# Scientific Computing Lecture 3

Differential equations Nikolay Koshev October 9, 2023



Skolkovo Institute of Science and Technology

#### Structure of the lecture

#### Differential equations

- Differential operator
- Cauchy problems
- Boundary-value problems
- Eigenvalue problems
- Partial differential equations of the first order
- Classification of second order PDEs



# Scientific Computing Lecture 3

Part 1: Differential operator Nikolay Koshev October 9, 2023

# Skoltech

Skolkovo Institute of Science and Technology

### Differential equations

# Many Physical Problems can be described with use of Differential Equations

- A differential equation is an equation for an unknown function of one or several variables that relates the values of the function itself and of its derivatives of various orders.
- Ordinary Differential Equation: Function has 1 independent variable.
- ▶ Partial Differential Equation: At least 2 independent variables.



### Differential operator

The operator equation for differential equations is often written as:

$$D\mathbf{u} = 0,$$

where D represents both mathematical model and the right-hand side; the function  ${\bf u}$  is an unknown function to be found.

#### **Examples**

Ordinary differential equation:

$$\frac{d\mathbf{u}}{dx} = f(x) \Leftrightarrow D \equiv \frac{d}{dx} - f(x).$$

Partial differential equation (Poisson equation, 2D):

$$\Delta \mathbf{u} = -f(\mathbf{u}, x, y); \quad \Delta \mathbf{u} = \frac{\partial^2 \mathbf{u}}{\partial x^2} + \frac{\partial^2 \mathbf{u}}{\partial y^2}; \quad D\mathbf{u} \equiv \Delta \mathbf{u} + f(\mathbf{u}, x, y)$$



### Differential operator

▶ Ordinary differential equation (ODE): Let  $x \in \mathbb{R}^1$ . ODE of the k-th order can be represented with the operator:

$$D\mathbf{u} = F(x, u(x), u'(x), u''(x), ..., u^{(k)}(x)).$$

Partial differential equation (PDE): Let  $\mathbf{x} \in \mathbb{R}^n \equiv (x_1, x_2, ..., x_n)$ . PDE of the k-th order can be represented with the following operator:

$$D\mathbf{u} = F\left(\mathbf{x}, \mathbf{u}(\mathbf{x}), \frac{\partial \mathbf{u}}{\partial x_1}, ..., \frac{\partial \mathbf{u}}{\partial x_n}, ..., \frac{\partial^2 \mathbf{u}}{\partial x_1^2}, ..., \frac{\partial^2 \mathbf{u}}{\partial x_n^2}, ..., \frac{\partial^{(k)} \mathbf{u}}{\partial x_1^{(k)}}, \frac{\partial^{(k)} \mathbf{u}}{\partial x_n^{(k)}}\right).$$



### Differential operator

- ▶ The function F(...) usually satisfies the regularity conditions: measurability, differentiability, continuity etc.
- ▶ **Linear equations:** If the function *F* depends linearly on **u** and its derivatives, then the ODE/PDE is linear.
- ▶ Quasi-linear equations: the function *F* linearly depends on higher-order derivatives.
- Non-linear equations: the most complicated case. The function F depends non-linearly on higher order (and, maybe, all other) derivatives.



#### ODE: Problems

Ordinary Differential Equation (ODE) of  $p^{th}$  order can be written in as follows:

$$F(x, u'(x), u''(x), ..., u^{(p)}(x)) \equiv D\mathbf{u} = 0.$$

Solve the equation above with respect to higher order derivative  $u^{(p)}$  to express it via all other derivatives:

$$u^{(p)}(x) = f(x, u, u', u'', ..., u^{(p-1)}).$$

We refer to the fact that with  $u^{(k)}(x) \equiv u_k(x)$  and  $u(x) = u_0(x)$ , the latter equation may be represented with a system of the first-order ODEs:

$$u_{k+1}(x) = u'_{k}(x), \quad 0 \le k \le p-2,$$
  
 $u'_{p}(x) = f(x, u_{0}, u_{1}, ..., u_{p-1}),$ 



# ODE: Three kinds of problems

Let  $\mathbf{u} = \{u_0, ..., u_p\}$ , and  $\mathbf{f} = \{f_1, ..., f_p\}$  By analogy, each system of differential equations of any order can be changed with equivalent system of the first-order ODE:

$$u'_{k} = f(x, u_0, u_1, ..., u_p),$$

Or, the vector form of the system above:

$$\mathbf{u}'(x) = \mathbf{f}(x, \mathbf{u}(x)), \quad \mathbf{u} = \{u_0, ..., u_p\}, \mathbf{f} = \{f_1, ..., f_p\}.$$

There are three kinds of problem related to the ODEs

- ► The Cauchy problem
- Boundary value problem
- ► Eigenvalues problem



# Scientific Computing Lecture 3

Part 2: Cauchy Problems Nikolay Koshev October 9, 2023

# Skoltech

Skolkovo Institute of Science and Technology

# ODE: The Cauchy problem

Vector form:

$$\begin{aligned} \mathbf{u}'(x) &= \mathbf{f}\big(x, \mathbf{u}(x)\big), \quad x \in [\xi, X] \\ \mathbf{u}(\xi) &= \eta, \quad \text{or } u_k(\xi) = \eta_k, 1 \leq k \leq p. \end{aligned}$$

Or common form:

$$F\left(x, u(x), u'(x), \dots, u^{(p)}(x)\right) \equiv D\mathbf{u} = 0.$$
$$u^{(k)}(\xi) = \eta_k, \quad 0 \le k < p$$

The conditions may be considered as the definition of some initial point  $(\xi, \eta_0, ..., \eta_{p-1})$  for an integral curve in (p+1-dimensional space  $(x, u_0, ..., u_p)$ .



Augustin-Louis Cauchy 1789-1857

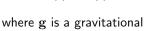


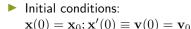
# The simplest Cauchy problem

Consider a ball with mass m. Let this ball be thrown; we need to calculate its trajectory.

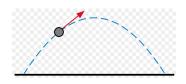
- ightharpoonup m the mass of the ball
- $\mathbf{x} \in L_2(\mathbb{R}^2) \times L_2(\mathbb{R}^2)$
- $\mathbf{v} \in L_2(\mathbb{R}^2) \times L_2(\mathbb{R}^2)$
- ► The governing equation:

$$\mathbf{x}''(t) = \mathbf{g},$$
  
 $\mathbf{x}'(t) = \mathbf{v}(t),$ 





acceleration.





# The simplest Cauchy problem

The solution of the problem can be easily found using integration:

$$\mathbf{v}(t) = Const + \int_{0}^{t} \mathbf{g} d\tau = \mathbf{v}_{0} + \mathbf{g}t,$$

$$\mathbf{x}(t) = Const + \int_{0}^{t} \mathbf{v}(\tau) d\tau = \mathbf{x}_{0} + \mathbf{v}_{0}t + \mathbf{g}t^{2}$$

Since in our case g has no x-component ( $\mathbf{g} = (0, -9.8)^T$ ), then, as we can see, the x component of the speed will be fixed. But...

- What if we have not only one body (ball)?
- ► What if we have air?
- ▶ What if we have other forces affecting the bodies?



# N-body problem

Consider a number of point-bodies  $B_i, i=1,...,N$ , defined with the coordinates  $x_i(t)$  at the moment of time t. Each body interacts with other bodies via the gravitation:

$$m_i \ddot{x}_i(t) = \sum_{j \neq i}^N G m_j m_i \frac{x_j(t) - x_i(t)}{|x_j(t) - x_i(t)|^3}, \quad 1 \leq i, j \leq N.$$

In order to state the problem, we also need to define the initial body positions with known coordinates  $\tilde{x}_i$ :

$$x_i(t_0) = \tilde{x}_i,$$

and the initial velocities with known velocities:

$$v_i(t_0) \equiv \dot{x}_i(t_0) = \tilde{v}_i.$$

Thus, we have the Cauchy problem.



# Solution of the Cauchy problem

The vector form:

$$\mathbf{u}'(x) = \mathbf{f}\big(x, \mathbf{u}(x)\big),$$
 
$$\mathbf{u}(\xi) = \eta, \quad \text{or } u_k(\xi) = \eta_k, 0 \le k \le p.$$

- 1. If the RHS  $\mathbf{f}(x, \mathbf{u})$  are continuous and bounded in some neighborhood of the initial point, then the solution  $\mathbf{u}(x)$  of a Cauchy problem exists (buy might be not unique!).
- 2. If the RHS additionally Lipschitz-continuous with respect to  $u_k$ , i.e. if  $\exists K$ :

$$|\mathbf{f}(x,\mathbf{u}_1) - \mathbf{f}(x,\mathbf{u}_2)| \le K|\mathbf{u}_1 - \mathbf{u}_2|,$$

then the solution  $\mathbf{u}(x)$  of a Cauchy problem is unique (**but might be not stable!**)

3. If, additionally to the conditions above, the RHS has continuous derivatives by all arguments up to  $q^{th}$  order, then the solution  $\mathbf{u}(x)$  of a Cauchy problem has continuous derivatives up to order q+1.



# Solving the Cauchy problem

- **Exact methods:** the solution of the Cuachy problem may be represented with elementary functions. Rarely applicable even for simple equations. For example, the solution of the equation  $u'(x) = x^2 + u^2(x)$  can not be presented with elementary functions (or just not reasonable).
- ▶ **Approximate methods:** the solution  $\mathbf{u}(x)$  is representable as a limit of some sequence  $y_n(x)$ , members of which can be represented with elementary functions or with quadratures. Applicable only for relatively simple linear problems.
- ▶ **Numerical solution:** applicable for a wide class of Cauchy problems. The problem must be well-conditioned.



# ODU: when the analytic solution is not reasonable

Consider the equation:

$$u'(x) = \frac{u - x}{u + x}$$

This equation can be integrated and we will find the equation for its solution:

$$\frac{1}{2}ln(x^2 + u^2) + arctg\frac{u}{x} = const.$$

However, solution of the latter equation is even more complicated than just numerical solution of the initial problem.



# ODU: ill-conditioned problem example

$$u'(x) = u - x$$
,  $0 \le x \le 100$ ,  $u(0) = 1$ .

The solution contains the constant being defined with the boundary condition

$$u(x;c) = 1 + x + ce^x,$$

which is equal to zero with the boundary condition u(0)=1, and to u(100)=101. However, even changing the boundary condition with  $u(0)=1+10^{-6}$  will change it to  $c=10^{-6}$ , which will result in  $u(100)=2.7\cdot 10^{37}$ .



# Methods of solving the Cauchy problems

- The Picard method.
- Poincare method (small parameter method).
- Runge-Kutta methods.
- Linear Multistep method.
- Implicit schemes.
- And many others...



# Runge-Kutta methods

- Numerical methods, being used with meshes.
- ▶ Explicit methods are commonly more accurate; there are some equations, however, which are not reasonable to solve with explicit schemes due to low stability/accuracy/performance. These equations called 'stiff equations'
  - Euler's method:
  - Runge-Kutta  $4^{th}$  order method;
- ▶ Implicit methods are less accurate and somewhat slower some ODE; however, due to better stability, implicit methods are capable for wider class of problems.



#### The Euler's method

The Cauchy problem:

$$u'(x) = f(x, u(x)), \quad \xi \le x \le X, \quad u(\xi) = \eta.$$

Let  $\{x_n, 0 \le n \le N\}$  be some mesh covering the interval  $[\xi, N]$ . Using the Taylor's series for the interval  $[x_n, x_{n+1}]$ :

$$u_{n+1} = u_n + h_n u'_n + \frac{1}{2} h_n^2 u''_n + \dots, \quad h_n = x_{n+1} - x_n$$

The equation states that u'(x)=f(x,u). Assuming the steps h are small and ignoring higher order derivatives multiplied with  $h_n^t, t>1$ , we obtain the Euler's method. The accuracy of the method is  $O(\max h_n)$ .



# $4^{th}$ order Runge-Kutta scheme

The Cauchy problem:

$$u'(x) = f(x, u(x)), \quad \xi \le x \le X, \quad u(\xi) = \eta.$$

Using the same mesh,  $\{u_n\}$  can be calculated as follows:

$$u_{n+1} = u_n + \frac{h}{6}(k_1 + 2k_2 + 2k_3 + k_4),$$

$$k_1 = f(x_n, u_n);$$

$$k_2 = f(x_n + \frac{h}{2}, u_n + \frac{h}{2}k_1);$$

$$k_3 = f(x_n + \frac{h}{2}, u_n + \frac{h}{2}k_2);$$

$$k_4 = f(x_n + h, y_n + hk_3).$$

The accuracy of the method is  $O(h^4)$ .



# Implicit Runge-Kutta scheme

The Cauchy problem:

$$u'(x) = f(x, u(x)), \quad \xi \le x \le X, \quad u(\xi) = \eta.$$

Using the same mesh,  $\{u_n\}$  can be calculated as follows:

$$u_{n+1} = u_n + h \frac{f(x_n, u_n) + f(x_{n+1}, u_{n+1})}{2}.$$

- ▶ The latter equation might not have a solution.
- ► The scheme will converge only with small *h*, which makes the method more computationally complicated.
- ▶ We need to bound the number of iterations.



# Scientific Computing Lecture 4

Part 3: Boundary-Value Problems Eigenvalue Problems Nikolay Koshev

October 9, 2023



Skolkovo Institute of Science and Technology

# Boundary value problems

#### The system of ODE:

$$u'_k(x) = f_k(x, u_1, u_2, ..., u_p), \quad 1 \le k \le p, \quad x \in [a, b]$$

In the Boundary Value problem, the conditions are being set both in points x=a and x=b.

- ▶ Possible with  $p \ge 2$ .
- Approximate methods: Fourier-based methods, Ritz methods, Galerkin methods.
- Numerical methods: Shooting method, Finite Differences.



# Boundary value problem

Simulation of loaded string (static deflection).

$$u''(x) = -f(x), \quad a \le x \le b, u(a) = u(b) = 0,$$

where f(x) is a deforming force.

Ballistic trajectory in space:

$$\mathbf{x}''(t) = \mathbf{g}, \quad \mathbf{x}(a) = \mathbf{x}_a, \mathbf{x}(b) = \mathbf{x}_b, |\mathbf{x}'(t)| = v_0.$$



# Shooting method

Consider the second order ODE:

$$u''(x) = f(x, u(x), u'(x)), \quad u(a) = u_a, u(b) = u_b.$$

Instead, we are looking for a parametrized solution u(x, d):

$$u''(x) = f(x, u(x), u'(x)), \quad u(a) = u_a, u'(a) = d.$$

If we find d such that  $u(b;d)=u_b$ , than the problem is solved. Thus:

The boundary value problem is being reduced to the set of Cauchy problems



# Shooting method

If the equation is linear: f(x,u,u') = p(x)u'(x) + q(x)u(x) + r(x), then

$$u(x) = u_1(x) + \frac{u_b - u_1(b)}{u_2(b)} u_2(x),$$

where  $u_1$  and  $u_2$  are the solutions of the following problems respectively:

$$u_1''(x) = f(x, u_1, u_1'), \quad u_1(a) = u_a, u_1'(a) = 0,$$
  
$$u_2''(x) = f(x, u_2, u_2') - r(x), \quad u_2(a) = 0, u_2'(a) = 1,$$

If the problem is non-linear, than we have to iterate in order to find a proper parameter d.



#### The Galerkin method

#### Consider the equation:

$$A(u(x)) = f(x), \quad a \le x \le b, u(a) = u_a, u(b) = u_b.$$

- Approximate the function  $u(x) \approx y_n(x) = \varphi_0(x) + \sum\limits_{k=1}^n c_k \varphi_k(x),$  where the continuous function  $\varphi_0(x)$  satisfies the boundary conditions:  $\varphi_0(a) = u_a, \varphi_0(b) = u_b;$  the functions  $\varphi_k(x), 1 \leq k < \infty$  are linearly independent and vanish at points a and b:  $\varphi_k(a) = \varphi_k(b) = 0.$
- ▶ The algebraic system of equations for the coefficients  $c_k$ :

$$\int_{a}^{b} \left( A(y_n(x)) - f(x) \right) \varphi_k(x) dx = 0, \quad 1 \le k \le n$$

ightharpoonup After definition of the coefficients  $c_k$ , we construct the approximate solution of the BV problem.



# Eigenvalue problems

#### The Eigenvalue problem:

$$\mathbf{u}'(x) = \mathbf{f}(x, \mathbf{u}; \lambda_r, r = 1, ..., q);$$
  
$$\mathbf{u} = \{u_1, u_2, ..., u_p\}, \quad \mathbf{f} = \{f_1, f_2, ..., f_p\}.$$

- Need to find both  $u_k$  and  $\lambda_r$ .
- Additional (boundary) conditions needed: p + q.
- ▶ The solution:  $u_k$  and  $\lambda_r$  eigenfunctions and eigenvalues.
- $\triangleright \lambda_k$  is also called **spectrum**.



# Oscillations of the string

The wave (D'Alembert) equation:

$$\left(\Delta - \frac{1}{v^2} \frac{\partial^2}{\partial t^2}\right) u = 0, \quad a \le x \le b, u(a) = 0; \quad u(b) = 0.$$

Assume the solution can be presented in a form  $u(x,t)=y(x)\cdot exp(i\omega t)$ . After substitution to the equation above, we obtain the following ODE:

$$\frac{d}{dx}\left(p(x)\frac{dy(x)}{dx}\right) = -k^2q(x)y(x), \quad y(a) = 0, y(b) = 0.$$

where  $k=\omega/v$  is a wavenumber,  $\omega$  - its frequency, and v is the speed.



### The Shooting method

Ignoring the right boundary condition, we have the Cauchy problem

$$\frac{d}{dx}\left(p(x)\frac{dy(x)}{dx}\right) = -k^2q(x)y(x), \quad y(a) = 0$$

the solution of which  $y(x;\boldsymbol{k})$  depends on the values of the parameter  $\boldsymbol{k}.$ 

The eigenvalues may be found then by changing the parameter k until the right boundary value be satisfied. Here the parameter k represents the proper (eigen) frequency of the string.



# Scientific Computing Lecture 3

Part 4: Partial differential equations of the first order Partial differential equations of the second order Nikolay Koshev

October 9, 2023



Skolkovo Institute of Science and Technology

# Partial differential equations of the first order

The Partial differential equation of the first order:

$$D\mathbf{u} = F\Big(\mathbf{x}, \mathbf{u}, \frac{\partial \mathbf{u}}{\partial x_1}, ..., \frac{\partial \mathbf{u}}{\partial x_n}\Big)$$

- N-body problem (Hamilton's equations)
- The Hopf equation.
- Eikonal equation.
- ► Molecular dynamics method.



# The N-body problem: potential energy

As we wrote before:

$$m_i \ddot{\mathbf{x}}_i = \sum_{j \neq i}^N G m_i m_j \frac{\mathbf{x}_j - \mathbf{x}_i}{|\mathbf{x}_j - \mathbf{x}_i|^3}, \quad 1 \leq i, j \leq N.$$
$$\mathbf{x}_i(t_0) = \tilde{\mathbf{x}}_i,$$
$$\mathbf{v}_i(t_0) \equiv \dot{\mathbf{x}}_i(t_0) = \tilde{\mathbf{v}}_i.$$

The equation, can be formulated in terms of the potential energy:

$$m_i \ddot{x}_i = \sum_{j \neq i}^N G m_i m_j \frac{x_j - x_i}{|x_j - x_i|^3} = -\frac{\partial U}{\partial \mathbf{x}_i},$$

where  $U(\mathbf{x})$  is self-potential energy:

$$U = -\sum_{1 \le i \le j \le n} \frac{Gm_i m_j}{|\mathbf{x}_j - \mathbf{x}_i|^3}.$$



# The Molecular Dynamics (MD) method

Microcanonical ensemble: the system is isolated from changes in moles (N), volume (V), and energy (E): adiabatic process with no heat exchange.

#### Newton's law of motion:

$$F(X) = -\nabla U(X) = M\dot{V}(t)$$
$$V(t) = \dot{X}(t).$$

#### Here:

- X are the particle coordinates.
- U(X) is a potential energy of the system (comes from pair or many-body potentials).
- ► M mass of particles.
- ightharpoonup V velocities of particles.



#### Some other examples

The eikonal equation:

$$|\nabla u(\mathbf{x})| = \frac{1}{f(\mathbf{x})}, \quad f(\mathbf{x}) > 0, \mathbf{x} \in \Omega, u(\mathbf{x})|_{\partial\Omega} = q(\mathbf{x});$$

► In electrostatics:

$$\mathbf{E} = -\nabla V$$
,

where V is electric potential, and  ${\bf E}$  is an electric field.

In fluid dynamics:

$$\mathbf{v} = \nabla \varphi,$$

where v is potential flow velocity, and  $\varphi$  is a potential.



### $2^{nd}$ order PDEs



#### Partial differential equations of second order

Most of PDE-driven problems can be described with PDEs of second order. The differential operator D can be represented as follows:

$$D\mathbf{u} = \sum_{i=1}^{n} \sum_{j=1}^{n} a_{ij}(\mathbf{x}) \frac{\partial^{2} \mathbf{u}}{\partial x_{i} \partial x_{j}} + F(\mathbf{x}, \mathbf{u}, \frac{\partial \mathbf{u}}{\partial x_{1}}, ..., \frac{\partial \mathbf{u}}{\partial x_{n}}) \equiv L\mathbf{u} + F(...).$$

The operator L is called the Principle Part and being defined with the matrix:

$$L\mathbf{u} = \begin{pmatrix} a_{11} & a_{12} & \dots & a_{1n} \\ a_{21} & \ddots & & \vdots \\ \vdots & & \ddots & \vdots \\ a_{n1} & \dots & \dots & a_{nn} \end{pmatrix} \begin{pmatrix} \frac{\partial^2 \mathbf{u}}{\partial x_1 \partial x_1} & \frac{\partial^2 \mathbf{u}}{\partial x_2 \partial x_1} & \dots & \frac{\partial^2 \mathbf{u}}{\partial x_n \partial x_1} \\ \frac{\partial^2 \mathbf{u}}{\partial x_1 \partial x_2} & \ddots & & \vdots \\ \vdots & & \ddots & \vdots \\ \frac{\partial^2 \mathbf{u}}{\partial x_1 \partial x_n} & \dots & \dots & \frac{\partial^2 \mathbf{u}}{\partial x_n \partial x_n} \end{pmatrix} = A_L H^T.$$

Here the matrix H is a Hessian.



#### General classification

We are able to find the eigenvalues  $\lambda$  of the matrix  $A_L$ :

$$A_L \mathbf{v} = \lambda \mathbf{v}$$

- ▶ Suppose we have found several eigenvalues  $\lambda_i$ .
- ▶ Elliptic equation:  $\lambda_i > 0$  or  $\lambda_i < 0$  for  $\forall i$ . All-positive of all-negative eigenvalues
- ▶ Parabolic equation:  $\exists k : \lambda_k = 0$  while  $\lambda_i > 0$  or  $\lambda_i < 0$  for  $\forall i \neq k$ . There is only one zero eigenvalue while all others have the same sign.
- ▶ **Hyperbolic equation:**  $\exists k: \lambda_k > 0$  while  $\lambda_i < 0 \forall i \neq k$  OR  $\exists k: \lambda_k > 0$  while  $\lambda_i < 0 \forall i \neq k$  The signs of eigenvalues are the same excepting one of them, which is non-zero and has contrary sign.



#### Elliptic equations

#### The equation

$$L\mathbf{u} + F(\dots) = 0 \tag{1}$$

is considered elliptic if the coefficient matrix  $\{a_{ij}\}$  has all-positive or all-negative eigenvalues.

- ► The solution is as smooth as the coefficients and boundary conditions allow.
- Well suited to describe static (for example, equilibrium states).
- Less suitable for dynamic processes.



#### Parabolic equations

#### The equation

$$L\mathbf{u} + F(\dots) = 0 \tag{2}$$

is considered parabolic if the coefficient matrix  $\{a_{ij}\}$  has one zero eigenvalue, while all others have the same sign (positive or negative)

- ▶ May be represented in form  $u_t = -L\mathbf{u}$ , where  $\mathbf{u}$  is an elliptic operator.
- Are well suited to describe the smoothly evolving processes such diffusion or heat transfers.
- ▶ The solution is generally smoother than the initial value.



#### Hyperbolic equations

#### The equation

$$L\mathbf{u} + F(\dots) = 0 \tag{3}$$

is considered parabolic if the coefficient matrix  $\{a_{ij}\}$  has eigenvalues, the signs of which are the same excepting one of them, which is non-zero and has contrary sign.

- May be represented with elliptic operator:  $L\mathbf{u} a^2 \frac{\partial^2 \mathbf{u}}{\partial t^2} = F(...)$ , where L is an elliptic operator.
- Hyperbolic equations solutions retain discontinuities of initial data.
- Well-suited to describe wave processes.



# Scientific Computing Lecture 3

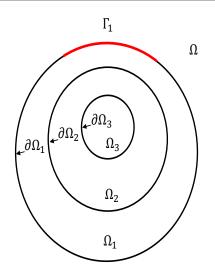
Part 5: Example 1. Cauchy problem for an elliptic PDE Nikolay Koshev

October 9, 2023



Skolkovo Institute of Science and Technology

#### The EEG Cauchy problem (Data mapping): the domain



#### Notations:

- $ightharpoonup \Omega \subset \mathbb{R}^3$  the whole domain.
- $\begin{array}{ll} & \Omega_k\,, k=1,...,N_\Omega-1 \text{ the outer} \\ \text{'sourceless' domains} \\ & (\sigma(\Omega_k)=\sigma_k=const). \end{array}$
- $\begin{array}{ll} & \Omega_{N_{\Omega}} \text{ the inner domain containing} \\ & \text{sources } \left(\sigma(\Omega_{N_{\Omega}}) = \sigma_{N_{\Omega}} = const\right) \end{array}$
- The accessible part of the boundary

$$\Gamma_k = \begin{cases} \Gamma_1, & k = 1 \\ \partial \Omega_k, & k \ge 2. \end{cases}$$

The rest of the boundary: Π<sub>k</sub>.

#### **Assumptions:**

- Current density volumetric distribution is non-zero only in the inner domain:  $\mathbf{J}: supp(\mathbf{J}) \in \Omega_{N\Omega}$
- $\begin{array}{c} \blacktriangleright \quad \text{Conductivities:} \\ \sigma_k \ll \sigma_{N_\Omega}, k=1,...,N_\Omega-1. \end{array}$



#### Data mapping: the Cauchy problem

Assume the function  $u(\mathbf{x}), \quad \mathbf{x} \in \Pi_1 \subset \partial \Omega_1$  to be known. Find the function  $U(\mathbf{x}), \quad \mathbf{x} \in \Omega_1$ , such that:

$$\Delta U = 0, \quad \mathbf{x} \in \Omega_1;$$

$$U(\mathbf{x}) = u(\mathbf{x}), \quad \mathbf{x} \in \Gamma_1;$$

$$\frac{\partial U}{\partial \mathbf{n}}(\mathbf{x}) = g(\mathbf{x}), \quad \mathbf{x} \in \Gamma_1,$$

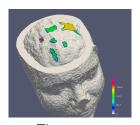
where

$$g(\mathbf{x}) = \begin{cases} 0, & \Gamma_1 \subset \partial \Omega, \\ \frac{\sigma_1}{\sigma_2} \mathbf{n} \cdot \nabla U(\mathbf{x}), & \Gamma_1 \subset \partial \Omega_2. \end{cases}$$

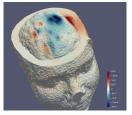
Whenever the function  $U(\mathbf{x})$  is found at every point  $\mathbf{x} \in \Omega_1$ , we obviously have the potential on the rest of a boundary  $\Pi_1 = \partial \Omega_1 n \Gamma_1$ .



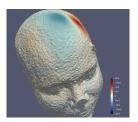
#### Numerical results: regular currents



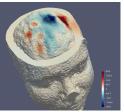
The sources



Real potential (brain)



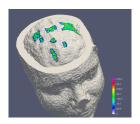
Potential on the scalp



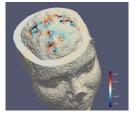
Reconstruction result



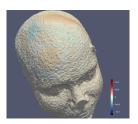
#### Numerical results: Random currents in each point



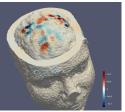
The sources



Real potential (brain)



Potential on the scalp



Reconstruction result



# Scientific Computing Lecture 3

Part 5: Example 2. Cauchy problem for the parabolic PDE Nikolay Koshev

October 9, 2023

# Skoltech

Skolkovo Institute of Science and Technology

#### Parabolic equations example: heat transfer equation

The quasilinear heat transfer equation (statement of 1D Forward Problem (simulation))

$$u_t = u_{xx} + aS(u) + F(x,t), \quad (x,t) \in [0,1]^2,$$
  
 $u(x,0) = f(x),$   
 $u(0,t) = g(t),$   
 $u(1,t) = p(t).$ 



#### The ill-posed Cauchy problem for heat transfer

The ill-posed Cauchy problem for heat transfer equation.

Suppose both initial condition f(x) and the left boundary condition g(t) to be unknown, and the function  $q(t) = \partial u/\partial x(1,t)$  to be known:

$$u_t = u_{xx} + aS(u) + F(x,t), \quad (x,t) \in [0,1]^2,$$
  
 $u(1,t) = p(t), \quad \frac{\partial u}{\partial x}(1,t) = q(t).$ 

Find u(x,t).

The problem is non-linear (quasi-linear), which leads that the cost functional has local minima



#### Coefficient problem for a heat-transfer equation

Let  $\Omega \subset \mathbb{R}^n$  be some domain,  $\partial \Omega$  - its boundary, and  $\Gamma \subset \partial \Omega$  is some part of this boundary. Consider the problem:

$$u_t = \Delta u + a(\mathbf{x})u, \quad \mathbf{x} \in \Omega$$
  
 $u(\mathbf{x})|_{\Gamma} = p(\mathbf{x}); \quad \partial_n u(\mathbf{x})|_{\Gamma} = q(x)$ 

Assuming  $u(\mathbf{x})>0, \mathbf{x}\in\Omega$  introduce the notation v=ln(u) (so that  $u=e^v$ ). Thus,

$$v_t u = u\Delta v + u(\nabla v)^2 + a(\mathbf{x})u, \quad v|_{\Gamma} = \ln(p).$$

Dividing the equation by u, we obtain:

$$v_t = \Delta v + (\nabla v)^2 + a(\mathbf{x}).$$

Now, differentiating the latter equation on t, we just exclude the coefficient a(x)!



## Thank you for your attention!



## **QUESTIONS**



#### Questions (Lec 3)

- Write down the forms of differential operator for Ordinary Differential Equation and for Partial Differential Equation. Explain the difference between them.
- Explain the meaning of terms: Linear differential equation, Quasi-linear differential equation, Non-linear differential equation. Provide examples of these equations.
- lacktriangle How the ODE of  $k^{th}$  order can be replaced with a system of ODE of the first order?
- List three kinds of problems for differential equations. Explain each of them in a few words.
- What is the Cauchy problem? Describe it with an example.
- What is the boundary value problem? Describe it with an example.
- What is the eigenvalue problem? Describe it with an example.
- ▶ When the Cauchy problem is well-posed?
- List three general approaches for solution of the Cauchy problem. Explain in few words advantages and disadvantages of them.
- What is explicit and implicit schemes for solving the Cauchy problem? Explain advantages and disadvantages for all of them.
- Explain the idea of Euler's method.
- Explain the idea of Shooting method.
- Write at least three real-world examples of problems leading to the partial differential equations of the first order. Analyze your examples in terms of linearity.
- Lan the Cauchy problem be stated for the partial differential equation? If yes, give an example of it.
- Describe the main idea of statement of N-body problem.
- Describe the main idea of the Molecular dynamics method.

