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OPTOMEKATRONİK UYGULAMA VE ARAŞTIRMA MERKEZİ

AO bench design
Zemax model and optimization
post-PDR design

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	Name	October 10, 2017	....
Released	Name	May 30, 2018	SIGNATURE if you want to import a scanned one
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FMV IŞ IK ÜNİ VERSİ TESİ OPTOMEKATRONİ K UYGULAMA VE ARAŞ TIRMA MERKEZİ İ KTİ SADİ İ Ş LETMESİ Meş rutiyet Mahallesi
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Issue	Date	Section/Page affected	Reason/Remarks
1.0	October 10, 2017	All	New
2.0	May 18, 2018	All	Review of all the design after decisions and Paolo S

Acronyms

aO	active Optics
AO	Adaptive Optics
CCD	Charge Coupled Device
CMOS	Complementary Metal Oxide Semiconductor
DAG	Dogu Anadolu Gozlemevi (East Anatolian Observatory)
DM	Deformable Mirror
FCL	Field Correction Lens
FoV	Field-of-View
FWHM	Full Width at Half Maximum
NCPA	Non Common Path Aberrations
NGS	Natural Guide Star
OTF	Optical Transfer Function
P2V	Peak-to-Valley
PSF	Point Spread Function
PSF-R	PSF Reconstruction
PWFS	Pyramid WFS
RMS	Root Mean Square
SH-WFS	Shack-Hartmann WFS
SNR	Signal-to-Noise Ratio
TT	Tip-Tilt
WFS	Wavefront Sensor
WFE	Wavefront Error

1 Scope of this document

This document presents the DAG adaptive optics (AO) optical train design and optimization.

2 Introduction

There are several mandatory points we have to handle :

- the exit pupil has to be imaged on the DM;
- the DM diameter is fixed to the DM-468 from ALPAO; that means $\varnothing$33 mm;
- the exits pupil has to be imaged on the TT mirror for the P-WFS modulation;
- the beam has to converge on the pyramid apex;
- the angle of the beam that arrives onto the pyramid apex is calculated in order to get a diffraction limited PSF size of 2 times the pyramid roof;
- the exit pupil has to be imaged onto the EMCCD detector
- the beam footprint diameter onto the detector has to be defined accordingly to the oversampling criterion (explained section 3.3).

We use off axis parabolas (OAPs) for all our design to reduce the aberrations (compared to lenses). We have added one more constraint on the design which is to have parallel beam arriving or leaving each off axis parabola in order to limit the spherical aberrations.

The fixed parameters are resumed here :

- exit pupil diameter : $\varnothing_{\mathrm{ExtP}} = 727.4046$ mm;
- distance exit pupil to focal plane : $\overline{\mathrm{ExP\text{-}FP}} = 10338.74$ mm;
- DM diameter : $\varnothing_{\mathrm{DM}} = 33.0$ mm;
- pixel size of the Nuvuu EMCCD AO : PxSize $= 24\ \mu$m

In order to design the AO bench we started with the telescope model. Indeed, we would try to compensate for the field of curvature introduced by the mirrors of the telescope with the off-axis parabolas (OAP) of the AO. The Zemax model of the telescope shows a curvature radius of about 1255 mm.

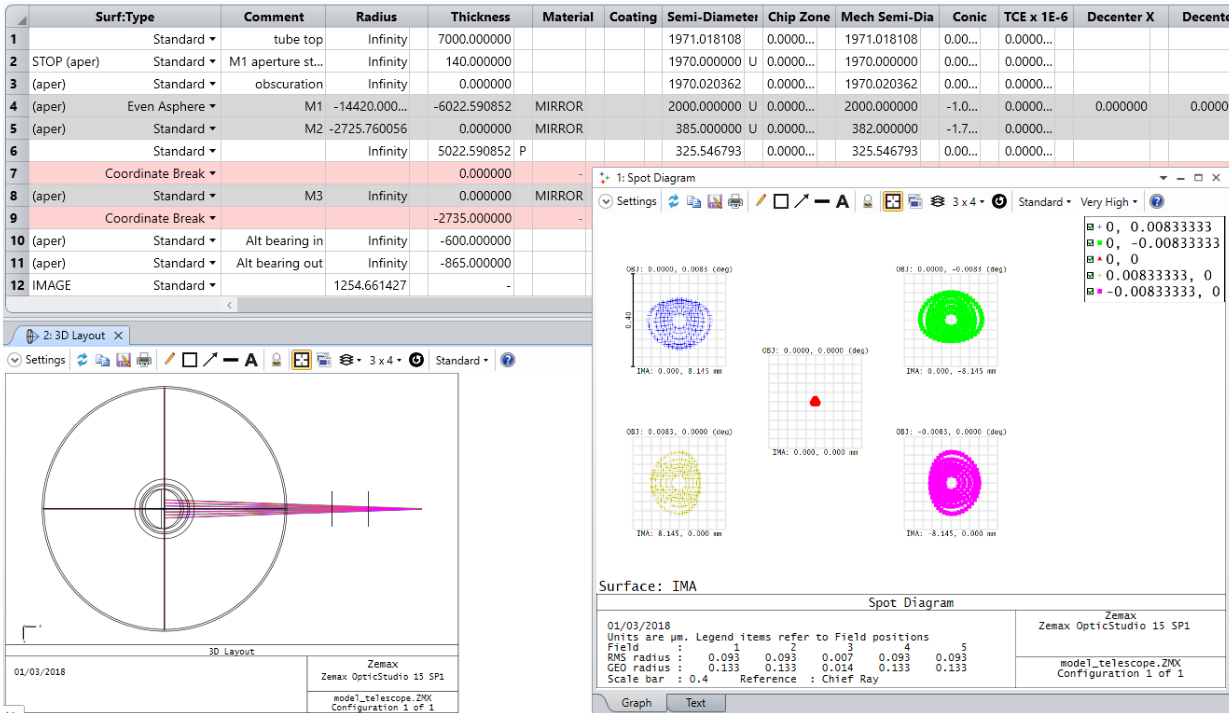

Figure 1: Zemax model of the telescope

3 Design development and optimization

A preliminary design have been set for the PDR and the main points are summarized at each optimization step of this report.

3.1 Imaging the exit pupil of the telescope onto the DM

The first branch consists on an imaging system of the telescope pupil onto the DM. The beam has to be collimated and the footprint has to take the entire clear aperture of the DM (for the DM-468 from ALPAO $\varnothing_{\mathrm{DM}} = 33$ mm).
The beam on the DM has to be reflected with a certain angle otherwise the reflected beam comes back on itself. The maximum acceptable angle can be calculated considering the position error of the beam per actuator on the DM. We can align the beam with an error of $1/10^{th}\Lambda$ (Λ the actuator pitch = 1.5 mm for the ALPAO DM-468). In order for the projected beam diameter to be no more than 10% smaller than the DM diameter on both side, the tilt of the DM must be no more than α (figure 2) :

$$\begin{aligned}
\alpha &= \arccos\left(1 - \frac{1}{5}\frac{\Lambda}{\varnothing_{\mathrm{DM}}}\right) \\
\alpha &= \arccos\left(1 - \frac{1}{5}\frac{1.5}{33}\right) \\
\alpha &= 7.73° \qquad (1) \\
2\alpha &= 15.46°
\end{aligned}$$

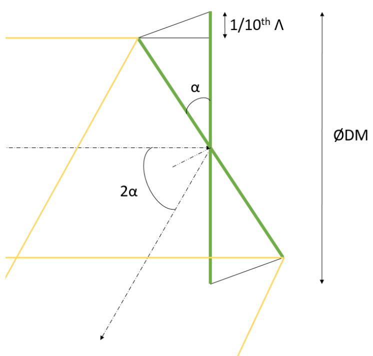

Figure 2: Sketch of the input and output beam depending on the tilt angle of the DM

The input beam diameter is then :

$$\varnothing_{\text{DM-beam}} = \varnothing_{\text{DM}} \cos \alpha = 33 \times \cos(7.73) = 32.7 \text{ mm} \tag{2}$$

In order to have a collimated beam on the DM with a diameter of $\varnothing_{\text{DM-beam}}$, the focal length of the OAP0 is calculated by the following sequence of equations. To visualize the parameters needed and the context we can look at figure 3. The with an angle θ with the corresponding reference coordinates $(0, x', y')$. The ray ΔCR corresponds to the chief ray, $\Delta\alpha u$ is the "upper" marginal ray and $\Delta\alpha b$ is the "bottom" marginal ray. The origin is placed at the focal plane of the telescope.

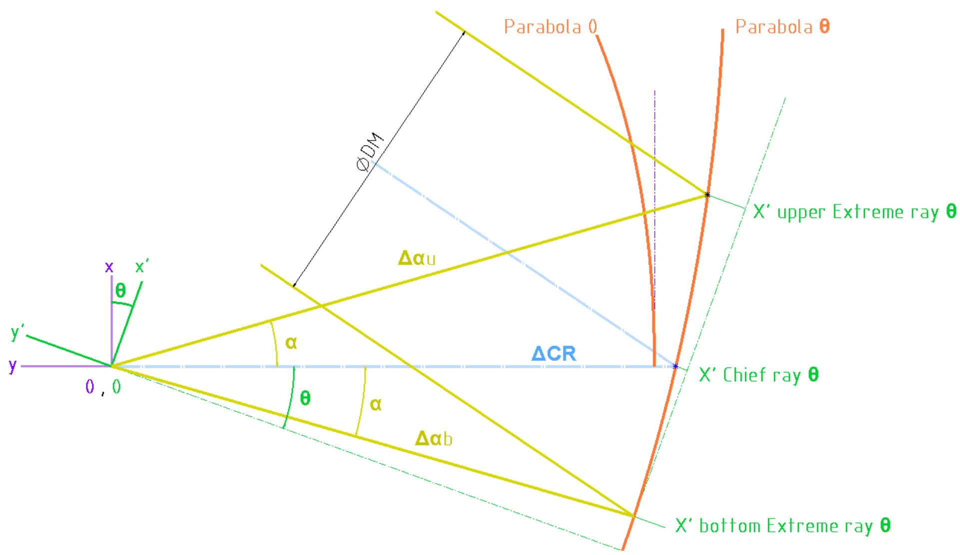

Figure 3: Ray tracing of a tilted parabola

In the reference coordinates $(0, x', y')$, we have (with PFL the Parental Focal Length of the OAP):

- equation of the parabola θ :

$$y' = \frac{x'^2}{4\,\text{PFL}} - \text{PFL} \tag{3}$$

- equation of $\Delta\alpha u$:

$$y' = -\frac{1}{\tan(\theta + \alpha)} x' \tag{4}$$

- equation of $\Delta\alpha b$:

$$y' = -\frac{1}{\tan(\theta-\alpha)}x' \tag{5}$$

Determination of $x'_{\text{Ex}\theta\ u}$ (intersection between parabola θ and $\Delta\alpha u$) :
(3) = (4)

$$x'_{\text{Ex}\theta\ u} = 2\,\text{PFL}\left(-\frac{1}{\tan(\theta+\alpha)} + \sqrt{\frac{1}{\tan^2(\theta+\alpha)}+1}\right) \tag{6}$$

Determination of $x'_{\text{Ex}\theta\ b}$ (intersection between parabola θ and $\Delta\alpha b$) :
(3) = (5)

$$x'_{\text{Ex}\theta\ b} = 2\,\text{PFL}\left(-\frac{1}{\tan(\theta-\alpha)} + \sqrt{\frac{1}{\tan^2(\theta-\alpha)}+1}\right) \tag{7}$$

In order to have the OAP0 output beam diameter equal to the DM clear aperture we have :

$$x'_{\text{Ex}\theta\ u} - x'_{\text{Ex}\theta\ b} = \varnothing_{\text{DM-beam}} \tag{8}$$

The parental focal length (see appendix A) of OAP0 is :

$$\text{PFL} = \frac{1}{2}\frac{\varnothing_{\text{DM-beam}}}{-\frac{1}{\tan(\theta+\alpha)}+\sqrt{\frac{1}{\tan^2(\theta+\alpha)}+1}+\frac{1}{\tan(\theta-\alpha)}-\sqrt{\frac{1}{\tan^2(\theta-\alpha)}+1}} \tag{9}$$

Using the equation (24) described in appendix A we can transform the parental focal length into the effective focal length.

For the OAP0 the angle α_{OAP0} is

$$\alpha_{\text{OAP0}} = \arctan\left(\frac{\varnothing_{\text{ExP}/2}}{\overline{\text{ExP-FP}}}\right) \tag{10}$$

The numerical application of (10) with $\varnothing_{\text{ExP}} = 727.4046$ mm and $\overline{\text{ExP-FP}} = 10338.74$ mm gives $\alpha_{\text{OAP0}} = 2.015°$, the aperture demi angle of the beam arriving on the first OAP.

After the discussion we had with Paolo Spano, we took his advice into account and choose the parabola tilt angle according to the following precept : $\theta = 35°$ is a critical choice and $\theta > 35°$ is a nightmare (manufacturing and alignment). We should take an tilt angle

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smaller than 30° to be realistic and the smaller the best. However, we are limited by the space around the focal plane because of the field correction lens (represented by the bloc on figure 4). Depending on where it is going to be placed we would have to increase the angle of OAP0. **An email has been sent to AMOS who is going to design the FCL to know if it is possible to have it before the focal plane.**

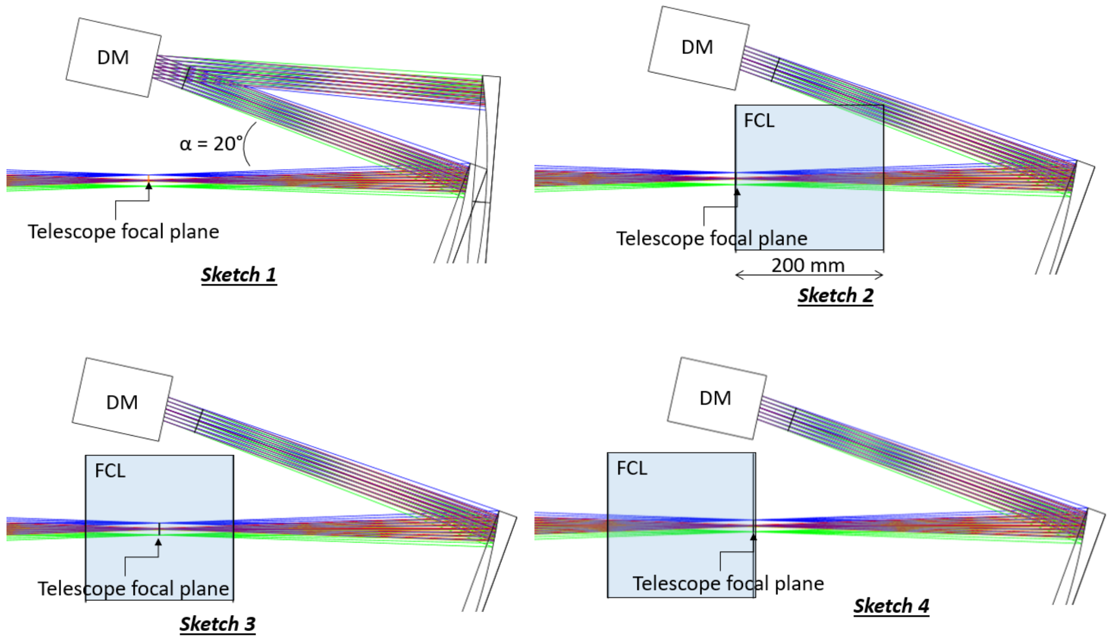


Figure 4: Field correction lens position from the focal plane of the telescope

We can set the parabola tilt angle to $\theta = 20°$ in order to keep this angle as small as possible to limit aberrations.

Using these inputs the focal length of the OAP0 is (appendix D): $\mathrm{PFL_{OAP0}} = 443.13\,\mathrm{mm}$. This value is too specific so we round it to[1] :

$$\mathrm{PFL_{OAP0}} = 445\,\mathrm{mm}$$

[1]In Zemax, we enter the $\mathrm{PFL_{OAP0}}$ and not the effective focal length in the coordinate break surface because using it, the translation is done before the tilt angle. This is why if we enter the translation in a surface before the coordinate break we cannot use the same length (we would then write the $\mathrm{EFL_{OAP0}}$)

The image of the exit pupil given by OAP0 gives the position of the DM. Its position p_i relative to the OAP0 vertex is given by (we use Gauss law) :

$$p_i = \frac{p_o \, \mathrm{PFL}_{\mathrm{OAP0}}}{p_o + \mathrm{PFL}_{\mathrm{OAP0}}} \tag{11}$$

$$p_i = \frac{-(10338.74 + 445)445}{-10338.74} \tag{12}$$

$$p_i = 464\,\mathrm{mm} \tag{13}$$

The Zemax "pick up pupil position" macro is used to get the perfect DM position wrt the OAP0. The DM is then at $p_i = 493.08$ mm. Even if we understand that we cannot align the DM reflecting surface with this accuracy, we keep it in the Zemax model and we will work on that during the tolerancing analysis[2].
The image of the pupil through a tilted OAP is also tilted (see [1]). This creates pupil aberrations.

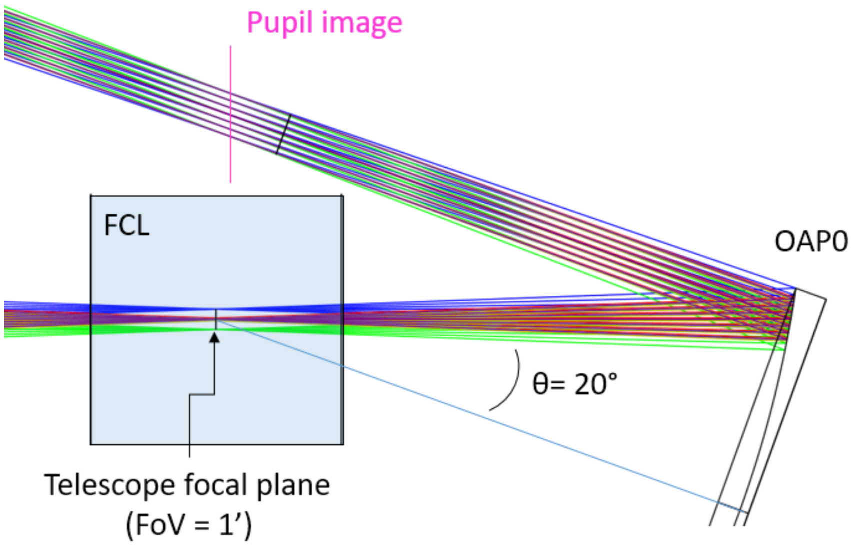
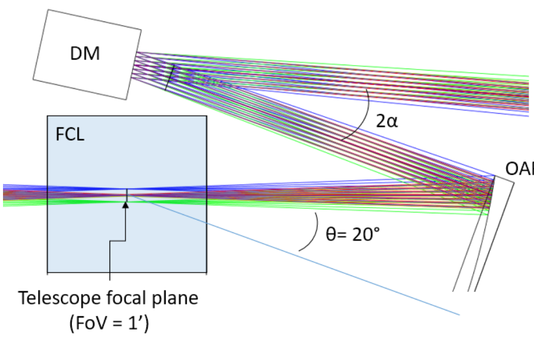


Figure 5: Tilted image of the pupil

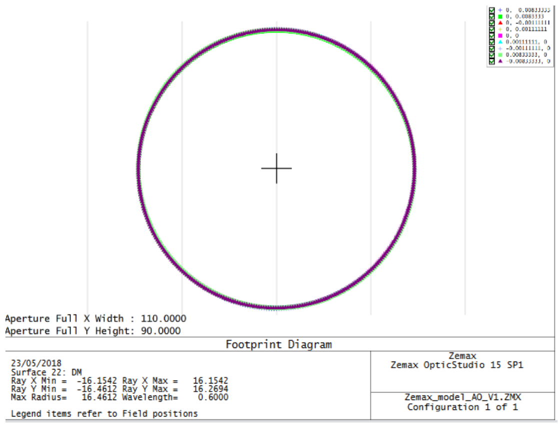
Here we want to place the DM at this pupil position. When we work with OAPs the image stop after an OAP is not perpendicular to the optical axis. The OAP introduce a tilt angle

[2]We always keep all the digit during an AO design in order to see what are the best results we can obtain with an ideal system.

of the image plane. If we tilt the DM in the opposite position of this OAP-introduced tilt angle we add pupil aberrations whereas if we tilt it the same direction we can compensate for it.



(a) Layout



(b) Zemax footprint diagram of the beam on the DM

Figure 6: Zemax model of the 1rst part of the AO design

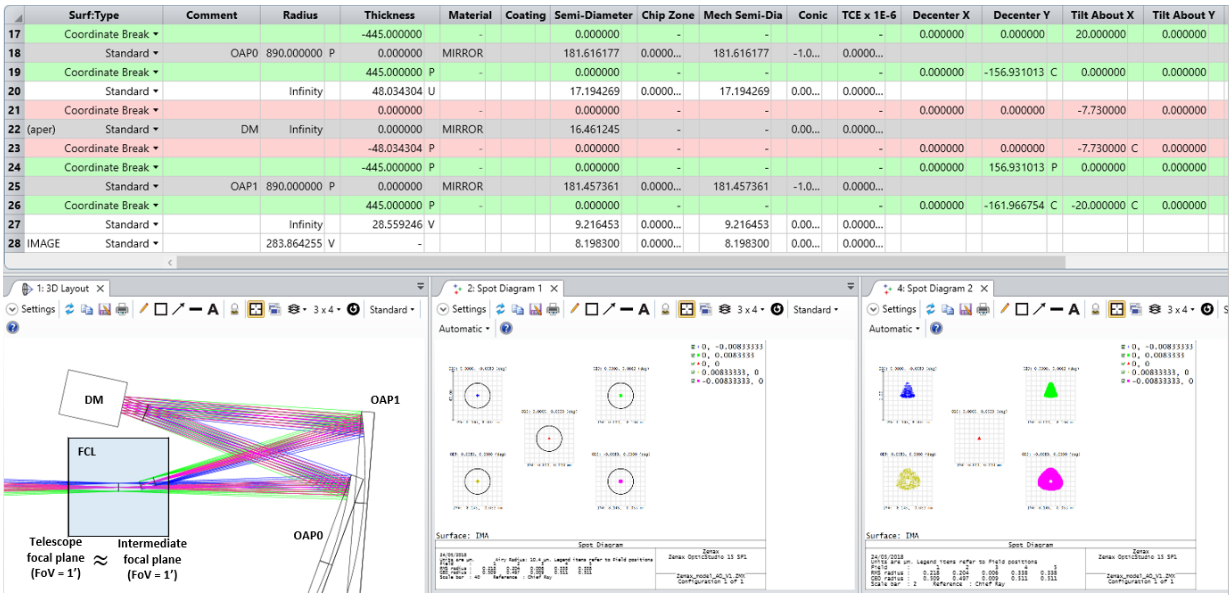
The beam footprint on the DM shows that all the field fits inside the clear aperture. The maximum diameter is equal to 32.7306 mm ($<$ 33 mm). The angle of the DM is set to -7.73° (equation (1)),and not 7.73°as explain above.

3.2 Imaging the pupil on the TT modulation mirror

3.2.1 Create an image of the pupil

We need to image the pupil on the TT modulation mirror. Moreover, we know that to compensate for the introduced aberration from OAP0 we can use an OAP1 with the same focal length. The F/D ratio would stay the same as the telescope one which is favourable for the science path.

Figure 7 shows that to model this combination we use the ”chief ray” solve on Zemax (Lens data) to keep the coordinates following the beam path (using the on-axis ray). OAP1 is placed at the exact inverse position of OAP0 in order to compensate for its aberrations so the intermediate focal plane is near the telescope focal plane.



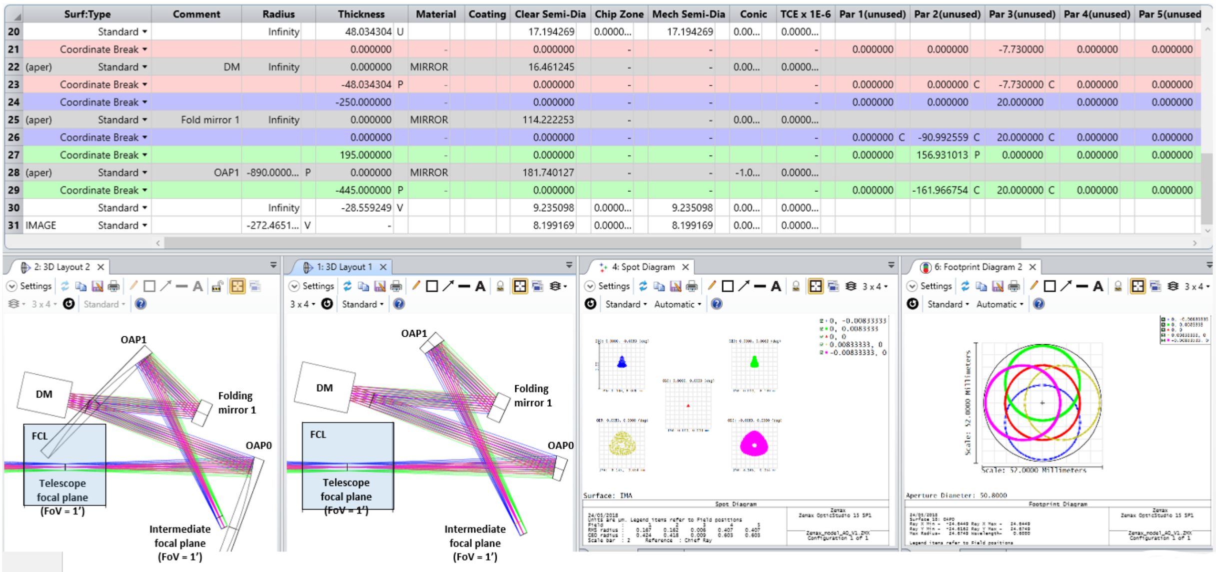
	Surf:Type	Comment	Radius	Thickness	Material	Coating	Semi-Diameter	Chip Zone	Mech Semi-Dia	Conic	TCE x 1E-6	Decenter X	Decenter Y	Tilt About X	Tilt About Y
17	Coordinate Break			-445.000000	-		0.000000	-	-		-	0.000000	0.000000	20.000000	0.000000
18	Standard	OAP0	890.000000 P	0.000000	MIRROR		181.616177	0.0000...	181.616177	-1.0...	0.0000...				
19	Coordinate Break			445.000000 P	-		0.000000	-	-		-	0.000000	-156.931013 C	0.000000	0.000000
20	Standard		Infinity	48.034304 U			17.194269	0.0000...	17.194269	0.00...	0.0000...				
21	Coordinate Break			0.000000	-		0.000000	-	-		-	0.000000	0.000000	-7.730000	0.000000
22 (aper)	Standard	DM	Infinity	0.000000	MIRROR		16.461245	-	-	0.00...	0.0000...				
23	Coordinate Break			-48.034304 P	-		0.000000	-	-		-	0.000000	0.000000	-7.730000 C	0.000000
24	Coordinate Break			-445.000000 P	-		0.000000	-	-		-	0.000000	156.931013 P	0.000000	0.000000
25	Standard	OAP1	890.000000 P	0.000000	MIRROR		181.457361	0.0000...	181.457361	-1.0...	0.0000...				
26	Coordinate Break			445.000000 P	-		0.000000	-	-		-	0.000000	-161.966754 C	-20.000000 C	0.000000
27	Standard		Infinity	28.559246 V			9.216453	0.0000...	9.216453	0.00...	0.0000...				
28 IMAGE	Standard		283.864255 V	-			8.198300	0.0000...	8.198300	0.00...	0.0000...				

Figure 7: Zemax model from the telescope beam to the intermediate focal plane

We use an optimization (smallest spot radius) to go to the best focal plane (surface # 76 thickness is set as variable) and in order to consider the field curvature we set the radius of the image plane variable. The optimization using a merit function (Type : P2V, Criteria : Spot Radius, Reference : Chief Ray) is done on the image surface to get the smallest spot radius of the central beam (FoV = 0'). We can see that the spot diameter varies between 0.511 μm and 0.009 μm (the center of the FoV) but when we look the spots with the Airy disk we can say that the beam is well focused.
However, the radius of the image curvature is smaller than the telescope output field of curvature so we did not compensate for that, instead we increase it (we will see how we could work on that or not).

3.2.2 Bend the beam

We do not want the beam to go back in the field correction lens and we would like to place the dichroic nearby the intermediate focal plane. In order to do that properly, we fold the beam after the DM with a flat mirror.



	Surf:Type	Comment	Radius	Thickness	Material	Coating	Clear Semi-Dia	Chip Zone	Mech Semi-Dia	Conic	TCE x 1E-6	Par 1(unused)	Par 2(unused)	Par 3(unused)	Par 4(unused)	Par 5(unused
20	Standard		Infinity	48.034304 U			17.194269	0.0000...	17.194269	0.00...	0.0000...					
21	Coordinate Break			0.000000	-		0.000000	-	-		-	0.000000	0.000000	-7.730000	0.000000	0.000000
22 (aper)	Standard	DM	Infinity	0.000000	MIRROR		16.461245	-	-	0.00...	0.0000...					
23	Coordinate Break			-48.034304 P	-		0.000000	-	-		-	0.000000	0.000000 C	-7.730000 C	0.000000	0.000000
24	Coordinate Break			-250.000000	-		0.000000	-	-		-	0.000000	0.000000	20.000000	0.000000	0.000000
25 (aper)	Standard	Fold mirror 1	Infinity	0.000000	MIRROR		114.222253	-	-	0.00...	0.0000...					
26	Coordinate Break			0.000000	-		0.000000	-	-		-	0.000000 C	-90.992559 C	20.000000 C	0.000000	0.000000
27	Coordinate Break			195.000000	-		0.000000	-	-		-	0.000000	156.931013 P	0.000000	0.000000	0.000000
28 (aper)	Standard	OAP1	-890.0000... P	0.000000	MIRROR		181.740127	-	-	-1.0...	0.0000...					
29	Coordinate Break			-445.000000 P	-		0.000000	-	-		-	0.000000	-161.966754 C	20.000000 C	0.000000	0.000000
30	Standard		Infinity	-28.559249 V			9.235098	0.0000...	9.235098	0.00...	0.0000...					
31 IMAGE	Standard		-272.4651... V	-			8.199169	0.0000...	8.199169	0.00...	0.0000...					

Figure 8: Zemax model from the telescope beam to the intermediate focal plane with a fold mirror after the DM

Figure 8 shows the new design with a folding mirror tilted at 20°. For more clarity we apply the real aperture size of the OAPs on layout 1. We can see on the footprint diagram that we can use two OAPs of 2 inches of diameter. The maximum diameter of the beam footprint (for a FoV = 1') is about 49.9 mm. However, we should be careful and verify during the tolerance analysis that when we move the OAP in the alignment precision range we can do, the entire beam stays reflected on the OAP (no vignetting). Concerning the folding mirror, we can also use a diameter of 2 inches it is positioned at 298 mm from the DM (same as above, we have to check during the tolerancing that we do not have vignetting when moving around this position).

After redoing the optimization of the intermediate focal plane position and the image radius of curvature we found spot diameters vary from 0.009 μm (on-axis) to 0.603 μm (0.5') across the FoV (see figure 9).

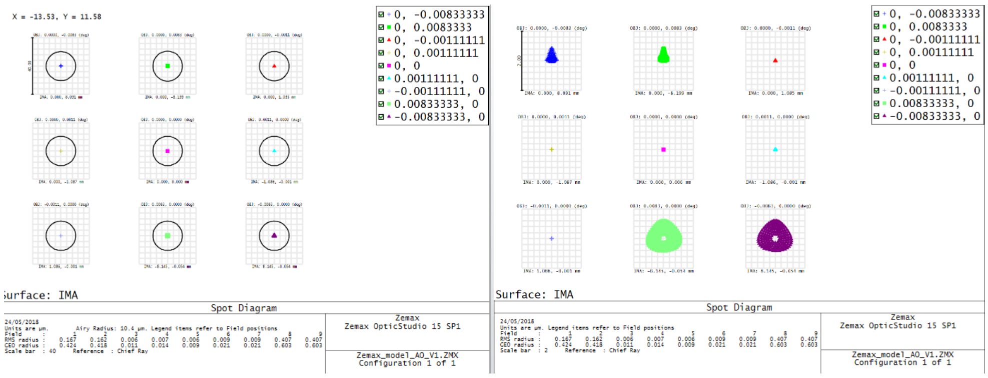


Figure 9: Spot diagram at the intermediate focal plane for the AO FoV and the WFS path FoV

3.2.3 Imaging the pupil on the TT modulation mirror

The selection of the star in the FoV is made using an XY stage where all the WFS path would be mounted (as we have a compact WFS path). Now, as the OAP0 and OAP1 are the same, this optical configuration creates an image of the pupil of the same diameter at the same distance from the FP as the entrance pupil.

It is so not possible to use it to place the TT modulation mirror which has to be about 10 mm in diameter and at a reasonable distance from FP. FOR these reasons....................................... we take a TT modulation mirror of ⌀0.5 inch mounted on the fast tip-tilt platform S-331 from PI. We want to have a beam of about 10 mm in diameter at the pupil to place the TT stage.

We have also to note that we would like to introduce a dichroic membrane nearby the intermediate focal plane.

We know that we need to arrive on the pyramid roof (P-WFS) with a corresponding :

$$F\# = 2\frac{\text{pyramid roof}}{\lambda} \tag{14}$$

Taking a pyramid roof of about 20 μm according to Jean-Pierre Veran (private communication) at $\lambda = 0.6\ \mu$m we have :

$$F\# = 66.7 \tag{15}$$

We round this value to F# = 60 (we need to investigate the pyramid roof size that can be manufactured).

We want then a vergence that takes the beam with a F# of 14 and brings it to 60. The TT modulation mirror is placed in this converging beam at the pupil position. At first approximation, we can do a geometrical dimensioning. We use a drawing software and apply the constraints to calculate which focal length and distance from the intermediate focal plane we need. In order to build this ray tracing properly we calculate the following parameters :

$$F\#_{\text{input}} = \frac{\overline{\text{ExP-FP}}}{\varnothing_{\text{ExtP}}}$$
$$F\#_{\text{output}} = \frac{f_{\text{lens}}}{\varnothing_{\text{beam on the lens}}}$$
$$\gamma_{\text{input}} = \frac{1}{2F\#_{\text{input}}}$$
$$\gamma_{\text{output}} = \frac{1}{2F\#_{\text{output}}}$$

The image of the pupil is nearby the foci (at 0.06 mm) and has to be about 5 mm in diameter. Taking the angles and the beam diameter on the TT modulation mirror we find the focal length of the lens and the distance from the intermediate focal plane described figure 10.

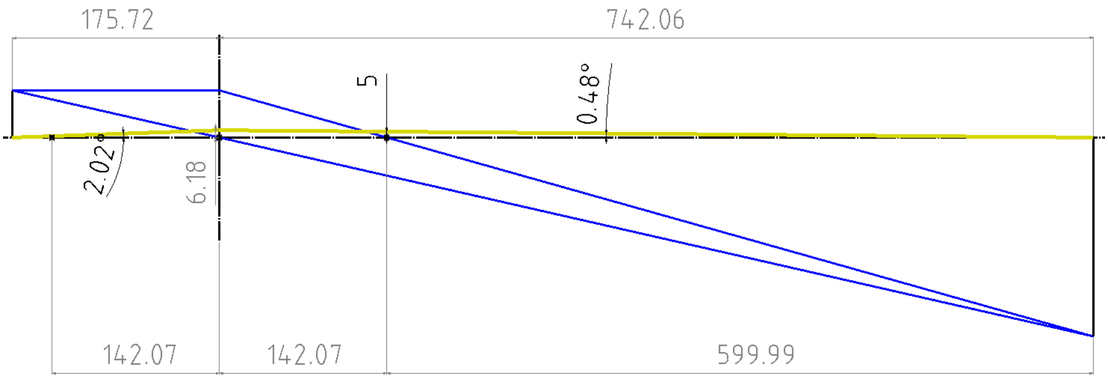


Figure 10: Ray tracing using SolidWorks to approximate the converging lens before the TT modulation mirror

We introduce these results in Zemax using a paraxial lens (figure 11). We can see that the beam diameter is about 5 mm as expected.

	Surf:Type	Comment	Radius	Thickness	Material	Coating	Semi-Diameter	Chip Zone	Mech Semi-Dia	Conic	TCE x 1E-6	Focal Length
30	Standard		Infinity	-28.559248			9.235098	0.0000...	9.235098	0.00...	0.0000...	
31 (aper)	Standard	Intermediate focal plane	Infinity	-175.000000			8.203247	0.0000...	1.100000	0.00...	0.0000...	
32	Paraxial			-138.493007 U			14.086406	-	-		0.0000...	142.000000
33	Coordinate Break			0.000000	-		0.000000	-	-		-	0.000000
34	Standard	TT mirror	Infinity	0.000000	MIRROR		5.095547	0.0000...	5.095547	0.00...	0.0000...	
35	Coordinate Break			614.537595 V	-		0.000000	-	-		-	0.000000
36 IMAGE	Standard		274.506411 V	-			35.413742	0.0000...	35.413742	0.00...	0.0000...	

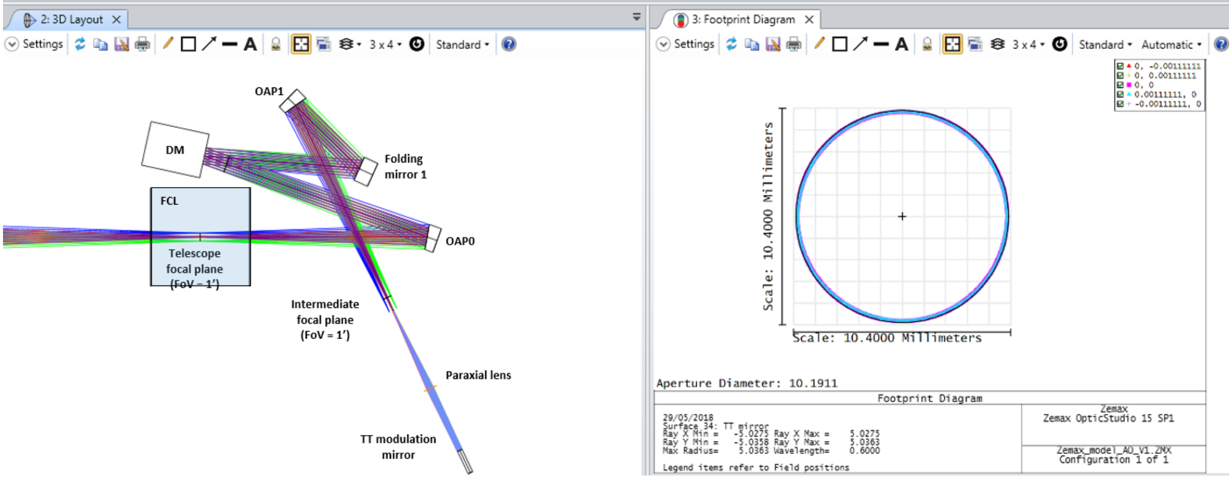


Figure 11: Zemax model from the telescope beam to the TT modulation mirror using a paraxial lens

We want to tilt the TT modulation mirror to send the beam in a convenient direction. We need to add the ADC nearby a pupil plane, so if we want to use the pupil where the TT modulation mirror is placed we need some space around it to insert the ADC. We can tilt the TT modulation mirror with an angle of about 20°. The radius of the beam footprint on the TT modulation mirror measured on the Zemax model is 5.32 mm at maximum (which is smaller than 0.5 inch).
Moreover, the focal length of 142 mm is not a common value so we can change it to 140 mm.

We use the optimization tool to find the best focal plane (distance and radius of curvature at the image plane). As we are going to move all the WFS path with a XY stage to pick up the star, we can determine the beam footprint diameter on the paraxial lens with the on-axis case. We can see that the F# of the beam arriving on the apex of the pyramid (figure 11) is F# $= \frac{136.59+579.88}{7.1884*2} = 58.5$ which is smaller than the F#fixed before at 60. We can adjust it moving a little bit the paraxial lens from the intermediate focal plane. We find the final solution (figure 12) with F# $= \frac{136.59+597.35}{6.0897*2} = 60.3$.

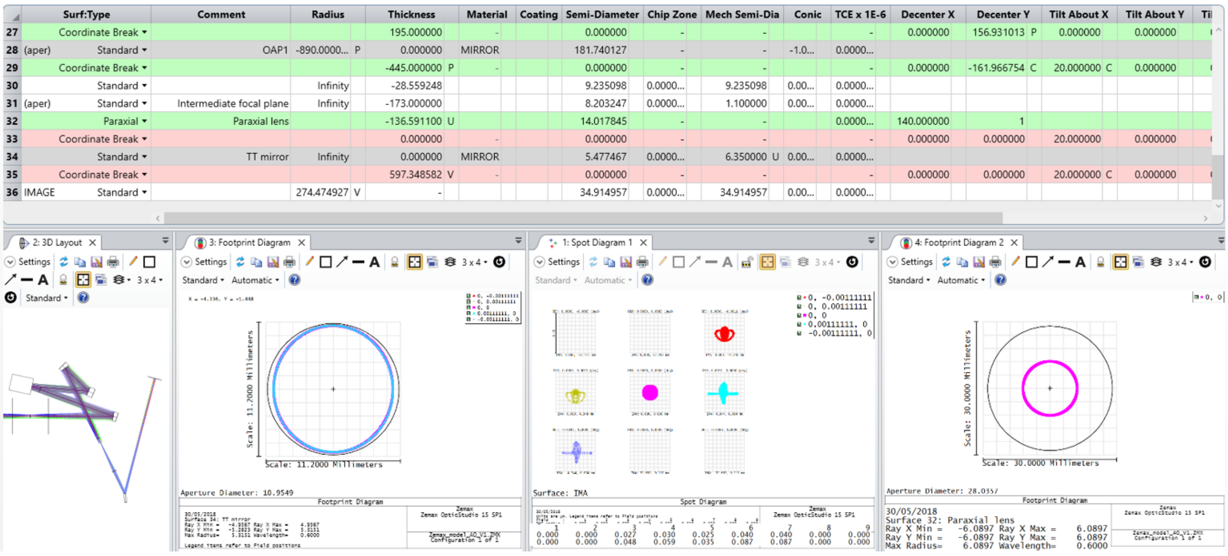


Figure 12: Zemax model from the telescope beam to the apex of the pyramid WFS.

Actually, we could play with the angle of the TT modulation mirror to place the ADC closer to the pupil plane. This would be investigate according to the ADC design.

3.3 Image the pupil on the CCD detector

The detector is the Nuvuu EMCCD with 128 pixels and a pixel size of 24 μ m. We have to sample the beam with at least 1 pixel per actuator. In order to relax the alignment problems we can oversample the beam. According to Jean-Pierre Veran advice an oversampling of 1.5 should be enough. The ALAPO DM-468 has 22 actuators across the clear aperture diameter (so 23 pitches). Then we can calculate the beam diameter on the detector :

$$\varnothing_{\text{CCD}} = \#_{\text{actuator across } \varnothing} \times \text{PxSize} \times \text{Oversampling factor} \quad (16)$$

$$\varnothing_{\text{CCD}} = 22 \times 24\ \mu\text{m} \times 1.5 \quad (17)$$

$$\varnothing_{\text{CCD}} = 0.792\ \text{mm} \quad (18)$$

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The pupil has to be images on the camera. The TT mirror (pupil plane) image through the OAP3 is at the infinity. So when we image it through the relay lens the pupil plane arrives at the focal plane of the relay lens. Moreover, in order to get the beam size on the camera, the focal length of the relay lens (f_{LR}, see appendix B) has to be :

$$f_{\mathrm{LR}} = \mathrm{F\#} \times \varnothing_{\mathrm{CCD}} \tag{19}$$

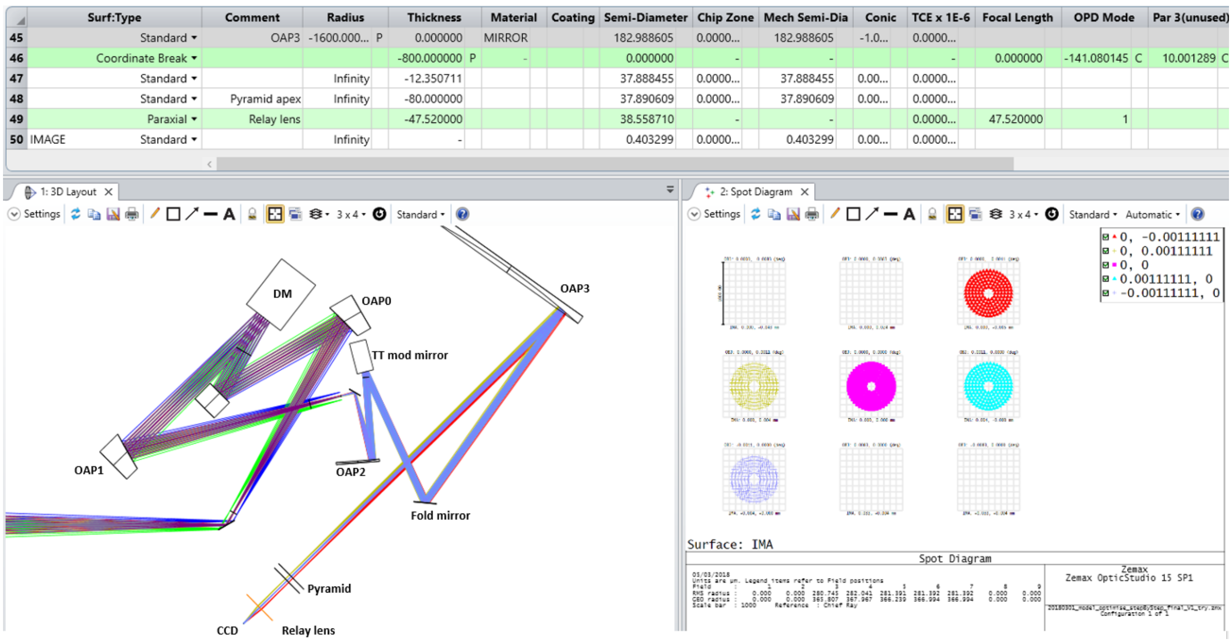
$$f_{\mathrm{LR}} = 60.26 \times 0.792 \tag{20}$$

$$f_{\mathrm{LR}} = 47.73 \text{ mm} \tag{21}$$

surface 26 thickness variable : jouer avec cette valeur pour obtenir un diametre de 0.792mm sur la CCD = 0.396mm en rayon

STOOOOOOOOOOOOOOOOOOOOOOOOOOOOOOOOPPPPPPPPPPPPPPPP

The remaining optics have been implemented in the Zemax model just to have an estimation of the space no optimization has been done (we can already see that the diameter is not precisely of $\varnothing_{\mathrm{CCD}}$).



	Surf:Type	Comment	Radius	Thickness	Material	Coating	Semi-Diameter	Chip Zone	Mech Semi-Dia	Conic	TCE x 1E-6	Focal Length	OPD Mode	Par 3(unused)
45	Standard	OAP3	-1600.000... P	0.000000	MIRROR		182.988605	0.0000...	182.988605	-1.0...	0.0000...			
46	Coordinate Break			-800.000000 P	-		0.000000	-	-		-	0.000000	-141.080145 C	10.001289 C
47	Standard		Infinity	-12.350711			37.888455	0.0000...	37.888455	0.00...	0.0000...			
48	Standard	Pyramid apex	Infinity	-80.000000			37.890609	0.0000...	37.890609	0.00...	0.0000...			
49	Paraxial	Relay lens		-47.520000			38.558710	-	-		0.0000...	47.520000	1	
50 IMAGE	Standard		Infinity	-			0.403299	0.0000...	0.403299	0.00...	0.0000...			

Figure 13: Zemax model of the bench until the CCD camera

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3.4 Atmopsheric Dispersion Compensator (ADC)

We are designing an ADC to compensate the atmospheric dispersion (see appendixC). The geometric parameters and the glass are not set yet this is why we do not have implemented it in the Zemax model for now. However, we can already say that the ADC would probably be inserted between OAP2 and OAP3 because the beam is collimated.

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A OAPs focal lengths relations

This section describe the relation between the parental and the effective focal length for an off-axis parabola. The sketch figure 14 describes the situation with :

- EFL the effective focal length
- PFL the parental focal length
- (x_S, y_S) the coordinate of the ray intersection with the parabola
- the parabola of equation

$$y = \frac{x^2}{4\text{PFL}} - \text{PFL} \tag{22}$$

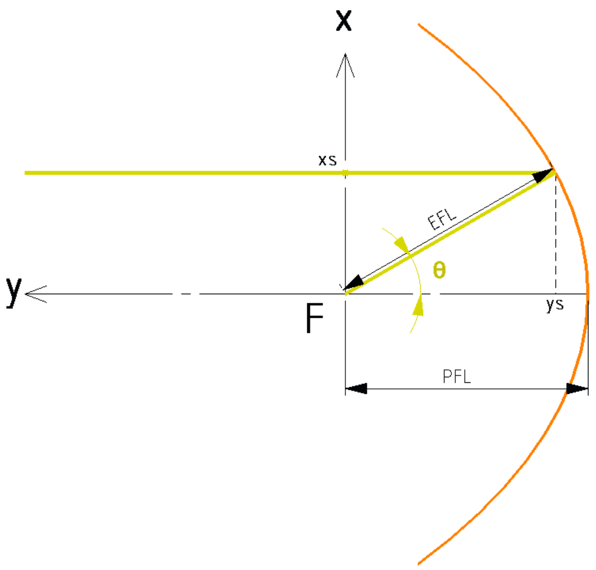

Figure 14: Sketch of a ray reflected on an OAP

The relation between EFL and PFL can be determined with the following set of equations.

$$\begin{aligned} x_S &= \text{EFL}\sin\theta \\ y_S &= \text{EFL}\cos\theta \end{aligned} \tag{23}$$

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Using (22) and (23) we have :

$$
\begin{aligned}
4\mathrm{PFL}\left(y_S+\mathrm{PFL}\right) &= x_S^2 \\
4\mathrm{PFL}\left(\mathrm{EFL}\cos\theta+\mathrm{PFL}\right) &= \left(\mathrm{EFL}\sin\theta\right)^2 \\
4\mathrm{PFL}^2-4\mathrm{PFL}\,\mathrm{EFL}\cos\theta-\mathrm{EFL}^2\sin^2\theta &= 0
\end{aligned}
$$

$$
\begin{aligned}
\mathrm{PFL}_{1,2} &= \frac{4\mathrm{EFL}\cos\theta \pm \sqrt{\left(-4\mathrm{EFL}\cos\theta\right)^2+16\mathrm{EFL}^2\sin^2\theta}}{2*4} \\
&= \frac{1}{2}\left[\mathrm{EFL}\cos\theta \pm \sqrt{\mathrm{EFL}^2\cos^2\theta+\mathrm{EFL}^2\sin^2\theta}\right]
\end{aligned}
$$

The final equation is :

$$
2\mathrm{PFL}=\mathrm{EFL}\left(1+\cos\theta\right) \tag{24}
$$

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B Focal length of the relay lens between the pyramid and the CCD

Gauss :

$$\begin{aligned} \frac{1}{p_i} - \frac{1}{p_o} &= \frac{1}{f_{LR}} \\ 1 - \frac{f_{LR}}{p_i} &= -\frac{f_{LR}}{p_o} \end{aligned} \tag{25}$$

Thales :

$$\begin{aligned} \frac{\varnothing_{\mathrm{CCD}}}{\varnothing_{\mathrm{L}}} &= \frac{|p_i| - |f_{LR}|}{|p_i|} \\ \frac{\varnothing_{\mathrm{CCD}}}{\varnothing_{\mathrm{L}}} &= 1 - \frac{f_{LR}}{p_i} \end{aligned} \tag{26}$$

Mixing (25) and (26) :

$$\frac{\varnothing_{\mathrm{CCD}}}{\varnothing_{\mathrm{L}}} = \frac{f_{LR}}{p_o} \tag{27}$$

In another hand :

$$\mathrm{F\#} = \frac{f_{\mathrm{OAP3}}}{\varnothing_{\mathrm{TT}}} = \frac{p_o}{\varnothing_L} \tag{28}$$

Mixing (27) and (28) :

$$f_{\mathrm{LR}} = \mathrm{F\#}\ \varnothing_{\mathrm{CCD}} \tag{29}$$

C ADC design

C.1 Scope

The previous report explains the different possibilities to configure an ADC. Here we are presenting the design we are going to implement for the case of the DAG telescope AO. This ADC will be for visible wavelength. The range is limited by the camera Nuvu spectral range and the star spectrum studied. We are taking only the Nuvu camera bandwidth into account for a first iteration.

C.2 Amici principle

As said in the previous report, the Amici prisms are commonly used in the ADC systems. The ADCs are composed by a 2-doublet design which mixes two Amici prisms. This kind of prism is an alliance of two pieces of glass which have different dispersion (different refractive indexes). The materials must have the same refractive number for a mean wavelength so that at this frequency the incident and emergent rays are parallel (zero-deviation at λ_{mean}). The two prisms can be rotated around the optical axis in order to change the dispersion and compensate it for each wavelength. Shorter or longer wavelength than the middle one are deflected in opposite directions. When the angle between them is 180° the dispersion is reduced at its minimum.

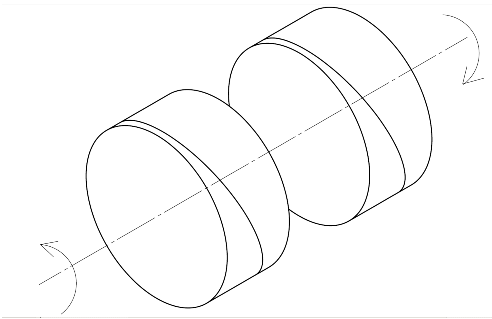

Figure 15: 2-doublet Amici prisms design

Another design can be a three-glass Amici prisms, which is called triplet-design. It consists of the insertion of a anomalous dispersion glass between the two firstly introduced surfaces.

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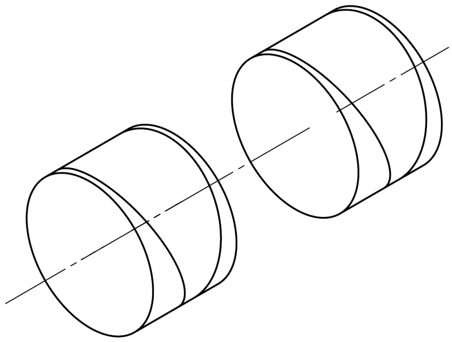

Figure 16: 2-triplet Amici prisms design [2]

The triplet-design seems to be appropriate for our case because it gives the best performance. According to Kopon thesis [2], the doublet corrects only the first order of chromatism while the triplet acts on both primary and secondary chromaticism aberration.

The system is made only by plane surfaces so that a plane front which go through the prisms should leave them as a plane front. In a collimated beam with a very narrow field of view the prisms combined thickness can be larger than the beam diameter [3].
We can also use vertex with a certain radius in a converging beam. This way we can change the F# leaving the ADC.
The doublet glass has to be designed in order to have a thermal expansion rate as close as possible between each glass to not break with the temperature changes. The internal reflection is also a parameter that we have to take into account because we want the ADC to transmit as much as possible light.

ADC design In order to design our ADC, we need to collect some information about its position in the AO bench,the dispersion of the atmosphere for our parameters, the field of view (FoV) and wavelength bandwidth it has to work in.

C.3 Model of the atmosphere refraction

C.3.1 Atmosphere refraction index

First, we need to model to atsmophere dispersion (refractive index) depending on the wavelength. The simulation of the atmosphere refractive index is based on Ciddor's approach [4] which is a compilation of all previous equations for the visible and near infrared. The following set of equations are used to model the atmosphere refraction :

$$
\begin{aligned}
10^8\left(n_{as}-1\right) &= k_1/\left(k_0-\sigma^2\right)+k_3/\left(k_2-\sigma^2\right) && (30)\\
\left(n_{asx}-1\right) &= \left(n_{as}-1\right)\left[1+0.534\times 10^{-6}\left(x_c-450\right)\right] && (31)\\
10^8\left(n_{ws}-1\right) &= 1.022\times\left(\omega_0+\omega_1\sigma^2+\omega_2\sigma^4+\omega_3\sigma^6\right) && (32)\\
n_{final} &= \left(\rho_a/\rho_{axs}\right)\left(n_{axs}-1\right)+\left(\rho_\omega/\rho_{\omega s}\right)\left(n_{\omega s}-1\right) && (33)
\end{aligned}
$$

where the parameters are defined as follow :

- the wave number
  $\sigma = 2\pi/\lambda\ [\mu\mathrm{m}^{-1}]$
- Constants involved in the standard phase and group refractivities of dry air [4]
  $k_0 = 238.0182,\ k_1 = 5792105,\ k_2 = 57.362,\ k_3 = 167917\ [\mu\mathrm{m}^{-2}]$
- n_{as} the refractive index of standard air at T=20°C, 101325 Pa, 0 % humidity, 450 ppm of CO_2
- for now, I took the concentration of CO_2 in the air of x_c = 450 ppm (standard) so equation (31) becomes $n_{asx} = n_{as}$ but we can change that easily changing the x_c parameter in the code
- n_{asx} the refractive index of x_c ppm of CO_2 at T=20°C, 101325 Pa, 0 % humidity
- Constants involved in the standard phase and group refractivities of water vapor [4]
  $\omega_0 = 295.235\,[\mu\mathrm{m}^{-2}],\ \omega_1 = 2.6422\,[\mu\mathrm{m}^{-2}],\ \omega_2 = -0.032380\,[\mu\mathrm{m}^{-4}],\ \omega_3 = 0.004028\,[\mu\mathrm{m}^{-6}]$
- $n_{\omega s}$ the refractive index of water vapour at T=20°C, 1333 Pa
- ρ_a [kg/m^3] the humid air density calculated equation (34)
- ρ_{axs} [kg/m^3] the density of dry air at standard conditions calculated equation (35)
- ρ_ω [kg/m^3] the density of water vapour calculated equation (36)
- $\rho_{\omega s}$ [kg/m^3] the density of water vapour at standard conditions calculated equation (37)

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Calcul of the humid air density

$$\rho_a = \frac{P\,M_a}{Z\,R\,T}\left[1 - x_V\left(1 - \frac{M_V}{M_a}\right)\right] \tag{34}$$

$$P = P_0\left(1 - \frac{\Delta T\,H}{T}\right)$$

$$M_a = (28.9635 + 12.011(xCO2 - 0.0004)) \times 10^{-3}$$

$$Z = 1 - \frac{P}{T}\left[a_0 + a_1\,t + a_2\,T^2 + (b_0 + b_1\,t) * x_v + (c_0 + c_1\,t) * x_v^2\right] + \frac{p^2}{T^2}\left(d + e * xv^2\right)$$

with :

- $P_0 = 1.01325 \times 10^5$ [Pa] the normal pressure at altitude 0 m;
- $\Delta T = 0.0065$ [K] the vertical gradient of temperature (0.65K for 100 m) [5]
- $T_0 = -10$ [°C] mean temperature from DAG-AWOS1
- $H = 3170$[m] Karakaya altitude
- M_a [kg/mol] the density of dry air with $xCO2 = 0.0004$ [6]
- $M_V = 18.01528 \times 10^{-3}$ [kg/mol] the mole mass of water
- $R = 8.314510$ [J/mol/K] the molar gas constant
- Z the compressibility with t the temperature in [°C] and the following constants and parameters [6] :

  - $a0 = 1.58123 \times 10^{-6}$ [K*Pa^{-1}], $a1 = -2.9331 \times 10^{-8}$ [Pa^{-1}], $a2 = 1.1043 \times 10^{-10}$ [(K*Pa)$^{-1}$], $b0 = 5.707 \times 10^{-6}$ [K*Pa^{-1}], $b1 = -2.051 \times 10^{-8}$ [Pa^{-1}], $c0 = 1.9898 \times 10^{-4}$ [K*Pa^{-1}] $c1 = -2.376 \times 10^{-6}$ [Pa^{-1}], $d = 1.83 \times 10^{-11}$ [K^{2}Pa^{-2}], $e = -0.765 \times 10^{-8}$ [K^{2}Pa^{-2}]
  - x_v =RH $f\frac{P_{sv}}{P}$ mole fraction of water vapour, RH the relative humidity (taken here as 0.8)
  - $f = \alpha + \beta P + \gamma T^2$ increasing factor ($\alpha = 1.00062$ [-], $\beta = 3.14 \times 10^{-8}$ [Pa], $\gamma = 5.6 \times 10^{-7}$ [K^{-2}])
  - $Psv = \exp\left(AT^2 + BT + C + \frac{D}{T}\right)$ the saturation vapour pressure of moist air ($A = 1.2378847 \times 10^{-5}$ [K^{-2}], $B = -1.9121316 \times 10^{-2}$ [K^{-1}], $C = 33.93711047$ [-], $D = -6.4341645 \times 10^3$ [K])

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Calcul of density of dry air at standard conditions

$$\rho_{axs} = \frac{P_0}{R_{gas} T_0} \tag{35}$$

with $R_{gas} = 287.05$ [J/kg/K] [7].

Calcul of the water vapour density

$$\rho_\omega = \frac{P_w M_V}{RT} \tag{36}$$

With :

- $M_V = 18.01528 \times 10^{-3}$ [kg/mol] the mole mass of water
- P_w [mb] partial pressure of water vapour (mmHg2Pa = 133.322365 and $Psat$ [mmHg] valid between -50°C and 200°C)

$$\begin{aligned} Psat &= \exp\left(46.784 - \frac{6435}{T} - 3.868 \log(T)\right) \\ Pw &= Psat * RH * \text{mmHg2Pa} \end{aligned}$$

Calcul of the water vapour density at standard conditions

$$\rho_{\omega s} = \frac{P_{w0} M_V}{RT_0} \tag{37}$$

With :

- $M_V = 18.01528 \times 10^{-3}$ [kg/mol] the mole mass of water
- P_{w0} [mb] partial pressure of water vapour (mmHg2Pa = 133.322365 and $Psat0$ [mmHg] valid between -50°C and 200°C)

$$\begin{aligned} Psat0 &= \exp\left(46.784 - 6435/(T0) - 3.868 \log(T0)\right) \\ P_{w0} &= Psat0 * RH * \text{mmHg2Pa} \end{aligned}$$

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C.3.2 Atmosphere refraction equation

The refraction of the atmosphere is computed with the equation [8]:

$$R(\lambda, z) = \kappa\,(n(\lambda)-1)\,(1-\beta)\tan(z) - \kappa\,(n(\lambda)-1)\left(\beta - \frac{(n(\lambda)-1)}{2}\right)\tan^3(z) \tag{38}$$

with

- $\beta = 0.001254\left(\frac{T(K)}{273.15}\right)$ the effective height of the observatory above the surface of the earth [8]
- $\kappa = 1$ for a spherical Earth surface [8] or instrumental correction no more useful [9]

This equation (38) is valid only for zenith angle ¡ 75°. We can even neglect the second term for zenith angles ¡ 65° [10]. The limit of zenith angle of our computation is set to 70°so the entire equation is set.

C.3.3 Glasses for the ADC

The choice of glass depends on the range of zenith angles, the wavelength interval, the maximum size of the blanks and the cost [11]). Most of the glass couple in the Amici conception are flint/crown pairs plus an anomalous dispersion glass inserted for the triplet design.
The glass choice is made in a data base built from Schott and Ohara catalogues. The data are taken from [12] where we can download glasses properties from many different suppliers. The data are sorted to extract Sellmeier coefficients [13] in order to compute the refractive index of each glass depending on the wavelength with the equation (39).

$$n^2(\lambda) = 1 + \frac{B_1\lambda^2}{\lambda^2 - C_1} + \frac{B_2\lambda^2}{\lambda^2 - C_2} + \frac{B_3\lambda^2}{\lambda^2 - C_3} \tag{39}$$

with B_i and C_i Sellmeier coefficients for λ in μm [13].

The final glass data base with the refractive index reported is generated for three wavelengths : le smallest, the highest and the mean wavelength. These are for our case the Nuvu camera [14] bandwidth limits $\lambda_{min} = 0.3\,[\mu m]$ and $\lambda_{max} = 1.0\,[\mu m]$ and the maximum quantum efficiency corresponding wavelength $\lambda_{mean} = 0.6\,[\mu m]$.

C.3.4 Beam propagation through the ADC

In order to compute de dispersion through the ADC a geometrical and refraction set of equations have been implemented. The sequence of these equations is listed bellow. I will follow the light on the figure 17 to describe each step.

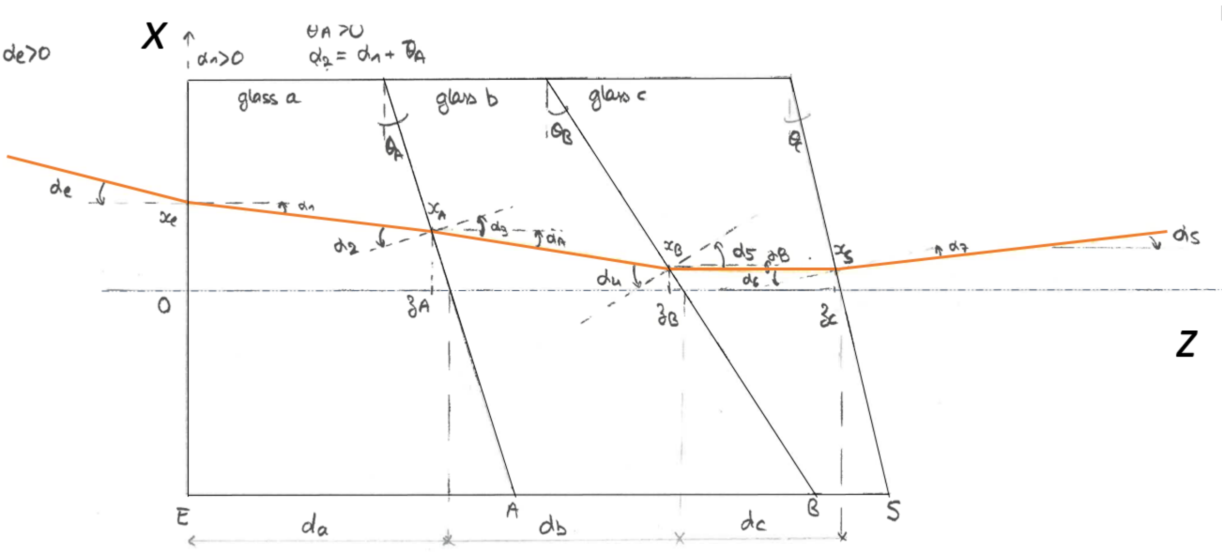

Figure 17: Sketch of the refraction through the ADC

The refractive index of the I-th glass is $n_I(\lambda_i)$. To simplify the notation, we just write n_I. We apply the sign convention described in the course material [15].

At the entrance we have the vector $\begin{bmatrix} x_E \\ z_E = 0 \\ n_0 \sin \alpha_E \end{bmatrix}$

At the entrance 0E, we have a refraction $\begin{bmatrix} x_E \\ z_E = 0 \\ n_A \sin \alpha_1 \end{bmatrix}$

which gives :

$\alpha_1 = \arcsin \left(\frac{n_0}{n_A} \sin \alpha_E \right)$

At the first interface AB, we arrive with the coordinates $\begin{bmatrix} x_A \\ z_A \\ n_A \sin(\alpha_2) \end{bmatrix}$

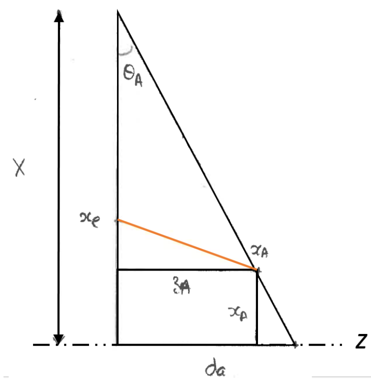

Figure 18: Decomposition of the parameters at an interface

We can determine this vector using the following equations from figure 18 :

$X = \frac{d_A}{\tan\theta_A}$; $\frac{X - x_A}{X} = \frac{z_A}{d_A}$; $z_A = \frac{x_E - x_A}{\tan\alpha_1}$

Then we find :

$$\begin{cases} x_A & = & \frac{x_E - d_A \tan\alpha_1}{1 - \tan\alpha_1 \tan\theta_A} \\ z_A & = & \frac{x_E - x_A}{\tan\alpha 1} \\ \alpha_2 & = & \alpha_1 + \theta_A \end{cases}$$

At the interface AB, we have a refraction $\begin{bmatrix} x_A \\ z_A \\ n_B \sin\alpha_3 \end{bmatrix}$

which gives :

$\alpha_3 = \arcsin\left(\frac{n_A}{n_B}\sin\alpha_2\right)$

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At the interface BC, we arrive with the coordinates $\begin{bmatrix} x_B \\ z_B \\ n_B \sin(\alpha_4) \end{bmatrix}$

We can determine this vector using the following equations from figure 18 :

$X = \frac{x_A}{\tan\alpha_A}$; $\frac{x_B}{X-Q-d_B+m} = \tan\alpha_A$; $m = d_A + d_B - z_B$; $m = x_B \tan\theta_B$; $Q = d_A - z_A$; $\alpha_A = \alpha_3 + \theta_A$

Then we find :

$$\begin{cases} x_B &= \frac{x_A - \tan\alpha_A(d_A+d_B-z_A)}{1-\tan\alpha_A \tan\theta_B} \\ z_B &= d_A + d_B - x_B \tan\theta_B \\ \alpha_4 &= \alpha_A + \theta_B \end{cases}$$

At the interface BC, we have a refraction $\begin{bmatrix} x_B \\ z_B \\ n_C \sin\alpha_5 \end{bmatrix}$

which gives :

$\alpha_5 = \arcsin\left(\frac{n_B}{n_C}\sin\alpha_4\right)$

At the interface C0, we arrive with the coordinates $\begin{bmatrix} x_C \\ z_C \\ n_C \sin(\alpha_6) \end{bmatrix}$

We can determine this vector using the following equations from figure 18 :

$X = \frac{x_B}{\tan\alpha_B}$; $\frac{x_C}{X-Q-d_C+m} = \tan\alpha_B$; $m = d_A + d_B - z_B$; $m = x_C \tan\theta_C$; $Q = d_A + d_B - z_B$; $\alpha_B = \alpha_5 + \theta_B$

Then we find :

$$\begin{cases} x_C &= \frac{x_B - \tan\alpha_B(d_A+d_B+d_C-z_B)}{1-\tan\alpha_B \tan\theta_C} \\ z_B &= d_A + d_B + d_C - x_C \tan\theta_C \\ \alpha_6 &= \alpha_B + \theta_C \end{cases}$$

At the interface C0, we have a refraction $\begin{bmatrix} x_C \\ z_C \\ n_0 \sin\alpha_7 \end{bmatrix}$

which gives :

$\alpha_7 = \arcsin\left(\frac{n_C}{n_0}\sin\alpha_6\right)$

The output angle with respect to the optical axis $\alpha_S = \alpha_7 - \theta_C$ This α_S is the dispersion angle of the prism : $R_{prism} = \alpha_S$.

C.3.5 The metric definition

Now we have computed the dispersion of the atmosphere and the ADC. We want the smallest total dispersion so we investigate all ADC configurations in order to minimize it. The metric we use comes from [10] to calculate the efficiency of our prism combination :

$$\mathrm{Eff}\,(\text{prism parameters}) = \sum_{\lambda_i} \left(R_{prism}\left(\lambda_i; \text{prism parameters}\right) - R_{atm}\left(\lambda_i\right)\right)^2 \tag{40}$$

C.3.6 Internal reflection

In parallel of the dispersion computation, we calculate the total internal reflection of the ADC. If the refractive index of two joint glasses are too different then we will loose a lot of intensity reflected on the interface. The external faces of the ADC will be coated. The determination of the total internal reflectivity is developed below.
The parameters are :

- $\alpha_Z = [\alpha_E; \alpha_1; \alpha_A; \alpha_B; \alpha_S]$ [rad] array of the rays angle wrt the optical axis
- $\alpha_I = [\alpha_1; \alpha_2; \alpha_3; \alpha_4; \alpha_5; \alpha_6; \alpha_7]$ [rad] array of the rays angles wrt the normal the vertex (refraction angle)
- $n = [n_0; n_A; n_B; n_C; n_0]$ medium refractive index
- $d = [z_A, z_B - z_A, z_C - z_B]$ [mm] distances on the optical axis between each medium change on the ray path

At the 1rst interface AB, we have :

$$r[1] = \frac{n[3]\cos\alpha_I[2] - n[2]\cos\alpha_I[3]}{n[3]\cos\alpha_I[2] + n[2]\cos\alpha_I[3]} \tag{41}$$

Taking the recursive initialization term $U(1) = r(1)$ we have for p = 2:

$$r[p] = \frac{n[p+1]\cos(\alpha_I[p]) - n[p]\cos(\alpha_I[p+1])}{n[p+1]\cos(\alpha_I[p]) + n[p]\cos(\alpha_I[p+1])} \tag{42}$$

$$U[p] = \frac{U[p-1] + r[p]\exp\left(\frac{2\pi}{\lambda}\cos(\alpha_Z[p])\ d[p-1]\right)}{1 + U[p-1]\,r[p]\exp\left(\frac{2\pi}{\lambda}\cos(\alpha_Z[p])\ d[p-1]\right)} \tag{43}$$

Using equation (43) we can calculate the reflection coefficient for $(p+1)$ number of glasses. In our case we have only 3 glasses so the total reflection (without taking external faces into account) is $U[2]$.

C.4 ADC design optimization

The efficiency of a combination (glass, angle, thickness) is computed by the function *Refraction_calculs_geometriques_20180302.m*. The *fmincon* Matlab algorithm is used to find minimum of this constrained nonlinear multivariable function. The parameters are the geometrical coefficients for one set of glass. When the optimum is found for a combination of glass, the efficiency and internal reflection are given as outputs.

The global matrix is made from two for-loops that look for all glasses from the catalogues Schott and Ohara. We decide to have the same glass type for the external wedges in order to simplify the anti-reflection coating process. This would give us the possibility to coat the prism at once and so reduce the cost. The refractive index of the middle wedge glass is chosen larger than the middle one to save effort and cost on the coating. In the optimization process on Matlab, when this refractive index (for λ_{mean} is larger than external one the loop stops. The outputs are set to 20'000 (arbitrarily large to oust them).
The wavelength for which the refraction is calculated are taken from the Nuvu camera quantum efficiency curve : $\lambda_{\text{min}} = 300$ nm, $\lambda_{\text{mean}} = 600$ nm and $\lambda_{\text{max}} = 1000$ nm. The zenith angle is set to 70°.

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D DimensionnementOAPs.py

```
# -*- coding: utf-8 -*-
"""
Created on Fri May 18 10:27:43 2018
This code compute the focal length of the OAPs and the image position.
@author: audrey.bouxin
"""
import numpy as np
import matplotlib.pyplot as plt

" Parameters"
DEG2RAD = np.pi/180;
theta_OAP0   = 20*DEG2RAD; #[rad]  pour l'OAP0
theta_OAP1 = theta_OAP0;
theta_OAP2   = 20*DEG2RAD; #[rad]  pour l'OAP2
#theta_OAP3   = 10*DEG2RAD;  #[rad]  pour l'OAP3
ExP2FP = 10338.74;        #[mm]
dExP   = 727.4046;        #[mm]
DDM    = 32.7;            #[mm]

# for OAP0
alpha = np.arctan(dExP*0.5/ExP2FP);  #[rad]
PFL_OAP0 =(DDM)/2/(-1/np.tan(theta_OAP1+alpha)+(1/(np.tan(theta_OAP1+alpha))
    **2+1)**0.5+1/np.tan(theta_OAP1-alpha)-(1/(np.tan(theta_OAP1-alpha))
    **2+1)**0.5);    #[mm] pour l'OAP0
EFL_OAP0 = 2*PFL_OAP0/(1+np.cos(theta_OAP1));
print('PFL_OAP0 (estim) : ', PFL_OAP0)


PFL_OAP0 = 445.0;
print('PFL_OAP0 : ', PFL_OAP0)
EFL_OAP0 = 2*PFL_OAP0/(1+np.cos(theta_OAP1))
print('EFL_OAP0 : ', EFL_OAP0)
po = -(PFL_OAP0+ExP2FP);
pi = po*PFL_OAP0/(po+PFL_OAP0);
print('pi (DM)  : ', pi)
```

```
# # pour l'OAP1
fOAP1 = PFL_OAP0

 # pour l'OAP2
print('------------------------------OAP2 : ')
TTangle = 10*DEG2RAD;
DTT = 10#*np.cos(TTangle);                  #[mm]
fOAP2 = DTT/DDM*fOAP1;        #[-]
alpha = np.arctan(DTT*0.5/fOAP2);      #[rad]
PFL_OAP2 = (DTT/2)/2/(-1/np.tan(theta_OAP2+alpha)+(1/(np.tan(theta_OAP2+
    alpha))**2+1)**0.5+1/np.tan(theta_OAP2)-(1/(np.tan(theta_OAP2))**2+1)
    **0.5)  #[mm] pour l'OAP2
EFL_OAP2 = 2*PFL_OAP2/(1+np.cos(theta_OAP2))
print('PFL_OAP2 : ', PFL_OAP2)
#
# # Pour l'OAP3
# pyr_roof = 20;   #[um]
# lambda    = 0.6;  #[um]
# fnum = 60;#2*pyr_roof/lambda;
# alpha = atan(1/(fnum*2));          #[rad] pour l'OAP3
# DTT = 25.11;
# PFL_OAP3 = (DTT/2)/2/(-1/tan(theta_OAP3+alpha)+(1/(tan(theta_OAP3+alpha))
    **2+1)**0.5+1/tan(theta_OAP3)-(1/(tan(theta_OAP3))**2+1)**0.5)    #[mm]
    pour l'OAP3
# EFL_OAP3 = 2*PFL_OAP3/(1+cos(theta_OAP3))
#
```

../DimensionnementOAPs.py