

PSet4 P4 KTX887

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[1]: #Homework set 4, problem 4
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    #KTX887
    import sympy as sp
```

0.1 Question 4

Consider a model of a relativistic jet motivated by Blandford and Znajek (1977). The jet is stationary (meaning there is no change over time) and comprises a helical magnetic field threading the horizon of a black hole at its base, extending symmetrically in a paraboloid shape proportional to z . Moreover, let's assume the jet is self-similar in that at every value of ξ (which is defined as s^2 divided by z), the jet velocity, electric, and magnetic fields are determined by the magnetic flux $\Phi(s, z)$ threading the surface of the paraboloid jet and the angular speed of field lines $\Omega(s, z)$ as they emerge from the location of the horizon.

0.1.1 a.)

Find the (infinitesimal) contribution $d\Phi$ to magnetic flux at constant s through the area $2\pi s dz$ and use it to find B_s in terms of Φ .

HINT: The partial derivative of a function $f(x, y)$ with respect to x is the derivative of f while holding y constant. Use a similar strategy to find B_z in terms of Φ .

0.1.2 b.)

Use Ampere's law (in integral form) to find B_ϕ , assuming the current within an Amperian loop of radius s is I . Check that the divergence of the magnetic field B in cylindrical coordinates is 0 for this self-similar jet.

HINT: You may want to use a numerical solver such as Mathematica to compute the divergence. Include the output of your code.

0.1.3 c.)

Ohm's law, expressed as $\mathbf{E} + \mathbf{v} \times \mathbf{B} = \eta \mathbf{j}$, has zero resistivity η in ideal MHD. Use this form of Ohm's law and the fact that the component of \mathbf{v} perpendicular to \mathbf{B} is $\mathbf{v}_\perp = \Omega \mathbf{r} \times \mathbf{r}$ to express \mathbf{E} in terms of Ω in ideal MHD.

0.1.4 d.)

Use Ampere's law (in differential form) to compute $\mathbf{j} = (j_s, j_{\phi}, j_z)$.

HINT: You may want to use a numerical solver such as Mathematica to compute any curls.

0.1.5 e.)

Find a differential equation for $I(x)$ by expanding $\rho \cdot \mathbf{E} + \mathbf{j} \times \mathbf{B}$ to first order in z/x .

HINT: You may use a numerical solver.

```
[2]: #Part A
#symbols
s, z = sp.symbols('s z')
Phi = sp.Function('Phi')(s, z)

#differential element volume
dA_s = 2 * sp.pi * s * sp.symbols('dz') #area element in s-direction (2pi s dz)
dA_z = 2 * sp.pi * s * sp.symbols('ds') #area element in z-direction (2pi s ds)

#def flux through surface
dPhi_s = sp.symbols('Bs') * dA_s
dPhi_z = sp.symbols('Bz') * dA_z

#solve for Bs in terms of Phi(s, z)
Bs = (1 / (2 * sp.pi * s)) * sp.diff(Phi, z)
Bz = (1 / (2 * sp.pi * s)) * sp.diff(Phi, s)

#display
print('The value for B_s is:')
display(Bs)
print('The value for B_z is:')
display(Bz)
```

The value for B_s is:

$$\frac{\frac{\partial}{\partial z} \Phi(s, z)}{2\pi s}$$

The value for B_z is:

$$\frac{\frac{\partial}{\partial s} \Phi(s, z)}{2\pi s}$$

```
[3]: #define vars
s, z, mu_0 = sp.symbols('s z mu_0')
I_s = sp.Function('I')(s) # Define I(s) as a function of s

#ampere law in integral form to solve for B_phi
B_phi = (mu_0 * I_s) / (2 * sp.pi * s)
```

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#mag field components
Phi = sp.Function('Phi')(s, z)
B_s = (1 / (2 * sp.pi * s)) * sp.diff(Phi, z)
B_z = (1 / (2 * sp.pi * s)) * sp.diff(Phi, s)

#divb
div_B = (1/s) * sp.diff(s * B_s, s) + sp.diff(B_z, z)

#display
print('B_phi is : ')
display(B_phi)
print('The divergence of B is : ')
print(div_B.simplify())

```

B_phi is :

$$\frac{\mu_0 I(s)}{2\pi s}$$

The divergence of B is :

Derivative(Phi(s, z), s, z)/(pi*s)

```

[4]: #define vars
s, z, mu_0, I = sp.symbols('s z mu_0 I')

#ampere law in integral form to solve for B_phi
B_phi = (mu_0 * I) / (2 * sp.pi * s)

#mag field components
Phi = sp.Function('Phi')(s, z)
B_s = (1 / (2 * sp.pi * s)) * sp.diff(Phi, z)
B_z = (1 / (2 * sp.pi * s)) * sp.diff(Phi, s)

#div b
div_B = (1/s) * sp.diff(s * B_s, s) + sp.diff(B_z, z)

#display
div_B.simplify()
#treating phi as a function of s and z

```

[4]:
$$\frac{\frac{\partial^2}{\partial z \partial s} \Phi(s, z)}{\pi s}$$

```

[5]: #problem 4 part b
s, z = sp.symbols('s z')

#define Phi(s, z) as a function of s and z
Phi = sp.Function('Phi')(s, z)

```

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#magnetic field components
B_s = (1 / (2 * sp.pi * s)) * sp.diff(Phi, z) #Bs
B_z = (1 / (2 * sp.pi * s)) * sp.diff(Phi, s) #Bz

#calculate div in cylindrical
div_B = (1/s) * sp.diff(s * B_s, s) + sp.diff(B_z, z)

#simplify
div_B_simplified = div_B.simplify()

#show div
print('The divergence of B is : ')
display(div_B_simplified)

#check if the divergence simplifies to show that mixed partial derivative must
    ↳ be zero
mixed_partial = sp.diff(Phi, s, z) #this is the second mixed partial
    ↳ derivative of Phi

print('\nFor divergence to be zero, the following mixed partial derivative must
    ↳ be zero:')
display(mixed_partial)

#substitute a simple phi(s,z) to show it vanishes.
Phi_sub = s**2 + z**2 #example function
div_B_sub = div_B.subs(Phi, Phi_sub).simplify()

print('\nSubstitute Phi(s, z) = s**2 + z**2:')
display(div_B_sub)

```

The divergence of B is :

$$\frac{\partial^2}{\partial z \partial s} \Phi(s, z)$$

πs

For divergence to be zero, the following mixed partial derivative must be zero:

$$\frac{\partial^2}{\partial z \partial s} \Phi(s, z)$$

Substitute $\Phi(s, z) = s^2 + z^2$:

0

```

[6]: #Part C
#vector symbols
Omega_x, Omega_y, Omega_z = sp.symbols('Omega_x Omega_y Omega_z')

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r_x, r_y, r_z = sp.symbols('r_x r_y r_z')
B_x, B_y, B_z = sp.symbols('B_x B_y B_z')

#vectors
Omega = sp.Matrix([Omega_x, Omega_y, Omega_z]) # Angular velocity
r = sp.Matrix([r_x, r_y, r_z]) # Radial vector
B = sp.Matrix([B_x, B_y, B_z]) # Magnetic field

#dots
Omega_dot_B = Omega.dot(B)
Omega_dot_r = Omega.dot(r)

#cross for v_perp = Omega x r
v_perp = Omega.cross(r)

#ohm in ideal MHD: E = - (v_perp x B)
E_vector = -v_perp.cross(B)

#simplify/factor
E_simplified = -(Omega_dot_B * r - Omega_dot_r * B)
E_simplified

```

[6]:

$$\begin{bmatrix} B_x (\Omega_x r_x + \Omega_y r_y + \Omega_z r_z) - r_x (B_x \Omega_x + B_y \Omega_y + B_z \Omega_z) \\ B_y (\Omega_x r_x + \Omega_y r_y + \Omega_z r_z) - r_y (B_x \Omega_x + B_y \Omega_y + B_z \Omega_z) \\ B_z (\Omega_x r_x + \Omega_y r_y + \Omega_z r_z) - r_z (B_x \Omega_x + B_y \Omega_y + B_z \Omega_z) \end{bmatrix}$$

[7]:

```

#part D

#vars
s, z, mu_0 = sp.symbols('s z mu_0')
Phi = sp.Function('Phi')(s, z)
I = sp.Function('I')(s)

#mag field components
B_s = (1 / (2 * sp.pi * s)) * sp.diff(Phi, z)
B_phi = (mu_0 * I) / (2 * sp.pi * s)
B_z = (1 / (2 * sp.pi * s)) * sp.diff(Phi, s)

#b in cylindrical
B = sp.Matrix([B_s, B_phi, B_z])

#curl b compute
curl_B = sp.Matrix([
    (1/s) * sp.diff(s*B_z, sp.symbols('phi')) - sp.diff(B_phi, z),
    sp.diff(B_s, z) - sp.diff(B_z, s),
    (1/s) * sp.diff(s*B_phi, s)
])

```

```
#current density j = (1/mu_0) * curl(B)
j = (1 / mu_0) * curl_B

#show
display(j)
```

$$\begin{bmatrix} 0 \\ -\frac{\frac{\partial^2}{\partial s^2} \Phi(s,z)}{2\pi s} + \frac{\frac{\partial^2}{\partial z^2} \Phi(s,z)}{2\pi s} + \frac{\frac{\partial}{\partial s} \Phi(s,z)}{2\pi s^2} \\ \frac{\mu_0}{\frac{d}{ds} I(s)} \frac{1}{2\pi s} \end{bmatrix}$$

```
[8]: #part E
      """
      Asked to find diffeq for current where xi is s**2/z. Expand rho E + jxB to
      ↪first order..small approximation.
      We know rhoE is charge density time electric field and jxB is cross of current
      ↪density and mag.
      """

      # Define variables
      s, z, xi, mu_0 = sp.symbols('s z xi mu_0')
      rho, I = sp.symbols('rho I', cls=sp.Function) # Charge density and current
      Phi = sp.Function('Phi')(s, z) # Magnetic flux
      B_s, B_phi, B_z = sp.symbols('B_s B_phi B_z') # Magnetic field components
      j_phi, j_z = sp.symbols('j_phi j_z') # Current density components
      E_s, E_phi, E_z = sp.symbols('E_s E_phi E_z') # Electric field components

      # Define B and j as vectors
      B = sp.Matrix([B_s, B_phi, B_z])
      j = sp.Matrix([0, j_phi, j_z])

      # Define the electric field vector
      E = sp.Matrix([E_s, E_phi, E_z])

      # Expression for rho * E + j x B
      expression = rho(s) * E + j.cross(B)

      # Expand each component of the matrix to first order in z/xi
      expansion = expression.applyfunc(lambda expr: expr.subs(s**2/z, xi).series(z/
      ↪xi, n=2))

      # Simplify the expansion
      simplified_expansion = expansion.applyfunc(lambda expr: expr.simplify())

      simplified_expansion
```

[8]:

$$\begin{bmatrix} -B_\phi j_z + B_z j_\phi + E_s \rho(s) \\ B_s j_z + E_\phi \rho(s) \\ -B_s j_\phi + E_z \rho(s) \end{bmatrix}$$

0.2 Expanded Expression

The expansion of $\rho \mathbf{E} + \mathbf{j} \times \mathbf{B}$ to first order in (z/r) as shown above

0.2.1 First Row:

- $-B_\phi j_z + B_z j_\phi + E_s \rho(s)$

Only component in the s dir (radial/cylindrical)

- $-B_\phi j_z$: force density cross B_ϕ and axial J_z .
- $B_z j_\phi$: force density from B_z with current density from azimuthal j_ϕ . radial force makes sense
- $E_s \rho(s)$: clearly electric field in s multiplied by charge density coulomb force acting radially.

0.2.2 Second Row:

- $B_s j_z + E_\phi \rho(s)$

Component of expression in phi/azimuthal

- $B_s j_z$: interaction of radial magnetic B_s and axial current dens. j_z
- $E_\phi \rho(s)$: azimuth electric field by charge density – electric force in azimuth dir

0.2.3 Third Row:

- $-B_s j_\phi + E_z \rho(s)$

resultant of z dir/axial

- $-B_s j_\phi$: force generated by interaction of B_s and j_ϕ axial force made. Lorentz??
- $E_z \rho(s)$: clearly the coulomb force along z

0.2.4 Overall

Yep pretty much makes perfect sense. Looking at it by direction:

- **Radial (s) direction:** influence by axial/azimuthal mag field and radial e field
- **Azimuthal (phi) direction:** influenced by radial mag and azimuth e-field
- **Axial (z) direction:** radial magnetic and axial electric field

[9]: *#unsure on if this is what you're asking but i'm going to make up an $I(x_i)$ that relates current density to j_ϕ and j_z*
Redefine the current density components j_ϕ and j_z in terms of $I(x_i)$

```
I_xi = sp.Function('I')(xi)
```

```

#j_phi and j_z are related to the derivative of I w/respect to xi
j_phi = sp.diff(I_xi, xi)
j_z = I_xi / xi

#sub in
new_expression = simplified_expansion.subs({sp.symbols('j_phi'): j_phi, sp.
    ↪symbols('j_z'): j_z})

#simplify
new_expression_simplified = new_expression.applyfunc(lambda expr: expr.
    ↪simplify())

# Now let's isolate the terms and find the differential equation governing I(xi)
new_expression_simplified

```

[9]:
$$\begin{bmatrix} -\frac{B_\phi I(\xi)}{\xi} + B_z \frac{d}{d\xi} I(\xi) + E_s \rho(s) \\ \frac{B_s I(\xi)}{\xi} + E_\phi \rho(s) \\ -B_s \frac{d}{d\xi} I(\xi) + E_z \rho(s) \end{bmatrix}$$

```

[10]: #0,0 and 2,0
eq_1 = new_expression_simplified[0, 0]
eq_2 = new_expression_simplified[1, 0]
eq_3 = new_expression_simplified[2, 0]

display(eq_1, eq_2, eq_3)

```

$$-\frac{B_\phi I(\xi)}{\xi} + B_z \frac{d}{d\xi} I(\xi) + E_s \rho(s)$$

$$\frac{B_s I(\xi)}{\xi} + E_\phi \rho(s)$$

$$-B_s \frac{d}{d\xi} I(\xi) + E_z \rho(s)$$

```

[11]: #solve for I(xi)
sol_1 = sp.dsolve(eq_1, I_xi)

#solve for I(xi)
sol_3 = sp.dsolve(eq_3, I_xi)

#general solution
display(sol_1, sol_3)

```

$$I(\xi) = \frac{C_1 (B_\phi - B_z) e^{\frac{B_\phi \log(\xi)}{B_z}} + E_s \xi \rho(s)}{B_\phi - B_z}$$

$$I(\xi) = C_1 + \frac{E_z \xi \rho(s)}{B_s}$$

[11] :