

# Documentation

Lukas Unglehrt

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## 1 Shallow water equations

### 1.1 Conservation form

The shallow water equations in conservation form are

$$\frac{\partial h}{\partial t} + \frac{\partial(hu)}{\partial x} + \frac{\partial(hv)}{\partial y} = 0 \quad (1a)$$

$$\frac{\partial(hu)}{\partial t} + \frac{\partial(hu^2)}{\partial x} + \frac{\partial(huv)}{\partial y} = -\frac{g}{2} \frac{\partial h^2}{\partial x} - gh \frac{\partial z}{\partial x} - \frac{h}{\rho} \tau_{b,x} \quad (1b)$$

$$\frac{\partial(hv)}{\partial t} + \frac{\partial(huv)}{\partial x} + \frac{\partial(hv^2)}{\partial y} = -\frac{g}{2} \frac{\partial h^2}{\partial y} - gh \frac{\partial z}{\partial y} - \frac{h}{\rho} \tau_{b,y} \quad (1c)$$

We abbreviate this system as

$$\frac{\partial \mathbf{U}}{\partial t} + \frac{\partial \mathbf{F}(\mathbf{U})}{\partial x} + \frac{\partial \mathbf{G}(\mathbf{U})}{\partial y} = \begin{bmatrix} 0 \\ -gh \frac{\partial z}{\partial x} - \frac{h}{\rho} \tau_{b,x} \\ -gh \frac{\partial z}{\partial y} - \frac{h}{\rho} \tau_{b,y} \end{bmatrix} \quad (2a)$$

with the conservative variables

$$\mathbf{U} = \begin{bmatrix} h \\ hu \\ hv \end{bmatrix} \quad (2b)$$

and the fluxes

$$\mathbf{F}(\mathbf{U}) = \begin{bmatrix} hu \\ hu^2 + \frac{1}{2}gh^2 \\ huv \end{bmatrix} = \begin{bmatrix} U_2 \\ \frac{U_2^2}{U_1} + \frac{1}{2}gU_1^2 \\ \frac{U_2 U_3}{U_1} \end{bmatrix} \quad (2c)$$

and

$$\mathbf{G}(\mathbf{U}) = \begin{bmatrix} hu \\ huv \\ hv^2 + \frac{1}{2}gh^2 \end{bmatrix} = \begin{bmatrix} U_3 \\ \frac{U_2 U_3}{U_1} \\ \frac{U_3^2}{U_1} + \frac{1}{2}gU_1^2 \end{bmatrix}. \quad (2d)$$

### 1.2 Quasi-linear form

Using the chain rule of differentiation, we can rewrite the system as

$$\frac{\partial \mathbf{U}}{\partial t} + \frac{\partial \mathbf{F}(\mathbf{U})}{\partial \mathbf{U}} \frac{\partial \mathbf{U}}{\partial x} + \frac{\partial \mathbf{G}(\mathbf{U})}{\partial \mathbf{U}} \frac{\partial \mathbf{U}}{\partial y} = \begin{bmatrix} 0 \\ -gh \frac{\partial z}{\partial x} - \frac{h}{\rho} \tau_{b,x} \\ -gh \frac{\partial z}{\partial y} - \frac{h}{\rho} \tau_{b,y} \end{bmatrix} \quad (3a)$$

where the flux Jacobians result as

$$\frac{\partial \mathbf{F}(\mathbf{U})}{\partial \mathbf{U}} = \begin{bmatrix} 0 & 1 & 0 \\ -\frac{U_2^2}{U_1^2} + gU_1 & 2\frac{U_2}{U_1} & 0 \\ -\frac{U_2 U_3}{U_1^2} & \frac{U_3}{U_1} & \frac{U_2}{U_1} \end{bmatrix} = \begin{bmatrix} 0 & 1 & 0 \\ gh - u^2 & 2u & 0 \\ -uv & v & u \end{bmatrix} \quad (3b)$$

and

$$\frac{\partial \mathbf{G}(\mathbf{U})}{\partial \mathbf{U}} = \begin{bmatrix} 0 & 0 & 1 \\ -\frac{U_2 U_3}{U_1^2} & \frac{U_3}{U_1} & \frac{U_2}{U_1} \\ -\frac{U_3^2}{U_1^2} + gU_1 & 0 & 2\frac{U_3}{U_1} \end{bmatrix} = \begin{bmatrix} 0 & 0 & 1 \\ -uv & v & u \\ gh - v^2 & 0 & 2v \end{bmatrix}. \quad (3c)$$

We can find out about the speed and direction of information propagation by computing the eigenvalue decomposition of the flux Jacobians. The eigenvalue decompositions read with  $c = \sqrt{gh}$ :

$$\frac{\partial \mathbf{F}(\mathbf{U})}{\partial \mathbf{U}} = \mathbf{V}^{-1} \begin{bmatrix} u & 0 & 0 \\ 0 & u + c & 0 \\ 0 & 0 & u - c \end{bmatrix} \mathbf{V} \text{ with } \mathbf{V} = \begin{bmatrix} 0 & 1 & 1 \\ 0 & u + c & u - c \\ 1 & v & v \end{bmatrix} \quad (4a)$$

and

$$\frac{\partial \mathbf{G}(\mathbf{U})}{\partial \mathbf{U}} = \mathbf{W}^{-1} \begin{bmatrix} v & 0 & 0 \\ 0 & v + c & 0 \\ 0 & 0 & v - c \end{bmatrix} \mathbf{W} \text{ with } \mathbf{W} = \begin{bmatrix} 0 & 1 & 1 \\ 1 & u & u \\ 0 & v + c & v - c \end{bmatrix} \quad (4b)$$

## 2 A Finite-Volume discretization

### 2.1 The discrete conservation law

For discretization, we choose an equidistant Cartesian grid with a collocated arrangement of variables. As primitive fields, we choose the conservative variables  $h$ ,  $hu$  and  $hv$ . We integrate the shallow water equations (2) over a cell  $[x_i - \frac{\Delta x}{2}; x_i + \frac{\Delta x}{2}] \times [y_j - \frac{\Delta y}{2}; y_j + \frac{\Delta y}{2}]$  and we obtain:

$$\begin{aligned} & \int_{x_i - \frac{\Delta x}{2}}^{x_i + \frac{\Delta x}{2}} \int_{y_j - \frac{\Delta y}{2}}^{y_j + \frac{\Delta y}{2}} \frac{\partial \mathbf{U}}{\partial t} dy dx + \int_{y_j - \frac{\Delta y}{2}}^{y_j + \frac{\Delta y}{2}} [\mathbf{F}(\mathbf{U})]_{x_i - \frac{\Delta x}{2}}^{x_i + \frac{\Delta x}{2}} dy + \int_{x_i - \frac{\Delta x}{2}}^{x_i + \frac{\Delta x}{2}} [\mathbf{G}(\mathbf{U})]_{y_j - \frac{\Delta y}{2}}^{y_j + \frac{\Delta y}{2}} dx \\ &= \int_{x_i - \frac{\Delta x}{2}}^{x_i + \frac{\Delta x}{2}} \int_{y_j - \frac{\Delta y}{2}}^{y_j + \frac{\Delta y}{2}} \begin{bmatrix} 0 \\ -gh \frac{\partial z}{\partial x} - \frac{h}{\rho} \tau_{b,x} \\ -gh \frac{\partial z}{\partial y} - \frac{h}{\rho} \tau_{b,y} \end{bmatrix} dy dx. \end{aligned} \quad (5)$$

We now define the volume averaged quantities

$$\bar{\mathbf{U}}_{i,j} = \frac{1}{\Delta x \Delta y} \int_{x_i - \frac{\Delta x}{2}}^{x_i + \frac{\Delta x}{2}} \int_{y_j - \frac{\Delta y}{2}}^{y_j + \frac{\Delta y}{2}} \mathbf{U} dy dx, \quad (6)$$

for which we obtain a system of ordinary differential equations

$$\begin{aligned} & \frac{d\bar{\mathbf{U}}_{i,j}}{dt} \Delta x \Delta y + \int_{y_j - \frac{\Delta y}{2}}^{y_j + \frac{\Delta y}{2}} [\mathbf{F}(\mathbf{U})]_{x_i - \frac{\Delta x}{2}}^{x_i + \frac{\Delta x}{2}} dy + \int_{x_i - \frac{\Delta x}{2}}^{x_i + \frac{\Delta x}{2}} [\mathbf{G}(\mathbf{U})]_{y_j - \frac{\Delta y}{2}}^{y_j + \frac{\Delta y}{2}} dx \\ &= \int_{x_i - \frac{\Delta x}{2}}^{x_i + \frac{\Delta x}{2}} \int_{y_j - \frac{\Delta y}{2}}^{y_j + \frac{\Delta y}{2}} \begin{bmatrix} 0 \\ -gh \frac{\partial z}{\partial x} - \frac{h}{\rho} \tau_{b,x} \\ -gh \frac{\partial z}{\partial y} - \frac{h}{\rho} \tau_{b,y} \end{bmatrix} dy dx. \end{aligned} \quad (7)$$

Also, we define the numerical fluxes

$$\tilde{\mathbf{F}}_{i+\frac{1}{2},j} = \frac{1}{\Delta y} \int_{y_j - \frac{\Delta y}{2}}^{y_j + \frac{\Delta y}{2}} [\mathbf{F}(\mathbf{U})]_{x_i + \frac{\Delta x}{2}} dy \quad (8a)$$

$$\tilde{\mathbf{G}}_{i,j+\frac{1}{2}} = \frac{1}{\Delta x} \int_{x_i - \frac{\Delta x}{2}}^{x_i + \frac{\Delta x}{2}} [\mathbf{G}(\mathbf{U})]_{y_j + \frac{\Delta y}{2}} dx, \quad (8b)$$

so our system results as

$$\begin{aligned} & \frac{d\bar{\mathbf{U}}_{i,j}}{dt} \Delta x \Delta y + \left( \tilde{\mathbf{F}}_{i+\frac{1}{2},j} - \tilde{\mathbf{F}}_{i-\frac{1}{2},j} \right) \Delta y + \left( \tilde{\mathbf{G}}_{i,j+\frac{1}{2}} - \tilde{\mathbf{G}}_{i,j-\frac{1}{2}} \right) \Delta x \\ &= \int_{x_i - \frac{\Delta x}{2}}^{x_i + \frac{\Delta x}{2}} \int_{y_j - \frac{\Delta y}{2}}^{y_j + \frac{\Delta y}{2}} \begin{bmatrix} 0 \\ -gh \frac{\partial z}{\partial x} - \frac{h}{\rho} \tau_{b,x} \\ -gh \frac{\partial z}{\partial y} - \frac{h}{\rho} \tau_{b,y} \end{bmatrix} dy dx. \end{aligned} \quad (9)$$

This equation is a discrete conservation law for the quantities  $h$ ,  $hu$  and  $hv$  on the cell  $(i, j)$ . To obtain a numerical scheme, we need to approximate the source terms and the fluxes in terms of the cell-averaged quantities  $\bar{\mathbf{U}}_{i,j}$  and choose a time integration scheme.

## 2.2 Roe's scheme

The idea of ? is to approximate the non-linear conservation law with a