

# Forceconstant Calculator

The force constant calculator `FCCalculator` implements properties derived from a harmonic model for the potential-energy surface,

$$H(\mathbf{U}) = \frac{1}{2} \sum_{IJ, \alpha\beta} \Phi_{IJ}^{\alpha\beta} U_I^\alpha U_J^\beta, \quad (1)$$

where  $\Phi_{IJ}^{\alpha\beta}$  denote the force constant matrix elements

$$\Phi_{IJ}^{\alpha\beta} = \left. \frac{\partial^2 E}{\partial R_I^\alpha \partial R_J^\beta} \right|_{\mathbf{R}^0}, \quad (2)$$

i.e., the second-order derivative of the potential energy  $E$  w.r.t. to the atomic positions  $\mathbf{R} = \{R_I^\alpha\}$  at a given configuration  $\mathbf{R}^0$ . Here,  $I, J$  are the atom labels and  $\alpha, \beta$  denote Cartesian components.  $\mathbf{U} = \{U_I^\alpha\}$  denotes the atomic displacements w.r.t. to the reference position  $\mathbf{R}^0$ , i.e.,

$$\mathbf{U} = \mathbf{R} - \mathbf{R}^0. \quad (3)$$

## Properties derived from the harmonic model

### Forces

Forces are given as a derivative of the potential energy and can therefore be obtained analytically from Eq. (1):

$$F_I^\alpha = -\frac{\partial H}{\partial R_I^\alpha} = -\sum_{J, \beta} \Phi_{IJ}^{\alpha\beta} U_J^\beta. \quad (4)$$

This equations can be written as a matrix-vector product

$$\mathbf{F} = -P\mathbf{U}, \quad (5)$$

where  $\mathbf{F}$  and  $\mathbf{U}$  are the  $3N \times 1$  vectors representing atomic forces and displacements, and  $P$  is the  $3N \times 3N$  forceconstant matrix, where  $N$  is the number of atoms in the simulation. In python, this is achieved by

```
1 | forces = -(fc @ displacements.flatten()).reshape(displacements.shape)
```

### Potential energy

The energy is given by computing the displacements  $\mathbf{U}$  according to Eq. (3), and performing the matrix product given by Eq. (1), i.e., by scalar multiplication of the displacements  $\mathbf{U}$  with the forces  $\mathbf{F}$ ,

$$E^{\text{pot}} = -\frac{1}{2} \mathbf{F} \cdot \mathbf{U}. \quad (6)$$

In python:

```

1 # energies: [N, 3] * [N, 3] -> [N]
2 energies = -(displacements * forces).sum(axis=1) / 2
3 # energy: [N] -> [1]
4 energy = energies.sum()

```

As a byproduct, atomic energy contributions (`energies`) are computed, given by

$$E_I^{\text{pot}} = \frac{1}{2} \sum_{J,\alpha\beta} \Phi_{IJ}^{\alpha\beta} U_I^\alpha U_J^\beta . \quad (7)$$

## Stress

Stress is given in terms of a strain derivative of the potential energy

$$\sigma^{\alpha\beta} = \frac{1}{V} \frac{\partial H}{\partial \epsilon^{\alpha\beta}} , \quad (8)$$

where  $\epsilon^{\alpha\beta}$  describes a rotation-free, homogeneous straining of the atomic positions. The strain derivatives of the atomic displacements read [Knuth2015]

$$\frac{\partial U_I^\gamma}{\partial \epsilon^{\alpha\beta}} = \delta^{\alpha\gamma} U_I^\beta . \quad (9)$$

In turn, the harmonic stress tensor reads

$$\sigma^{\alpha\beta} = \frac{1}{V} \frac{\partial H}{\partial \epsilon^{\alpha\beta}} = \frac{1}{V} \underbrace{\sum_{IJ,\gamma} \Phi_{IJ}^{\alpha\gamma} U_I^\beta U_J^\gamma}_{\sigma_{\text{HA}}^{\alpha\beta}} + \sigma_{\text{QHA}}^{\alpha\beta} , \quad (10)$$

where the purely harmonic contribution is given by

$$\sigma_{\text{HA}}^{\alpha\beta} = \frac{1}{V} \sum_{IJ,\gamma} \Phi_{IJ}^{\alpha\gamma} U_I^\beta U_J^\gamma , \quad (11)$$

and the *quasi*harmonic contribution  $\sigma_{\text{QHA}}^{\alpha\beta}$  stems from the strain derivative of the force constants. The quasiharmonic contribution can be approximated by

$$\frac{1}{V} \frac{\partial H}{\partial \epsilon^{\alpha\beta}} \approx 3 \left( \frac{\partial H}{\partial V} \right)^{\alpha\beta} , \quad (12)$$

which is strictly true only in cubic systems. By denoting the quasiharmonic force constants, i.e., the volume derivative  $\partial \Phi / \partial V$  as  $\Phi'$ , we have

$$\sigma_{\text{QHA}}^{\alpha\beta} = \frac{3}{2} \sum_{IJ,\gamma} \Phi_{IJ}'^{\alpha\gamma} U_I^\beta U_J^\gamma . \quad (13)$$

## Virial stress

The virial stress in periodic systems is given by [Louwerse2006]

$$\sigma_{\text{vir}}^{\alpha\beta} = -\frac{1}{2V} \sum_{IJ} (R_I^\alpha - R_J^\alpha) F_{IJ}^\beta , \quad (14)$$

where  $F_{IJ}^\beta$  denotes the pairwise force between atom  $I$  and  $J$  fulfilling Newton's 3rd law,

$$F_{IJ}^\alpha = -F_{JI}^\alpha , \quad (15)$$

and sums to the atomic force  $F_I^\beta$ ,

$$F_I^\beta = \sum_J F_{IJ}^\beta . \quad (16)$$

This pairwise force is given in the harmonic approximation as

$$F_{IJ}^\beta = -\Phi_{IJ}^{\beta\gamma} U_J^\gamma + \Phi_{JI}^{\beta\gamma} U_I^\gamma , \quad (17)$$

where summing over repeated indices is implied. This definition complies with both Eq. (15) and (16), by using the acoustic sum rule  $\sum_J \Phi_{IJ}^{\alpha\beta} = 0$ . Using this pairwise force in the definition of virial stress, Eq. (14), we have

$$\sigma_{\text{vir}}^{\alpha\beta} = -\frac{1}{V} \sum_{IJ} (R_I^\alpha - R_J^\alpha) \Phi_{IJ}^{\beta\gamma} U_J^\gamma \quad (18)$$

$$\begin{aligned} &= -\frac{1}{V} \sum_{IJ} (R_I^{0\alpha} - R_J^{0\alpha}) \Phi_{IJ}^{\beta\gamma} U_J^\gamma \\ &\quad - \frac{1}{V} \sum_{IJ} (U_I^\alpha - U_J^\alpha) \Phi_{IJ}^{\beta\gamma} U_J^\gamma , \end{aligned} \quad (19)$$

where in the second row  $\mathbf{R} = \mathbf{R}^0 + \mathbf{U}$  was written out. However, the first contribution can be written as

$$V\sigma_{\text{vir},0}^{\alpha\beta} = -\sum_{IJ} (R_I^{0\alpha} - R_J^{0\alpha}) \Phi_{IJ}^{\beta\gamma} U_J^\gamma = \underbrace{\sum_{IJ} R_J^{0\alpha} \Phi_{IJ}^{\beta\gamma} U_J^\gamma}_{\text{acoustic sum rule}} - \sum_{IJ} R_I^{0\alpha} \Phi_{IJ}^{\beta\gamma} U_J^\gamma \equiv \sum_I R_I^{0\alpha} F_I^\beta , \quad (20)$$

which does not contribute to thermal averages because  $\langle F \rangle = 0$ .

## Remark on finite supercells

In finite supercells, care must be taken when computing the pair distance  $\mathbf{R}_I - \mathbf{R}_J$ . They should obey a minimum image convention (MIC). Since MIC is a non-continuous map as it involves wrapping of atoms by finite displacements, continuity of the map must be ensured, e.g., by averaging over MIC-equivalent pairs of atoms.

## Atomic virials

The harmonic virial stress can be decomposed into atomic contributions by writing

$$\sigma_{\text{vir}}^{\alpha\beta} = \sum_I \sigma_{\text{vir},I}^{\alpha\beta} , \quad (21)$$

with

$$\sigma_{\text{vir},I}^{\alpha\beta} = -\frac{1}{V} \sum_{IJ} (R_I^\alpha - R_J^\alpha) \Phi_{IJ}^{\beta\gamma} U_J^\gamma . \quad (22)$$

This quantity is related to the *virial heat flux*  $\mathbf{J}_{\text{vir}}$  by multiplying with the atomic velocities  $\dot{\mathbf{U}}$ ,

$$\mathbf{J}_{\text{vir}}^\alpha = \sum_{I,\gamma} \sigma_{\text{vir},I}^{\alpha\gamma} \dot{U}_I^\gamma . \quad (23)$$

## Literature

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- **Knuth2015**: F. Knuth, C. Carbogno, V. Atalla, V. Blum, and M. Scheffler, Comput. Phys. Commun. **190**, 33 (2015).
- **Louwerse2006**: M. J. Louwerse and E. J. Baerends, Chem. Phys. Lett. **421**, 138 (2006).