

# Notes

## On Scaled Equations of Nonlinear Radiative Transfer Problems and Computational Units

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## 1 Computational Units

Let us introduce the following scaling by means of special computational units:

- $t = \rho_t \hat{t}$ ,
- $I_g = \rho_I \hat{I}_g$ ,
- $B_g = \rho_B \hat{B}_g$ ,
- $\varepsilon = \rho_\varepsilon \hat{\varepsilon}$ .

## 2 Scaled Equations

The multigroup radiative transfer equations with scaling are given by

$$\frac{1}{\rho_t c} \frac{\partial \hat{I}_g}{\partial \hat{t}} + \mathbf{\Omega} \cdot \nabla \hat{I}_g + \kappa_{E,g}(T) \hat{I}_g = \frac{\rho_B}{\rho_I} \kappa_{B,g}(T) \hat{B}_g(T), \quad (1a)$$

$$\hat{I}_g \Big|_{\mathbf{r}=\mathbf{r}_\gamma} = \frac{1}{\rho_I} I_g^{in}, \quad \text{for } \mathbf{r}_\gamma \in \partial G \text{ and } \mathbf{\Omega} \cdot \mathbf{e}_n < 0, \quad \hat{t} \geq \frac{1}{\rho_t} t_0, \quad (1b)$$

$$\hat{I}_g \Big|_{\hat{t}=\rho_t^{-1} t_0} = \frac{1}{\rho_I} I_g^0, \quad \text{for } \mathbf{r} \in G \text{ and all } \mathbf{\Omega}, \quad (1c)$$

where

$$\hat{B}_g(T) = \frac{\hat{\sigma}_{R,g}(T)}{\pi} T^4 \quad (2)$$

and

$$\hat{\sigma}_{R,g}(T) = \frac{1}{\rho_B} \sigma_{R,g}(T). \quad (3)$$

The moments of the scaled intensity are

$$\hat{E}_g = \frac{1}{c} \int_{4\pi} \hat{I}_g d\Omega, \quad (4)$$

$$\hat{\mathbf{F}}_g = \int_{4\pi} \Omega \hat{I}_g d\Omega. \quad (5)$$

The scaled energy balance equation has the following form:

$$\frac{\rho_\varepsilon}{\rho_t} \frac{\partial \hat{\varepsilon}(T)}{\partial \hat{t}} = c \sum_{g=1}^{N_g} \left( \varkappa_{E,g}(T) \rho_I \hat{E}_g - \varkappa_{B,g}(T) \rho_B \hat{E}_g^{pl} \right), \quad (6)$$

$$\hat{\varepsilon}(T) = \hat{c}_v T, \quad (7)$$

$$T|_{\hat{t}=\rho_t^{-1}t_0} = T^0, \quad (8)$$

where

$$T^0 = T(t_0), \quad (9)$$

$$\hat{c}_v = \frac{1}{\rho_\varepsilon} c_v, \quad (10)$$

$$\hat{E}_g^{pl} = \frac{4\pi}{c} \hat{B}_g(T) = \frac{4\hat{\sigma}_{R,g}(T)}{c} T^4. \quad (11)$$

We note that

$$\hat{E} = \sum_{g=1}^{N_g} \hat{E}_g, \quad (12)$$

$$\hat{\mathbf{F}} = \sum_{g=1}^{N_g} \hat{\mathbf{F}}_g, \quad (13)$$

$$\hat{E}^{pl}(T) = \frac{4\hat{\sigma}_R}{c} T^4, \quad (14)$$

$$\hat{\sigma}_R = \frac{1}{\rho_B} \sigma_R. \quad (15)$$

### 3 A Particular Case of Computational Units

Consider the following case:

- $t = [shakes]$ ,
- $\rho_t = 1$ ,
- $\rho_B = \rho_I$ ,
- $\rho_\varepsilon = \rho_I$ ,

where, for instance,  $\rho_I = 10^{13}$ . Using *shakes* we obtain

- $c = 299.792458 \left[ \frac{cm}{sh} \right]$ .
- $h = 6.62613 \times 10^{-19} \text{ erg} \cdot \text{sh}$

- $\sigma_R = \frac{2\pi^5 k^4}{15c^2 h^3} = 5.67 \times 10^{-13} \frac{\text{erg}}{\text{cm}^2 \text{ sh } (\text{°K})^4} = 1.0266 \times 10^4 \frac{\text{erg}}{\text{cm}^2 \text{ sh } \text{eV}^4}$
- specific intensity  $I(\mathbf{x}, \mathbf{\Omega}, \nu, t) = ch\nu n(\mathbf{x}, \mathbf{\Omega}, \nu, t)$ :  $\left[ \frac{\text{erg}}{\text{cm}^2 \text{ sh } \text{Hz } \text{ster}} \right]$ ,
- monochromatic radiation energy  $E_\nu = E(\mathbf{x}, \nu, t) = \int_{4\pi} h\nu n d\mathbf{\Omega} = \frac{1}{c} \int_{4\pi} I d\mathbf{\Omega}$ :  $\left[ \frac{\text{erg}}{\text{cm}^3 \text{Hz}} \right]$ .

We now define the scaled Stefan-Boltzmann constant

$$\hat{\sigma}_R = \frac{1}{\rho_I} \sigma_R, \quad (16)$$

the scaled function

$$\hat{\sigma}_{R,g}(T) = \frac{1}{\rho_I} \sigma_{R,g}(T), \quad (17)$$

scaled heat capacity constant

$$\hat{c}_v = \frac{1}{\rho_I} c_v, \quad (18)$$

and scaled intensities at boundaries and at the initial moment of time

$$\hat{I}_g^{in} = \frac{1}{\rho_I} I_g^{in}, \quad \hat{I}_g^0 = \frac{1}{\rho_I} I_g^0. \quad (19)$$

The set of scaled equations (1) and (6) get the form of the original NRT equations, namely,

$$\frac{1}{c} \frac{\partial \hat{I}_g}{\partial \hat{t}} + \mathbf{\Omega} \cdot \nabla \hat{I}_g + \kappa_{E,g}(T) \hat{I}_g = \kappa_{B,g}(T) \hat{B}_g(T), \quad (20a)$$

$$\hat{I}_g \Big|_{\mathbf{r}=\mathbf{r}_\gamma} = \hat{I}_g^{in}, \quad \text{for } \mathbf{r}_\gamma \in \partial G \text{ and } \mathbf{\Omega} \cdot \mathbf{e}_n < 0, \quad \hat{t} \geq \rho_t^{-1} t_0, \quad (20b)$$

$$\hat{I}_g \Big|_{\hat{t}=\rho_t^{-1} t_0} = \hat{I}_g^0, \quad \text{for } \mathbf{r} \in G \text{ and all } \mathbf{\Omega}. \quad (20c)$$

$$\frac{\partial \hat{\varepsilon}(T)}{\partial \hat{t}} = c \sum_{g=1}^{N_g} \left( \kappa_{E,g}(T) \hat{E}_g - \kappa_{B,g}(T) \hat{E}_g^{pl} \right), \quad (20d)$$