

# 计算物理第八次作业

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## 1 题目 1: Poisson Equation

### 1.1 题目描述

Consider the Poisson equation

$$\nabla^2 \phi(x, y) = -\frac{\rho(x, y)}{\epsilon_0} \quad (1)$$

from electrostatics on a rectangular geometry with  $x \in [0, L_x]$  and  $y \in [0, L_y]$ . Write a program that solves this equation using the relaxation method. Test your program with:

$$(1) \quad \rho(x, y) = 0, \phi(0, y) = \phi(L_x, y) = \phi(x, 0) = 0, \phi(x, L_y) = 1V, L_x = 1m, L_y = 1.5m$$

$$(2) \quad \frac{\rho(x, y)}{\epsilon_0} = 1V/m^2, \phi(0, y) = \phi(L_x, y) = \phi(x, L_y) = \phi(x, 0) = 0, L_x = L_y = 1m$$

### 1.2 程序描述

本程序使用 Gauss-Seidel 方法求解二维静电势分布的 Poisson 方程 (Eq.1)。在空间上进行离散化, 得到离散的 Poisson 方程为

$$u_{i,j} = \frac{1}{4}(u_{i,j-1} + u_{i-1,j} + u_{i,j+1} + u_{i+1,j}) - \frac{h^2}{4}\rho_{i,j} \quad (2)$$

其中  $h$  为  $x, y$  方向实空间离散化的步长。Gauss-Seidel 方法每次迭代求解每个格点上取值的公式为

$$u_{i,j} = u_{i,j} + \frac{\omega}{4}(-\rho_{i,j} \cdot h^2 + u_{i,j-1} + u_{i-1,j} - 4u_{i,j} + u_{i,j+1} + u_{i+1,j}) \quad (3)$$

等号右边  $u$  的两个指标大于  $i, j$  的为上轮迭代的值, 小于  $i, j$  的为本轮迭代的值。对  $u$  矩阵在相邻两轮迭代间作差并求  $l-2$  范数, 当其小于指定值时认为算法收敛到了 Poisson 方程的解。

本程序源文件为 Poisson.py, 运行依赖 Python 第三方库 Numpy 和 Matplotlib。在终端进入当前目录, 使用命令 `python -u Poisson.py` 运行本程序。运行后绘制出如 Fig.1所示的两种电荷分布 + 边界条件下 Poisson 方程的解。

### 1.3 伪代码

求解 Eq.1的 Gauss-Seidel 方法的伪代码如 Alg.1所示。

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**Algorithm 1** Gauss-Seidel Method

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**Input:** Charge density distribution  $\rho$ , boundary condition  $\phi(x, 0), \phi(0, y), \phi(x, L_y), \phi(L_x, y)$ , tolerable error  $\epsilon$ , maximum iteration number  $M$ , real space discretization step  $h$ , and a initial guess  $u_0$  (with boundary condition satisfied).

**Output:** The solution of the 2-D Poisson equation.

```
1: normf  $\leftarrow$  0
2: norml  $\leftarrow$   $||u_0||$ 
3: iter_num  $\leftarrow$  0
4: while iter_num  $< M$  &  $|norm_l - norm_f| > \epsilon$  do
5:   normf  $\leftarrow$  norml
6:   for  $i \leftarrow 1, N_x - 1$  do
7:     for  $j \leftarrow 1, N_y - 1$  do
8:        $u_{i,j} \leftarrow u_{i,j} + \frac{\omega}{4}(-\rho_{i,j} \cdot h^2 + u_{i,j-1} + u_{i-1,j} - 4u_{i,j} + u_{i,j+1} + u_{i+1,j})$ 
9:     end for
10:  end for
11:  norml  $\leftarrow$   $||u||$ 
12:  iter_num  $\leftarrow$  iter_num + 1
13: end while
14: return  $u$ 
```

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## 1.4 输入输出示例

在  $h = 0.005, \epsilon = 1 \times 10^{-6}$  时, 程序输出两种电荷分布 + 边界条件下 Poisson 方程的解如 Fig.1所示。可见结果符合边界条件, 符合定性预期, 故知程序正确。

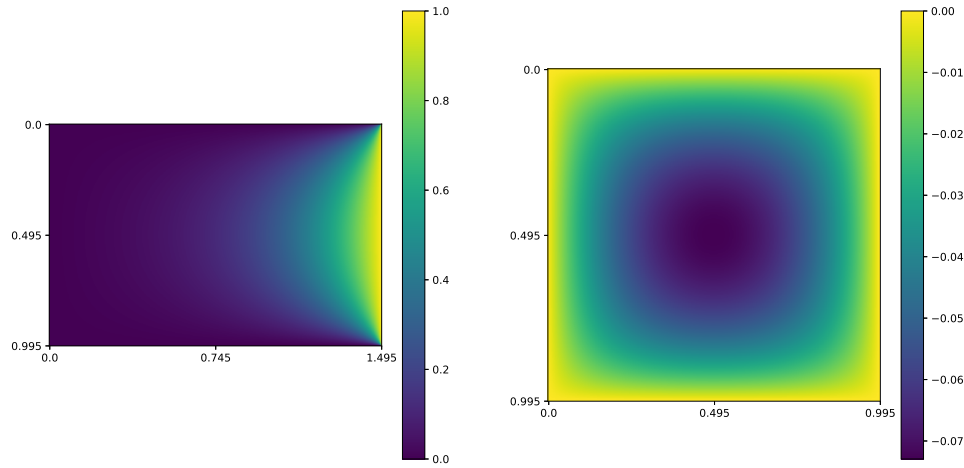


Figure 1: 两种电荷分布 + 边界条件下 Poisson 方程的解

## 2 题目 2: 求解含时 Schrodinger 方程

### 2.1 题目描述

Solve the time-dependent Schrodinger equation using both the Crank–Nicolson scheme and stable explicit scheme. Consider the one-dimensional case and test it by applying it to the problem of a square well with a Gaussian initial state coming in from the left.

$$\psi(x, 0) = \sqrt{\frac{1}{\pi}} e^{ik_0 x - \frac{(x - \xi_0)^2}{2}} \quad (4)$$

$$V(x) = \begin{cases} +\infty, & x < 0 \text{ or } x > a_3 \\ 0, & x < a_1 \text{ or } x > a_2 \\ -V_0, & a_2 < x < a_2 \end{cases}$$

### 2.2 程序描述

本程序分别通过 Crank Nicolson 方法和 Stable Explicit 方法求解含时 Schrodinger 方程。

$$i\hbar \frac{\partial \psi}{\partial t} = -\frac{\hbar^2}{2m} \frac{\partial^2 \psi}{\partial x^2} + V\psi$$

Crank Nicolson 方法

$$\begin{pmatrix} 1 + \frac{i\Delta t\hbar}{2m\Delta x^2} + \frac{i\Delta tV_j}{2\hbar} & -\frac{i\Delta t\hbar}{4m\Delta x^2} & & & \\ -\frac{i\Delta t\hbar}{4m\Delta x^2} & 1 + \frac{i\Delta t\hbar}{2m\Delta x^2} + \frac{i\Delta tV_j}{2\hbar} & -\frac{i\Delta t\hbar}{4m\Delta x^2} & & \\ & -\frac{i\Delta t\hbar}{4m\Delta x^2} & \ddots & \ddots & \\ & & \ddots & \ddots & -\frac{i\Delta t\hbar}{4m\Delta x^2} \\ & & & -\frac{i\Delta t\hbar}{4m\Delta x^2} & 1 + \frac{i\Delta t\hbar}{2m\Delta x^2} + \frac{i\Delta tV_j}{2\hbar} \end{pmatrix} \begin{pmatrix} \psi_0^{n+1} \\ \psi_1^{n+1} \\ \vdots \\ \psi_{N_x-1}^{n+1} \\ \psi_{N_x}^{n+1} \end{pmatrix} = \begin{pmatrix} 1 - \frac{i\Delta t\hbar}{2m\Delta x^2} - \frac{i\Delta tV_j}{2\hbar} & \frac{i\Delta t\hbar}{4m\Delta x^2} & & & \\ \frac{i\Delta t\hbar}{4m\Delta x^2} & 1 - \frac{i\Delta t\hbar}{2m\Delta x^2} - \frac{i\Delta tV_j}{2\hbar} & \frac{i\Delta t\hbar}{4m\Delta x^2} & & \\ & \frac{i\Delta t\hbar}{4m\Delta x^2} & \ddots & \ddots & \\ & & \ddots & \ddots & \frac{i\Delta t\hbar}{4m\Delta x^2} \\ & & & \frac{i\Delta t\hbar}{4m\Delta x^2} & 1 - \frac{i\Delta t\hbar}{2m\Delta x^2} - \frac{i\Delta tV_j}{2\hbar} \end{pmatrix} \begin{pmatrix} \psi_0^n \\ \psi_1^n \\ \vdots \\ \psi_{N_x-1}^n \\ \psi_{N_x}^n \end{pmatrix} \quad (5)$$

Stable Explicit 方法:

$$\begin{pmatrix} \psi_0^{n+1} \\ \psi_1^{n+1} \\ \vdots \\ \psi_{N_x-1}^{n+1} \\ \psi_{N_x}^{n+1} \end{pmatrix} = \begin{pmatrix} \psi_0^{n-1} \\ \psi_1^{n-1} \\ \vdots \\ \psi_{N_x-1}^{n-1} \\ \psi_{N_x}^{n-1} \end{pmatrix} + \begin{pmatrix} -\frac{2i\Delta t\hbar}{m\Delta x^2} - \frac{2i\Delta tV_j}{\hbar} & \frac{i\Delta t\hbar}{m\Delta x^2} & & & \\ \frac{i\Delta t\hbar}{m\Delta x^2} & -\frac{2i\Delta t\hbar}{m\Delta x^2} - \frac{2i\Delta tV_j}{\hbar} & \frac{i\Delta t\hbar}{m\Delta x^2} & & \\ & \frac{i\Delta t\hbar}{m\Delta x^2} & \ddots & \ddots & \\ & & \ddots & \ddots & \frac{i\Delta t\hbar}{m\Delta x^2} \\ & & & \frac{i\Delta t\hbar}{m\Delta x^2} & -\frac{2i\Delta t\hbar}{m\Delta x^2} - \frac{2i\Delta tV_j}{\hbar} \end{pmatrix} \begin{pmatrix} \psi_0^n \\ \psi_1^n \\ \vdots \\ \psi_{N_x-1}^n \\ \psi_{N_x}^n \end{pmatrix} \quad (6)$$

本程序源文件为 Schrodinger.py, 运行依赖 Python 第三方库 Numpy、math 和 Matplotlib. 在终端进入当前目录, 使用命令 `python -u Schrodinger.py` 运行本程序. 运行后输出 Fig.2所示图像

## 2.3 伪代码

Crank Nicolson 方法的伪代码如 Alg.2所示, Stable Explicit Scheme 的伪代码如 Alg.3所示,

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**Algorithm 2** Crank Nicolson method

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**Input:** The initial wavefunction  $\psi_0$ , the time step  $\Delta t$ , the time step number  $N_t$

**Output:** The time evolution of  $\psi$

```

1: for  $n \leftarrow 1, 2, \dots, N_t$  do
2:    $\psi_n \leftarrow (1 + \frac{i\Delta t \hat{H}}{2\hbar})^{-1} (1 - \frac{i\Delta t \hat{H}}{2\hbar}) \psi_{n-1}$ 
3:    $\psi_n(0), \psi_n(a) \leftarrow 0$ 
4: end for

```

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**Algorithm 3** Stable Explicit Scheme

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**Input:** The initial wavefunction  $\psi_0$ , the time step  $\Delta t$ , the time step number  $N_t$

**Output:** The time evolution of  $\psi$

```

1: for  $n \leftarrow 1, 2, \dots, N_t$  do
2:    $\psi_n \leftarrow \psi_{n-2} - 2i\hat{H}\Delta t\psi_{n-1}$ 
3:    $\psi_n(0), \psi_n(a) \leftarrow 0$ 
4: end for

```

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## 2.4 输入输出示例

程序中令  $m = \hbar = 1, V_0 = 10, a_1 = 75, a_2 = 125, a_3 = 200, k = 5, \xi_0 = 40$ , 两种方法的结果分别如 Fig.2所示。可以看到初始高斯波包在遇到势阱时分成反射和透射两部分, 符合定性预期, 两种方法结果也大致相符。

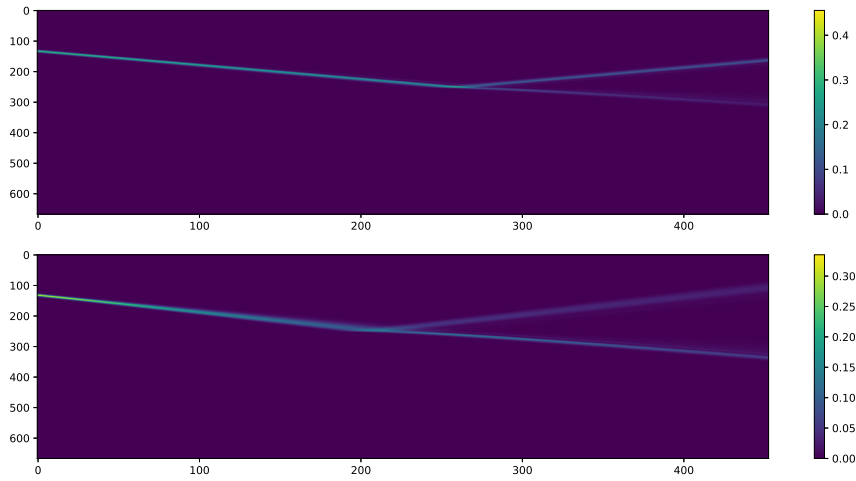


Figure 2: Crank Nicolson 方法 (上) 和 Stable Explicit Scheme (下) 的计算结果

### 3 题目 3: 稳定条件证明

#### 3.1 题目描述

Prove the stability condition( $\frac{c\Delta t}{\Delta x}$ ) of the explicit scheme of the 1D wave equation by performing Von Neumann stability analysis.

$$\frac{\partial^2 u}{\partial t^2} = c^2 \frac{\partial^2 u}{\partial x^2} \quad (7)$$

#### 3.2 证明过程

$$\begin{aligned} \left. \frac{\partial^2 u}{\partial t^2} \right|_{ij} &= \frac{1}{\Delta t^2} (u_{i,j+1} + u_{i,j-1} - 2u_{i,j}) \\ \left. \frac{\partial^2 u}{\partial x^2} \right|_{ij} &= \frac{1}{\Delta x^2} (u_{i+1,j} + u_{i-1,j} - 2u_{i,j}) \\ u_{i,j+1} + u_{i,j-1} - 2u_{i,j} &= \frac{c^2 \Delta t^2}{\Delta x^2} (u_{i+1,j} + u_{i-1,j} - 2u_{i,j}) = \alpha^2 (u_{i+1,j} + u_{i-1,j} - 2u_{i,j}) \\ \frac{(\xi - 1)^2}{\xi} &= \alpha^2 \frac{(e^{IK\Delta x} - 1)^2}{e^{IK\Delta x}} = 4\alpha^2 \sin^2 \frac{1}{2} K\Delta x \end{aligned} \quad (8)$$

当  $\alpha \leq 1$  时,  $|\xi| < 1$ , 所以稳定。