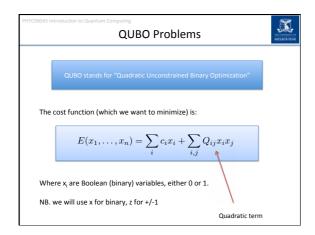


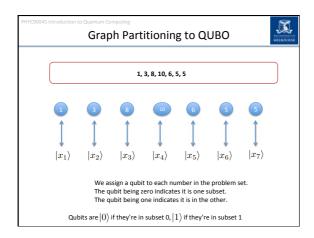
Recall Total Energy of the System
Consider a system that has an energy function:
$E = J_{12}z_1z_2 + J_{23}z_2z_3 + J_{13}z_1z_3 + B_1z_1 + B_2z_2 + B_3z_3$
where the $z_i$ are +/-1, and the J's and B's are specific parameters defining the particular problem at hand.
To get ready to map to a QC, we write the total energy as the operator "H" (which physicists would call the "Hamiltonian") on a system of qubits as a sum of these terms with $z_i \! \! \! \! \! \! \! \! \! \! \! \! \! \! \! \! \! \! \!$
$H = J_{12}Z_1Z_2 + J_{23}Z_2Z_3 + J_{13}Z_1Z_3 + B_1Z_1 + B_2Z_2 + B_3Z_3$
Pairwise interactions between qubits Bias on individual qubit

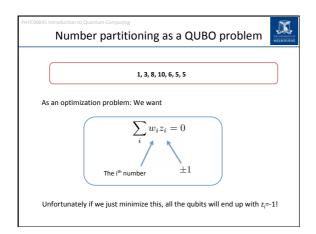
PHYC90045 Introduction to Q  Ma	pping the Spin Glass form to QC
Optimisation proble	ems can often be cast into an equivalent "spin glass" form:
-> convert to a conv	$E = \sum_{i  eq j} J_{ij} z_i z_j + \sum_i B_i z_i$ renient form to map onto a quantum computer:
	Ising coupling local "field" $H = \sum_{i \neq j} J_{ij} Z_i Z_j + \sum_i B_i Z_i$
	erators defined as per our definitions with eigenvalues +/-1 pped to binary variables 0/1)



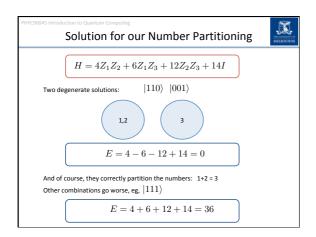
PHYC90045 Introduction to Quantum Computing	<b>P</b>
Binary to energy	DE CONTRACTOR DE METEROLIENE
Typically when we write such energy functions we write	e in terms of the Z variables:
$oxed{z_i}$	(lower case z)
But the binary variables in terms of 0 or 1:	
$\boxed{x_i = 0  \text{or}  x_i = 1}$	
Can convert between $\mathbf{x}_i$ and $\mathbf{Z}_i$ using:	
$x_i \to \frac{z_i + 1}{2}$	

PHYC9004	Example: Number partitioning
	Given a set, S, of numbers:
	1, 3, 8, 10, 6, 5, 5
	is there a partition of this set of numbers into two disjoint subsets R and S – R, such that the sum of the elements in both sets is the same?
	Yes (in this case): <b>{1, 8, 10}</b> and <b>{3, 6, 5, 5}</b>

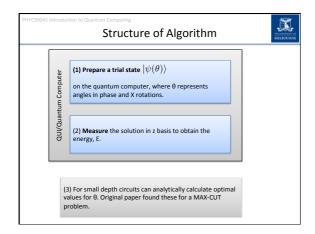


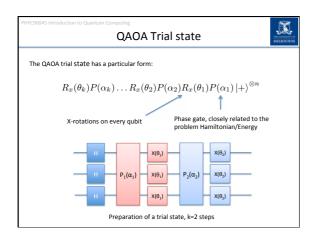


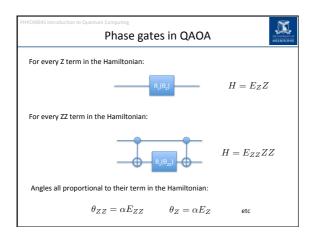
Number partitioning as a QUBO problem
But if we square, we should get a positive solution (or zero). We want to find the assignment of spins which has the minimum energy (ie. closest to zero):
$H = \left(\sum_i w_i Z_i\right)^2 = \sum_{i \neq j} 2w_i w_j Z_i Z_j + \sum_i w_i^2 I$ Coupling is the product of numbers
Eg. For the set <b>{1, 2, 3}</b> :
$H = 4Z_1Z_2 + 6Z_1Z_3 + 12Z_2Z_3 + 14I$ Finding minimum energy state will solve the problem!

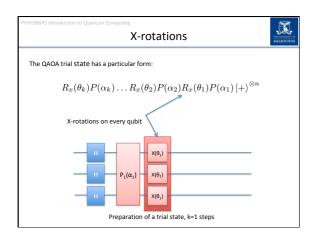


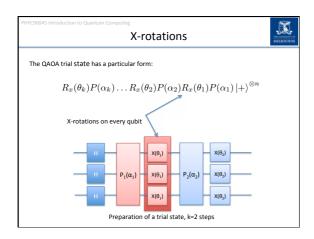




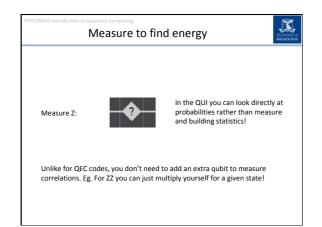


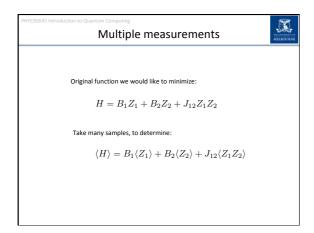


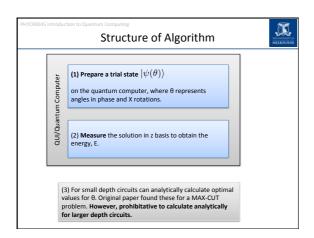


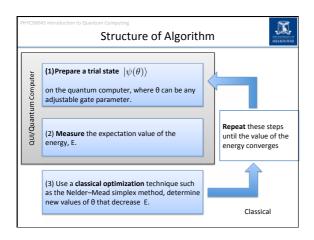


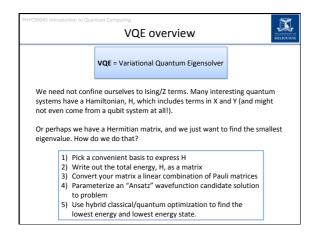
PHYC90045 Introduction to	Measuring the energy
The Hamiltonian	of a QUBO problem can be expressed as a sum of several terms:
	Ising coupling local "field" $H = \sum_{i \neq j} J_{ij} Z_i Z_j + \sum_i B_i Z_i$
For example:	$H = B_1 Z_1 + B_2 Z_2 + J_{12} Z_1 Z_2$
and the choice o	mation to the ground state (depending on the number of steps, k, f angles) will have a higher probability of measuring the minimum lowest energy state.

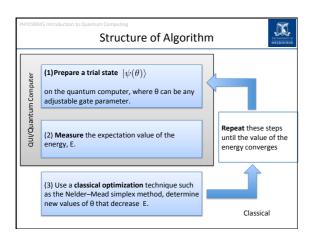


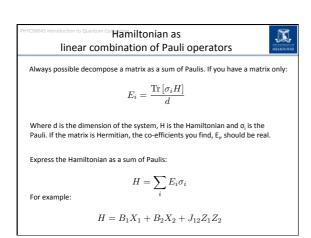












# Find the expectation value of H

We can express the expectation value of the energy as a sum of expectation values of the Paulis:

$$\langle H \rangle = \sum_{i} E_{i} \langle \sigma_{i} \rangle$$

For our example:

$$\langle H \rangle = B_1 \langle X_1 \rangle + B_2 \langle X_2 \rangle + J_{12} \langle Z_1 Z_2 \rangle$$

For a given trial state, these can be found directly from experiment (or through the QUI)

## Reminder: How to measure Paulis



Measure X:



Measure Y:





In the QUI you record probabilities rather than measure and building statistics!

Measure Z:



Unlike for QEC codes, you don't need to add an extra qubit to measure correlations. Eg. For ZZ you can just multiply yourself given the state which is  $\,$ measured!

### VQE: Jordan Wigner Transformation



A classic of physics from 1928, by Jordan and Wigner. You've got a system of qubits. You want to use it to simulate fermions (eg. someone's given you a chemistry problem involving electrons). Electrons do not behave the same as qubits.

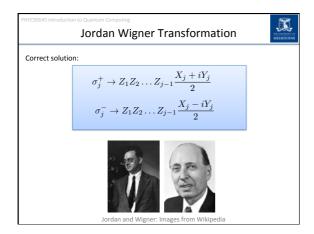
How do we do this?

$$\sigma_j^+ \to \frac{X_j + iY_j}{2} = \left[ \begin{array}{cc} 0 & 0 \\ 1 & 0 \end{array} \right]_j = f_j^+ \quad \text{Create a fermion at the j$^{\text{th}}$ site?}$$

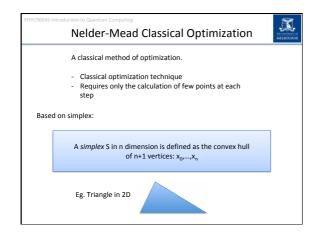
$$\sigma_j^- \to \frac{X_j - iY_j}{2} = \left[ \begin{array}{cc} 0 & 1 \\ 0 & 0 \end{array} \right]_j = f_j^- \quad \text{Destroy a fermion at the j$^{\text{th}}$ site?}$$

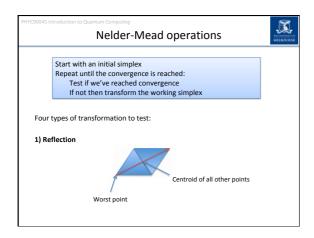
This is close but **WRONG!** The commutation relations between different sites are wrong (fermions anti-commute).  $\begin{bmatrix} f & f \end{bmatrix} = 0$ 

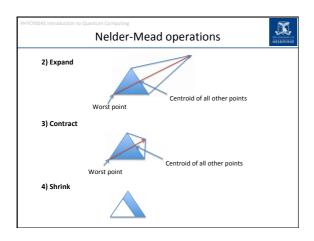
$$[f_j, f_k] = 0$$

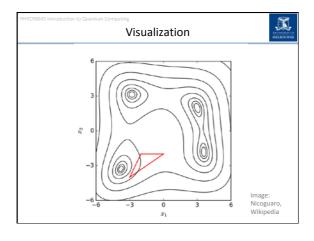


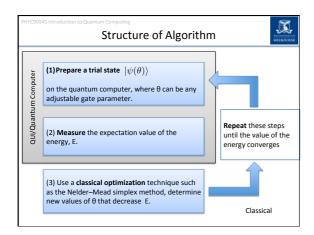
Pricking a VQE Ansatz	ELECCIANI
QAOA State	
Use a combination of X rotations and phase rotations.	
Adiabatic Methods	
Slowly vary the Hamiltonian, in a parameterized way, to obtain an approximation to the ground state. $\\$	
Coupled Cluster Methods	
First start in a (often unentangled) reference state, which can be calculated classically, eg. with mean field methods.	
Consider successively more complicated perturbations away from this reference state: First we consider just (parameterized) single qubit rotations away from the reference state. Then we consider unitaries with both single qubit rotations and two body interactions, then one, two and three body interactions. As with QAOA we can consider several rounds of interaction.	

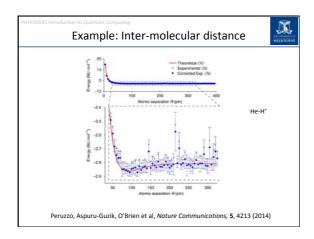












Week 10	MELBOLI
Lecture 19 Quantum Approximate Optimization Algorithm (QAOA),	
Variational Quantum Eigensolver (VQE), classical feedback	(
Lecture 20	
Exponentials, and Quantum Optimization	
Lab 10	