The Calderón Problem

on Riemannian Manifolds

Colin Roberts





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- Rephrase the problem in a geometrical way.
- Prove the problem in 2 dimensions using the boundary control method.
- What can and can't we do to generalize this method?

Section 1

Introduction

Subsection 1

Calderón Problem

In 1980, Alberto Calderón proposed a problem in his paper On an inverse

boundary value problem.

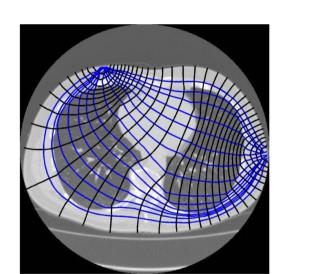
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- This problem sparked interest due to its usefulness in geophysical and medical imaging.



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<u>Idea:</u> Given a domain Ω with interior Ω^+ that we cannot probe, can we determine the conductivity γ matrix by studying the boundary $\partial\Omega$?

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- This defines the voltage-to-current map Λ so that $\Lambda(f) = h$.
- Can we determine the conductivity matrix γ from Λ ?

Section 2

The Calderón Problem on Riemannian Manifolds

Subsection 1

Preliminaries

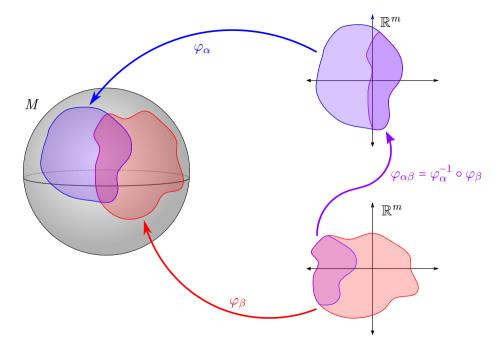


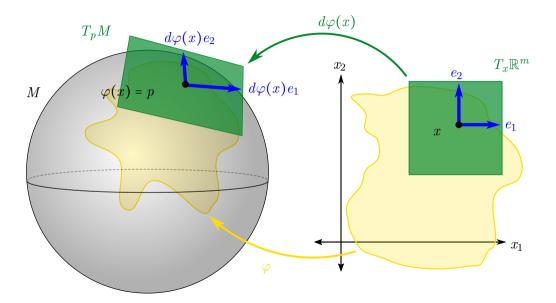
■ Smooth n-dimensional manifold: A space that locally looks like (is C^{∞} diffeomorphic to) an open subset of \mathbb{R}^n .

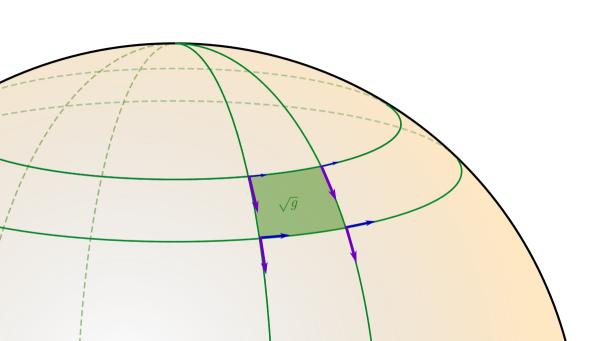
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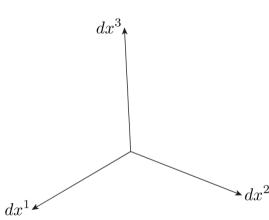
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- Exterior algebra: Differential forms with the wedge product \wedge .
- *Hodge Star*: Attached to the exterior algebra when we also have a Riemannian metric. Gives an isomorphism between k and n k-forms.



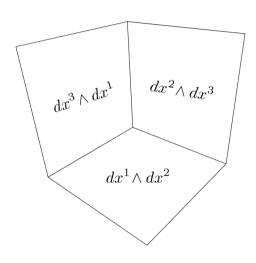




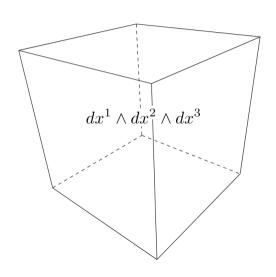
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2-Forms



3-Forms



•
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- *Dirac Operator*: $D = d + \delta$.
- Laplace-Beltrami Operator: $\Delta = d\delta + \delta d = D^2$ and in coordinates

$$\Delta f = \frac{1}{\sqrt{|g|}} \sum_{j=1}^{n} \sum_{k=1}^{n} \frac{\partial}{\partial x^{i}} \sqrt{|g|} g^{ij} \frac{\partial}{\partial x^{j}} f$$

Subsection 2

Raphrasing EIT Problem in a Geometrical Language

■ Let (unknown) connected Ω be a smooth Riemannian manifold with boundary $\partial\Omega$.

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■ Recover g from knowing Λ .

Subsection 3

Expected and Current Results

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So, g is a function of n variables that needs to be determined by the kernel λ which is 2n-2 variables.

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- $n \ge 3$ is overdetermined.

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Indeed, let $\tilde{g} = e^{2\phi}g$, then

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When n = 2, the extra term cancels.

1 Dimension

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- Can only know the total impedence between the two electrodes.

Isotropic Case

For g isotropic and $n \geq 3$ one can determine g from Λ . (Sylvester-Uhlmann 1987)

2 Dimensional Anisotropic

- \blacksquare Can recover g up to conformal class and can't do better.
- Proven by Lassas and Uhlmann in On Determining the Riemannian manifold from the Dirichlet to Neumann map.

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- Can determine the boundary C^{∞} -jet of g in Lee and Uhlmann's Determining anisotropic real-analytic conductivities by boundary measurements.
- For smooth manifolds, the anisotropic problem is open. The goal is to recover the metric up to isometry.

Section 3

Boundary Control Method in 2 Dimensions

Theorem

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Two 2-dimensional compact orientable manifolds with single common boundary are conformally equivalent iff their DN-maps coincide.

Belishev's The Calderón Problem for Two-Dimensional Manifolds by the BC-Method.



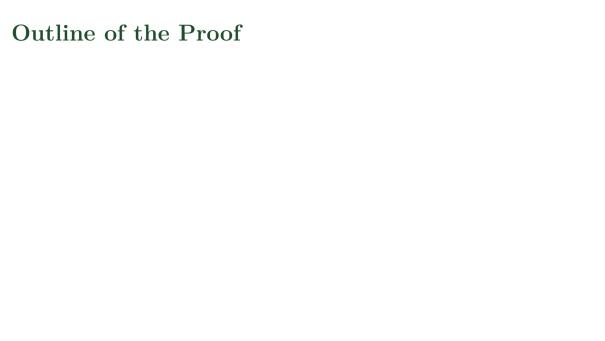
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- Gelfand transform relates an algebra \mathcal{A} to the algebra of continuous functions on the spectrum of that algebra, $C(\operatorname{spec}\mathcal{A})$.
- This gives us a way to realize Ω from functions defined on Ω that we have access to.



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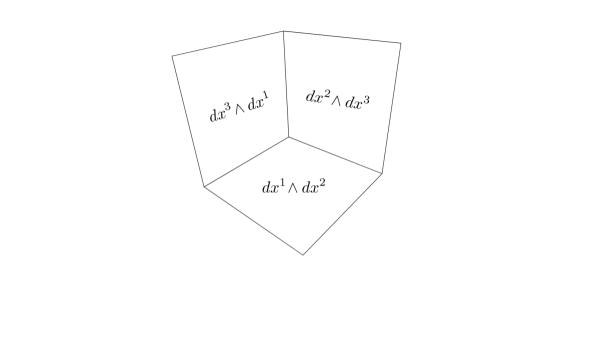
- From DN map, recover the algebra of holomorphic functions. (Lemma 1)
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- Represent the trace algebra with the DN map. (Lemma 3)
- Construct the manifold. (Theorem)

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- Define Λ by $\iota^*(\star du)$ for a harmonic u.
- Λ maps boundary k-forms to boundary n-k-1 forms.



Subsection 1

Lemma 1



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■ A function u satisfying $\Delta u = 0$ has a conjugate function v if and only if the trace $\iota^* u$ satisfies

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 $dim Ran \left[\Lambda + d\Lambda^{-1} d \right] = \beta_1(\Omega).$

Corollary

 Λ completely determines the topology of $\Omega.$



Proof

- Since Ω is a single connected component, $\beta_0(\Omega) = 1$.
- We have $\beta_1(\Omega)$ from before.
- Since Ω is a surface with boundary, $\beta_2(\Omega) = 0$.
- Since Ω is two dimensional, $\beta_n(\Omega) = 0$ for $n \geq 3$.

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- \blacksquare Let u be a 0-form and v as a 2-form.
- Then $\frac{\partial}{\partial \overline{z}}$ is given by $D = d + \delta$.
- Dw = 0 gives us the Cauchy-Riemann equations
- We call u and v conjugate by CREs.

Hilbert Transform

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- Define $\mathcal{H} = d\Lambda^{-1}$.

■ By Lemma 1 we can now create the algebra $\mathcal{A}(\Omega) \subset C(\Omega)$ from harmonic functions u with conjugates v by

$$\mathcal{A}(\Omega) \coloneqq \{w = u + iv\}.$$

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- This is analogous to having the Hodge star on a surface.

Subsection 2

Lemma 2

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- \mathcal{M} with this topology is called the *spectrum* spec \mathcal{A} .

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- A function algebra $A \subset C(X)$ is *generic* if ϵ is a homeomorphism.
- A generic algebra is (spatially) isomorphic to its Gelfand transform.

Lemma 2

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The algebra of holomorphic functions $\mathcal{A}(\Omega)$ is generic.



Importance

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- $\hat{\mathcal{A}}(\partial\Omega)$ is (spatially) isomorphic to $\mathcal{A}(\Omega)$ by taking the Gelfand transform of the trace.
- The lemma shows that $\epsilon: \Omega \to \operatorname{spec} \mathcal{A}(\Omega)$ is a homeomorphism, so we have determined Ω up to homeomorphism.

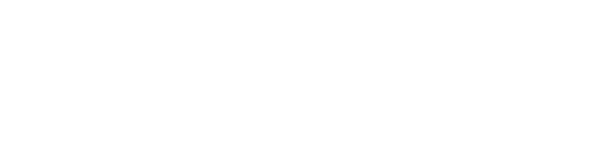


What's Left?

We can only have hope access to the trace algebra $\mathcal{A}(\partial\Omega)$. So we need to determine this to reach our goal.

Subsection 3

Lemma 3



Trace Algebra

Trace Algebra

■ The trace algebra $\mathcal{A}(\partial\Omega) \coloneqq \iota^* \mathcal{A}(\Omega)$ is isometrically isomorphic to $\mathcal{A}(\Omega)$ since

$$\|w\|_{\mathcal{A}(\Omega)} = \|\iota^* w\|_{\mathcal{A}(\partial\Omega)}$$

and since a holomorphic function is uniquely determined by its boundary values.

Lemma 3

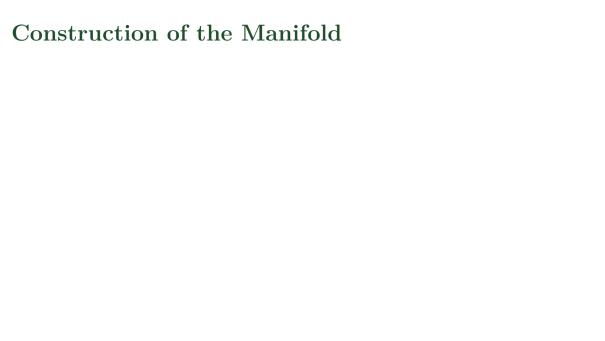
We have the representation

$$\mathcal{A}(\partial\Omega) = \operatorname{clos}_{C(\partial\Omega)}\{f+ih\},\,$$

where h is conjugate to f by \mathcal{H} .

Subsection 4

Proof of the Main Theorem



Following these steps yields a manifold (Ω, g) with the DN map Λ .

■ Step 1: We know $g|_{\partial\Omega}$ by Lee and Uhlmann, and thus we know \mathcal{H} and $C^{\infty}(\partial\Omega)$. This allows us to recover the trace algebra $\mathcal{A}(\partial\Omega)$ using the representation in Lemma 3.

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- Step 4: $\mathcal{A}(\Omega)$ gives us the complex structure on Ω by Lemma 1.
- Step 5: Equip Ω with a metric g conforming to this complex structure.

Section 4

Generalizing This Method

First Issue

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 \blacksquare No complex structure in higher dimensions.

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- $lue{}$ Use a Clifford algebra/calculus structure to replace \mathbb{C} .
- The tools of Clifford analysis allow us to recover a notion of holomorphicity known as *monogenicity*.
- We can recover a similar algebra (Hardy space) of monogenic functions.

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■ The quotient of the tensor algebra

$$C\ell(V,Q) = \bigoplus_{j=0}^{\infty} V^{\otimes j} / \langle v \otimes v - Q(v) \rangle.$$

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■ If Q = g is an inner product, this yields a geometric product on vectors $u, v \in C\ell(V, g)$

$$uv = g(u, v) + u \wedge v.$$

• Geometric product can be extended to multivectors.

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■ We can replace $d + \delta$ with the Dirac operator D

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- lacktriangle There are Cauchy integral and Hilbert transform type operators for D in arbitrary dimension.

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- The spectral theory for Belishev's solution required a commutative Banach algebra.
- The spectral theory for noncommutative Banach algebras is not as developed.
- There is still work to do here to find some way around this.

Section 5

Conclusion

■ Calderón proposed a useful and challenging problem for both theorists and practicioners.

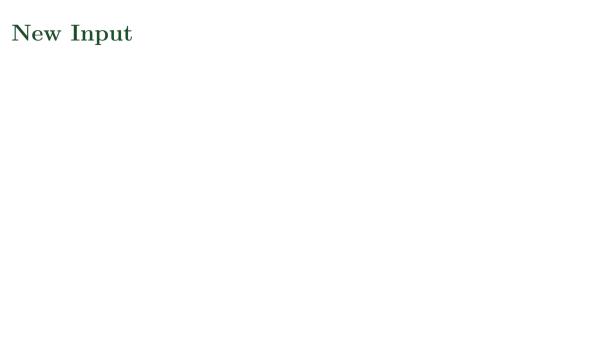
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- Advances have been made in both areas.
- Ideal results are still not yet obtained.

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- These issues remain if we try to naively generalize this approach.



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- We can still construct the same algebra of holomorphic functions.
- The relevant algebras are noncommutative for dimensions ≥ 3 .
- There are possibly other tools at our disposal that may be able to replace the loss of commutivity.

Thank you!