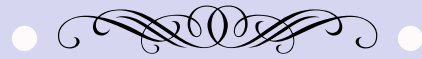


Dimensional Dependence of

Light

Interaction with Nanowires



**A Ph.D. defense by
Zhihuan Wang**

Advised by Dr. Bahram Nabet

Outline

- 1 Introduction
- 2 Crowding Effects
- 3 Experimental starting points
- 4 Monte Carlo Techniques
- 5 Modeling and Results
- 6 Extensions and Future Work

Protein Folding

Proteins

- Responsible for wide variety of biological function
- Linear polymer chains of amino acids

Folding

- Central problem: from primary sequence and physiological conditions determine the *unique* native state
- Misfolding and aggregation results in many clinical disorders including neurodegenerative disorders like Huntington's and Alzheimer's disease and diabetes.

Entropy: The Science of Counting

Equilibrium statistical mechanics

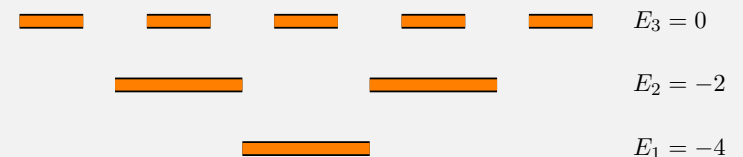
$$Z = \sum g(E_i) e^{-\beta E_i}$$

- The partition function gives a complete description of the system.

$$S = k \ln g(E)$$

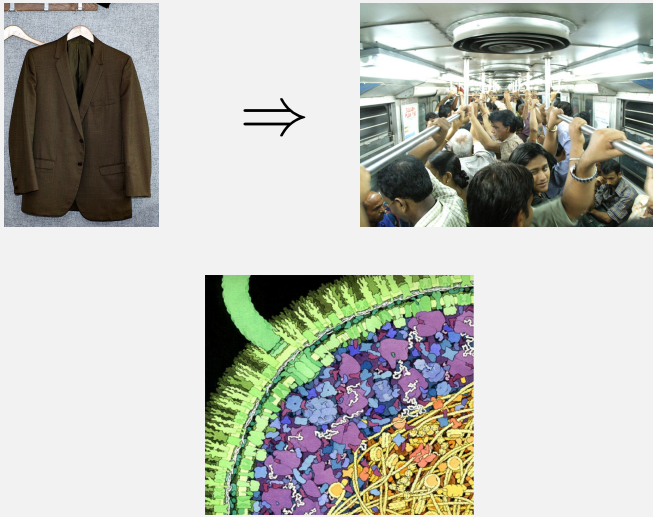
- Try to solve for the density of states directly!

Example:



This gives a DOS of $g(-4) = 1$, $g(-2) = 2$, $g(0) = 5$.

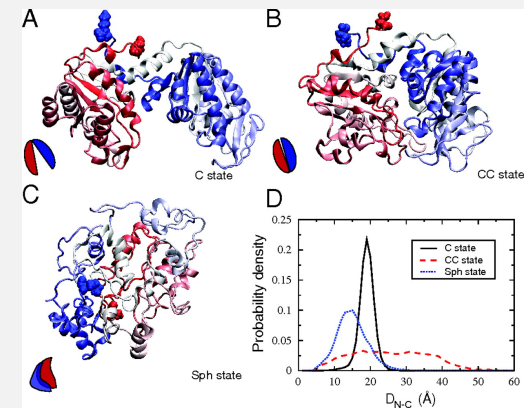
What is Crowding?



Crowding reduces the conformational state space of the peptide.

3 / 29

Crowding Induced Function Change

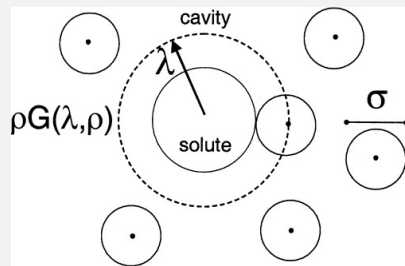


Crowding changes activity rates and function, and may change the native state

Dhar, Gruebele, Cheung, et. al., PNAS 107, 17586 (2010)

4 / 29

Hard-Sphere Experiments, Crowding Insights



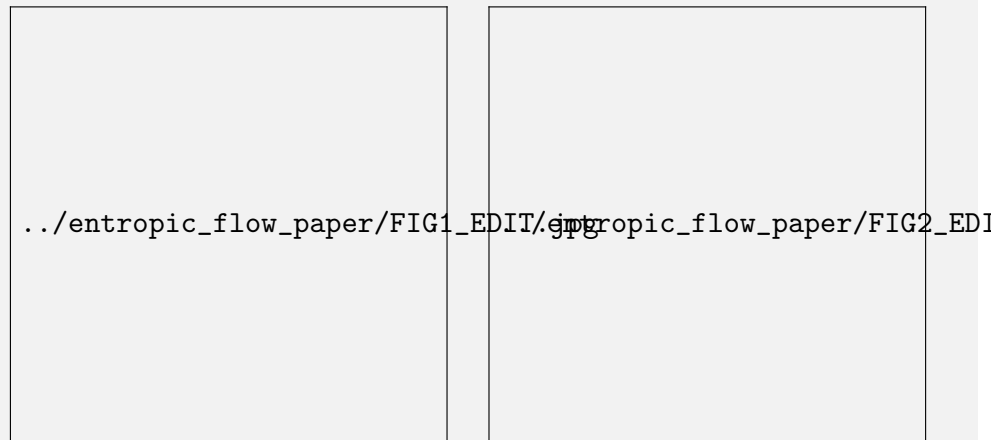
Picture from [Heying, 2004], depletion zone traced with dashed lines

Hard-sphere studies of the crowding effect have shown:

- **Changes of the conformations of the cosolute**
- **Motion towards minimal curvature**

5 / 29

Crowding Induced Cosolute Conformation



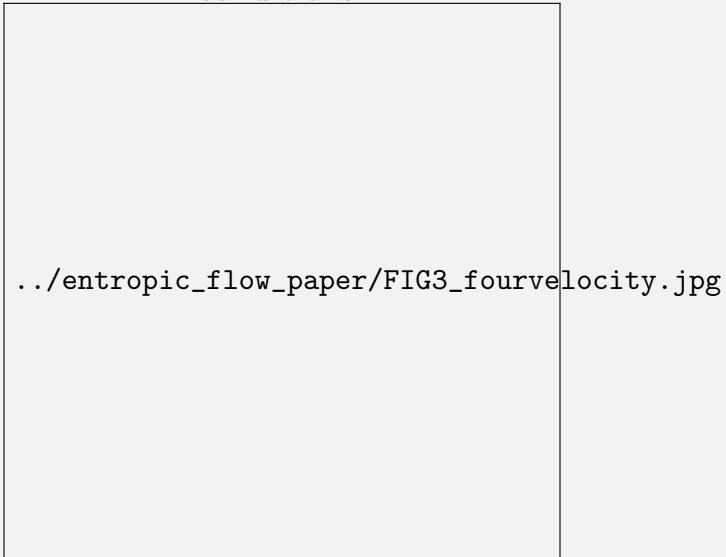
Hard-sphere experiments with fixed inner elliptical boundary shows standard radial density profiles

Hoppe and Yuan, Phys. Rev. E 80, 011404 (2009)

6 / 29

Crowding Induced Flow

Non-spherical objects generate flows from random initial conditions!



Four vortices of paired spins, net angular momentum is still conserved
Effect: Crowding induced shape change (also considered by Cheung)

7 / 29

Designed β -sheet peptides

Collection of experiments by Feng Gai at U. Penn.

Properties:

- β -sheet motif
- Low aggregation propensity
- Known native-state

- #1 RFSEV^D[PG]KKFITS^D[PG]KTYTEV^D[PG]KKILQ
- #2 RFIEV^D[PG]KKFITS^D[PG]KTYTE

^DP = D-Pro

Xu, Gai, et. al., Biochem.47, 2064 (2008)
Hudson and Anderson, Biopolymers.83(4) pg. 424-433 (2006)

9 / 29

Scaled Particle Theory

- Measures the cost of inserting a new particle into a solution
- Rigorous results known for binary hard-sphere solutions
- Extended to aspherical particles by Kihara
- Aspherical SPT for proteins by Minton

Kihara, Rev. Mod. Phys. 25, 831 (1953)
Minton, Method Enzymol. 295, 127 (1998)

8 / 29

Native state representation

Wild-type β -sheet peptide (trpzip-m1)

- #3 GEWTWAD[AT]KTWTWTE

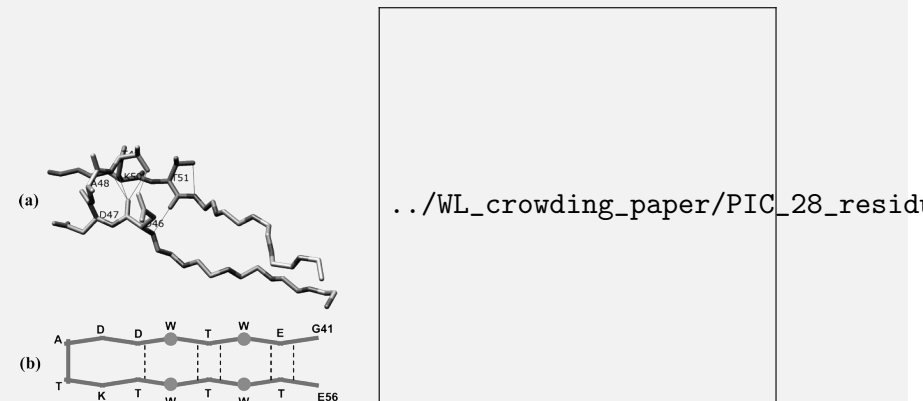
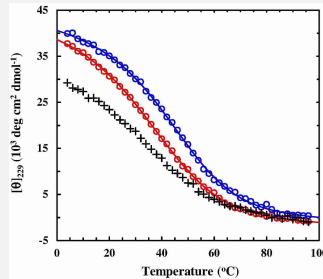


Figure: Native state of the wild-type 16-residue and designed 28-residue peptides represented schematically.

10 / 29

Why study these peptides?

- Thermal transition data
- Crowding data under equilibrium properties



- Isolated β -sheet experiments are relatively new
- Comparison between theoretical/experimental studies
- Detailed thermodynamic MD studies are feasible, but time consuming
- **What effect do crowders have on the folding process?**

11 / 29

Monte Carlo: How to sample

Most methods rely on Monte-Carlo, i.e. a selection method that is probabilistically determined. Any method must overcome two obstacles:

- Enthalpic barriers
- Entropic barriers (Levinthal's paradox)

The enthalpic barriers are usually directly observable from the Hamiltonian while the entropic ones are often more subtle.

12 / 29

Traditional Sampling Methods

The typical method uses a Metropolis-Hastings selection scheme:

$$P(A \rightarrow B) = \exp\left(\frac{\Delta E_{A,B}}{kT}\right)$$

This works well at finding local minimum - however it needs help getting out. Various models have been designed to get around this such as: Replica Exchange, Simulated Annealing, Reversible Jump Dynamics, etc...

These simulations only collect data at fixed temperatures; thermodynamic parameters like the free energy or the specific heat at arbitrary temperatures are time consuming

$$C_V = T \left(\frac{\partial S}{\partial T} \right)_{N,V}$$

13 / 29

Wang-Landau

- Originally designed to solve the Ising model
- Extremely effective when the cardinality of Ω (density of states) is small
- Can usually be made parallel trivially
- Works best when energy levels are discrete

F. Wang and D. P. Landau, Phys. Rev. Lett. 86(10), 2001

14 / 29

Wang-Landau: The equations

$$P(A \rightarrow B) = \frac{\Omega(E_A)}{\Omega(E_B)}$$

On accepting state:

$$\Omega(E_{ex}) \rightarrow c\Omega(E_{ex})$$

where c is a constant that is iteratively reduced (originally as):

$$c \rightarrow \sqrt{c}$$

Once $\Omega(E)$ has converged, we have the partition function *for all temperatures*

$$Z(T) = \sum \Omega(E_i) \exp(E_i/kT)$$

15 / 29

Wang-Landau

- Starts as a random walk across the state space
- As high entropic states are sampled they become less likely
- Random walk \rightarrow biased random-walk
- As it converges - detailed balance is obeyed
- **Flat histogram in energy space**

16 / 29

Difficult Questions

How to model?

- All-atom?
- Explicit water?
- Effective forces?
- Level of coarse-graining?

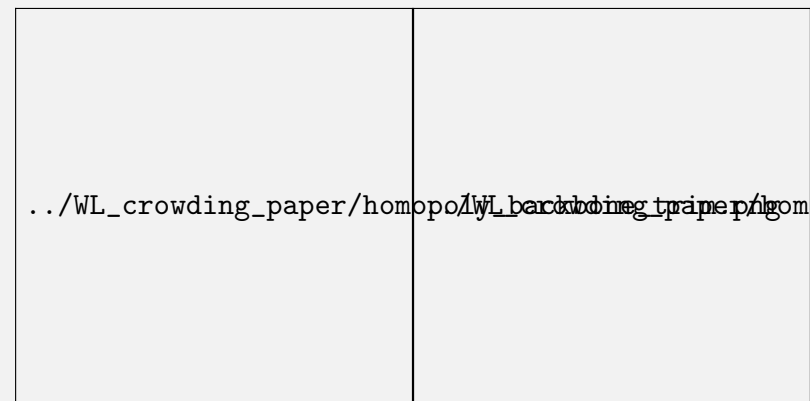
We use a face-centered cubic lattice Gō-model, with a Potts/Ising approximation for the dihedral angles and aspherical scaled-particle theory to model a crowded solution.

lets break that down...

17 / 29

Lattice Models

- Discretize the positions of the amino acids onto a fixed lattice
- Short-range interactions dominate (hydrophobicity, charge, etc)
- Lattice geometry can play a role on the pairing of amino acids

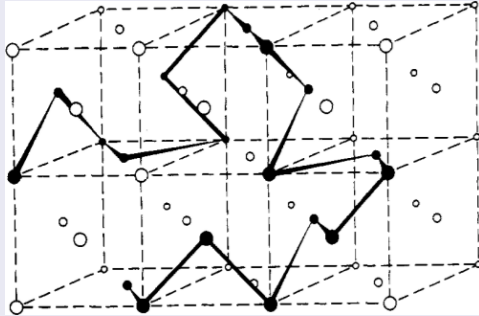


Chan and Dill, J. Chem. Phys. 118, 8106 (1994)

18 / 29

Lattice Choice

- High coordination number (FCC vs SC)
- Common motifs are better represented (alpha, beta)
- Still a crude approximation to an off-lattice model



Picture from [Downey, 1986]

Ising Model

- Analytical solution for 1D (Ising 1925), 2D (Onsager 1944) grids
- Each vertex on a graph has spin $\sigma \in \{-1, 1\}$, Hamiltonian is the sum over all adjacent spins

$$\mathcal{H} = -J \sum_{i,j \in \text{edges}} \sigma_i \sigma_j + h \sum_{k \in \text{vertices}} \sigma_k \delta_{1, \sigma_k}$$

- Gives true phase transitions for 2D infinite grids
- Finite graphs can be solved *exactly*

Gō-model & extensions

Gō-model

- Only native-state contacts contribute (all-or-nothing)

Extensions

- Extra degree of freedom, each bead now has a ‘spin’.
- Accounts for the missing oriental entropy of $\phi - \psi$ angles
- Only native-state contacts *and* aligned spins contribute, reducing the conformation to a simple graph

```
../WL_crowding_paper/homopoly./WL_crowding_paper/fig_example_
```

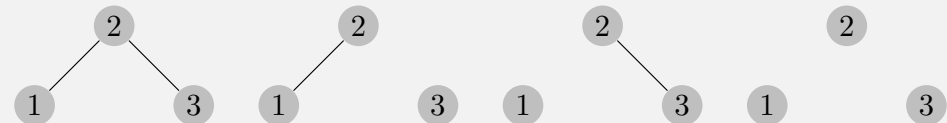
General solution to Potts/Ising model (with external fields!)

With Self-Similarity

- Can often be solved/approximated with RG
- e.g. n-dimensional lattices, Cayley/Bethe trees

General Solution

- Solved with a subgraph decomposition



$$Z = \sum_{A \subseteq G} \prod_{c \in C_A}^{k(A)} \sum_j^q \left(w_j^{n(c)} v_j^{e(c)} \right)$$

Model Hamiltonian

Hamiltonian counts aligned native contacts in close proximity

$$\mathcal{H}(\mathbf{c}, \mathbf{s}) = -J_+ k_+ + J_- k_-$$

$$k_+ = \sum_{i=1}^L \sum_{j=i+2}^L \omega_{ij} \mathbf{s}_i \mathbf{s}_j \mathbf{G}_{ij}$$

$$k_- = \sum_{i=1}^L \sum_{j=i+2}^L \omega_{ij} (1 - \mathbf{G}_{ij})$$

Free energy terms for crowding $\Delta\mu$ and entropic cost of alignment $\Delta\psi$

$$\mathcal{F}(\mathbf{c}, \mathbf{s}) = \mathcal{H}(\mathbf{c}, \mathbf{s}) - \beta \Delta\psi(\sigma) - \beta \Delta\mu(\mathbf{c}) \quad (1)$$

23 / 29

Fraction Folded

Measure of fraction folded well-described by model

../WL_crowding_paper/PLOT_all_experimental_fits-crop

24 / 29

Crowded Conditions → Folding Stability

../WL_crowding_paper/PLOT_trpzip_CV-crop.pdf

25 / 29

Crowded Conditions → Folding Instability

While the smaller peptide trpzip4-m1 displays crowding enhanced stability, the larger designed β -sheet peptides show (small) instability!

Simulations predict that the entropic effect due to the crowding of aspherical proteins may be destabilizing.

Why might this be so?

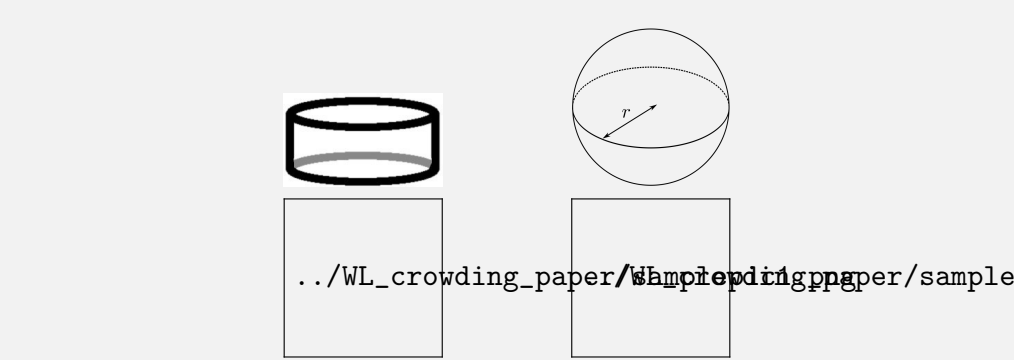
- Ensemble of states has to be considered
- Only possible with a full density of states calculation

26 / 29

Figure: Specific heat per residue count for the wild-type peptide trpzip4-m1 in the

Crowded Conditions \rightarrow Folding Instability

Consider the underlying geometry of our β -sheets



Change of native state conformation under crowding of 3-strand peptide

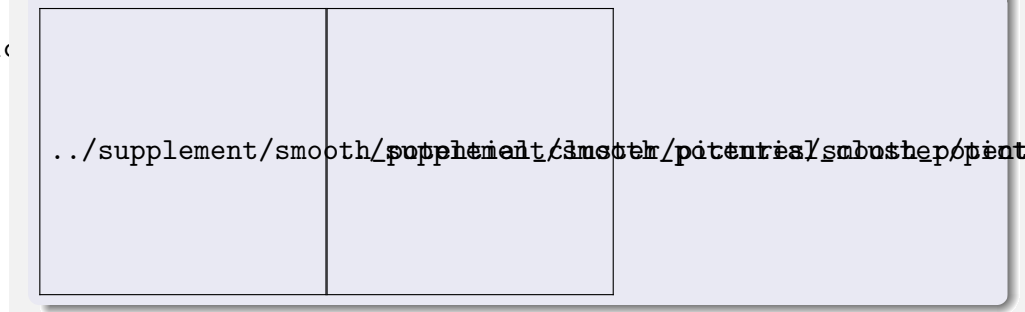
Trade-off between the **energetically favored native-state**, and an **entropically unfavorable conformation**.

Macrostate Clustering

Group **microstates** into **macrostates**

- Coarse-grain from kinetics, compute transition probabilities
- Transpose of eigenvectors suggest clustering

Example: Random walk in a three-well potential



Future Work

Where is the field going?

What are the important questions?

Nuanced Interactions

- Electrostatic Effects
- Hydrophobicity
- Hydrogen Bonds

Beyond thermodynamics

- Protein function and activity
- Aggregation models and kinetics
- Diffusion rates in crowded environments

Thank You.

Extra: Scaled Particle Theory Equations

Activity coefficient

$$\ln \gamma_i = -\ln(1 - (V)) + \frac{H_i(S) + S_i(H) + V_i(1)}{1 - (V)} + \frac{H_i^2(S)^2 + 2V_i(H)(S)}{2(1 - (V))^2} + \frac{V_i(H^2)(S)^2}{3(1 - (V))^3}$$

H is the Kihara support function, geometrical measure of ‘roundness’

For example:

$$H_{\text{sphereocylin}} = r\pi/4 + L/2$$



$$H_{\text{cylinder}} = r + L/4$$



We model the crowders, Ficoll 70, as sphereocylinders and the β -sheet peptides as cylinders to best match their native state geometry.

Fodeke and Minton, J. Phys. Chem. B 114, 10876 (2010)