

Clifford Analysis and a Noncommutative Gelfand Representation

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Overview

1 Introduction

2 Clifford analysis

3 Gelfand theory

4 Future work

5 Conclusions

Section 1

Introduction

Motivating problems

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- The *Calderón problem* replaces the medium with a manifold M , conductivity with g , and replaces the voltage-to-current map with the Dirichlet-to-Neumann operator Λ .

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- Do these functions also contain geometric information such as metric data?
- How much can we learn about M if our data is supported only on the boundary?

Subsection 1

Preliminaries

- *Clifford algebra* originated in 1878 with William Kingdon Clifford's work that extends Hermann Grassmann's *exterior algebra*.

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- *Clifford analysis* arrived in the 1980's due to Hestenes, Sobczyk, Sommen, Brackx, and Delenghe in order to enrich Élie Cartan's *differential forms*. See: [Hestenes, Sobczyk: 1984] and [Doran, Lasenby: 2003].

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$$\mathcal{T}(V) := \bigoplus_{j=0}^{\infty} V^{\otimes j} = \mathbb{F} \oplus (V \otimes V) \oplus (V \otimes V \otimes V) \oplus \dots.$$

- Form the *Clifford algebra* via a quotient

$$Cl(V, Q) := \mathcal{T}(V) / \langle \mathbf{v} \otimes \mathbf{v} - Q(\mathbf{v}) \rangle.$$

Geometric and exterior algebras

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$$\bigwedge(V) := Cl(V, 0).$$

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- The scalar part is symmetric: $\mathbf{u} \cdot \mathbf{v} = g(\mathbf{u}, \mathbf{v})$.
- The bivector part is antisymmetric: $\mathbf{u} \wedge \mathbf{v} = -\mathbf{v} \wedge \mathbf{u}$.

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E.g., $\mathbf{A}_r = \mathbf{v}_1 \wedge \cdots \wedge \mathbf{v}_r$.

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 - Elements of the even grade subalgebra, \mathcal{G}^+ , are called *spinors*.
- The most general elements are *multivectors* and are given by

$$A = \sum_{r=0}^n \langle A \rangle_r,$$

where $\langle A \rangle_r \in \mathcal{G}^r$ extracts the grade r part of A . So $\mathcal{G} = \bigoplus_{r=0}^n \mathcal{G}^r$.

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- The most important products are

$$A_r \cdot B_s := \langle A_r B_s \rangle_{|r-s|}$$

$$A_r \wedge B_s := \langle A_r B_s \rangle_{r+s}$$

$$A_r \rfloor B_s := \langle A_r B_s \rangle_{s-r}$$

$$A_r \lrcorner B_s := \langle A_r B_s \rangle_{r-s}$$

Reciprocals and reverses

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- The *reverse* of a multivector is extended linearly from the action on r -blades by

$$\mathbf{A}_r^\dagger = (\mathbf{v}_1 \wedge \cdots \wedge \mathbf{v}_r)^\dagger = \mathbf{v}_r \wedge \cdots \wedge \mathbf{v}_1.$$

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- Define the *multivector norm* by

$$|A| := \sqrt{(A, A)}.$$

Adjoint

Note the reverse acts as an adjoint by

$$(CA, B) = (A, C^\dagger B)$$

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- We define the *unit pseudoscalar* by

$$\boldsymbol{I} := \frac{1}{|\boldsymbol{\mu}|} \boldsymbol{\mu}.$$

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- If $|\mathbf{A}_r| = 1$, then \mathbf{A}_r is a *unit blade*.
- Unit r -blades correspond to r -dimensional subspaces so they correspond to points in $\text{Gr}(r, n)$.

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- Note $A_r^\perp \in \mathcal{G}^{n-r}$, like the Hodge star \star .

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- Both are grade preserving.

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 - Take the standard basis $\mathbf{e}_1, \mathbf{e}_2$, and define $\mathbf{B}_{12} = \mathbf{e}_1 \mathbf{e}_2$ and note $\mathbf{B}_{12}^2 = -1$.
Thus,

$$(u_1 + v_1 \mathbf{B}_{12})(u_2 + v_2 \mathbf{B}_{12}) = u_1 u_2 - v_1 v_2 + (u_1 v_2 + u_2 v_1) \mathbf{B}_{12}.$$

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- Right multiplication by \mathbf{B}_{12} rotates counter-clockwise by $\pi/2$.

Section 2

Clifford analysis

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- Retain the same naming scheme as before.

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$$\nabla_{\mathbf{u}} A_r = \langle \nabla_{\mathbf{u}} A_r \rangle_r.$$

- $\nabla_{\mathbf{u}}$ is compatible with dot and wedge since

$$\nabla_{\mathbf{u}}(A \cdot B) = (\nabla_{\mathbf{u}} A) \cdot B + A \cdot (\nabla_{\mathbf{u}} B)$$

$$\nabla_{\mathbf{u}}(A \wedge B) = (\nabla_{\mathbf{u}} A) \wedge B + A \wedge (\nabla_{\mathbf{u}} B).$$

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- Note $\nabla^2 = \Delta$, the Laplace-Beltrami operator.

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- Specifically,

$$\text{curl}(\mathbf{v}) = (\nabla \wedge \mathbf{v})^\perp$$

Differential forms

Differential forms

- Define the *r-dimensional directed measure*

$$dX_r := \mathbf{v}_{j_1} \wedge \cdots \wedge \mathbf{v}_{j_r} dx^{j_1} \cdots dx^{j_r}$$

where $1 \leq j_1 < \cdots < j_r \leq n$ and summation is implied.

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- Define an *r*-form a_r by

$$a_r = A_r \cdot dX_r^\dagger$$

where $A_r = \frac{1}{r!} a_{i_1 \dots i_r} \mathbf{v}^{i_1} \wedge \cdots \wedge \mathbf{v}^{i_r}$.

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- Refer to A_r the *multivector equivalent* of a_r .

Exterior algebra and calculus

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- Given r -forms a_r , b_r , and an s -form c_s , we have

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- The Hodge star on multivector equivalents is

$$\star a_r = (I^{-1} A_r)^\dagger \cdot dX_{n-r}^\dagger$$

Volume form

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- We integrate scalar fields A_0 on M by

$$\int_M A_0^\perp \cdot dX_n = \int_M A_0 \mu.$$

Multivector field inner product

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- $\ll A_r, B_s \gg = 0$ when $r \neq s$ so the L^2 -direct sum agrees with the grade based direct sum.

Boundary

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- The boundary volume form is

$$\mu_\partial := \mathbf{I}_\partial^{-1} \cdot dX_{n-1}$$

and we define

$$\ll A, B \gg_\partial := \frac{1}{\text{vol}(M)} \int_{\partial M} (A, B) \mu_\partial.$$

Multivector valued integrals

We can define a multivector valued integral on an oriented submanifold R by taking $A \in \mathcal{G}(M)$ and computing

$$\int_R A \mathbf{I}_R \mu_R.$$

Theorem (Hestenes, Sobczyk, 1984)

Let $A, B \in \mathcal{G}(M)$, then

$$\begin{aligned}\int_M \dot{A} \dot{\nabla} \mathbf{I} \mu &= \int_{\partial M} A \mathbf{I} \partial \mu \partial \\ \int_M \mathbf{I} \nabla B \mu &= \int_{\partial M} \mathbf{I} \partial B \mu \partial \\ \int_M \dot{A} \dot{\nabla} \mathbf{I} B \mu &= (-1)^n \int_M A \mathbf{I} \nabla B \mu + \int_{\partial M} A \mathbf{I} \partial B \mu \partial.\end{aligned}$$

Theorem

We have the Green's formula for the gradient

$$\ll A, \mathbf{I} \nabla B \gg = (-1)^n \ll \nabla A, \mathbf{I} B \gg + \ll A, \mathbf{I} \partial B \gg_{\partial} .$$

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- Let $f = u + v\mathbf{B} \in \mathcal{G}_2^+(\mathbb{R}^2)$ then $\nabla f = 0$ yields the Cauchy-Riemann equations

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$$\frac{\partial u}{\partial x} = \frac{\partial v}{\partial y}, \qquad \frac{\partial u}{\partial y} = -\frac{\partial v}{\partial x}.$$

- Define the *gradients*

$$\nabla \mathcal{G}(M) := \{\nabla A \mid A \in \mathcal{G}(M) \text{ and } A|_{\partial M} = 0\}.$$

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- Note,

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- Define the *Cauchy kernel* by $G(x, x') := E(x' - x)$.

Cauchy integral

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- Let $A \in \mathcal{M}(M)$, then we have the *Cauchy integral formula*

$$A(x) = (-1)^n \mathbf{I}^{-1} \int_{\partial M} G(x, x') \mathbf{I}_{\partial}(x') A(x') \mu_{\partial}(x').$$

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$$A(x) = (-1)^n \mathbf{I}^{-1} \int_{\partial M} G(x, x') \mathbf{I}_{\partial}(x') A(x') \mu_{\partial}(x').$$

- This uniquely determines a monogenic field from boundary values.

Lemma

Let $A \in \mathcal{M}(M)$ be such that $A|_{\partial M} = 0$. Then $A = 0$ on all of M .

Lemma

Fix a multivector field $A \in \mathcal{G}(M)$. If

$$\ll A, B \gg = 0$$

for all $B \in \mathcal{G}(M)$ with $B|_{\partial M} = 0$, then $A = 0$.

Theorem (Clifford-Hodge-Morrey Decomposition)

The space of multivector fields $\mathcal{G}(M)$ has the L^2 -orthogonal decomposition

$$\mathcal{G}(M) = \mathcal{M}(M) \oplus \mathbf{I}\nabla\mathcal{G}(M).$$

Proof.

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■ *Orthogonality:* Let $A \in \mathcal{M}(M)$ and $\mathbf{I}\nabla B \in \mathbf{I}\nabla\mathcal{G}(M)$ and note

$$\ll A, \mathbf{I}\nabla B \gg = (-1)^n \ll \nabla A, \mathbf{I}B \gg + \ll A, \mathbf{I}\partial B \gg = 0,$$

by the multivector Green's formula.

- Let $C \in \mathcal{G}(M)$ be in the orthogonal complement to $\mathbf{I}\nabla\mathcal{G}(M)$.

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$$0 = \langle\langle C, \mathbf{I}\nabla B \rangle\rangle = \langle\langle \nabla C_0, \mathbf{I}B \rangle\rangle .$$

- By the previous lemmas, it must be that $C_0 = 0$. Hence the orthogonal complement to $\mathbf{I}\nabla\mathcal{G}(M)$ is $\mathcal{M}(M)$.

Comparing to Hodge-Morrey

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- The Hodge-Morrey decomposition reads

$$\Omega^r(M) = \underbrace{\mathcal{E}_D^r(M)}_{\text{Im}(\nabla \wedge)} \oplus \underbrace{\mathcal{C}_N^r(M)}_{\text{Im}(\nabla \cdot)} \oplus \underbrace{\mathcal{H}^r(M)}_{\text{Ker}(\nabla)}.$$

via [Schwarz: 1995].

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via [Schwarz: 1995].

- Whereas the Clifford-Hodge-Morrey decomposition ignores specific grades

$$\mathcal{G}(M) = \mathcal{M}(M) \oplus \mathbf{I} \nabla \mathcal{G}(M).$$

Section 3

Gelfand theory

Open questions

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- Belishev and Vakulenko ask whether this is true in higher dimensions.
- We prove an analogous result for an arbitrary \mathbb{B} in \mathbb{R}^n .
- This approach can hopefully be used to prove the analogous result for any smooth orientable Riemannian manifold M .

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- The classical Gelfand representation shows the spectrum of $\mathcal{A}(M)$ is homeomorphic to M via the weak-* topology.
- Functions in $\mathcal{A}(M)$ determine the complex structure on M .
- Thus, we can find a g that is conformal with the complex structure.

Subsurface spinor fields

Subsurface spinor fields

- Let $\mathbf{B} \in \mathcal{G}(M)$ be a constant unit 2-blade, then $f_+ \in \mathcal{G}^+(M)$ satisfying

$$f_+ = P_{\mathbf{B}} \circ f_+ \circ P_{\mathbf{B}}$$

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- The space of monogenic subsurface spinors

$$\mathcal{A}_{\mathbf{B}}(M) = \{f_+ \in \mathcal{G}_{\mathbf{B}}^+(M) \mid \nabla f_+ = 0\}$$

is a commutative unital Banach algebra.

Functionals

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- Define the *spinor dual* $\mathcal{M}^*(M)$ as the continuous right \mathcal{G}_n^+ -module homomorphisms

$$\mathcal{M}^*(M) := \{l: \mathcal{M}^+(M) \rightarrow \mathcal{G}_n^+ \mid l(fs+g) = l(f)s+l(g), \forall f, g \in \mathcal{M}(M), s \in \mathcal{G}_n^+\}$$

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- Assert the weak-* topology on $\mathcal{M}^*(M)$ so that every $x \in M$ corresponds to a continuous map on $\mathcal{M}^+(M)$.

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- We show these are the only elements in the spectrum.

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- Note $z_{ij} \in \mathcal{A}_{\mathbf{B}_{ij}}(\mathbb{B}_{R,w})$.

Monogenic polynomials

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- Let σ be a permutation of $\{2, 3, \dots, n\}$, then

$$p_{j_2 \dots j_n}(x) = \frac{1}{j!} \sum_{\text{permutations}} z_{1\sigma(1)}(x) \cdots z_{1\sigma(j)}(x)$$

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- Collect these into the set of *monogenic polynomials*

$$\mathcal{M}^{\mathcal{P}}(\mathbb{B}_{R,w}) = \left\{ \sum_{j=0}^N \left(\sum_{\substack{j_2 \dots j_n \\ j_2 + \dots + j_n = j}} p_{j_2 \dots j_n} a_{j_2 \dots j_n} \right) \mid j_2 + \dots + j_n = j, N \in \mathbb{N}, a_{j_2 \dots j_n} \in \mathcal{G}_n^+ \right\}.$$

Lemma (Density)

The space $\mathcal{M}^{\mathcal{P}}(\mathbb{B}_{R,w})$ is dense in $\mathcal{M}^+(\mathbb{B}_{R,w})$.

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$$a_{j_2 \dots j_n} = \int_{\partial \mathbb{B}_{R,w}} \frac{\partial^j G(w, y)}{\partial y_2^{j_2} \dots \partial y_n^{j_n}} \nu(y) f(y) \mu_{\partial}(y),$$

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- Then

$$f(x) = \sum_{j=0}^{\infty} \left(\sum_{\substack{j_2 \dots j_n \\ j_2 + \dots + j_n = j}} p_{j_2 \dots j_n}(x - w) a_{j_2 \dots j_n} \right),$$

converges pointwise for $x \in \mathbb{B}_{R,w}$ by **[Ryan, 2004]**.

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- By linearity, we can note that for $\delta \in \mathfrak{M}(\mathbb{B}_{R,w})$

$$\delta(f(x)) = \sum_{j=0}^{\infty} \left(\sum_{\substack{j_2 \dots j_n \\ j_2 + \dots + j_n = j}} \delta(p_{j_2 \dots j_n}(x - w)) a_{j_2 \dots j_n} \right)$$

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- On each monogenic polynomial

$$\delta(p_{j_2 \dots j_n}(x)) = \frac{1}{j!} \sum_{\text{permutations}} \delta((z_{1\sigma(1)}(x)) \cdots \delta(z_{1\sigma(j)}(x)))$$

by the multiplicativity of δ .

Lemma (Point evaluation)

Let $\delta \in \mathfrak{M}(\mathbb{B}_{R,w})$ and $z_{ij} \in \mathcal{A}_{\mathbf{B}_{ij}}(\mathbb{B}_{R,w})$, then $\delta(z_{ij}) = z_{ij}(x^\delta)$ for some $x^\delta \in \mathbb{R}^n$.

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- The set of constants α and β are determined by n independent numbers, so we can say $\delta(z_{ij}) = z_{ij}(x^\delta)$ for some $x^\delta \in \mathbb{R}^n$.

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so this sequence not converge due to a singularity at x^δ .

- Hence, it must be that $x^\delta \in \mathbb{B}_{R,w}$ by continuity of δ .

Theorem (Noncommutative Gelfand representation)

For any $\delta \in \mathfrak{M}(\mathbb{B}_{R,w})$, there is a point $x^\delta \in \mathbb{B}_{R,w}$ such that $\delta(f) = f(x^\delta)$ for any $f \in \mathcal{M}(\mathbb{B}_{R,w})$. Given the weak- $$ topology on $\mathcal{M}^*(\mathbb{B}_{r,w})$, the map*

$$\gamma: \mathfrak{M}(\mathbb{B}_{R,w}) \rightarrow \mathbb{B}_{R,w}, \quad \delta \mapsto x^\delta$$

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- For $f \in \mathcal{M}^+(\mathbb{B}_{R,w})$ we have

$$f(\gamma(\delta_n)) = f(x^{\delta_n}) = \delta_n(f) = \gamma^{-1}(x^{\delta_n})(f).$$

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- Taking $n \rightarrow \infty$ shows γ and γ^{-1} are continuous so γ is a homeomorphism.

Section 4

Future work

Calderón problem

Calderón problem

- Let (M, g) be an unknown Riemannian manifold with known boundary ∂M . Consider the forward problem

$$\begin{cases} \Delta \omega = 0 & \text{in } M \\ \iota^* \omega = \phi & \text{on } \partial M \end{cases}$$

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- Define the *Dirichlet-to-Neumann map* on forms by $\Lambda\phi = \iota^*(\star d\omega)$.
- Question: Can we determine (M, g) from Λ ?

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- Solved in dimensions $n \geq 3$ when M is an analytic manifold [**Lassas, Taylor, Uhlmann: 2003**].
- The smooth cases is still unsolved.

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- For $n = 3$, the scalar potential u and magnetic bivector field b are two parts of a monogenic field $f = u + b$ due to Ohm's and Ampere's laws

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- If Λ can provide us $b|_{\partial M}$, then we can possibly reconstruct $\mathcal{M}^+(M)$.
- Given the algebraic structure of each $\mathcal{A}_B(M) \subset \mathcal{M}^+(M)$, can this be used to determine g ?

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- Can the magnetic impedance tomography problem can provide some extra insight on the EIT problem?
- The Hodge-Morrey decomposition is an instrumental tool for boundary value problems that, for example, allows one to show that Λ determines the Betti numbers of M [Belishev, Sharafutdinov: 2008].
- Can the Clifford-Hodge-Morrey decomposition can allow us to work on other related boundary inverse problems?

Section 5

Conclusions

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- This provides a new way to decompose fields on domains of \mathbb{R}^n and this can likely be generalized to arbitrary compact orientable Riemannian manifolds.

Conclusion

- We have utilized Clifford analysis to serve as a meaningful generalization of both the complex analysis and differential forms.
- This provides a new way to decompose fields on domains of \mathbb{R}^n and this can likely be generalized to arbitrary compact orientable Riemannian manifolds.
- Likewise, we have proven that the monogenic spinors contain a wealth of topological information.

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- We are continuing to work to apply our scheme to new models such as the Modular Arbitrary-Order Ocean-Atmospheric Model (MAOOAM).

Thank you!