

# Clifford Analysis and a Noncommutative Gelfand Representation

Colin Roberts

# Overview

**1** Introduction

**2** Clifford analysis

**3** Gelfand theory

**4** Future work

**5** Conclusions

# Section 1

## **Introduction**

# Motivating problems

- *Electrical Impedance Tomography (EIT)* asks whether one can determine the conductivity of a medium from the voltage-to-current map.
- The *Calderón problem* replaces the medium with a manifold  $M$ , conductivity with  $g$ , and replaces the voltage-to-current map with the Dirichlet-to-Neumann operator  $\Lambda$ .

# Other questions

- What topological information can we retrieve from functions on a manifold  $M$ ?
- Do these functions also contain geometric information such as metric data?
- How much can we learn about  $M$  if our data is supported only on the boundary?

## Subsection 1

### **Preliminaries**

- *Clifford algebra* originated in 1878 with William Kingdon Clifford's work that extends Hermann Grassmann's *exterior algebra*.
- *Clifford analysis* arrived in the 1980's due to Hestenes, Sobczyk, Sommen, Brackx, and Delenghe in order to enrich Élie Cartan's *differential forms*. See: [Hestenes, Sobczyk: 1984] and [Doran, Lasenby: 2003].

# Clifford algebras

Let  $V$  be a vector space over a field  $\mathbb{F}$  with quadratic form  $Q$ .

- Define the tensor algebra

$$\mathcal{T}(V) := \bigoplus_{j=0}^{\infty} V^{\otimes j} = \mathbb{F} \oplus (V \otimes V) \oplus (V \otimes V \otimes V) \oplus \dots.$$

- Form the *Clifford algebra* via a quotient

$$Cl(V, Q) := \mathcal{T}(V) / \langle \mathbf{v} \otimes \mathbf{v} - Q(\mathbf{v}) \rangle.$$



# Geometric and exterior algebras

- Given a (pseudo) inner product  $g$ , we set  $Q(\cdot) = g(\cdot, \cdot)$  and define a *geometric algebra*

$$\mathcal{G} := Cl(V, g).$$

- The *exterior algebra* is given by

$$\bigwedge(V) := Cl(V, 0).$$

# Algebra structure

Multiplication in  $\mathcal{G}$  is seen by looking at how  $\otimes$  acts in the quotient.

- Given  $\mathbf{u}, \mathbf{v} \in \mathcal{G}$  we can take the product

$$\mathbf{u}\mathbf{v} = \underbrace{\mathbf{u} \cdot \mathbf{v}}_{\text{scalar}} + \underbrace{\mathbf{u} \wedge \mathbf{v}}_{\text{bivector}}.$$

- The scalar part is symmetric:  $\mathbf{u} \cdot \mathbf{v} = g(\mathbf{u}, \mathbf{v})$ .
- The bivector part is antisymmetric:  $\mathbf{u} \wedge \mathbf{v} = -\mathbf{v} \wedge \mathbf{u}$ .

# Multivectors

- $\mathcal{G}$  is graded and of dimension  $2^n$ .
  - There are  $\binom{n}{r}$  elements in the space of grade- $r$  elements,  $\mathcal{G}^r$ , called  *$r$ -vectors*.
  - Those that are exterior products of  $r$  independent vectors are  *$r$ -blades*.  
E.g.,  $\mathbf{A}_r = \mathbf{v}_1 \wedge \cdots \wedge \mathbf{v}_r$ .
  - Elements of the even grade subalgebra,  $\mathcal{G}^+$ , are called *spinors*.
- The most general elements are *multivectors* and are given by

$$A = \sum_{r=0}^n \langle A \rangle_r,$$

where  $\langle A \rangle_r \in \mathcal{G}^r$  extracts the grade  $r$  part of  $A$ . So  $\mathcal{G} = \bigoplus_{r=0}^n \mathcal{G}^r$ .

# Algebraic Structure

- Extend the multiplication from vectors to multivectors.
- On homogeneous elements,

$$A_r B_s = \langle A_r B_s \rangle_{|r-s|} + \langle A_r B_s \rangle_{|r-s|+2} + \cdots + \langle A_r B_s \rangle_{r+s}$$

- The most important products are

$$A_r \cdot B_s := \langle A_r B_s \rangle_{|r-s|}$$

$$A_r \wedge B_s := \langle A_r B_s \rangle_{r+s}$$

$$A_r \rfloor B_s := \langle A_r B_s \rangle_{s-r}$$

$$A_r \lrcorner B_s := \langle A_r B_s \rangle_{r-s}$$

# Reciprocals and reverses

- Given any vector basis  $\mathbf{v}_i$ , define the *reciprocal vectors* by  $\mathbf{v}^i \cdot \mathbf{v}_j = \delta_j^i$ .
- The *reverse* of a multivector is extended linearly from the action on  $r$ -blades by

$$\mathbf{A}_r^\dagger = (\mathbf{v}_1 \wedge \cdots \wedge \mathbf{v}_r)^\dagger = \mathbf{v}_r \wedge \cdots \wedge \mathbf{v}_1.$$

# Inner product and norm

- Define the *multivector inner product* by

$$(A, B) := \langle A^\dagger B \rangle$$

which is bilinear, symmetric, and positive definite if  $g$  is positive definite.

- Define the *multivector norm* by

$$|A| := \sqrt{(A, A)}.$$

# Adjoint

Note the reverse acts as an adjoint by

$$(CA, B) = (A, C^\dagger B)$$

$$(AC, B) = (A, BC^\dagger).$$

# Pseudoscalars

- *Pseudoscalars* are the grade- $n$  elements. For example, the volume element

$$\boldsymbol{\mu} = \mathbf{v}_1 \wedge \cdots \wedge \mathbf{v}_n.$$

- We define the *unit pseudoscalar* by

$$\boldsymbol{I} := \frac{1}{|\boldsymbol{\mu}|} \boldsymbol{\mu}.$$



# Blades and subspaces

- If  $g$  is positive definite all blades are invertible [Chisholm: 2012].
- If  $|\mathbf{A}_r| = 1$ , then  $\mathbf{A}_r$  is a *unit blade*.
- Unit  $r$ -blades correspond to  $r$ -dimensional subspaces so they correspond to points in  $\text{Gr}(r, n)$ .

# Duality

- Given any multivector  $A$ , we can take its *dual*

$$A^\perp := A\mathbf{I}^{-1}.$$

- Note  $A_r^\perp \in \mathcal{G}^{n-r}$ , like the Hodge star  $\star$ .

# Projection and rejection

- The *projection* of  $B$  into a subspace  $\mathbf{A}_r$  by

$$P_{\mathbf{A}_r}(B) := B \rfloor \mathbf{A}_r \mathbf{A}_r^{-1}$$

- The *rejection* by

$$R_{\mathbf{A}_r}(B) := B \wedge \mathbf{A}_r \mathbf{A}_r^{-1}.$$

- Both are grade preserving.

# Examples

- Define  $\mathcal{G}_{p,q}$  by letting  $\mathbf{e}_i^2 = -1$  for  $i = 1, \dots, p$  and  $\mathbf{e}_i^2 = +1$  otherwise.
- **Claim:**  $\mathbb{H}$  arises naturally as the even subalgebra  $\mathcal{G}_3^+ := \mathcal{G}_{0,3}^+$ .
- **Claim:**  $\mathbb{C}$  arises naturally as the even subalgebra  $\mathcal{G}_2^+ := \mathcal{G}_{0,2}^+$ .
  - Take the standard basis  $\mathbf{e}_1, \mathbf{e}_2$ , and define  $\mathbf{B}_{12} = \mathbf{e}_1 \mathbf{e}_2$  and note  $\mathbf{B}_{12}^2 = -1$ .  
Thus,

$$(u_1 + v_1 \mathbf{B}_{12})(u_2 + v_2 \mathbf{B}_{12}) = u_1 u_2 - v_1 v_2 + (u_1 v_2 + u_2 v_1) \mathbf{B}_{12}.$$

- Right multiplication by  $\mathbf{B}_{12}$  rotates counter-clockwise by  $\pi/2$ .

## Section 2

# Clifford analysis

# Multivector Fields

- Let  $M$  be a smooth, compact, connected, and oriented  $n$ -dimensional Riemannian manifold with metric  $g$ .
- Idea: Form the Clifford algebras on tangent spaces.
  - Each  $Cl(T_p M, g_p)$  is a *geometric tangent space* which we glue together to form

$$Cl(TM, g) := \bigsqcup_{p \in M} Cl(T_p M, g_p).$$

- The space of *(smooth) multivector fields* is

$$\mathcal{G}(M) := \{C^\infty\text{-smooth sections of } Cl(TM, g)\}.$$

- Retain the same naming scheme as before.

# Multivector derivative

On  $M$ , take the unique Levi-Civita connection  $\nabla$  and covariant derivative  $\nabla_{\mathbf{u}}$ .

- $\nabla_{\mathbf{u}}$  is extended to multivectors and is grade preserving [Schindler: 2018],

$$\nabla_{\mathbf{u}} A_r = \langle \nabla_{\mathbf{u}} A_r \rangle_r.$$

- $\nabla_{\mathbf{u}}$  is compatible with dot and wedge since

$$\nabla_{\mathbf{u}}(A \cdot B) = (\nabla_{\mathbf{u}} A) \cdot B + A \cdot (\nabla_{\mathbf{u}} B)$$

$$\nabla_{\mathbf{u}}(A \wedge B) = (\nabla_{\mathbf{u}} A) \wedge B + A \wedge (\nabla_{\mathbf{u}} B).$$

# Gradient

- Define the *gradient* (or *Dirac operator*) locally by

$$\nabla = \sum_{i=1}^n \mathbf{v}^i \nabla_{\mathbf{v}_i}$$

- $\nabla$  acts as a vector in  $\mathcal{G}(M)$  and obeys the Leibniz rule

$$\nabla(AB) = \dot{\nabla} \dot{A} B + \dot{\nabla} A \dot{B}.$$

- Note  $\nabla^2 = \Delta$ , the Laplace-Beltrami operator.



# Example

- In  $\mathcal{G}_3(\mathbb{R}^3)$ ,  $\nabla_{\mathbf{e}_i}$  is the partial derivative.
- Take a vector field  $\mathbf{v}$ , then

$$\nabla \mathbf{v} = \underbrace{\nabla \cdot \mathbf{v}}_{\text{divergence}} + \underbrace{\nabla \wedge \mathbf{v}}_{\text{curl}}.$$

- Specifically,

$$\text{curl}(\mathbf{v}) = (\nabla \wedge \mathbf{v})^\perp$$

# Differential forms

- Define the *r-dimensional directed measure*

$$dX_r := \mathbf{v}_{j_1} \wedge \cdots \wedge \mathbf{v}_{j_r} dx^{j_1} \cdots dx^{j_r}$$

where  $1 \leq j_1 < \cdots < j_r \leq n$  and summation is implied.

- Define an *r-form*  $\alpha_r$  by

$$\alpha_r = A_r \cdot dX_r^\dagger$$

where  $A_r = \frac{1}{r!} \alpha_{i_1 \dots i_r} \mathbf{v}^{i_1} \wedge \cdots \wedge \mathbf{v}^{i_r}$ .

- Refer to  $A_r$  the *multivector equivalent* of  $\alpha_r$ .

# Exterior algebra and calculus

- Given an  $r$ - and  $s$ -form  $\alpha_r$  and  $\beta_s$  we have

$$\alpha_r + \beta_s = (A_r + B_s) \cdot dX_r^\dagger, \quad \alpha_r \wedge \beta_s = (A_r \wedge B_s) \cdot dX_{r+s}^\dagger.$$

- The exterior derivative on multivector equivalents is

$$d\alpha_r = (\nabla \wedge A_r) \cdot dX_{r+1}^\dagger$$

- The Hodge star on multivector equivalents is

$$\star \alpha_r = (\mathbf{I}^{-1} A_r)^\dagger \cdot dX_{n-r}^\dagger$$

# Volume form

- The *volume form* on  $M$  is given in local coordinates by

$$\mu = \sqrt{|g|} dx^1 \cdots dx^n = \mathbf{I}^{-1} \cdot dX_n$$

- We integrate scalar fields  $A_0$  on  $M$  by

$$\int_M A_0^\perp \cdot dX_n = \int_M A_0 \mu.$$

# Multivector field inner product

- We define the  $L^2$ -inner product on multivector fields by

$$\ll A, B \gg := \frac{1}{\text{vol}(M)} \int_M (A, B) \mu$$

- This realizes the  $r$ -form inner product

$$\int_M \alpha_r \wedge \star \beta_r = \int_M \langle A_r^\dagger B_r \rangle \mu = \text{vol}(M) \ll A, B \gg$$

- $\ll A_r, B_s \gg = 0$  when  $r \neq s$  so the  $L^2$ -direct sum agrees with the grade based direct sum.

# Boundary

- On  $\partial M$ , the boundary pseudoscalar  $\mathbf{I}_\partial$  induces the boundary normal

$$\boldsymbol{\nu} = \mathbf{I}_\partial^\perp.$$

- The boundary volume form is

$$\mu_\partial := \mathbf{I}_\partial^{-1} \cdot dX_{n-1}$$

and we define the *multivector field inner product*

$$\ll A, B \gg_\partial := \frac{1}{\text{vol}(M)} \int_{\partial M} (A, B) \mu_\partial.$$

# Multivector valued integrals

We can define a multivector valued integral on an oriented submanifold  $R$  by taking  $A \in \mathcal{G}(M)$  and computing

$$\int_R A \mathbf{I}_R \mu_R.$$

## Theorem (Hestenes, Sobczyk, 1984)

*Let  $A, B \in \mathcal{G}(M)$ , then*

$$\begin{aligned}\int_M \dot{A} \dot{\nabla} \mathbf{I} \mu &= \int_{\partial M} A \mathbf{I} \partial \mu \partial \\ \int_M \mathbf{I} \nabla B \mu &= \int_{\partial M} \mathbf{I} \partial B \mu \partial \\ \int_M \dot{A} \dot{\nabla} \mathbf{I} B \mu &= (-1)^n \int_M A \mathbf{I} \nabla B \mu + \int_{\partial M} A \mathbf{I} \partial B \mu \partial.\end{aligned}$$



## Theorem

*We have the Green's formula for the gradient*

$$\ll A, \mathbf{I} \nabla B \gg = (-1)^n \ll \nabla A, \mathbf{I} B \gg + \ll A, \mathbf{I} \partial B \gg_{\partial} .$$

*Proof.* Fix  $A^{\dagger}, B \in \mathcal{G}(M)$  and note that

$$\begin{aligned} \int_M A^{\dagger} \mathbf{I} \nabla B \mu &= (-1)^n \int_M \dot{A}^{\dagger} \dot{\nabla} \mathbf{I} B \mu + \int_{\partial M} A^{\dagger} \mathbf{I} \partial B \mu_{\partial} \\ &= (-1)^n \int_M (\nabla A)^{\dagger} \mathbf{I} B \mu + \int_{\partial M} A^{\dagger} \mathbf{I} \partial B \mu_{\partial} . \end{aligned}$$

Take the scalar part and divide by  $\text{vol}(M)$  to find

$$\ll A, \mathbf{I} \nabla B \gg = (-1)^n \ll \nabla A, \mathbf{I} B \gg + \ll A, \mathbf{I} \partial B \gg_{\partial} .$$

# Special fields

- Define the *monogenic fields*

$$\mathcal{M}(M) := \{A \in \mathcal{G}(M) \mid \nabla A = 0\}.$$

- Let  $f = u + v\mathbf{B} \in \mathcal{G}_2^+(\mathbb{R}^2)$  then  $\nabla f = 0$  yields the Cauchy-Riemann equations

$$\frac{\partial u}{\partial x} = \frac{\partial v}{\partial y}, \qquad \frac{\partial u}{\partial y} = -\frac{\partial v}{\partial x}.$$

- Define the *gradients*

$$\nabla \mathcal{G}(M) := \{\nabla A \mid A \in \mathcal{G}(M) \text{ and } A|_{\partial M} = 0\}.$$

# Cauchy kernel

- For the remainder, take  $M$  imbedded in  $\mathbb{R}^n$  with  $n \geq 2$ .
- Define the vector field

$$E(x) := \frac{1}{S_n} \frac{x}{|x|^n}$$

where  $S_n$  is the surface area of the unit ball.

- Note,

$$\nabla E(x) = -\dot{E}(x)\dot{\nabla} = \delta(x).$$

- Define the *Cauchy kernel* by  $G(x, x') := E(x' - x)$ .

# Cauchy integral

- Let  $A \in \mathcal{M}(M)$ , then we have the *Cauchy integral formula*

$$A(x) = (-1)^n \mathbf{I}^{-1} \int_{\partial M} G(x, x') \mathbf{I}_{\partial}(x') A(x') \mu_{\partial}(x').$$

- This uniquely determines a monogenic field from boundary values.

### Lemma

*Let  $A \in \mathcal{M}(M)$  be such that  $A|_{\partial M} = 0$ . Then  $A = 0$  on all of  $M$ .*

### Lemma

*Fix a multivector field  $A \in \mathcal{G}(M)$ . If*

$$\ll A, B \gg = 0$$

*for all  $B \in \mathcal{G}(M)$  with  $B|_{\partial M} = 0$ , then  $A = 0$ .*

## Theorem (Clifford-Hodge-Morrey Decomposition)

*The space of multivector fields  $\mathcal{G}(M)$  has the  $L^2$ -orthogonal decomposition*

$$\mathcal{G}(M) = \mathcal{M}(M) \oplus \mathbf{I}\nabla\mathcal{G}(M).$$

*Proof.*

■ *Orthogonality:* Let  $A \in \mathcal{M}(M)$  and  $\mathbf{I}\nabla B \in \mathbf{I}\nabla\mathcal{G}(M)$  and note

$$\ll A, \mathbf{I}\nabla B \gg = (-1)^n \ll \nabla A, \mathbf{I}B \gg + \ll A, \mathbf{I}\partial B \gg = 0,$$

by the multivector Green's formula.

- Let  $C \in \mathcal{G}(M)$  be in the orthogonal complement to  $\mathbf{I}\nabla\mathcal{G}(M)$ .
- Use the Cauchy integral formula, construct a monogenic field  $\tilde{C}$  from  $C|_{\partial M}$  and note  $C = \tilde{C} + C_0$  where  $C_0|_{\partial M} = 0$ .
- Note,

$$0 = \langle\langle C, \mathbf{I}\nabla B \rangle\rangle = \langle\langle \nabla C_0, \mathbf{I}B \rangle\rangle .$$

- By the previous lemmas, it must be that  $C_0 = 0$ . Hence the orthogonal complement to  $\mathbf{I}\nabla\mathcal{G}(M)$  is  $\mathcal{M}(M)$ .



# Comparing to Hodge-Morrey

- The Hodge-Morrey decomposition reads

$$\Omega^r(M) = \underbrace{\mathcal{E}_D^r(M)}_{\text{Im}(\nabla \wedge)} \oplus \underbrace{\mathcal{C}_N^r(M)}_{\text{Im}(\nabla \lrcorner)} \oplus \underbrace{\mathcal{H}^r(M)}_{\text{Ker}(\nabla)}.$$

via [Schwarz: 1995].

- Whereas the Clifford-Hodge-Morrey decomposition ignores specific grades

$$\mathcal{G}(M) = \mathcal{M}(M) \oplus \mathbf{I} \nabla \mathcal{G}(M).$$

## Section 3

### **Gelfand theory**

# Open questions

- In [Belishev: 2003], we see an algebraic proof for the 2-dimensional Calderón problem.
- In [Belishev, Vakulenko: 2017], we see a proof for a noncommutative Gelfand representation using quaternion fields for a ball  $\mathbb{B}$  in  $\mathbb{R}^3$ .
- Belishev and Vakulenko ask whether this is true in higher dimensions.
- We prove an analogous result for an arbitrary  $\mathbb{B}$  in  $\mathbb{R}^n$ .
- This approach can hopefully be used to prove the analogous result for any smooth orientable Riemannian manifold  $M$ .

## 2-dimensional BC method

The boundary control (BC) method is implemented in [Belishev: 2003] in the following manner.

- Determine the algebra  $\mathcal{A}(M)$  of holomorphic functions on  $M$  from continuous function algebra on the boundary  $\mathcal{A}(\partial M)$  using  $\Lambda$ .
- The classical Gelfand representation shows  $\mathcal{A}(M)$  is homeomorphic to  $M$  via the weak- $*$  topology.
- Functions in  $\mathcal{A}(M)$  determine the complex structure on  $M$ .
- Thus, we can find a  $g$  that is conformal with the complex structure.

# Subsurface spinor fields

- Let  $\mathbf{B} \in \mathcal{G}(M)$  be a constant unit 2-blade, then  $f_+ \in \mathcal{G}^+(M)$  satisfying

$$f_+ = P_{\mathbf{B}} \circ f_+ \circ P_{\mathbf{B}}$$

is a *subsurface spinor field*. Let  $\mathcal{G}_{\mathbf{B}}^+(M)$  denote the space such fields.

- The space of monogenic subsurface spinors

$$\mathcal{A}_{\mathbf{B}}(M) = \{f_+ \in \mathcal{G}_{\mathbf{B}}^+(M) \mid \nabla f_+ = 0\}$$

is a commutative unital Banach algebra.

# Functionals

- Define the *spinor dual*  $\mathcal{M}^*(M)$  as the continuous right  $\mathcal{G}_n$ -module homomorphisms

$$\mathcal{M}^*(M) := \{l: \mathcal{M}^+(M) \rightarrow \mathcal{G}_n^+ \mid l(fs+g) = l(f)s + l(g), \forall f, g \in \mathcal{M}(M), s \in \mathcal{G}_n^+\}$$

and refer to the elements as *spin functionals*.

- Assert the weak-\* topology on  $\mathcal{M}^*(M)$  so that every  $x \in M$  corresponds to a continuous map on  $\mathcal{M}^*(M)$ .

# Characters

- Define the algebra  $\mathbb{A}_{\mathbf{B}}$  to be the algebra generated by 1 and  $\mathbf{B}$ .
- The *spinor spectrum*  $\mathfrak{M}(M)$  is the set of algebra homomorphisms

$$\mathfrak{M}(M) := \{ \delta \in \mathcal{M}^*(M) \mid \delta(f) \in \mathbb{A}_{\mathbf{B}}, \delta(fg) = \delta(f)\delta(g), \forall f, g \in \mathcal{A}_{\mathbf{B}}(M), \mathbf{B} \in \text{Gr}(2, n) \}$$

and refer to the elements as *spin characters*.

- One example of such characters are point evaluations  $\delta(f) = f(x^\delta)$ .
- We show these are the only elements in the spectrum.

## $z$ analogs

- Take the standard basis for  $\mathbb{R}^n$ , and consider  $M = \mathbb{B}_{R,w}$ .
- Let  $\mathbf{B}_{ij} = \mathbf{e}_i \mathbf{e}_j$ , and define

$$z_{ij} = x_j - x_i \mathbf{B}_{ij}$$

- Note  $z_{ij} \in \mathcal{A}_{\mathbf{B}_{ij}}(\mathbb{B}_{R,w})$ .



# Monogenic polynomials

- Let  $\sigma$  be a permutation of  $\{2, 3, \dots, n\}$ , then

$$p_{j_2 \dots j_n}(x) = \frac{1}{j!} \sum_{\text{permutations}} z_{1\sigma(1)}(x) \cdots z_{1\sigma(j)}(x)$$

is a monogenic homogeneous polynomial of degree  $j$ .

- Collect these into the set of *monogenic polynomials*

$$\mathcal{M}^{\mathcal{P}}(\mathbb{B}_{R,w}) = \left\{ \sum_{j=0}^N \left( \sum_{\substack{j_2 \dots j_n \\ j_2 + \dots + j_n = j}} p_{j_2 \dots j_n} a_{j_2 \dots j_n} \right) \mid j_2 + \dots + j_n = j, N \in \mathbb{N}, a_{j_2 \dots j_n} \in \mathcal{G}_n \right\}.$$

## Lemma (Density)

*The space  $\mathcal{M}^{\mathcal{P}}(\mathbb{B}_{R,w})$  is dense in  $\mathcal{M}(\mathbb{B}_{R,w})$ .*

*Proof sketch.*

- Let  $f \in \mathcal{M}(\mathbb{B}_{R,w})$  and use the Cauchy integral formula to define the coefficients  $a_{j_2 \dots j_n} \in \mathcal{G}_n^+$  by

$$a_{j_2 \dots j_n} = \int_{\partial \mathbb{B}_{R,w}} \frac{\partial^j G(w, y)}{\partial y_2^{j_2} \dots \partial y_n^{j_n}} \nu(y) f(y) \mu_{\partial}(y),$$

- Then

$$f(x) = \sum_{j=0}^{\infty} \left( \sum_{\substack{j_2 \dots j_n \\ j_2 + \dots + j_n = j}} p_{j_2 \dots j_n}(x - w) a_{j_2 \dots j_n} \right),$$

converges pointwise for  $x \in \mathbb{B}_{R,w}$  by **[Ryan, 2004]**.

# Idea

- By linearity, we can note that for  $\delta \in \mathfrak{M}(\mathbb{B}_{R,w})$

$$\delta(f(x)) = \sum_{j=0}^{\infty} \left( \sum_{\substack{j_2 \dots j_n \\ j_2 + \dots + j_n = j}} \delta(p_{j_2 \dots j_n}(x-w)) a_{j_2 \dots j_n} \right)$$

- On each monogenic polynomial

$$\delta(p_{j_2 \dots j_n}(x)) = \frac{1}{j!} \sum_{\text{permutations}} \delta((z_{1\sigma(1)}(x)) \cdots \delta(z_{1\sigma(j)}(x)))$$

by the multiplicativity of  $\delta$ .

## Lemma (Point evaluation)

Let  $\delta \in \mathfrak{M}(\mathbb{B}_{R,w})$  and  $z_{ij} \in \mathcal{A}_{\mathbf{B}_{ij}}(\mathbb{B}_{R,w})$ , then  $\delta(z_{ij}) = z_{ij}(x^\delta)$  for some  $x^\delta \in \mathbb{R}^n$ .

*Proof sketch.*

- We have  $\delta(z_{ij}) = \alpha_{ij} + \beta_{ij}\mathbf{B}_{ij}$ .
- Note  $z_{ij}\mathbf{B}_{ji} = -z_{ji}$  and  $z_{ij} = z_{kj} + z_{ik}\mathbf{B}_{kj}$  yield the relationships

$$\alpha_{ji} = -\beta_{ij} \quad \alpha_{ij} = \alpha_{kj} \quad \beta_{ij} = \beta_{ik} \quad \alpha_{ik} = -\beta_{kj}.$$

- The set of constants  $\alpha$  and  $\beta$  are determined by  $n$  independent numbers, so we can say  $\delta(z_{ij}) = z_{ij}(x^\delta)$  for some  $x^\delta \in \mathbb{R}^n$ .

## Lemma (Identification)

Let  $f \in \mathcal{M}(\mathbb{B}_{R,w})$ , then  $\delta(f) = f(x^\delta)$  for some  $x^\delta \in \mathbb{B}_{R,w}$ .

*Proof.*

- Fix  $\delta \in \mathfrak{M}(\mathbb{B}_{R,w})$  and suppose  $x^\delta \notin \mathbb{B}_{R,w}$ .
- Take a sequence  $x_n \rightarrow x^\delta$  with  $x_n \notin \mathbb{B}_{R,w}$ .
- Define  $G_n(x) := G(x, x_n)\mathbf{e}_1 \in \mathcal{M}^+(\mathbb{B}_{R,w})$ .
- Note,

$$\lim_{n \rightarrow \infty} \delta(G_n) = \lim_{n \rightarrow \infty} G_n(x^\delta)$$

so this sequence not converge due to a singularity at  $x^\delta$ .

- Hence, it must be that  $x^\delta \in \mathbb{B}_{R,w}$  by continuity of  $\delta$ .

## Theorem (Noncommutative Gelfand representation)

*For any  $\delta \in \mathfrak{M}(\mathbb{B}_{R,w})$ , there is a point  $x^\delta \in \mathbb{B}_{R,w}$  such that  $\delta(f) = f(x^\delta)$  for any  $f \in \mathcal{M}(\mathbb{B}_{R,w})$ . Given the weak- $*$  topology on  $\mathcal{M}^*(\mathbb{B}_{r,w})$ , the map*

$$\gamma: \mathfrak{M}(\mathbb{B}_{R,w}) \rightarrow \mathbb{B}_{R,w}, \quad \delta \mapsto x^\delta$$

*is a homeomorphism.*

*Proof.*

- The lemmas show that  $\gamma: \mathfrak{M}(\mathbb{B}_{R,w}) \rightarrow \mathbb{B}_{R,w}$  is bijective.
- To see that  $\gamma$  is a homeomorphism, take a sequence  $\delta_n \rightarrow \delta$  in  $\mathfrak{M}(\mathbb{B}_{R,w})$ .
- For  $f \in \mathcal{M}^+(\mathbb{B}_{R,w})$  we have

$$f(\gamma(\delta_n)) = f(x^{\delta_n}) = \delta_n(f) = \gamma^{-1}(x^{\delta_n})(f).$$

- Taking  $n \rightarrow \infty$  shows  $\gamma$  and  $\gamma^{-1}$  are continuous so  $\gamma$  is a homeomorphism.

## Section 4

### **Future work**



# Calderón problem

- Question: Let  $(M, g)$  be an unknown Riemannian manifold with known boundary  $\partial M$ . Consider the forward problem

$$\begin{cases} \Delta\omega = 0 & \text{in } M \\ \iota^*\omega = \phi & \text{on } \partial M \end{cases}$$

- Define the *Dirichlet-to-Neumann map* on forms by  $\Lambda\phi = \iota^*(\star d\omega)$ .
- Can we determine  $(M, g)$  from  $\Lambda$ ?

# Calderón problem

- This problem is equivalent to the electrical impedance tomography problem in dimension 3.
- The problem has been solved in dimension  $n = 2$  [Belishev: 2003].
- Solved in dimensions  $n \geq 3$  when  $M$  is an analytic manifold [Lassas, Taylor, Uhlmann: 2003].
- The smooth cases is still unsolved.

# Thoughts

- Even monogenic fields have harmonic components. Given a harmonic  $r$ -vector  $A_r$ , can we reconstruct a monogenic multivector containing  $A_r$ ?
- For  $n = 3$ , the scalar potential  $u$  and magnetic bivector field  $b$  are two parts of a monogenic field  $f = u + b$  due to Ohm's and Ampere's laws

$$-\nabla \wedge u = \mathbf{j} = \nabla \rfloor b.$$

- If  $\Lambda$  can provide us  $b|_{\partial M}$ , then we can possibly reconstruct  $\mathcal{M}^+(M)$ .
- Given the algebraic structure of each  $\mathcal{A}_B(M) \subset \mathcal{M}^+(M)$ , can this be used to determine  $g$ ?

# Other inverse problems

- Can the magnetic impedance tomography problem can provide some extra insight on the EIT problem?
- The Hodge-Morrey decomposition is an instrumental tool for boundary value problems that, for example, allows one to show that  $\Lambda$  determines the Betti numbers of  $M$  [Belishev, Sharafutdinov: 2008].
- Can the Clifford-Hodge-Morrey decomposition can allow us to work on other related boundary inverse problems?
- These problems could include spacetime problems where the metric  $g$  is of mixed signature.

## Section 5

# Conclusions

# Conclusion

- We have utilized multivector fields to serve as a meaningful generalization of both the complex numbers and differential forms.
- This provides a new way to decompose fields on domains of  $\mathbb{R}^n$  and this can likely be generalized to arbitrary compact orientable pseudo-Riemannian manifolds.
- Likewise, we have proven that the monogenic fields contain a wealth of topological information and this information is supported on the boundary by the Cauchy integral formula.

# Data Assimilation

- Over the past two years I have also worked with a team on developing new techniques for data assimilation.
- We have submitted *Model and Data Reduction for Data Assimilation: Particle Filters Employing Projected Forecasts and Data with Application to a Shallow Water Model*.
- We are continuing to work to apply our scheme to new models such as the Modular Arbitrary-Order Ocean-Atmospheric Model (MAOOAM).