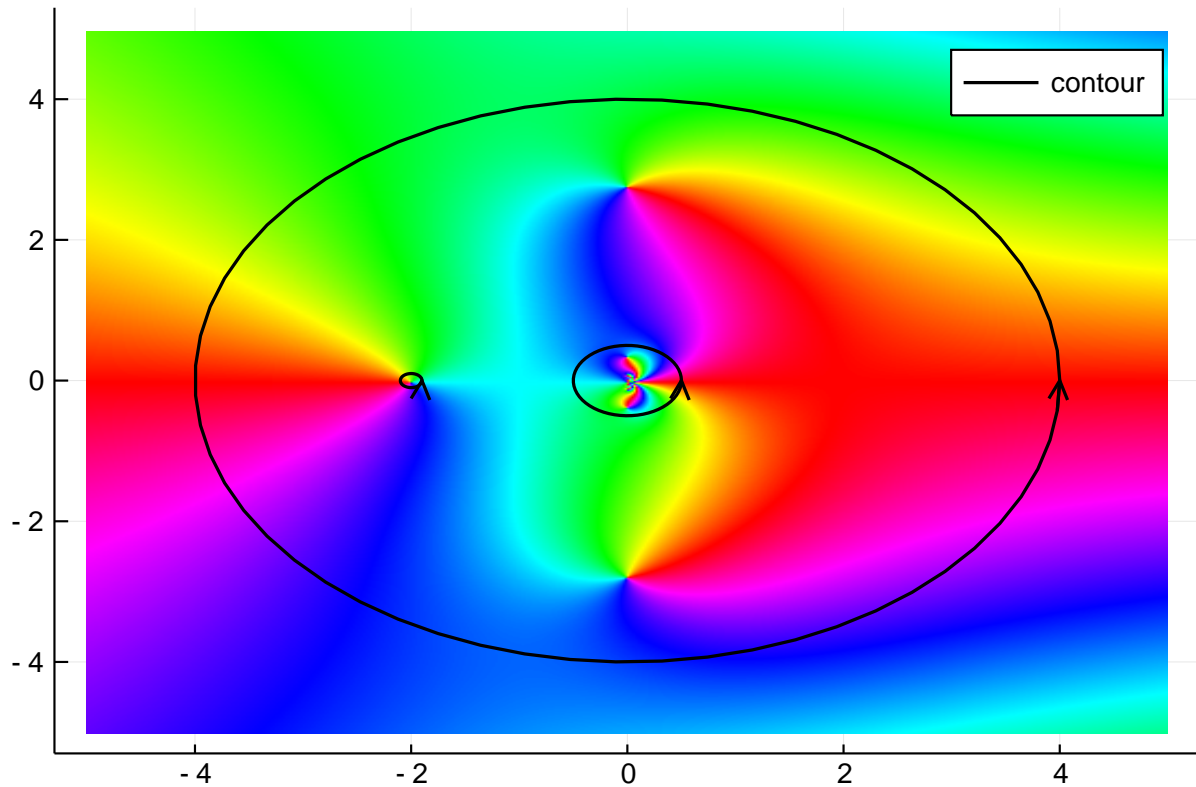
$$f(z) = \frac{1}{2\pi i} \left[\sum_{k=1}^g \oint_{\gamma_k} + \oint_{\gamma_\infty} \right] \frac{f(\zeta)}{\zeta - z} d\zeta$$

Here are is an example. Consider

$$f(z) = (e^{1/z} + e^z)/(z(z+2))$$

which has an essential singularity at 0 and ∞ and a simple pole at 2. We can recover f from contours around each singularity:

```
using ApproxFun, ComplexPhasePortrait, Plots
f = z -> (exp(1/z) + exp(z))/(z*(z+2))
Γ = Circle(0.0, 4.0) ∪ Circle(0.0,0.5,false) ∪ Circle(-2.0,0.1,false)
phaseplot(-5..5, -5..5, f)
plot!(Γ; color=:black, label=:contour, arrow=true, linewidth=1.5)
```



Cauchy's integral formula is still valid:

```

ζ = Fun(Γ)
z = 2.0+1.0im
sum(f.(ζ)/(ζ - z))/(2π*im), f(z)

(0.8671607060038513 + 0.10261889457156066im, 0.8671607060038514 + 0.1026188
9457156062im)

```

Now we specialise to the case where we have a rational function

$$r(z) = \frac{p(z)}{q(z)}$$

where p, q are both polynomials. This is analytic everywhere apart from the roots of q , which we enumerate $\lambda_1, \dots, \lambda_g$. If we integrate over negatively oriented circles around each root:

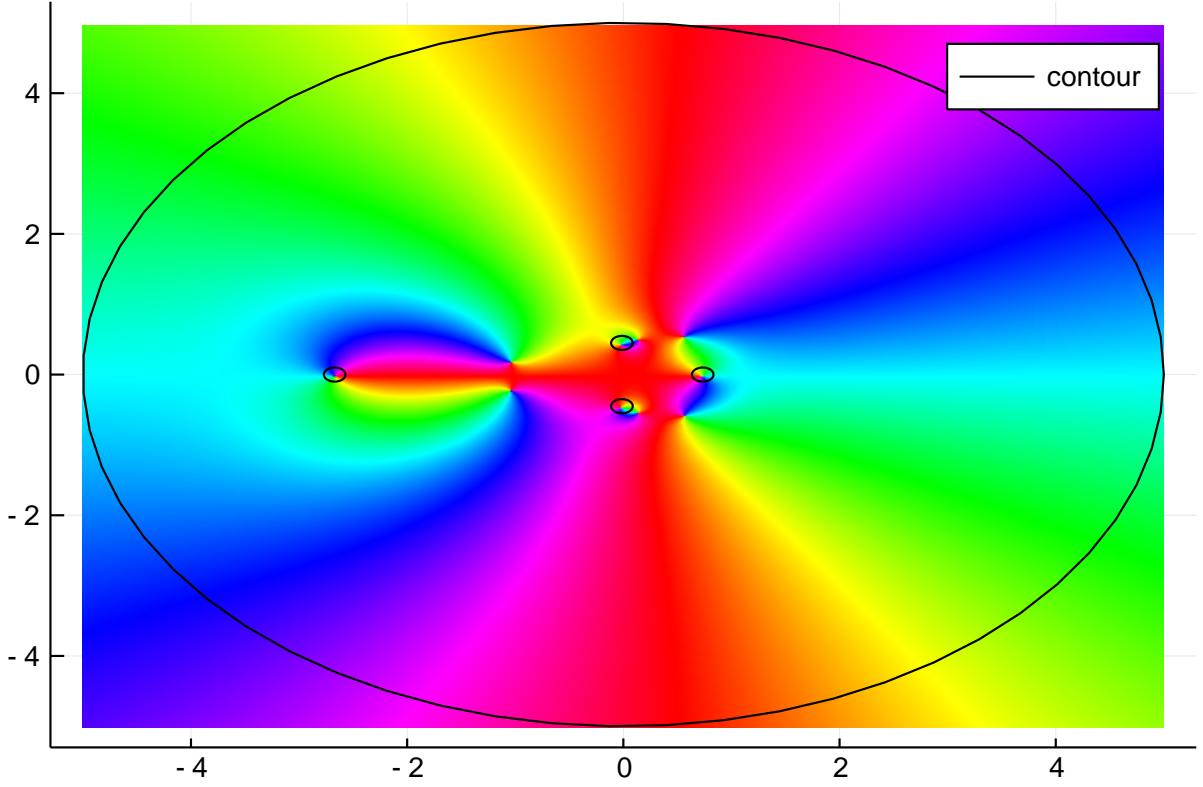
```

n = 7
m = 5
p = Fun(Taylor(), randn(n))
q = Fun(Taylor(), randn(m))
λ = complexroots(q)

Γ = Circle(0.0, 5.0)
for λ in λ
    global Γ
    Γ = Γ ∪ Circle(λ, 0.1, false)
end
r = z -> extrapolate(p,z)/extrapolate(q,z)

phaseplot(-5..5, -5..5, r)
plot!(Γ; color=:black, label=:contour)

```



we recover the function:

```

ζ = Fun(Γ)
z = 2.0+2.0im
sum(r.(ζ)/(ζ - z))/(2π*im) , r(z)

(4.372717213418816 - 46.349833892497706im, 4.372717213418801 - 46.349833892
497685im)

```

But now we can use the residue theorem to simplify the integrals!

Near the j th root we have the Laurent series

$$r(z) = r_{-N_j}^j (z - \lambda_j)^{-N_j} + \cdots + r_{-1}^j (z - \lambda_j)^{-1} + r_0 + r_1(z - \lambda_j) + \cdots$$

where N_j is the order of the zero of $q(z)$ at λ_j .

Then it follows that

$$\frac{1}{2\pi i} \oint_{\gamma_j} \frac{r(\zeta)}{z - \zeta} d\zeta = r_{-N_j}^j (z - \lambda_j)^{-N} + \cdots + r_{-1}^j (z - \lambda_j)^{-1}$$

for z outside the contour γ_j .

Similarly, for the contour around infinity γ_∞ , if we have the Laurent series

$$r(z) = \cdots + r_{-1}^\infty z^{-1} + r_0^\infty + r_1^\infty z + \cdots + r_{N_0}^\infty z^{N_0}$$

where N_∞ is the degree of $p(z)$ divided by the degree of $q(z)$. Then we have

$$\frac{1}{2\pi i} \oint_{\gamma_\infty} \frac{r(\zeta)}{z - \zeta} d\zeta = r_0^\infty + r_1^\infty z + \cdots + r_{N_0}^\infty z^{N_0}.$$

Thus we have the expansion summing over the behaviour near each singularity that holds for all z :

$$r(z) = \sum_{k=0}^{N_\infty} r_k^\infty z^k + \sum_{j=1}^d \sum_{k=-N_j}^{-1} r_k^j (z - \lambda_j)^k$$

Example When we only have simple poles and no polynomial growth at ∞ , this has a simple form in terms of residues:

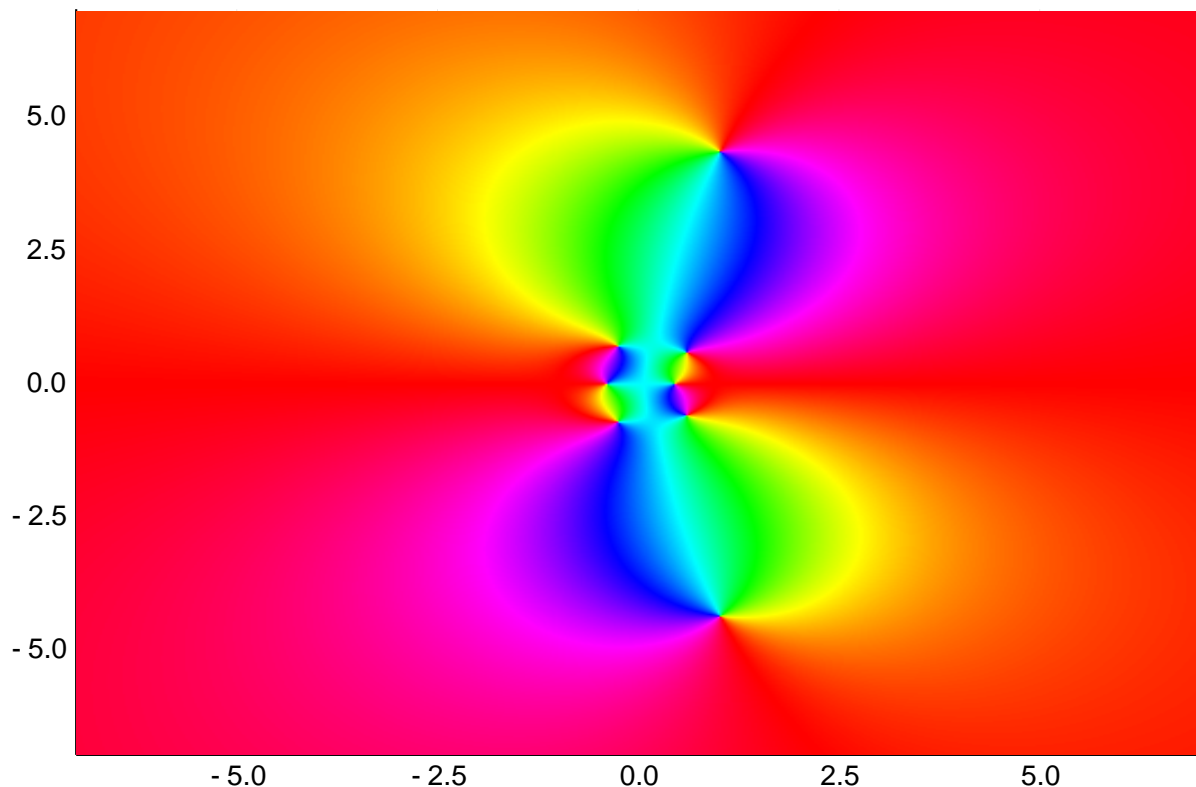
$$r(z) = r(\infty) + \sum_{j=1}^d (z - \lambda_j)^{-1} \operatorname{Res}_{z=\lambda_j} r(z)$$

Here we demonstrate it on a random polynomial:

```
n = 5
m = 5
p = Fun(Taylor(), randn(n))
q = Fun(Taylor(), randn(m))
λ = complexroots(q)

r = z -> extrapolate(p,z)/extrapolate(q,z)

phaseplot(-7..7, -7..7, r)
```



This constructs `r_2` as the partial fraction expansion of `r`:

```
res = extrapolate.(p,λ)./extrapolate.(q',λ)
r∞ = p.coefficients[n]/q.coefficients[m]

r_2 = z -> r∞ + sum(res.*(z .- λ).^(-1))
```

```

z = 0.1+0.2im
r(z) - r.2(z) # we match to high accuracy

0.0 + 8.153200337090993e-16im

```

1.2 Recovering analytic functions

We now consider the above approach for 2 examples with branch cuts.

Example 1

Consider $\phi(z) = \log(z-1) - \log(z+1)$. For $x < -1$ the branch cuts cancel and we have

$$\phi_+(x) = \lim \phi(x+i\epsilon) = \log_+(x-1) - \log_+(x+1) = \log|x-1| + i\pi - \log|x+1| - i\pi = \log(1-x) - \log(-1-x).$$

Similarly

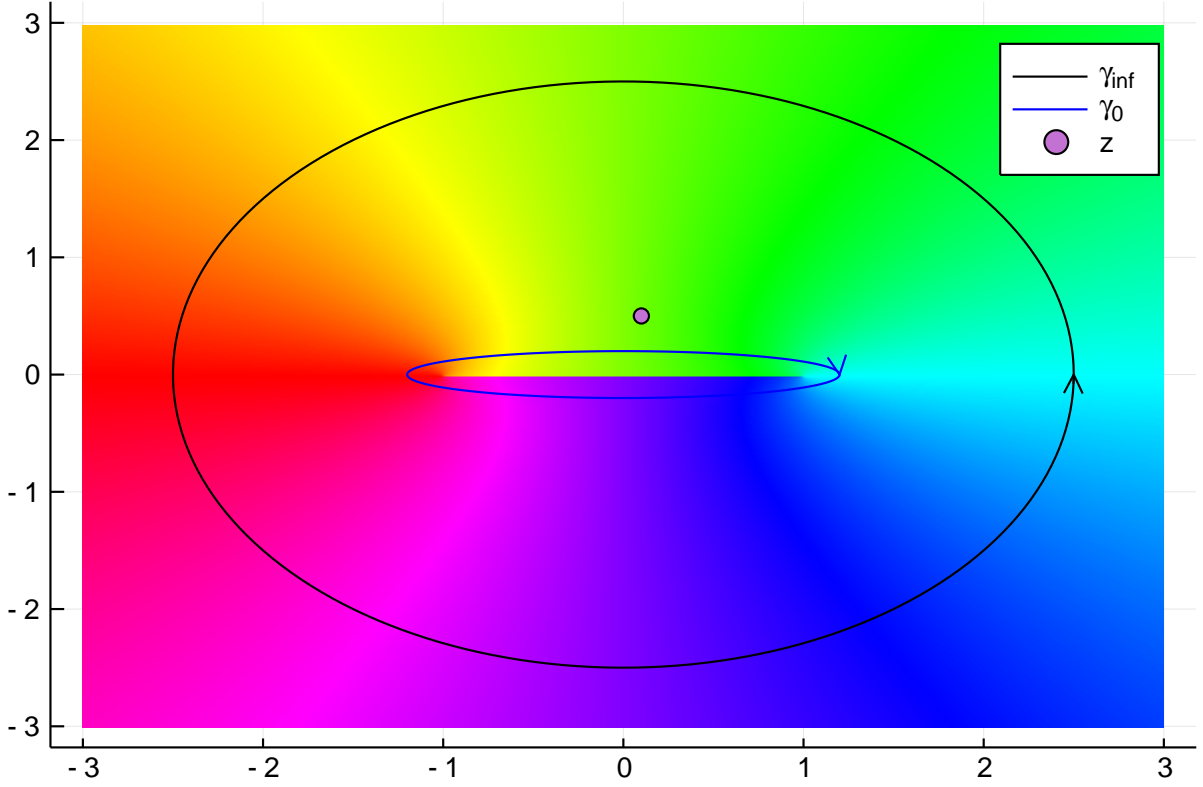
$$\phi_-(x) = \log(1-x) - \log(-1-x) = \phi_+(x)$$

i.e., we are continuous on the branch cut (with $\phi(x) := \phi_+(x)$) and therefore analytic. Thus $\phi(z)$ is analytic off $[-1, 1]$ which can be seen clearly from a phase portrait. Using the corollary above we can recover f from integrating over two contours: γ_∞ surrounding ∞ and γ_0 surrounding the branch cut, with z in-between:

```

φ = z -> log(z-1) - log(z+1)
phaseplot(-3..3, -3..3, φ)
θ = range(0, 2π; length=200)
plot!(2.5cos.(θ), 2.5sin.(θ); color=:black, label="\gamma_\infty", arrow=true)
plot!(1.2cos.(θ), 0.2sin.(-θ); color=:blue, label="\gamma_0", arrow=true)
scatter!([0.1], [0.5]; label="z")

```



That is, we have

$$\phi(z) = \frac{1}{2\pi i} \left[\oint_{\gamma_0} + \oint_{\gamma_{\infty}} \right] \frac{\phi(\zeta)}{\zeta - z} d\zeta$$

Note that $\phi(z)$ is analytic in a neighbourhood of ∞ and has weaker-than-pole growth: it can only grow at worse like $\log z$. Thus it must be analytic at ∞ . To determine the constant we use a simple Taylor series argument, evaluating on the real axis to simplify rules of logarithms: as $x \rightarrow \infty$ we have

$$\phi(x) = \log(x-1) - \log(x+1) = \log x + \log(1-1/x) - \log x - \log(1+1/x) \rightarrow 0,$$

that is $\phi(\infty) = 0$, and in fact $\phi(z) = O(z^{-1})$ as $\phi(z^{-1})$ is analytic at zero and therefore has a converging Taylor series. It follows from Cauchy's theorem (exterior) that

$$\oint_{\gamma_{\infty}} \frac{\phi(\zeta)}{\zeta - z} d\zeta = 0$$

as the integrand decays like $O(\zeta^{-2})$.

We are left with the integral on γ_0 . We can think of it as a rectangular contour with contours $[-1 - \epsilon - i\epsilon, -1 - \epsilon + i\epsilon, 1 + \epsilon + i\epsilon, 1 + \epsilon - i\epsilon]$. Letting $\epsilon \rightarrow 0$, on the contour above $\phi(z)$ tends to

$$\lim_{\epsilon \rightarrow 0} \phi(x + i\epsilon) = \phi_+(x)$$

and similar to the contour below. Since ϕ only has logarithmic singularities this limit can be done safely. Thus we end up with the expression

$$\phi(z) = \frac{1}{2\pi i} \oint_{\gamma_0} \frac{\phi(\zeta)}{\zeta - z} d\zeta = \frac{1}{2\pi i} \int_{-1}^1 \frac{\phi_+(x) - \phi_-(x)}{x - z} dx = \int_{-1}^1 \frac{1}{x - z} dx.$$

Example 2

We repeat the above procedure with $\phi(z) = \sqrt{z-1}\sqrt{z+1}$. Again this is analytic off $[-1, 1]$ and we can express it as integrals over γ_0 and γ_∞ . Now it grows like z at ∞ ,

$$\phi(z) = z + O(z^{-1}),$$

hence we have (as above)

$$\frac{1}{2\pi i} \oint_{\gamma_\infty} \frac{\phi(\zeta)}{\zeta - z} d\zeta = z.$$

The integral over the contour γ_0 can be collapsed. On the jump $-1 < x < 1$ we have

$$\phi_+(x) = \sqrt{x-1}_+ \sqrt{x+1} = i\sqrt{|x-1|}\sqrt{x+1} = i\sqrt{1-x}\sqrt{x+1} = i\sqrt{1-x^2}$$

while $\phi_-(x) = -\phi_+(x) = -i\sqrt{1-x^2}$. We thus have

$$\begin{aligned} \phi(z) &= z + \frac{1}{2\pi i} \oint_{\gamma_0} \frac{\phi(\zeta)}{\zeta - z} d\zeta = z + \frac{1}{2\pi i} \int_{-1}^1 \frac{\phi_+(x) - \phi_-(x)}{x - z} dx \\ &= z + \frac{1}{\pi} \int_{-1}^1 \frac{\sqrt{1-x^2}}{x - z} dx. \end{aligned}$$