Quantum simulation on a random tensor network

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1 Differentiating single step time evolution

The m site Rydberg Hamiltonian is

$$H_{\text{Rydberg}} = \sum_{i,j=1,i>j}^{m} \frac{C}{|r_i - r_j|^6} n_i n_j + \Omega(t) \sum_{i=1}^{m} \frac{1}{2}\sigma_i^x + \Delta(t) \sum_{i=1}^{m} n_i \qquad (1)$$

For simplicity, we consider the following general representation of a time-space dependent Hamiltonian with k terms

$$H = \sum_{k=1}^{K} c_k O_k \qquad (2)$$

where c_k can be dependent on a set of parameters like locations $r_1, r_2, \ldots, r_m$, and pulses $\Omega(t)$ and $\Delta(t)$.

1.1 The ODE version

In each step of the ODE solver, it performs the following update

$$|\psi'\rangle = (1 - iH\Delta t)|\psi\rangle \qquad (3)$$

We derive the backward rules for the gradients by inspecting the following equations

$$\begin{aligned}
\overline{\mathcal{L}}\delta\mathcal{L} &= \overline{|\psi'\rangle} \circ \delta|\psi'\rangle \\
&= \sum_k \overline{c_k}\delta c_k + \overline{|\psi\rangle} \circ \delta|\psi\rangle + \overline{\Delta t}\delta\Delta t
\end{aligned} \qquad (4)$$

where $\circ$ is the Hadamard product applied on real numbers, note a complex number in computer is composed of two real numbers. The above euqations has a more elegant linear algebra version as the following.

$$\begin{aligned}
\overline{\mathcal{L}}\delta\mathcal{L} &= \overline{\langle\psi'|}\delta|\psi'\rangle \\
&= \sum_k \overline{c_k}\delta c_k + \overline{\langle\psi|}\delta|\psi\rangle + \overline{\Delta t}\delta\Delta t
\end{aligned} \qquad (5)$$

where we have used $\langle\psi|$ to represent the hermitian conjugate of $|\psi\rangle$.

$$\delta|\psi'\rangle = -i\sum_k \delta c_k O_k \Delta t|\psi\rangle - iH\delta\Delta t|\psi\rangle + (1 - iH\Delta t)\delta|\psi\rangle \tag{6}$$

By observing Eq. (**??**) and Eq. (**??**), one can see

$$\overline{\langle\psi|} = \overline{\langle\psi'|}(1 - iH\Delta t) \tag{7}$$

$$\overline{c_k} = \Re\left[-i\Delta t\overline{\langle\psi'|}O_k|\psi\rangle\right] \tag{8}$$

$$\overline{\Delta t} = \Re\left[-i\overline{\langle\psi'|}H|\psi\rangle\right] \tag{9}$$

After a step, a normalization procedure might be called on the wave functions, this is trivial so that we do not discuss it at this stage.

1.2 The `expmv` version

To differentiate the time evolution directly, one can use the Taylor expansion

$$|\psi'\rangle = e^{-iHt}|\psi\rangle$$
$$= \sum_{n=0}^{\infty} \frac{(-it)^n H^n}{n!}|\psi\rangle \tag{10}$$

Similarly, we have

$$\delta|\psi'\rangle = e^{-iHt}\delta|\psi\rangle + \sum_{n=0}^{\infty} \frac{(-it)^n \delta(H^n)}{n!}|\psi\rangle + \left(e^{-iH(t+\delta t)} - e^{-iHt}\right)|\psi\rangle \tag{11}$$

$$\overline{\langle\psi|} = \overline{\langle\psi|}e^{-iHt}\delta \tag{12}$$

$$\overline{t} = \overline{\langle\psi'|} - iHe^{-iHt}|\psi\rangle \tag{13}$$

$$\overline{c_k} = \sum_n \frac{(-it)^n}{n!} \sum_{p=0}^{n-1} \overline{\langle\psi'|}H^p O_k H^{n-p-1}|\psi\rangle \tag{14}$$

2 How to differentiate an ODE solver

2.1 The adjoint state method (or neural ODE)

Since the time evolution is reversible, one can reverse it by doing inverse time evolution (or the adjoint state method [**?**, **?**]).


Algorithm 1: The continuous adjoint state method

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input : parameters θ, initial time t_0, ending time t_n, final state s_n and its adjoint $\overline{s_n}$

output: $\overline{s_0}$, $\overline{\theta}$

1 function aug_dynamics$((s, a, \text{-}), t, \theta)$

2 $q = f(s, t, \theta)$ # the augmented dynamics

3   **return** $(q, -a^T \frac{\partial q}{\partial s}, -a^T \frac{\partial q}{\partial \theta})$

4 end

5 $S_0 = (s_n, \overline{s_n}, 0)$ # initial state of the augmented dynamics

6 $(s_0, \overline{s_0}, \overline{\theta}) = \text{ODESolve}(\text{aug\_dynamics}, S_0, \theta, t_n, t_0)$ # integrate the augmented dynamics in the reversed time

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2.2 The treeverse algorithm (or optimal checkpointing)

For the cases reversibility is not guarented, one can use treeverse algorithm [**?**, **?**]. Peter is a road maintainer. He is assigned a job by his male willed employer to paint a one-way road to blue in the reversed order (do not ask me why!). Lucky enough, Peter picked up some teleportation magics in Hogwarts. In the following illustration, we represent the road as N grids arranged in one dimension.
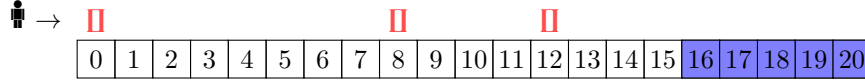
Peter can take the following actions.

1. paint one block s_i if he is in i-th block and s_{i+1} is painted,

2. move in one direction, i.e. $i \to i + 1$,

3. setup a teleportation gate (marked with red symbols) at where he locates. He can at most create $\delta \leq N$ gates at the same time. When this upper limit is reached, he must distroy an existing teleportation gate to create a new one.

4. teleport himself to any existing teleportation gate.

Given $N = 10000$, $\delta = 10$, can you please help Peter design a scheme so that he can drive the least?

The optimal solution is described by the binomial function $\eta(d, t) \equiv \frac{(d+t)!}{d!t!}$.

1. In the first forward sweep, Peter uses teleportation gates to devide the road into δ segments of size $\eta(d, \tau - 1), d = \delta, \delta - 1, \ldots, 1$ plus one extra grid. Note $\sum_{d=0}^{\delta} \eta(d, \tau - 1) = \eta(\delta, \tau)$.

2. Sweep over the last block, use up the remaining checkpoints and paint the last grid. If Peter has paint all grids in a block, destroy the last checkpoint.

3. In the t-th forward sweep of a block, Peter uses teleportation gates to devide the block into d segments of size $\eta(k, \tau - t), k = d, d-1, \ldots, 1$ plus one extra grid, where d is equal to δ minus the number of checkpoints already used.

References

[1] TreeverseAlgorithm.jl. https://github.com/GiggleLiu/TreeverseAlgorithm.jl.

[2] Ricky T. Q. Chen, Yulia Rubanova, Jesse Bettencourt, and David K Duvenaud. Neural ordinary differential equations. In *Advances in Neural Information Processing Systems*, volume 31. Curran Associates, Inc., 2018.

[3] Andreas Griewank. Achieving logarithmic growth of temporal and spatial complexity in reverse automatic differentiation. *Optimization Methods and software*, 1(1):35–54, 1992.

[4] R.-E. Plessix. A review of the adjoint-state method for computing the gradient of a functional with geophysical applications. *Geophysical Journal International*, 167(2):495–503, 11 2006.

Algorithm 2: The Treeverse algorithm

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input : State cache $S = \{0 : s_0\}$, the initial adjoint $\overline{s_n} \equiv \frac{\partial \mathcal{L}}{\partial s_n}$, the maximum number of checkpoints δ, the number of scan τ, the starting location of current block $\beta = 0$, the end point of current block $\phi = n$, and the bisection point of current block $\sigma = 0$

output: Back-propagated adjoint $\overline{s_0} \equiv \frac{\partial \mathcal{L}}{\partial s_0}$

1 function treeverse(S, $\overline{s_\phi}$, δ, τ, β, σ, ϕ)

2 | **if** $\sigma > \beta$ **then**

3 | | $\delta = \delta - 1$

4 | | $s = S[\beta]$ # `load initial state` s_β

5 | | **for** $j = \beta, \beta + 1, ..., \sigma - 1$ **do**

6 | | | $s_{j+1} = f_j(s_j)$ # `compute` s_σ

7 | | **end**

8 | | $S[\sigma] = s_\sigma$

9 | **end**

10 | # let κ be the division point, call the treeverse algorithm recursively

11 | **while** $\tau > 0$ **and** $\kappa = \text{mid}(\delta, \tau, \sigma, \phi) < \phi$ **do**

12 | | $\overline{s_\kappa} = $ treeverse(S, $\overline{s_\phi}$, δ, τ, σ, κ, ϕ)

13 | | $\tau = \tau - 1$

14 | | $\phi = \kappa$

15 | **end**

16 | $\overline{s_\sigma} = \overline{f_\sigma}(\overline{s_{\sigma+1}}, s_\sigma)$ # `back propagate the gradient`

17 | **if** $\sigma > \beta$ **then**

18 | | remove($S[\sigma]$)   # `remove state` s_σ `from cache`

19 | **end**

20 | **return** $\overline{s_\sigma}$

21 end

22 function mid(δ, τ, σ, ϕ)

      # choose the bisection point

23 | $\kappa = \lceil (\delta\sigma + \tau\phi)/(\tau + \delta) \rceil$

24 | **if** $\kappa \geq \phi$ **and** $\delta > 0$ **then**

25 | | $\kappa = \max(\sigma + 1, \phi - 1)$

26 | **end**

27 end