

Numerical implementation of Neutron star cooling

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- Relativistic stellar structure



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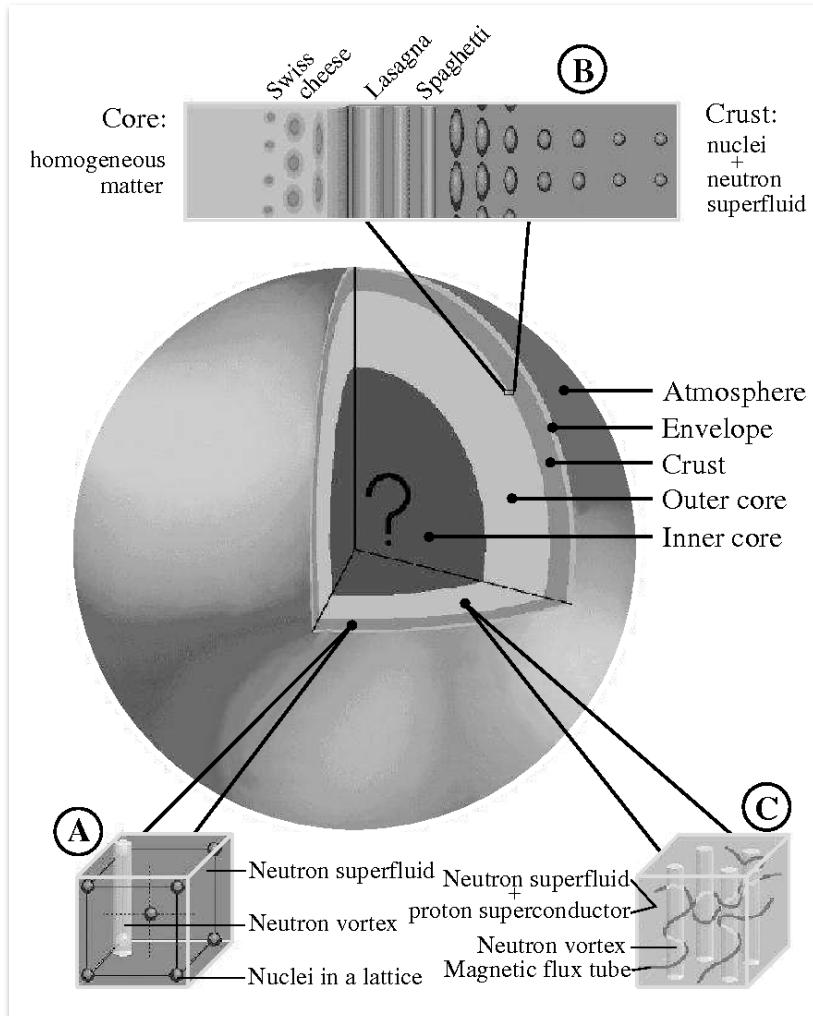


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- Relativistic stellar structure
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Basic structure of Neutron star

Basic structure of Neutron star



- **Atmosphere (10cm)**

Determines the thermal radiation (the spectrum). Of upmost importance for interpretation of X-ray (and optical) observation. However it has NO effect on the thermal evolution of the star.

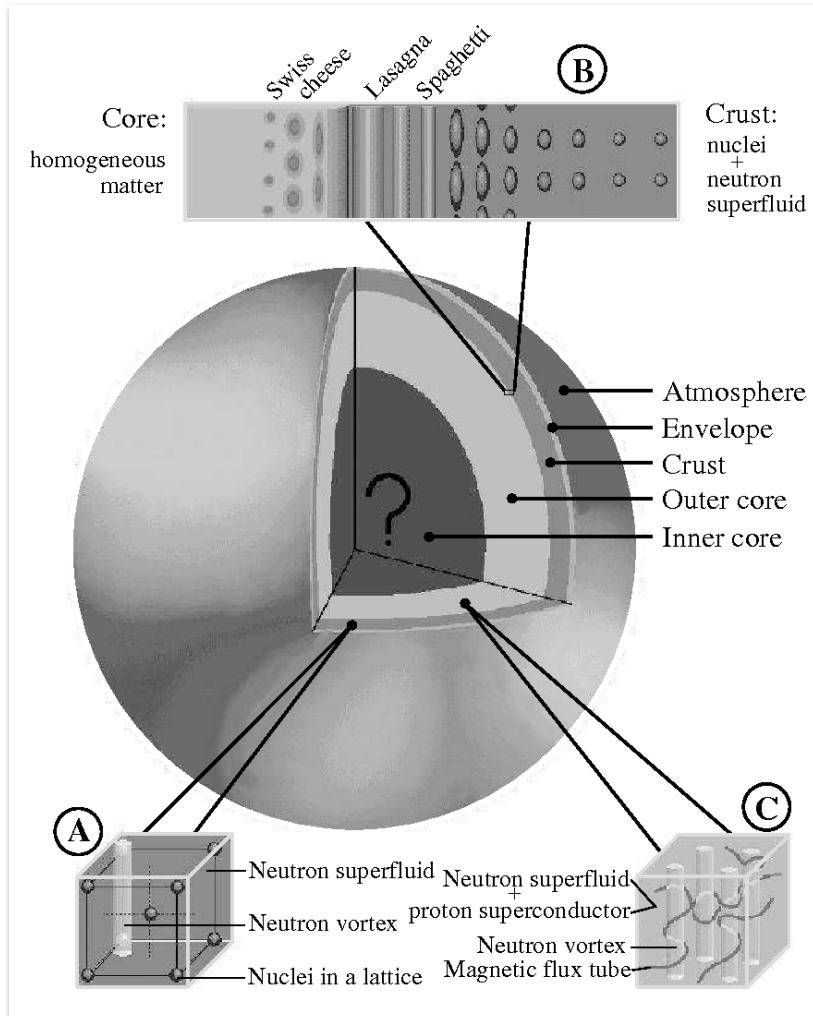
- **Envelope (100m)**

Contains a huge temperature gradient - extremely important for the cooling, strongly affected by magnetic fields and the presence of *polluting* light elements

- **Crust (1km)**

Composed of nuclei, quasi-free neutrons. The nuclei are arranged into a crystalline lattice. At bottom of crust, *pasta phases* of nuclear matter forms mantle.

Basic structure of Neutron star



- **Outer Core (10km)**

Contains most of the stellar mass, composed of neutrons with an admixture of protons & leptons - electrons, muons ($npe\mu$ matter)

- **Inner Core (? km)**

Exists for massive neutron star, compositions & properties are not well known. There are some proposed theoretical models:

1. **hyperonization of matter** – the appearance of various hyperons (first of all, Λ^- and Σ^- hyperons – $(npe\mu\Lambda\Sigma)$ matter)
2. **pion condensation** – formation of a Bose condensate of collective interactions with the properties of π -mesons.
3. **kaon condensation** – formation of a similar condensate of K -mesons.



Basic structure of neutron star

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 - Describe relativistic stellar structure



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 - **Tolman-Oppenheimer-Volkoff equations**
 - Describe relativistic stellar structure
 - **Thermal evolution equations**
 - Describe stellar evolution

Relativistic stellar structure



Tolman-Oppenheimer-Volkoff Equation

- For static, non-rotating and spherical symmetric star, metric is given as

$$ds^2 = -e^{2\Phi(r)}dt^2 + e^{2\Lambda(r)}dr^2 + r^2d\Omega^2$$



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- For isolated star, the metric must reduce to the Schwarzschild metric at outside of the star

$$ds^2 = - \left(1 - \frac{2M}{r}\right) dt^2 + \frac{1}{1 - \frac{2M}{r}} dr^2 + r^2 d\Omega^2$$



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- At interior of the star, let define new metric function for convenience

$$e^{-2\Lambda(r)} = 1 - \frac{2m(r)}{r}$$

The quantity $m(r)$ can be interpreted as the **mass interior to radius r** .



Tolman-Oppenheimer-Volkoff Equation

- Now, consider perfect fluid matter

$$T_{\mu\nu} = (\rho + P)u_\mu u_\nu + Pg_{\mu\nu}, \quad u = e^{-\Phi} \frac{\partial}{\partial t}$$



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- Rewrite: (index-free notation)

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- From Einstein equation, we can get two equations: [Specific calculation with SageManifolds](#)

$$(0, 0) \Rightarrow \frac{dm(r)}{dr} = 4\pi r^2 \rho$$

$$(1, 1) \Rightarrow \frac{d\Phi}{dr} = \frac{m(r) + 4\pi r^3 P}{r^2 \left(1 - \frac{2m(r)}{r}\right)}$$



Tolman-Oppenheimer-Volkoff Equation

- From energy conservation, we can get one more equation:

$$\nabla_\mu T^{\mu\nu} = 0 \Rightarrow (\rho + P) \frac{d\Phi}{dr} + \frac{dP}{dr} = 0$$



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- From energy conservation, we can get one more equation:

$$\nabla_\mu T^\mu{}_\nu = 0 \Rightarrow (\rho + P) \frac{d\Phi}{dr} + \frac{dP}{dr} = 0$$

- These three fundamental equations are called **Tolman-Oppenheimer-Volkoff Equation**

$$\begin{aligned}\frac{dm}{dr} &= 4\pi r^2 \rho \\ \frac{d\Phi}{dr} &= \frac{m + 4\pi r^3 P}{r^2 \left(1 - \frac{2m}{r}\right)} \\ \frac{dP}{dr} &= -\frac{(\rho + P)(m + 4\pi r^3 P)}{r^2 \left(1 - \frac{2m}{r}\right)}\end{aligned}$$



Polytropic Equation of State

- Neutron stars are well modeled by polytropes

$$P = K \rho_0^\Gamma$$

where

K : polytropic gas constant

ρ_0 : rest mass density

$\rho_0 = m_u n_B$ (for neutron star)

m_u : nucleon mass - atomic mass unit

n_B : baryonic number density

$\rho = \rho_0(1 + \epsilon)$

ϵ : internal energy density per unit mass

$$\Gamma \equiv 1 + \frac{1}{n}$$

n : polytropic index



Polytropic Equation of State

- By 1st law of thermodynamics with isentropic assumption, we can get

$$P = n_B^2 \frac{\partial(\rho/n_B)}{\partial n_B}$$



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$$P = n_B^2 \frac{\partial(\rho/n_B)}{\partial n_B}$$

- Substitute polytropic EoS:

$$\rho = \rho_0 + \frac{P}{\Gamma - 1} = \left(\frac{P}{K} \right)^{\frac{1}{\Gamma}} + \frac{\rho}{\Gamma - 1}$$



Solving TOV equation

- Total running terms

$$\frac{dm}{dr} = 4\pi r^2 \rho$$

$$\frac{dP}{dr} = -\frac{(\rho + P)(m + 4\pi r^3 P)}{r^2 \left(1 - \frac{2m}{r}\right)}$$

$$\frac{d\Phi}{dr} = \frac{m + 4\pi r^3 P}{r^2 \left(1 - \frac{2m}{r}\right)}$$

$$\frac{dN_B}{dr} = \frac{4\pi r^2 n_B}{\sqrt{1 - \frac{2m}{r}}}$$

$$\rho_0 = \left(\frac{P}{K}\right)^{\frac{1}{\Gamma}}$$

$$\rho = \rho_0 + \frac{P}{\Gamma - 1} = \left(\frac{P}{K}\right)^{\frac{1}{\Gamma}} + \frac{P}{\Gamma - 1}$$

$$n_B = \frac{\rho_0}{m_u} = \frac{1}{m_u} \left(\frac{P}{K}\right)^{\frac{1}{\Gamma}}$$



Solving TOV equation

- Total boundary conditions

$$m(0) = 0$$

$$\rho_0(0) = \rho_c$$

$$P(0) = P_c = K\rho_c^\Gamma$$

$$n_B(0) = n_c = \frac{\rho_c}{m_u}$$

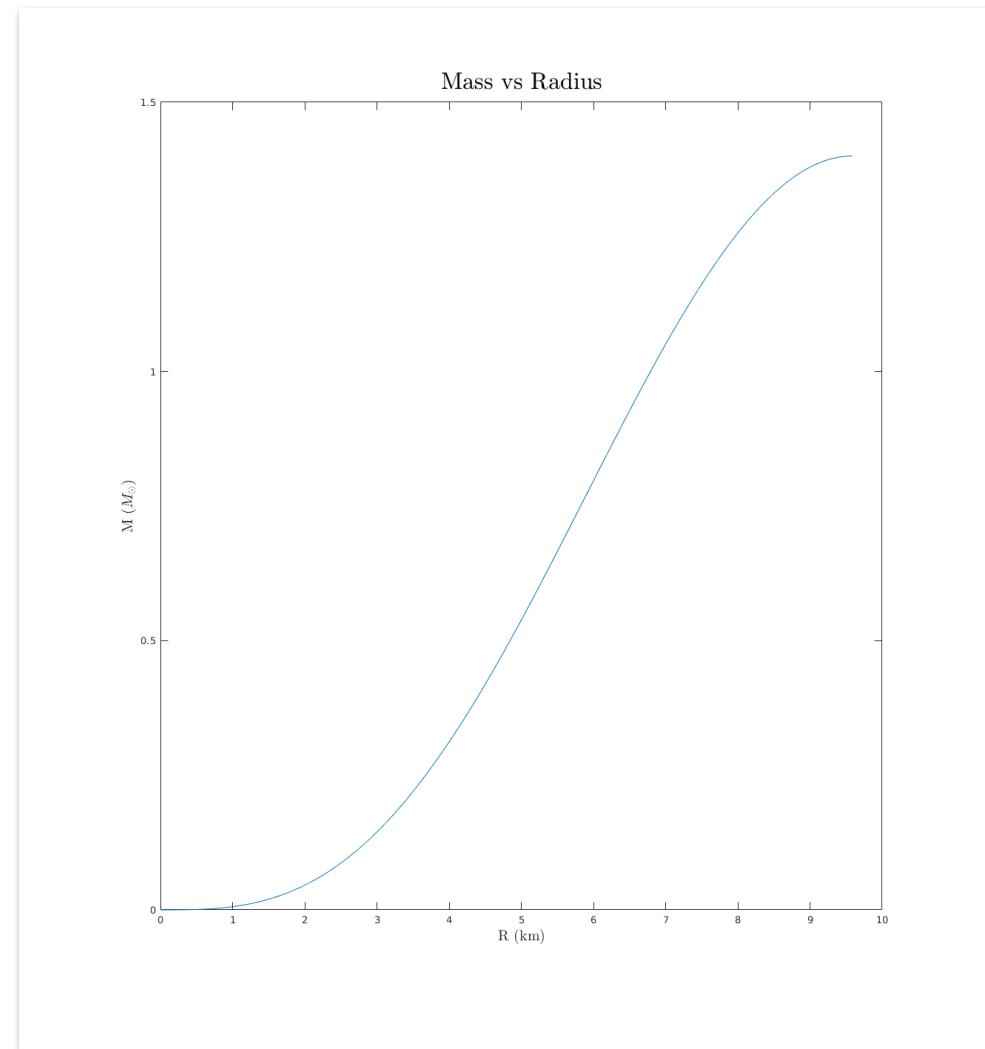
$$P(R) = 0$$

$$\Phi(R) = \frac{1}{2} \ln \left(1 - \frac{2M}{R} \right)$$

Results ($K = 100, \Gamma = 2, \rho_c = 1.28 \times 10^{-3}$)

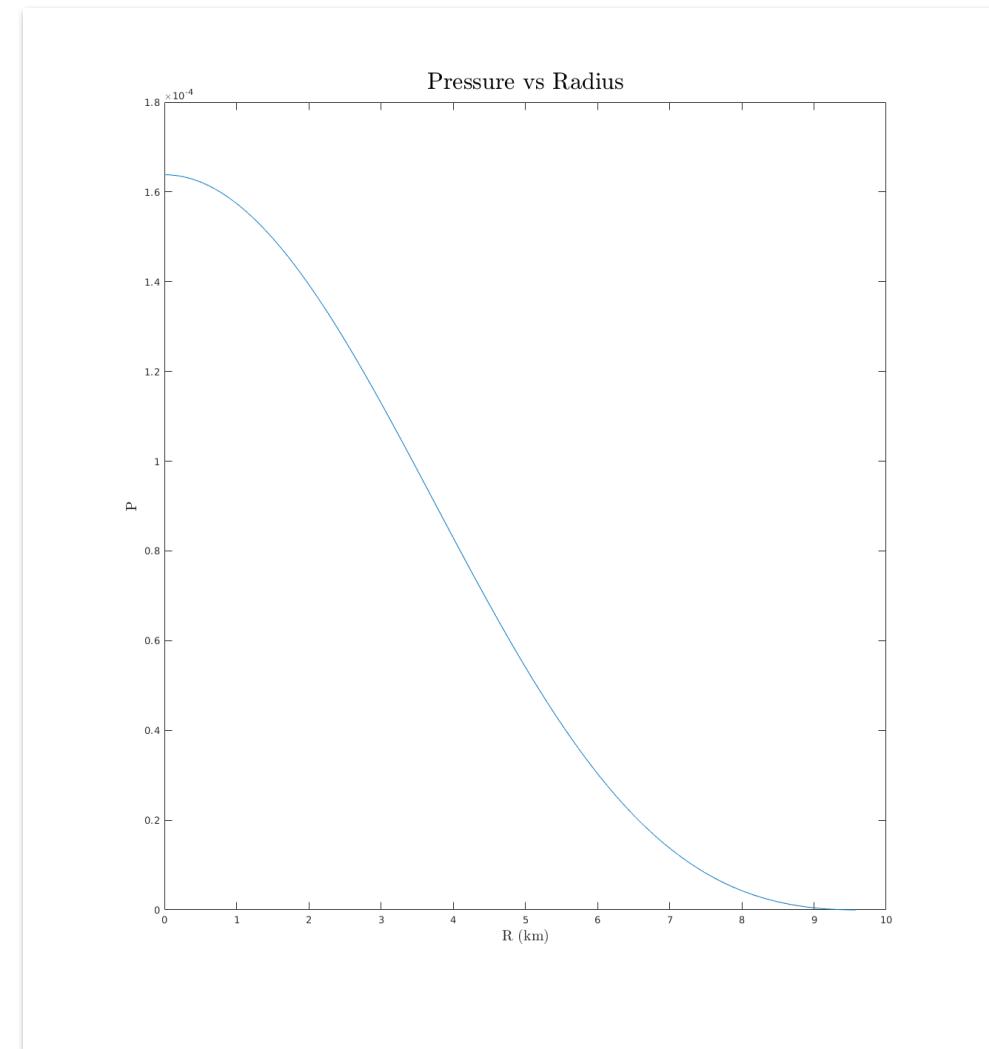
$$R = 9.5850\text{km}$$

$$M = 1.4002M_{\odot}$$



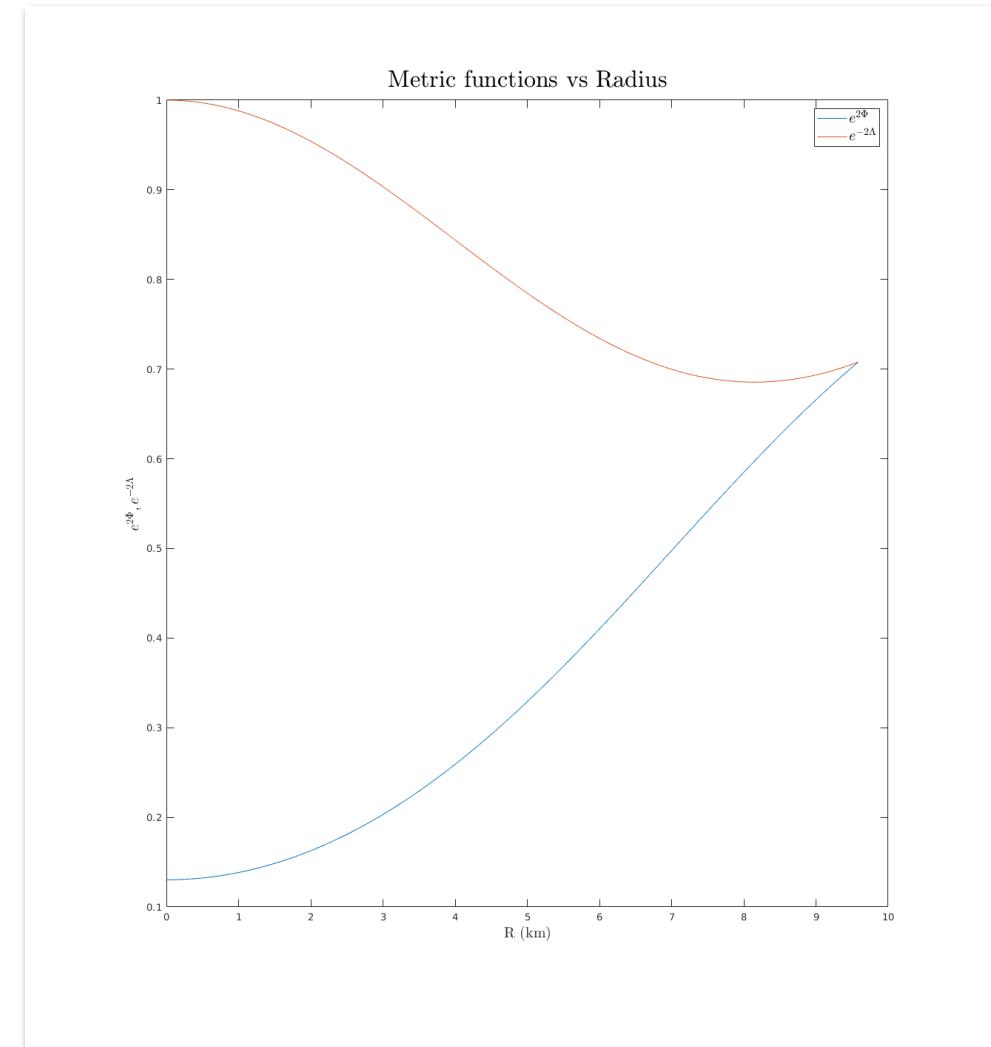
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$$P(0) = P_c, \quad P(R) = 0$$



Results ($K = 100, \Gamma = 2, \rho_c = 1.28 \times 10^{-3}$)

$$e^{2\Phi(R)} = e^{-2\Lambda(R)}$$



Implicit scheme to solve stiff ODE



Brief view - Thermal evolution equations

- Energy balance

$$\frac{d(L e^{2\Phi})}{dr} = - \frac{4\pi r^2 e^\Phi}{\sqrt{1-2m/r}} \left(C_v \frac{dT}{dt} + e^\Phi (Q_\nu - Q_h) \right)$$



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$$\frac{d(T e^\Phi)}{dr} = - \frac{1}{\lambda} \frac{L e^\Phi}{4\pi r^2 \sqrt{1-2m/r}}$$



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$$\frac{d(T e^\Phi)}{dr} = - \frac{1}{\lambda} \frac{L e^\Phi}{4\pi r^2 \sqrt{1-2m/r}}$$

where

- C_v : Specific heat
- L : Diffusive luminosity at r
- λ : Thermal conductivity
- Q_ν : Neutrino emissivity
- Q_h : Heating late



Brief view - Thermal evolution equations

- Let's rewrite these equations with red-shifted functions & Lagrangian coordinate a

$$\mathcal{T} = e^\Phi T, \quad \mathcal{L} = e^{2\Phi} L, \quad da = \frac{4\pi r^2 n_B dr}{\sqrt{1 - 2m/r^2}}$$



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$$\frac{d\mathcal{T}}{dt} = F \left(\mathcal{T}, \frac{d\mathcal{L}}{da} \right) \quad \mathcal{L} = G \left(\mathcal{T}, \frac{d\mathcal{T}}{da} \right)$$



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- Our choice is **Backward Euler** - Most robust method (but high cost)



Brief view - Backward Euler

- Suppose that there is non-autonomous ODE system

$$\mathbf{Y}' = f(t, \mathbf{Y})$$



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- I have achieved to reduce high cost of Backward Euler with AD & Rust (High performance programming language)

TODO



Current Progress

Symbolic

- Solve TOV from spherical symmetric metric & polytropic EM Tensor



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Numeric

- Implement Explicit RK4
- Apply ERK4 to TOV eq
- Implement Implicit Method
 - Backward Euler Method



TODO

Numeric

- Implement more implicit methods
 - 4th order Gauss-Legendre (GL4)
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 - Focus on Cas A (Cassiopeia A)
 - Modified Urca
 - Bremstrahlung
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Interface

- Implement one-click (or type) installation for any OS
- Various language API (Python or even Mathematica?)

Thank you