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## 1 Picard's Iteration Method

This method is used to find successive approximations to the unknown solution of a first-order initial value problem. Consider the IVP:

$$\frac{dy}{dx} = f(x, y) ; \quad y(x_0) = y_0 \quad (1)$$

This is equivalent to the following integral equation:

$$y(x) = y_0 + \int_{x_0}^x f(t, y(t)) dt \quad (2)$$

Integrating (1), we get

$$\begin{aligned} \int_{x_0}^x \frac{dy}{dx} dx &= \int_{x_0}^x f(t, y(t)) dt \\ \implies y(x) &= y_0 + \int_{x_0}^x f(t, y(t)) dt \end{aligned}$$

$\therefore$  (1)  $\implies$  (2). Now, differentiating (2), we get

$$\frac{dy}{dx} = 0 + f(x, y(x)) = f(x, y).$$

Also,

$$y(x_0) = y_0 + \int_{x_0}^{x_0} f(t, y(t)) dt = y_0,$$

$\therefore$  (2)  $\implies$  (1). Since  $y(t)$  is unknown, we can't use that in the integral, so we use  $y(t) = y_0$  in the integral to get a function  $y_1(x)$ , i.e.

$$y_1(x) = y_0 + \int_{x_0}^x f(t, y_0) dt.$$

Then we use  $y(t) = y_1(t)$  in the integral to get the next approximation:

$$y_2(x) = y_0 + \int_{x_0}^x f(t, y_1(t)) dt.$$

In general,

$$y_{n+1}(x) = y_0 + \int_{x_0}^x f(t, y_n(t)) dt.$$

$y_1(x), y_2(x), \dots, y_n(x), \dots$  are called the Picard iterations. Note that all  $y_n(x)$  satisfy the initial condition  $y_n(x_0) = y_0$ . But none of them may satisfy the ODE. However, under some conditions on  $f(x, y)$ ,  $y_n(x)$  converges to the unique solution  $y(x)$  to the IVP.

**Example 1.** Consider the IVP  $\frac{dy}{dx} = x + y$ ;  $y(0) = 0$ . Note that this can be solved easily to get  $y(x) = e^x - x - 1$ . Let's find the Picard iterations. Here,  $f(x, y) = x + y$  &  $x_0 = 0, y_0 = 0$ .

$$\begin{aligned} y_1(x) &= y_0 + \int_{x_0}^x f(t, y_0) dt \\ &= 0 + \int_0^x f(t, 0) dt \\ \therefore y_1(x) &= \int_0^x t dt = \frac{x^2}{2}. \end{aligned}$$

$$\begin{aligned} y_2(x) &= y_0 + \int_{x_0}^x f(t, y_1(t)) dt \\ &= 0 + \int_0^x (t + y_1(t)) dt \\ &= \int_0^x \left( t + \frac{t^2}{2} \right) dt = \frac{x^2}{2} + \frac{x^3}{6}. \end{aligned}$$

$$\begin{aligned} \text{Similarly, } y_3(x) &= 0 + \int_0^x f(t, y_2(t)) dt \\ &= \int_0^x \left( t + \frac{t^2}{2} + \frac{t^3}{6} \right) dt = \frac{x^2}{2} + \frac{x^3}{6} + \frac{x^4}{24}. \end{aligned}$$

$$\begin{aligned} \text{By induction, } y_n(x) &= \frac{x^2}{2!} + \frac{x^3}{3!} + \dots + \frac{x^{n+1}}{(n+1)!}, \quad n \geq 1, \\ &= e^x - x - 1. \end{aligned}$$

In practice, we can use  $y_n(x)$  for large enough  $n$  to approximate the solution  $y(x)$  to the given IVP when we are not able to find the actual solution.

## 2 Existence and Uniqueness of First-Order IVP

Some examples:

1. The IVP  $|y'| + |y| = 0$  ;  $y(0) = 1$  has no solutions because  $|y'| + |y| = 0 \implies |y| = 0 \implies y = 0 \implies y(0) = 0 \neq 1$ .
2. The IVP  $y' = 2x$  ;  $y(0) = 1$  has a unique solution  $y = x^2 + 1$ .
3. Consider the IVP  $xy' = y - 1$  ;  $y(0) = 1$ , then  $y = 1 + cx$  is a solution for every  $c \in \mathbb{R}$ .

Hence, an IVP may have no solutions, a unique solution, or more than one solution.

### **Theorem 1. (Sufficient conditions for existence of solutions)**

Consider the IVP:  $\frac{dy}{dx} = f(x, y)$  ;  $y(x_0) = y_0$ . Suppose the function  $f(x, y)$  is continuous on a rectangle  $R = \{(x, y) : |x - x_0| \leq a, |y - y_0| \leq b\}$ . Assume  $|f(x, y)| \leq k \forall (x, y) \in R$  for some constant  $k$ . Then the IVP has at least one solution  $y(x)$ . Also, the solution must be defined for  $x \in (x_0 - \alpha, x_0 + \alpha)$ , where  $\alpha = \min\{a, \frac{b}{k}\}$ .

**Remark 1.** The  $\alpha$  that we get in the theorem need not be the maximum possible for a given IVP.

**Example 2.** Consider the IVP  $\frac{dy}{dx} = 1 + y^2$  ;  $y(0) = 0$ . Let's find the maximum possible  $\alpha$  given by Theorem 1. Here,  $f(x, y) = 1 + y^2$  is continuous everywhere. So, we can take  $a$  and  $b$  to be as large as we want.

If  $R = \{(x, y) : |x| \leq a, |y| \leq b\}$ , then

$$|f(x, y)| = 1 + y^2 \leq 1 + b^2 \forall (x, y) \in R.$$

So,  $k = 1 + b^2$ , then  $\frac{b}{k} = \frac{b}{1+b^2} \leq \frac{1}{2}$ ,

$$\therefore \alpha = \min \left\{ a, \frac{b}{k} \right\} \leq \frac{b}{k} \leq \frac{1}{2}.$$

So, the maximum value  $\alpha$  that we get using the existence theorem is  $\alpha = \frac{1}{2}$ .

$\therefore$  Solution  $y(x)$  must be defined in the interval  $(-\frac{1}{2}, \frac{1}{2})$ .

But, by solving the IVP, we get  $y = \tan x$ . The solution  $y = \tan x$  is defined on  $(-\frac{\pi}{2}, \frac{\pi}{2})$  which is bigger than  $(-\alpha, \alpha)$ .

Next question is about the uniqueness of solution. We will see that under an additional condition we can guarantee the uniqueness and existence of solution to the IVP.

**Definition 1. (*Lipschitz condition*)** A function  $f(x, y)$  of two variables is said to satisfy the "Lipschitz condition" on a region  $R \subseteq \mathbb{R}^2$  if  $\exists$  a constant  $M$  such that

$$|f(x, y_1) - f(x, y_2)| \leq M|y_1 - y_2|$$

for all  $(x, y_1), (x, y_2) \in R$ .

**Remark 2.** If  $f(x, y)$  is continuous on  $R$  and  $\frac{\partial f}{\partial y}$  is also continuous on  $R$ , then  $f$  satisfies the Lipschitz condition.

*Proof.* Let  $(x, y_1), (x, y_2) \in R$ . Then by the mean value theorem,

$$f(x, y_1) - f(x, y_2) = \frac{\partial f}{\partial y}(x, y^*) (y_1 - y_2)$$

for some  $y^*$  between  $y_1$  &  $y_2$ .

$$\implies |f(x, y_1) - f(x, y_2)| = \left| \frac{\partial f}{\partial y}(x, y^*) \right| |y_1 - y_2|,$$

since  $\frac{\partial f}{\partial y}$  is continuous on  $R$ ,  $\exists M$  such that  $\left| \frac{\partial f}{\partial y} \right| \leq M$  on  $R$ .

$$\implies |f(x, y_1) - f(x, y_2)| \leq M|y_1 - y_2|.$$

$\therefore f$  satisfies the Lipschitz condition. □

The converse of the above remark is not true in general.

**Example 3.** Let  $f(x, y) = x + |\sin y|$ . Then  $f$  satisfies the Lipschitz condition on  $\mathbb{R}^2$ , because

$$\begin{aligned} |f(x, y_1) - f(x, y_2)| &= ||\sin y_1| - |\sin y_2|| \\ &\leq |\sin y_1 - \sin y_2| \\ &\leq |y_1 - y_2|. \end{aligned}$$

But,  $\frac{\partial f}{\partial y}$  does not exist on the  $x$ -axis.

**Theorem 2. (*Existence and uniqueness*)** Consider the IVP :

$$\frac{dy}{dx} = f(x, y); \quad y(x_0) = y_0.$$

Suppose  $f(x, y)$  is continuous on a rectangle  $R = \{(x, y) : |x - x_0| \leq a, |y - y_0| \leq b\}$  and  $f$  satisfies the Lipschitz condition on  $R$ . Then the IVP has a unique solution  $y(x)$  defined on some open interval containing  $x = x_0$ .



**Example 4.** (Non-uniqueness) Consider  $\frac{dy}{dx} = \sqrt{|y|}$  ;  $y(0) = 0$ . Clearly,  $y \equiv 0$  is a solution to the above IVP. Let

$$y = \begin{cases} \frac{x^2}{4}, & x \geq 0 \\ -\frac{x^2}{4}, & x < 0 \end{cases}$$

$$\implies \frac{dy}{dx} = \begin{cases} \frac{x}{2}, & x \geq 0 \\ -\frac{x}{2}, & x < 0 \end{cases}$$

$$\text{Also, } \sqrt{|y|} = \sqrt{\frac{x^2}{4}} = \frac{|x|}{2} = \begin{cases} \frac{x}{2}, & x \geq 0 \\ -\frac{x}{2}, & x < 0 \end{cases},$$

$$\therefore \frac{dy}{dx} = \sqrt{|y|}.$$

Also,  $y(0) = 0$ ,  $\therefore y(x)$  is also a solution to the IVP.

**Example 5.** Consider  $\frac{dy}{dx} = \sqrt{|y|}$  ;  $y(x_0) = y_0$ . If  $y_0 \neq 0$ , then  $f(x, y) = \sqrt{|y|}$  satisfies the Lipschitz condition on a small enough rectangle containing  $(x_0, y_0)$ . Hence, the IVP has a unique solution if  $y_0 \neq 0$ .

**Example 6.** Consider the IVP  $y \frac{dy}{dx} = x$  ;  $y(0) = \beta$ . Find the values of  $\beta$  for which the IVP has a unique solution, more than one solution, or no solution.

**Solution** Here,  $f(x, y) = \frac{x}{y}$  is continuous for  $(x, y)$  except on the line  $y = 0$ . Also,  $\frac{\partial f}{\partial y} = -\frac{x}{y^2}$  is continuous except when  $y = 0$ . If  $\beta \neq 0$ , then the existence-uniqueness theorem guarantees a unique solution to the IVP. Here, we can find the solution as follows:

$$y \frac{dy}{dx} = x$$

$$\implies \int y dy = \int x dx$$

$$\implies \frac{y^2}{2} = \frac{x^2}{2} + c$$

$$\because y(0) = \beta \implies c = \frac{\beta^2}{2}$$

$$\therefore y^2 = x^2 + \beta^2 \implies y = \pm \sqrt{x^2 + \beta^2}.$$

We have already seen that if  $\beta \neq 0$ , then the IVP has a unique solution. Since  $y = \sqrt{x^2 + \beta^2} \implies y(0) = |\beta|$ ,  $\therefore$  if  $\beta > 0$ , then  $y = \sqrt{x^2 + \beta^2}$  is a

solution but  $y = -\sqrt{x^2 + \beta^2}$  is not a solution to the IVP. Similarly, if  $\beta < 0$ , then  $y = -\sqrt{x^2 + \beta^2}$  is a solution to the IVP but  $y = \sqrt{x^2 + \beta^2}$  is not a solution to the IVP.

For  $\beta = 0$ , the existence-uniqueness theorem does not give us any information. However,  $y = x$  and  $y = -x$  are clearly solutions to the IVP:  $y \frac{dy}{dx} = x$  ;  $y(0) = 0$ . ■

**Exercise 1.** Consider the IVP:

$$(x^2 - 4x) \frac{dy}{dx} = (2x - 4)y ; y(x_0) = y_0.$$

Find  $(x_0, y_0)$  for which the IVP has no solution, a unique solution, or more than one solution.

### 3 Laplace Transforms

**Definition 2.** Let  $f : [0, \infty) \rightarrow \mathbb{R}$  be any function. We define the Laplace transform of  $f$  as

$$F(s) = (Lf)(s) = \int_0^{\infty} e^{-st} f(t) dt,$$

provided the improper integral converges.

**Example 7.** Let  $f(t) = 1 \quad \forall t$ , then

$$\begin{aligned} F(s) &= \int_0^{\infty} e^{-st} \cdot 1 dt = \lim_{b \rightarrow \infty} \int_0^b e^{-st} dt \\ &= \lim_{b \rightarrow \infty} \frac{1}{s} (1 - e^{-bs}) \\ &= \frac{1}{s} \quad \text{for } s > 0, \end{aligned}$$

$$\therefore L(1)(s) = \frac{1}{s} \quad \text{for } s > 0.$$

**Example 8.** Let  $f(t) = e^{at}$ , then

$$\begin{aligned} F(s) &= \int_0^{\infty} e^{-st} e^{at} dt \\ &= \int_0^{\infty} e^{-(s-a)t} dt \\ &= \frac{1}{s-a} \quad \text{for } s > a, \end{aligned}$$

$$\therefore L(e^{at})(s) = \frac{1}{s-a} \quad \text{for } s > a.$$

**Example 9.** Let  $f(t) = t^n$ ,  $n \in \mathbb{N}$ , then

$$\begin{aligned} F(s) &= \int_0^{\infty} e^{-st} t^n dt \\ &= t^n \left. \frac{e^{-st}}{-s} \right|_0^{\infty} - \int_0^{\infty} n t^{n-1} \frac{e^{-st}}{-s} dt \\ &= 0 + \frac{n}{s} \int_0^{\infty} t^{n-1} e^{-st} dt, \quad \text{if } s > 0, \\ &= \frac{n}{s} L(t^{n-1})(s) \end{aligned}$$

$$\therefore L(t^n)(s) = \frac{n}{s} L(t^{n-1})(s).$$

By induction,

$$L(t^n)(s) = \frac{n!}{s^{n+1}}, \quad s > 0.$$

### 3.1 Properties of Laplace Transforms

#### 1. Linearity:

$$L(af(t) + bg(t))(s) = aL(f(t))(s) + bL(g(t))(s).$$

*Proof.*

$$\begin{aligned}
 L(af(t) + bg(t))(s) &= \int_0^{\infty} e^{-st}(af(t) + bg(t))dt \\
 &= a \int_0^{\infty} e^{-st}f(t)dt + b \int_0^{\infty} e^{-st}g(t)dt \\
 &= aL(f)(s) + bL(g)(s).
 \end{aligned}$$

□

Using this, we can find the Laplace transform of any polynomial. For example,

$$\begin{aligned}
 L(t^3 - 2t^2 + t + 3)(s) &= L(t^3)(s) - 2L(t^2)(s) + L(t)(s) + 3L(1)(s) \\
 &= \frac{3!}{s^4} - 2\left(\frac{2!}{s^2}\right) + \frac{1!}{s^2} + 3\cdot\frac{1}{s}.
 \end{aligned}$$

### Laplace transforms of hyperbolic sine and hyperbolic cosine functions:

We know that

$$\begin{aligned}
 \sinh(at) &:= \frac{e^{at} - e^{-at}}{2} \\
 \therefore L(\sinh(at))(s) &= \frac{1}{2} [L(e^{at}) - L(e^{-at})] \\
 &= \frac{1}{2} \left[ \frac{1}{s-a} - \frac{1}{s+a} \right] \\
 \therefore L(\sinh(at))(s) &= \frac{a}{s^2 - a^2} \quad \text{if } s > |a|.
 \end{aligned}$$

Similarly,

$$\begin{aligned}
 \cosh(at) &:= \frac{e^{at} + e^{-at}}{2} \\
 \therefore L(\cosh(at))(s) &= \frac{1}{2} [L(e^{at}) + L(e^{-at})] \\
 &= \frac{1}{2} \left[ \frac{1}{s-a} + \frac{1}{s+a} \right] \\
 \therefore L(\cosh(at))(s) &= \frac{s}{s^2 - a^2} \quad \text{if } s > |a|.
 \end{aligned}$$

### Laplace transforms of sine and cosine functions:

Let  $f(t) = \cos(\omega t)$ , then

$$\begin{aligned} F(s) &= \int_0^{\infty} e^{-st} \cos(\omega t) dt \\ &= \frac{e^{-st}}{-s} \cos(\omega t) \Big|_0^{\infty} - \int_0^{\infty} \frac{e^{-st}}{-s} (-\omega) \sin(\omega t) dt \\ &= \frac{1}{s} - \frac{\omega}{s} \int_0^{\infty} e^{-st} \sin(\omega t) dt, \quad \text{if } s > 0, \\ &= \frac{1}{s} - \frac{\omega}{s} \left[ \frac{e^{-st}}{-s} \sin(\omega t) \Big|_0^{\infty} - \int_0^{\infty} \frac{e^{-st}}{-s} \omega \cos(\omega t) dt \right] \\ &= \frac{1}{s} - \frac{\omega^2}{s^2} F(s) \\ \implies \left( 1 + \frac{\omega^2}{s^2} \right) F(s) &= \frac{1}{s} \\ \therefore F(s) &= L(\cos(\omega t))(s) = \frac{s}{s^2 + \omega^2}. \end{aligned}$$

Similarly,

$$L(\sin(\omega t))(s) = \frac{\omega}{s^2 + \omega^2}.$$

### 3.2 Inverse Laplace Transforms

**Definition 3.** If  $F(s)$  is the Laplace transform of a function  $f(t)$ , then  $f(t)$  is said to be the inverse Laplace transform of  $F(s)$ , and is denoted by  $L^{-1}(F)(t)$ .

**Example 10.** From example 9,

$$\begin{aligned} L^{-1} \left( \frac{1}{s^n} \right) &= \frac{1}{(n-1)!} L^{-1} \left( \frac{(n-1)!}{s^n} \right) \\ &= \frac{t^{n-1}}{(n-1)!} \quad \text{for } n \in \mathbb{N}. \end{aligned}$$

**Example 11.**

$$L^{-1} \left( \frac{1}{s^2 + 4} \right) = \frac{1}{2} L^{-1} \left( \frac{2}{s^2 + 2^2} \right) = \frac{1}{2} \sin(2t).$$

### 3.3 Some More Properties of Laplace Transforms

#### s-Shifting Property of Laplace Transforms:

Let  $F(s) = L(f)(s) = \int_0^{\infty} e^{-st} f(t) dt$ , then

$$\begin{aligned} F(s-a) &= \int_0^{\infty} e^{-(s-a)t} f(t) dt \\ &= \int_0^{\infty} e^{-st} (e^{at} f(t)) dt \\ &= L(e^{at} f(t))(s) \end{aligned}$$

$$\implies L(e^{at} f(t))(s) = L(f)(s-a).$$

Also,

$$L^{-1}(F(s-a)) = e^{at} f(t),$$

$$\therefore L^{-1}(F(s-a))(t) = e^{at} L^{-1}(F(s))(t).$$

**Example 12.** Find  $L^{-1}\left(\frac{1}{s^2+2s+5}\right)$ .

**Solution**

$$\frac{1}{s^2+2s+5} = \frac{1}{(s+1)^2+2^2}$$

Let  $F(s) = \frac{1}{s^2+2s+5}$ , then from example 11,

$$\begin{aligned} F(s) &= \frac{1}{2} L(\sin(2t))(s) \\ \implies F(s+1) &= \frac{1}{2} L(e^{-t} \sin(2t))(s) \\ &= L\left(\frac{1}{2} e^{-t} \sin(2t)\right)(s) \end{aligned}$$

$$\therefore L^{-1}\left(\frac{1}{s^2+2s+5}\right) = L^{-1}(F(s+1))(t) = \frac{1}{2} e^{-t} \sin(2t).$$

■

### Derivatives of Laplace Transforms

Let  $F(s) = L(f)(s) = \int_0^{\infty} e^{-st} f(t) dt$ , then

$$\begin{aligned} F'(s) &= \frac{d}{ds} \left[ \int_0^{\infty} e^{-st} f(t) dt \right] \\ &= \int_0^{\infty} \frac{\partial}{\partial s} [e^{-st} f(t)] dt \\ &= \int_0^{\infty} -te^{-st} f(t) dt \\ &= \int_0^{\infty} e^{-st} [-tf(t)] dt \\ &= L(-tf(t))(s) \end{aligned}$$

$$\therefore F'(s) = L(-tf(t))(s) \quad \text{or} \quad L(tf(t))(s) = -F'(s).$$

$$\begin{aligned} \therefore L(t^2 f(t))(s) &= L(t(tf(t)))(s) \\ &= -\frac{d}{ds} [L(tf(t))(s)] \\ &= -\frac{d}{ds} (-F'(s)) \\ &= F''(s). \end{aligned}$$

In general,

$$L(t^n f(t))(s) = (-1)^n F^{(n)}(s) \quad \text{for } n = 1, 2, 3, \dots$$

**Example 13.** Find  $L(t \cos \omega t)$ .

**Solution** We know that  $L(\cos \omega t) = \frac{s}{s^2 + \omega^2} = F(s)$ , therefore,

$$\begin{aligned} L(t \cos \omega t) &= -F'(s) \\ &= -\frac{d}{ds} \left( \frac{s}{s^2 + \omega^2} \right) \\ &= \frac{s^2 - \omega^2}{(s^2 + \omega^2)^2}. \end{aligned}$$

■

## Laplace Transform of Derivatives

$$\begin{aligned}L(f')(s) &= \int_0^{\infty} e^{-st} f'(t) dt \\&= e^{-st} f(t) \Big|_0^{\infty} - \int_0^{\infty} -s e^{-st} f(t) dt \\&= 0 - f(0) + s \int_0^{\infty} e^{-st} f(t) dt \\&= -f(0) + sF(s) \\ \therefore L(f')(s) &= sL(f)(s) - f(0).\end{aligned}$$

$$\begin{aligned}\implies L(f'')(s) &= sL(f')(s) - f'(0) \\&= s[sL(f)(s) - f(0)] - f'(0) \\ \therefore L(f'')(s) &= s^2L(f)(s) - sf(0) - f'(0)\end{aligned}$$

In general,

$$L(f^{(n)})(s) = s^n L(f)(s) - s^{n-1} f(0) - s^{n-2} f'(0) - \dots - f^{(n-1)}(0) \quad \forall n \in \mathbb{N}.$$

## Laplace Transform of Integrals

Let

$$g(t) = \int_0^t f(\tau) d\tau,$$

then  $g'(t) = f(t)$  and  $g(0) = 0$ .

$$\begin{aligned}\therefore L(g')(s) &= L(f)(s) \\ \implies sL(g)(s) - g(0) &= F(s) \\ \implies L(g)(s) &= \frac{F(s)}{s} \\ \therefore L\left(\int_0^t f(\tau) d\tau\right) &= \frac{F(s)}{s}.\end{aligned}$$

**Example 14.** Find  $L^{-1}\left(\frac{1}{s(s^2+\omega^2)}\right)$ .



**Solution** Let  $F(s) = \frac{1}{s^2 + \omega^2} = \frac{1}{\omega} \left( \frac{\omega}{s^2 + \omega^2} \right) = \frac{1}{\omega} L(\sin \omega t)(s)$ ,

$$\implies f(t) = \frac{1}{\omega} \sin(\omega t)$$

$$\begin{aligned} \therefore L^{-1} \left( \frac{F(s)}{s} \right) &= \int_0^t f(\tau) d\tau \\ &= \int_0^t \frac{1}{\omega} \sin(\omega \tau) d\tau \end{aligned}$$

$$\therefore L^{-1} \left( \frac{1}{s(s^2 + \omega^2)} \right) = \frac{1}{\omega^2} (1 - \cos \omega t).$$

■

### 3.4 Application of Laplace Transforms to Solve IVPs

**Example 15.** Solve  $y'' - y = t$ ;  $y(0) = 1$ ,  $y'(0) = 1$ .

**Solution** Taking the Laplace transform,

$$\begin{aligned} L(y'') - L(y) &= L(t) \\ \implies [s^2 L(y) - sy(0) - y'(0)] - L(y) &= \frac{1}{s^2} \\ \implies (s^2 - 1)L(y)(s) - s - 1 &= \frac{1}{s^2} \\ \implies (s^2 - 1)L(y)(s) &= \frac{1}{s^2} + s + 1 \\ \implies L(y)(s) &= \frac{1}{s^2(s^2 - 1)} + \frac{s + 1}{s^2 - 1} \\ &= \frac{1}{s^2 - 1} - \frac{1}{s^2} + \frac{1}{s - 1} \end{aligned}$$

Taking the inverse Laplace transform  $L^{-1}$ ,

$$\begin{aligned} y(t) &= \sinh t - t + e^t \\ &= \frac{3}{2}e^t - \frac{1}{2}e^{-t} - t. \end{aligned}$$

■

### 3.5 Heaviside Function (Unit Step Function)

Define

$$u(t) = \begin{cases} 0, & t \leq 0 \\ 1, & t > 0 \end{cases}$$

So, for any  $a > 0$ ,

$$\begin{aligned}u_a(t) &= \begin{cases} 0, & t \leq a \\ 1, & t > a \end{cases} \\ &= u(t-a).\end{aligned}$$

$$\begin{aligned}\therefore L(u(t-a)) &= \int_0^{\infty} e^{-st} u(t-a) dt \\ &= \int_a^{\infty} e^{-st} dt \\ &= \frac{e^{-as}}{s}.\end{aligned}$$

Let

$$\tilde{f}(t) = f(t-a)u(t-a) = \begin{cases} 0 & \text{if } t \leq a \\ f(t-a) & \text{if } t > a \end{cases}$$

Then

$$\begin{aligned}L(\tilde{f})(s) &= \int_0^{\infty} e^{-st} \tilde{f}(t) dt \\ &= \int_a^{\infty} e^{-st} f(t-a) dt \\ &= e^{-as} \int_a^{\infty} e^{-s(t-a)} f(t-a) dt \\ &= e^{-as} \int_0^{\infty} e^{-s\tau} f(\tau) d\tau \\ &= e^{-as} L(f)(s)\end{aligned}$$

$$\therefore L(f(t-a)u(t-a)) = e^{-as} L(f).$$

**Hat Function:**

For  $0 \leq a < b$ ,

$$H(t) := \begin{cases} 1, & a < t < b \\ 0, & \text{otherwise} \end{cases}$$

i.e.

$$H(t) = u(t - a) - u(t - b).$$

**Example 16.** Let

$$f(t) = \begin{cases} 2 & , \quad 0 < t < 1 \\ \frac{t^2}{2} & , \quad 1 < t < \frac{\pi}{2} \\ \cos t & , \quad t > \frac{\pi}{2} \end{cases}.$$

Find  $L(f)(s)$ .

**Solution** First we express  $f(t)$  in terms of the Heaviside function.

$$\begin{aligned} f(t) &= \begin{cases} 2 & , \quad 0 < t < 1 \\ \frac{t^2}{2} & , \quad 1 < t < \frac{\pi}{2} \\ \cos t & , \quad t > \frac{\pi}{2} \end{cases} \\ &= 2[1 - u(t - 1)] + \frac{t^2}{2} \left[ u(t - 1) - u\left(t - \frac{\pi}{2}\right) \right] + \cos t \, u\left(t - \frac{\pi}{2}\right) \\ \therefore L(f)(s) &= 2(L(1) - L(u(t - 1))) + \frac{1}{2}L(t^2 u(t - 1)) - \frac{1}{2}L\left(t^2 u\left(t - \frac{\pi}{2}\right)\right) \\ &\quad + L\left(\cos t \, u\left(t - \frac{\pi}{2}\right)\right) \end{aligned} \quad (1)$$

Now,

$$\begin{aligned} t^2 u(t - 1) &= [(t - 1) + 1]^2 u(t - 1) \\ &= (t - 1)^2 u(t - 1) + 2(t - 1)u(t - 1) + u(t - 1) \\ \therefore L(t^2 u(t - 1)) &= L((t - 1)^2 u(t - 1)) + 2L((t - 1)u(t - 1)) + L(u(t - 1)) \\ &= e^{-s}L(t^2) + 2e^{-s}L(t) + \frac{e^{-s}}{s} \\ &= e^{-s}\frac{2}{s^3} + 2e^{-s}\frac{1}{s^2} + \frac{e^{-s}}{s} \end{aligned}$$

Similarly,

$$\begin{aligned} t^2 u\left(t - \frac{\pi}{2}\right) &= \left[\left(t - \frac{\pi}{2}\right) + \frac{\pi}{2}\right]^2 u\left(t - \frac{\pi}{2}\right) \\ &= \left(t - \frac{\pi}{2}\right)^2 u\left(t - \frac{\pi}{2}\right) + \pi\left(t - \frac{\pi}{2}\right) u\left(t - \frac{\pi}{2}\right) + \frac{\pi^2}{4} u\left(t - \frac{\pi}{2}\right) \\ \therefore L\left(t^2 u\left(t - \frac{\pi}{2}\right)\right) &= e^{-\frac{\pi}{2}s}\frac{2}{s^3} + \pi e^{-\frac{\pi}{2}s}\frac{1}{s^2} + \frac{\pi^2}{4} e^{-\frac{\pi}{2}s}\frac{1}{s} \end{aligned}$$

Also,

$$\begin{aligned}
 \cos t \, u\left(t - \frac{\pi}{2}\right) &= \sin\left(\frac{\pi}{2} - t\right) u\left(t - \frac{\pi}{2}\right) \\
 &= -\sin\left(t - \frac{\pi}{2}\right) u\left(t - \frac{\pi}{2}\right) \\
 \therefore L\left(\cos t \, u\left(t - \frac{\pi}{2}\right)\right) &= -L\left(\sin\left(t - \frac{\pi}{2}\right) u\left(t - \frac{\pi}{2}\right)\right) \\
 &= -e^{-\frac{\pi}{2}s} L(\sin t)(s) \\
 &= -e^{-\frac{\pi}{2}s} \frac{1}{s^2 + 1}.
 \end{aligned}$$

■

### 3.6 Dirac Delta Function

For  $k \in \mathbb{N}$ , let

$$f_k(t - a) = \begin{cases} k & , \quad a \leq t \leq a + \frac{1}{k} \\ 0 & , \quad \text{otherwise} \end{cases}$$

Note that

$$\int_0^{\infty} f_k(t - a) dt = 1 \quad \forall k.$$

Define the Dirac-delta function as

$$\begin{aligned}
 \delta(t - a) &= \lim_{k \rightarrow \infty} f_k(t - a) \\
 &= \begin{cases} \infty & , \quad t = a \\ 0 & , \quad t \neq a \end{cases}
 \end{aligned}$$

**Properties:**

$$1. \int_0^{\infty} \delta(t - a) dt = 1.$$

2. Note that

$$\begin{aligned}
 f_k(t-a) &= \begin{cases} k & , \quad a \leq t \leq a + \frac{1}{k} \\ 0 & , \quad \text{otherwise} \end{cases} \\
 &= k \left[ u(t-a) - u\left(t - \left(a + \frac{1}{k}\right)\right) \right] \\
 \therefore L(f_k(t-a)) &= k \left[ L(u(t-a)) - L\left(u\left(t - \left(a + \frac{1}{k}\right)\right)\right) \right] \\
 &= k \left[ \frac{e^{-as}}{s} - \frac{e^{-(a+\frac{1}{k})s}}{s} \right] \\
 &= e^{-as} \left[ \frac{1 - e^{-\frac{s}{k}}}{\frac{s}{k}} \right] \\
 \therefore L(\delta(t-a)) &= \lim_{k \rightarrow \infty} L(f_k(t-a)) \\
 &= \lim_{k \rightarrow \infty} e^{-as} \left[ \frac{1 - e^{-\frac{s}{k}}}{\frac{s}{k}} \right] \\
 &= e^{-as} \\
 \therefore L(\delta(t-a)) &= e^{-as}.
 \end{aligned}$$

**Example 17.** Solve the IVP:

$$y'' + 3y' + 2y = \delta(t-1), \quad y(0) = 0, \quad y'(0) = 0.$$

**Solution** Taking the Laplace transform, we get

$$L(y'') + 3L(y') + 2L(y) = L(\delta(t-1))$$

$$\implies [s^2Y(s) - sy(0) - y'(0)] + 3[sY(s) - y(0)] + 2Y(s) = e^{-s}$$

$$\implies (s^2 + 3s + 2)Y(s) = e^{-s}$$

$$\begin{aligned}
 \implies Y(s) &= e^{-s} \cdot \frac{1}{(s+1)(s+2)} \\
 &= \frac{e^{-s}}{s+1} - \frac{e^{-s}}{s+2} \\
 \implies y(t) &= L^{-1} \left[ \frac{e^{-s}}{s+1} \right] - L^{-1} \left[ \frac{e^{-s}}{s+2} \right] \\
 &= L^{-1}[e^{-s}L(e^{-t})(s)] - L^{-1}[e^{-s}L(e^{-2t})(s)] \\
 &= e^{-(t-1)}u(t-1) - e^{-2(t-1)}u(t-1)
 \end{aligned}$$

$$\therefore y(t) = \begin{cases} 0 & , \quad t < 1 \\ e^{-(t-1)} - e^{-2(t-1)} & , \quad t > 1 \end{cases}.$$

■

### 3.7 Convolution

**Question:** Is  $L(fg) = L(f)L(g)$ ?

**Example 18.** Let  $f = 1$ ;  $g = 1$ , then  $L(f) = \frac{1}{s} = L(g)$ . Also,  $L(fg) = L(1) = \frac{1}{s}$ ,

$$\therefore L(fg) \neq L(f)L(g).$$

**Example 19.** Let  $f = e^t$ ;  $g = 1$ , then  $L(f) = \frac{1}{s-1}$ ;  $L(g) = \frac{1}{s}$ , but  $L(fg) = L(e^t) = \frac{1}{s-1} \neq L(f)L(g)$ .

If  $L^{-1}(F(s)) = f(t)$  &  $L^{-1}(G(s)) = g(t)$ , then we have seen by the above examples that  $L^{-1}(F(s)G(s)) \neq f(t)g(t)$ . Let

$$F(s) = \int_0^{\infty} e^{-s\tau} f(\tau) d\tau$$

$$G(s) = \int_0^{\infty} e^{-su} g(u) du$$

Put  $u = t - \tau$ , then  $du = dt$ .

$$\implies G(s) = \int_{\tau}^{\infty} e^{-s(t-\tau)} g(t-\tau) dt.$$

Multiplying, we get

$$\begin{aligned} F(s)G(s) &= \int_0^{\infty} e^{-s\tau} f(\tau) d\tau \int_{\tau}^{\infty} e^{-s(t-\tau)} g(t-\tau) dt \\ &= \int_0^{\infty} \left( e^{-s\tau} f(\tau) e^{s\tau} \int_{\tau}^{\infty} e^{-st} g(t-\tau) dt \right) d\tau \\ &= \int_0^{\infty} f(\tau) \int_{\tau}^{\infty} e^{-st} g(t-\tau) dt d\tau \end{aligned} \quad (*)$$

This is a double integral over the region  $R$  given below :

$$\begin{aligned} R &= \{(t, \tau) : \tau \leq t < \infty, 0 < \tau < \infty\} \\ &= \{(t, \tau) : 0 \leq \tau \leq t, 0 \leq t < \infty\} \end{aligned}$$

Changing the order of integrals in (\*), we get

$$\begin{aligned} F(s)G(s) &= \int_0^\infty e^{-st} \left( \int_0^t f(\tau)g(t-\tau)d\tau \right) dt \\ &= L((f * g)(t))(s), \end{aligned}$$

where

$$(f * g)(t) := \int_0^t f(\tau)g(t-\tau)d\tau.$$

**Definition 4.**

$$(f * g)(t) := \int_0^t f(\tau)g(t-\tau)d\tau$$

is called the convolution of  $f$  and  $g$ .

**Exercise 2.** Use the above formula for  $f * g$  and show that  $L(f * g) = L(f)L(g)$ . Hence,

$$L^{-1}(F(s)G(s)) = L^{-1}(F(s)) * L^{-1}(G(s)).$$

**Example 20.** Find  $L^{-1}\left(\frac{1}{(s^2+\omega^2)^2}\right)$ .

**Solution** We know that  $L^{-1}\left(\frac{1}{s^2+\omega^2}\right) = \frac{1}{\omega} \sin(\omega t)$ ,

$$\begin{aligned} \therefore L^{-1}\left(\frac{1}{(s^2+\omega^2)^2}\right) &= \frac{1}{\omega} \sin(\omega t) * \frac{1}{\omega} \sin(\omega t) \\ &= \frac{1}{\omega^2} \int_0^t \sin(\omega \tau) \sin \omega(t-\tau) d\tau \\ &= \frac{1}{2\omega^2} \int_0^t [\cos(2\omega \tau - \omega t) - \cos(\omega t)] d\tau \\ \implies L^{-1}\left(\frac{1}{(s^2+\omega^2)^2}\right) &= \frac{1}{2\omega^3} [\sin(\omega t) - \omega t \cos(\omega t)]. \end{aligned}$$

■

### Properties of Convolution:

1.  $f * g = g * f$ .

*Proof.*  $(f * g)(t) = \int_0^t f(\tau)g(t - \tau)d\tau$ .

Put  $t - \tau = u \implies \tau = t - u \implies d\tau = -du$ , and when  $\tau = 0$ ,  $u = t$ ;  
when  $\tau = t$ ,  $u = 0$ ,

$$\begin{aligned}\therefore (f * g)(t) &= \int_t^0 f(t - u)g(u)(-du) \\ &= \int_0^t g(u)f(t - u)du \\ &= (g * f)(t).\end{aligned}$$

□

2.  $f * (g + h) = f * g + f * h$ .

3.  $(f * g) * h = f * (g * h)$ .

4.  $f * 0 = 0$ , where 0 denotes the zero function.

5.  $f * 1 \neq f$  (for example let  $f(t) = t$ ).

6.  $(f * f)(t)$  need not be a non-negative function. Let  $f(t) = \sin t$ , then

$$(f * f)(t) = \frac{1}{2}[\sin t - t \cos t].$$

**Example 21.** Solve  $y'' + 3y' + 2y = r(t)$ ,  $y(0) = 0, y'(0) = 0$ , where

$$r(t) = \begin{cases} 1 & , \quad 1 < t < 2 \\ 0 & , \quad \text{otherwise} \end{cases}.$$

**Solution Method 1:** Observe that  $r(t) = u(t - 1) - u(t - 2)$ . Take the Laplace transform and find  $Y(s)$ . Then take the inverse Laplace transform



to get  $y(t)$ .

**Method 2:** (Using convolution) Taking the Laplace transform, we get

$$\begin{aligned} L(y'') + 3L(y') + 2L(y) &= L(r) \\ \implies [s^2Y(s) - sy(0) - y'(0)] + 3[sY(s) - y(0)] + 2Y(s) &= R(s) \\ \implies (s^2 + 3s + 2)Y(s) &= R(s) \\ \implies Y(s) &= \frac{1}{(s+1)(s+2)} \cdot R(s) \end{aligned}$$

Now,

$$\frac{1}{(s+1)(s+2)} = \frac{1}{s+1} - \frac{1}{s+2} = L(e^{-t} - e^{-2t}).$$

Taking the inverse Laplace transform, we get

$$\begin{aligned} y(t) &= (r * q)(t); \quad \text{where } q(t) = e^{-t} - e^{-2t}, \\ &= \int_0^t r(\tau)q(t-\tau)d\tau \end{aligned}$$

For  $0 < t < 1$ ,  $r(t) = 0$ ,

$$\therefore y(t) = \int_0^t 0 \cdot q(t-\tau)d\tau = 0.$$

For  $1 < t < 2$ ,

$$\begin{aligned} y(t) &= \int_1^t 1 \cdot (e^{-(t-\tau)} - e^{-2(t-\tau)}) d\tau \\ &= \frac{1}{2} - e^{-(t-1)} + \frac{1}{2}e^{-2(t-1)}. \end{aligned}$$

For  $t > 2$ ,

$$\begin{aligned} y(t) &= \int_1^2 1 \cdot (e^{-(t-\tau)} - e^{-2(t-\tau)}) d\tau \\ &= e^{-(t-2)} - \frac{e^{-2(t-2)}}{2} - e^{-(t-1)} + \frac{e^{-2(t-1)}}{2}. \end{aligned}$$

■